



Threshold resummation for Drell-Yan production: theory and phenomenology

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HP².3rd

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In collaboration with:

Stefano Forte, Giovanni Ridolfi

Plan of the talk:

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- for which values of $\tau = \frac{Q^2}{s}$ is resummation important?

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- two prescriptions for resummation
 - the minimal prescription
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 - subleading terms
- new phenomenological results
 - rapidity distributions at NNLO + NNLL

For which τ is resummation important?

$z \sim 1$: logarithmic enhancement \rightarrow resummation of $\frac{\log^k(1-z)}{1-z}$

$$\sigma(\tau) = \int_{\tau}^1 \frac{dz}{z} \mathcal{L}\left(\frac{\tau}{z}\right) \hat{\sigma}(z), \quad \tau = \frac{Q^2}{s}, \quad z = \frac{Q^2}{\hat{s}}$$

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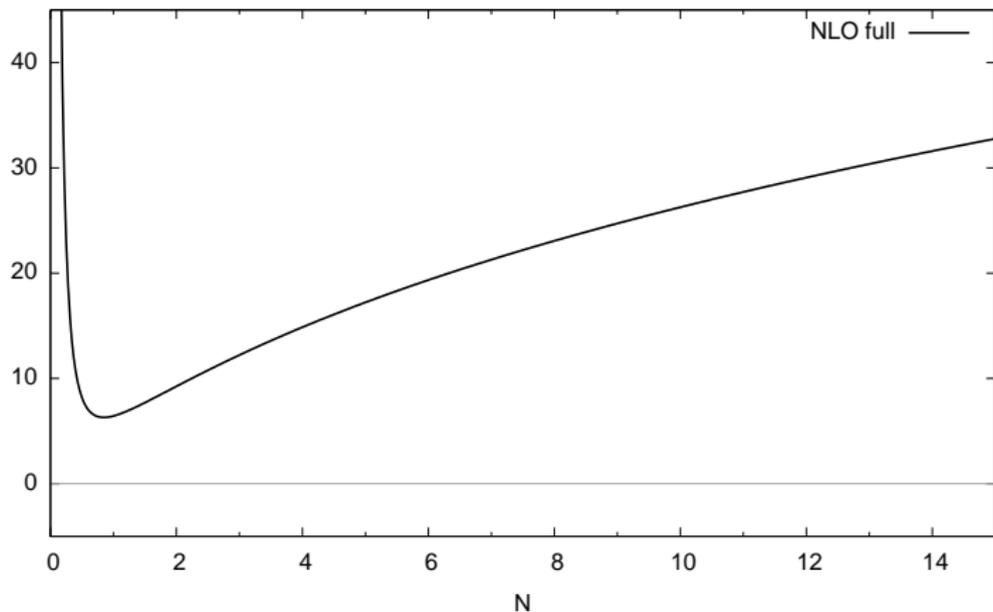
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N -space analysis
and saddle point argument

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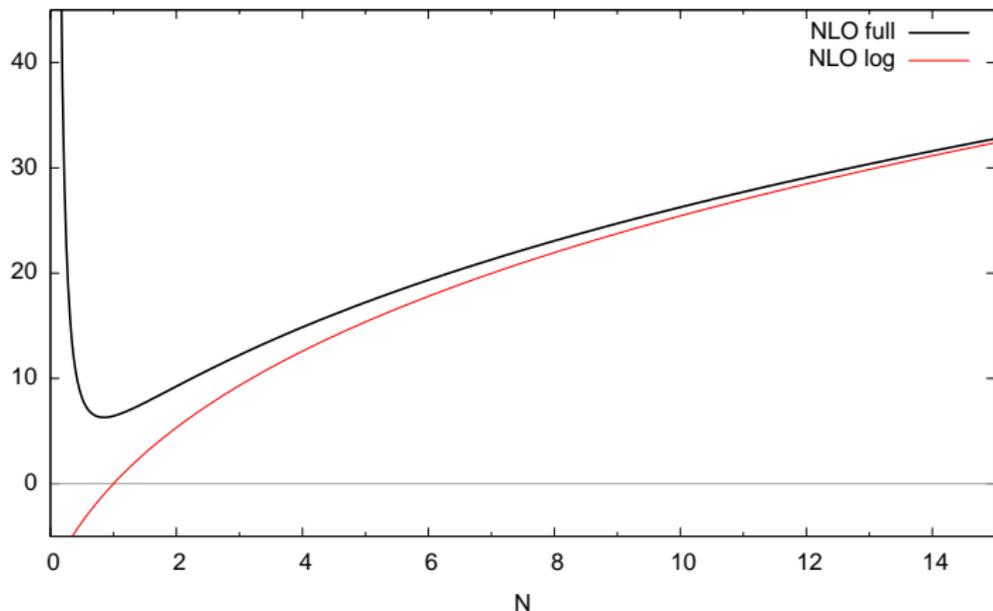
Drell-Yan $q\bar{q}$ at NLO in N -space

$$\frac{\alpha_s}{\pi} 4C_F \left\{ \left[\frac{\log(1-z)}{1-z} \right]_+ - \frac{\log \sqrt{z}}{1-z} - \frac{1+z}{2} \log \frac{1-z}{\sqrt{z}} + \left(\frac{\pi^2}{12} - 1 \right) \delta(1-z) \right\}$$



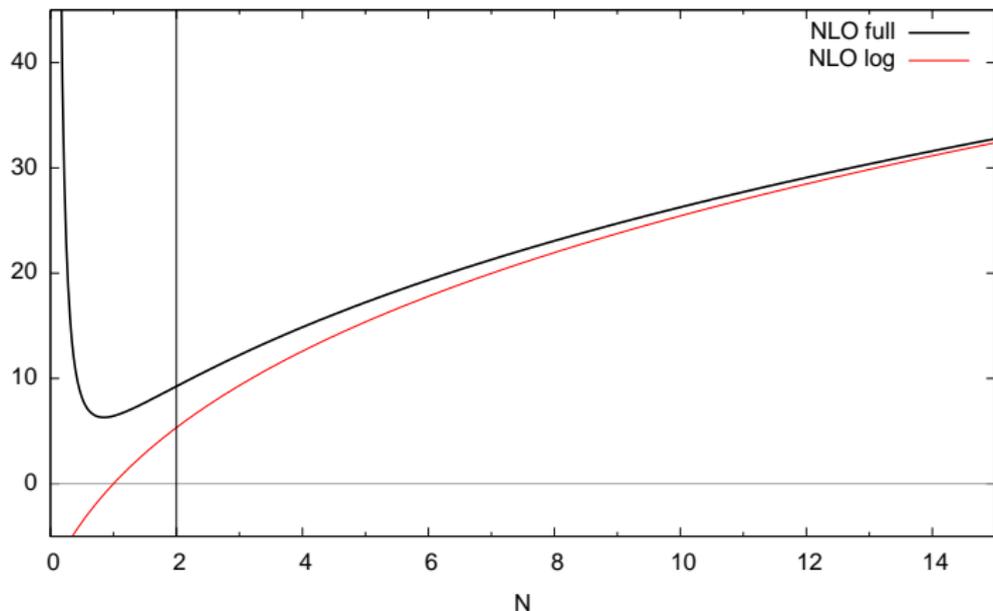
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For $N \gtrsim 2$ more than 50% of the NLO is given by the log term

Saddle point argument

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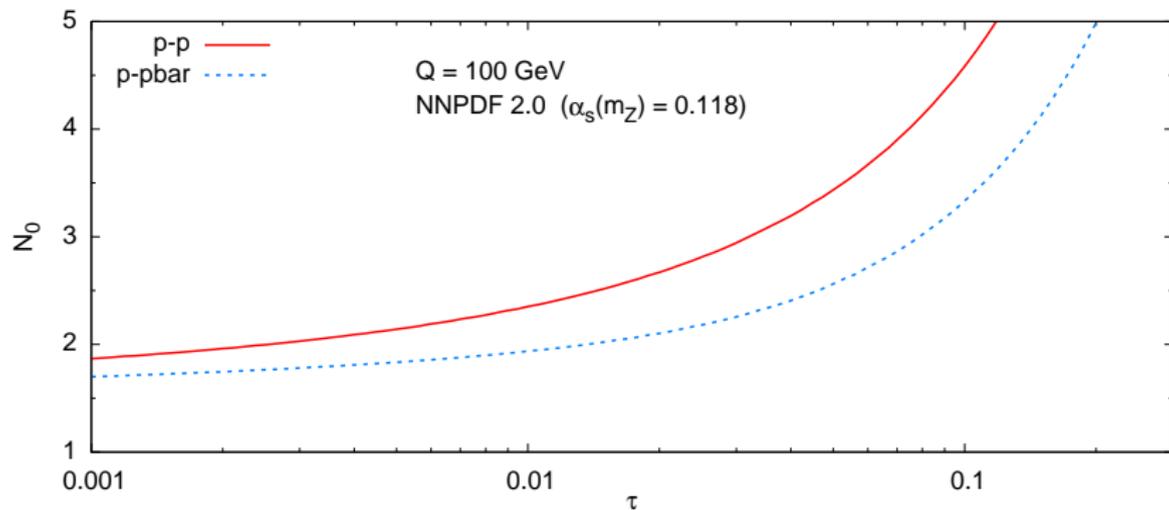
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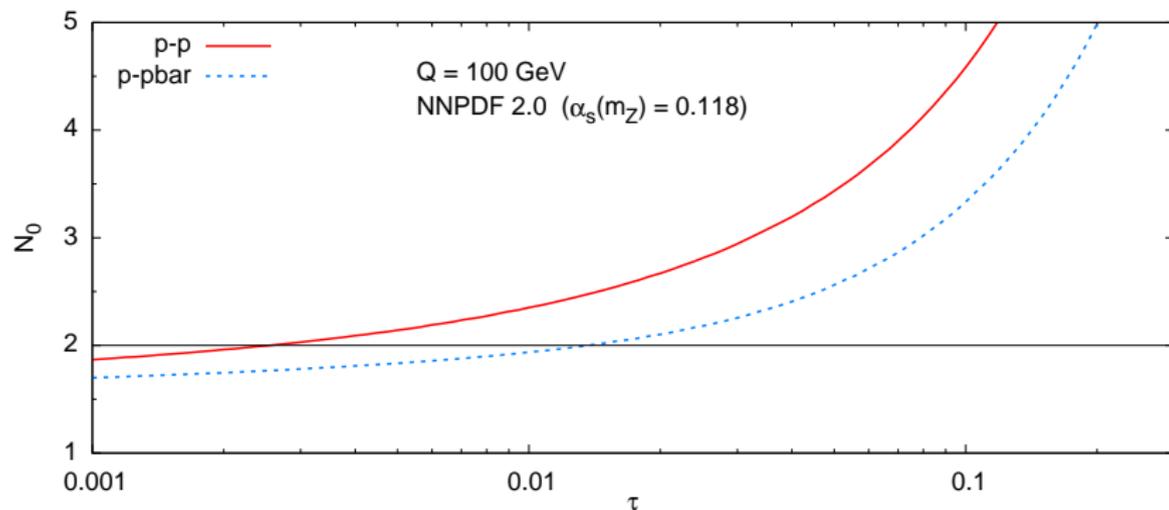
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How small?

Saddle point N_0 vs τ

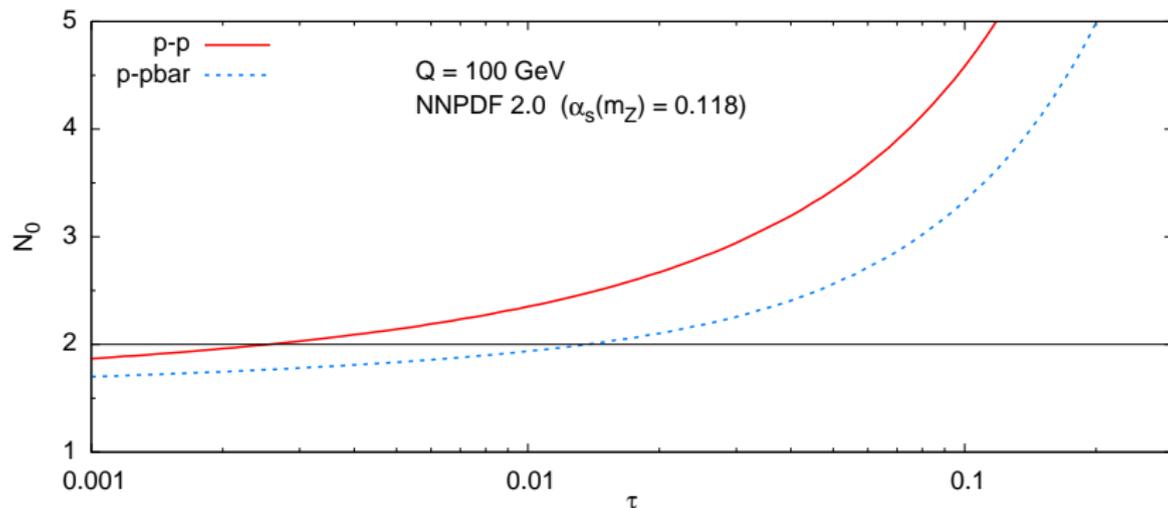


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Much smaller than expected!

Resummation

Resummation is performed in N -space $(L = 2\beta_0\alpha_s \log \frac{1}{N})$

$$\hat{\sigma}^{\text{res}}(N) = g_0(\alpha_s) \exp \left[\frac{1}{\alpha_s} g_1(L) + g_2(L) + \alpha_s g_3(L) + \alpha_s^2 g_4(L) + \dots \right]$$

known up to g_4 ($N^3\text{LL}$): S.Moch, J.A.M.Vermaseren, A.Vogt (hep-ph/0506288)

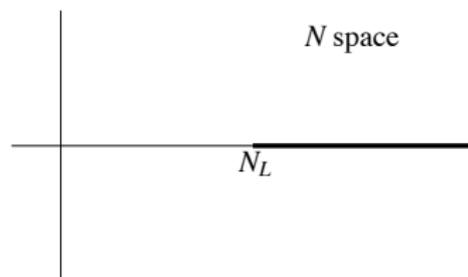
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Branch cut due to the Landau singularity for $N > N_L = \exp \frac{1}{2\beta_0\alpha_s}$



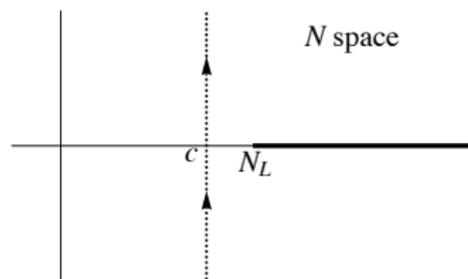
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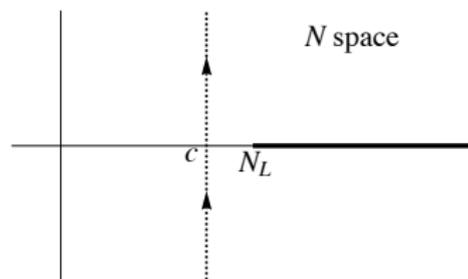
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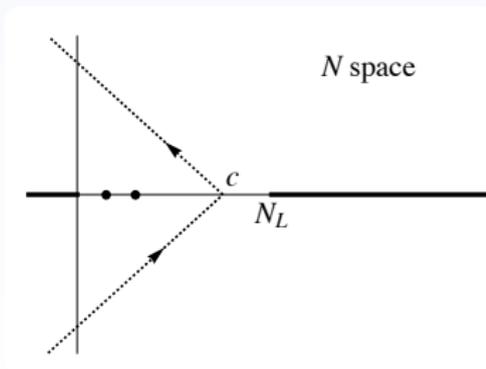
The Mellin inverse does not exist

Minimal prescription

S.Catani, M.L.Mangano, P.Nason, L.Trentadue (hep-ph/9604351)

$$\sigma_{\text{MP}}(\tau) = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} dN \tau^{-N} \mathcal{L}(N) \hat{\sigma}^{\text{res}}(N)$$

with $c < N_L = \exp \frac{1}{2\beta_0\alpha_s}$, as in the figure.



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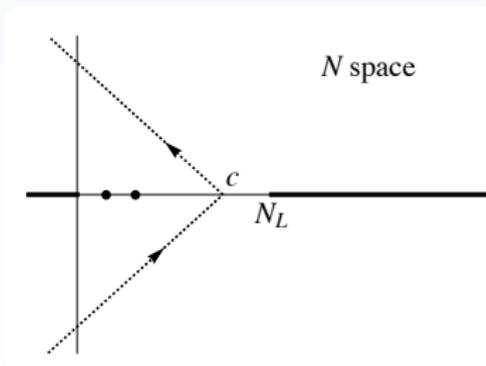
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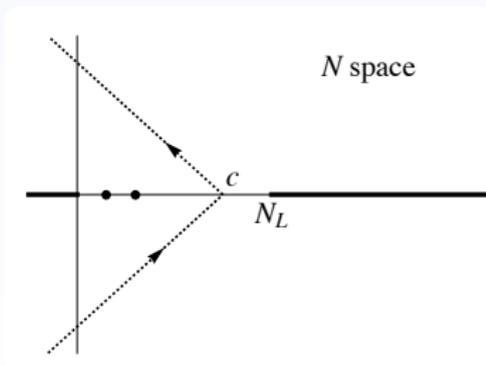
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- exact for invertible functions



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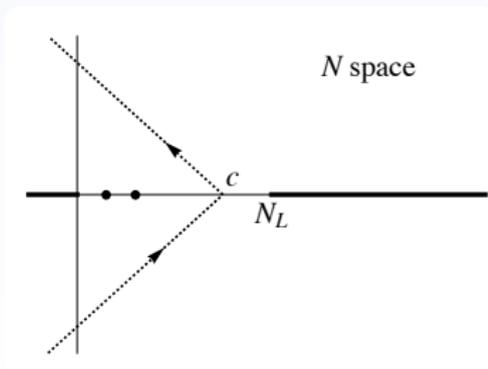
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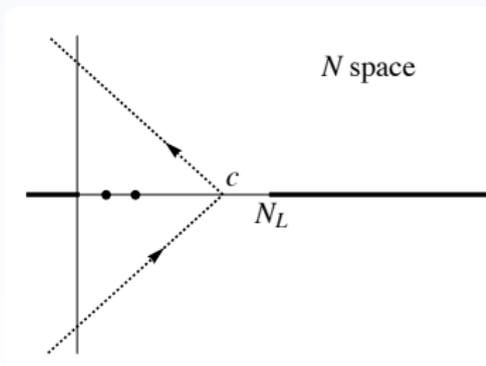
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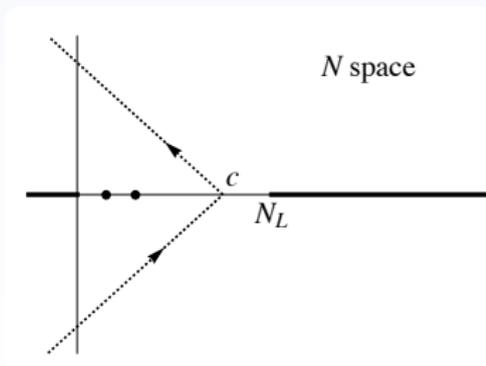
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But...

- a non-physical region of the parton cross-section contributes
- difficult numerical implementation



Minimal prescription: non-physical contribution

$$\sigma_{\text{MP}}(\tau) = \int_{\tau}^{+\infty} \frac{dz}{z} \mathcal{L}\left(\frac{\tau}{z}\right) \hat{\sigma}_{\text{MP}}(z)$$

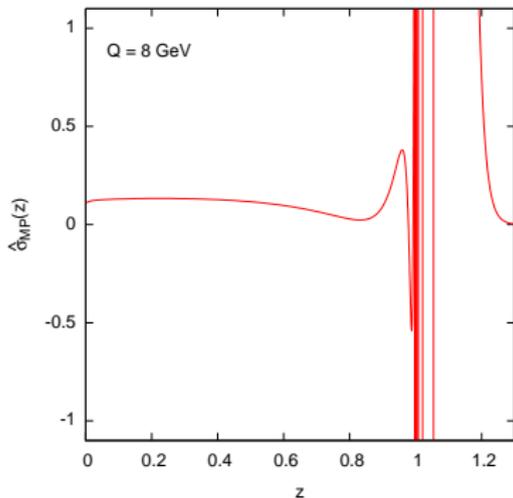
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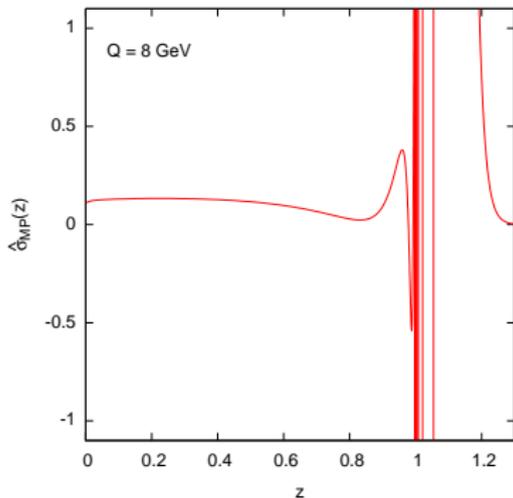
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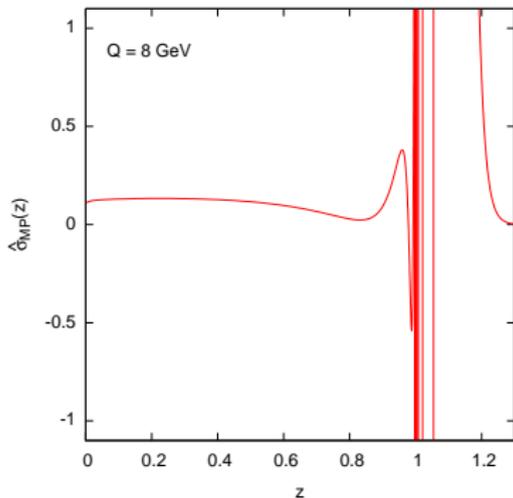
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Need for $\mathcal{L}(N)$, for values of N where the Mellin transform of $\mathcal{L}(x)$ does not converge



$$\hat{\sigma}^{\text{res}}(N)$$

Borel prescription (1)

$$\hat{\sigma}^{\text{res}}(N) = \sum_{k=1}^{\infty} h_k(\bar{\alpha}) \bar{\alpha}^k \quad \log^k \frac{1}{N} \quad , \quad \bar{\alpha} = 2\beta_0\alpha_s$$

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Treat the divergent series $\mathcal{M}^{-1}(\hat{\sigma}^{\text{res}}(N))$ with Borel method:*

* S.Forte, G.Ridolfi, J.Rojo, M.Ubiali (hep-ph/0601048); R.Abbate, SF, GR (hep-ph/0707.2452); MB, SF, GR (hep-ph/0807.3830); MB, SF, GR (coming soon)

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$$\mathcal{M}^{-1}[\hat{\sigma}^{\text{res}}(N)] = \sum_{k=1}^{\infty} h_k(\bar{\alpha}) \bar{\alpha}^k \mathcal{M}^{-1}\left[\log^k \frac{1}{N}\right], \quad \bar{\alpha} = 2\beta_0\alpha_s$$

Treat the divergent series $\mathcal{M}^{-1}(\hat{\sigma}^{\text{res}}(N))$ with Borel method:*

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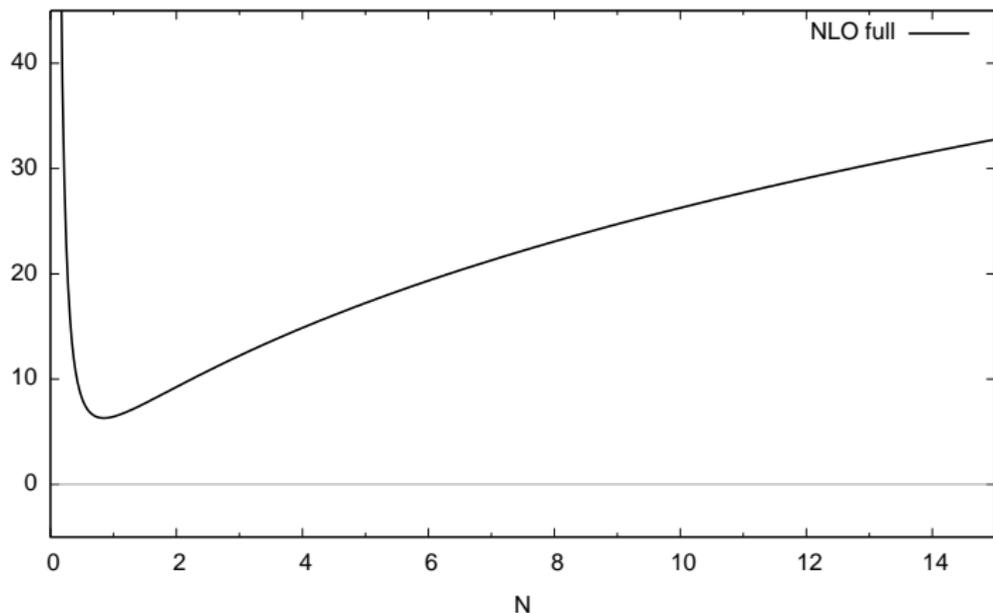
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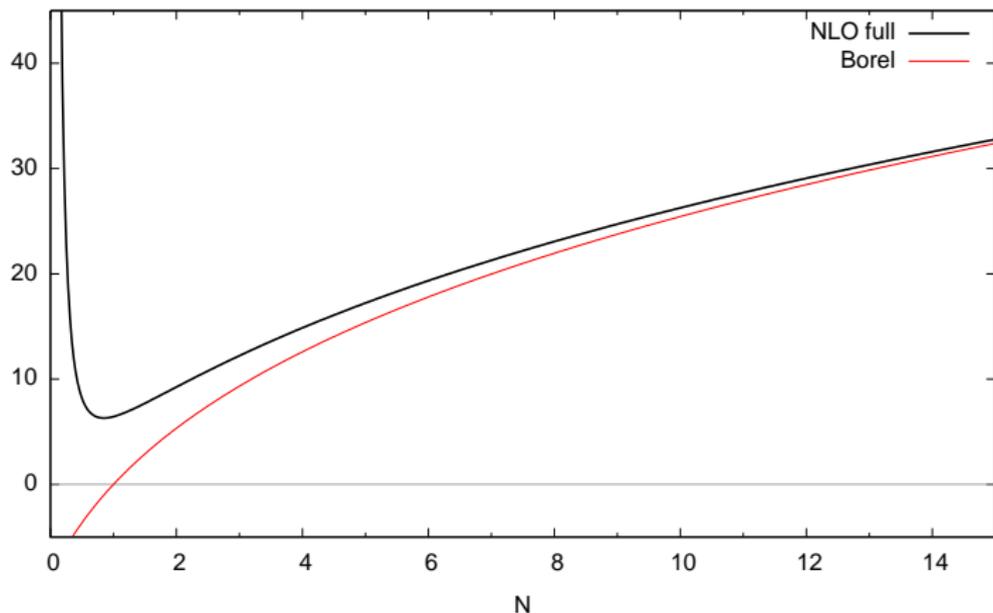
Comparison with fixed order: Drell-Yan $q\bar{q}$ at NLO

$$\frac{\alpha_s}{\pi} 4C_F \left\{ \left[\frac{\log(1-z)}{1-z} \right]_+ - \frac{\log \sqrt{z}}{1-z} - \frac{1+z}{2} \log \frac{1-z}{\sqrt{z}} + \left(\frac{\pi^2}{12} - 1 \right) \delta(1-z) \right\}$$



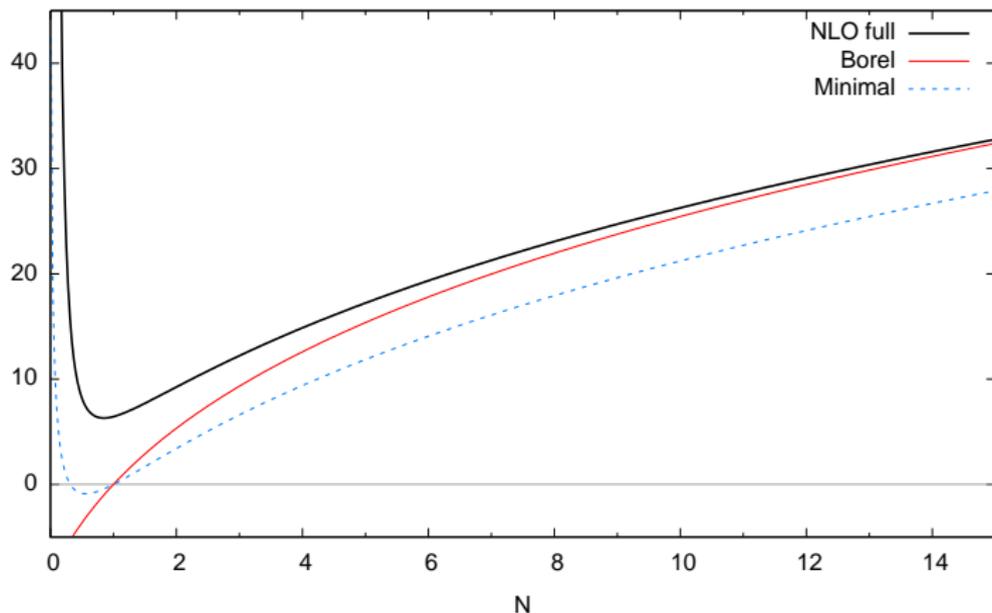
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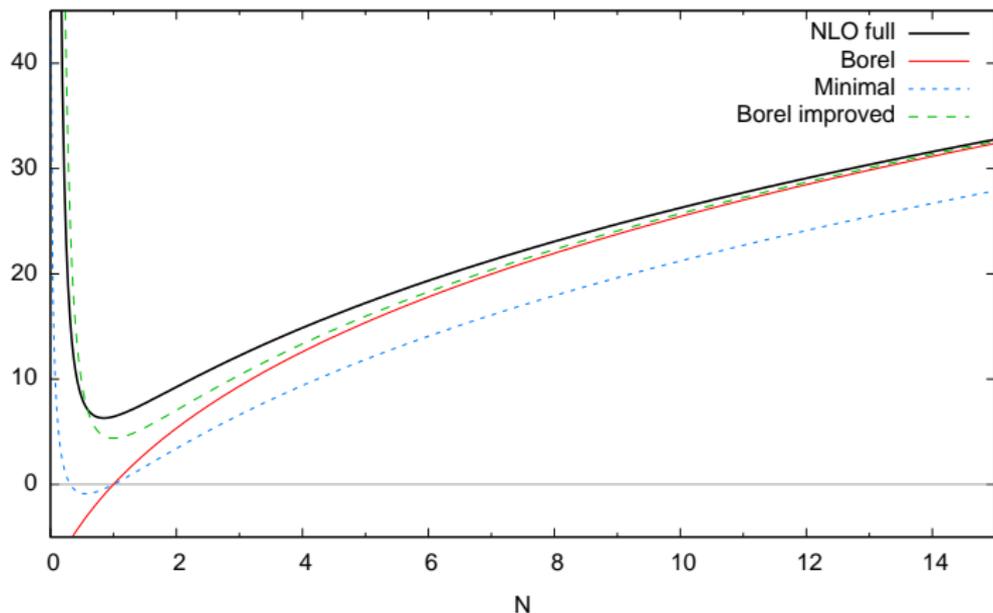
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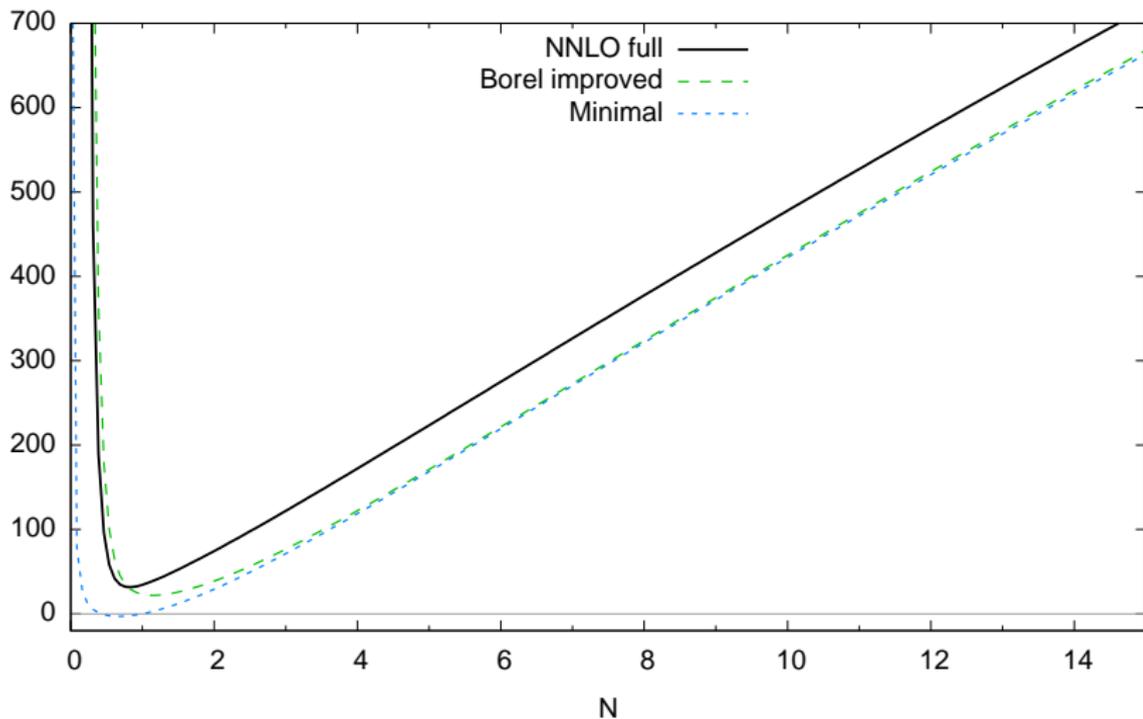
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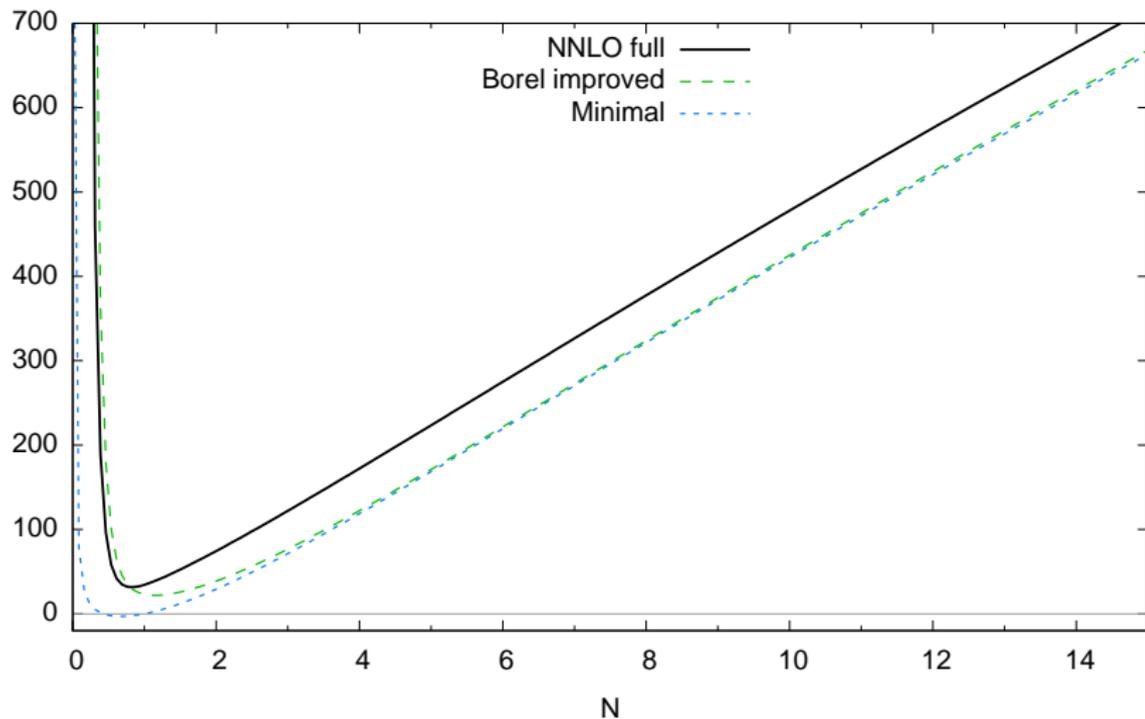
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Discrepancy due to terms like $\log^k(1-z) \rightarrow \frac{\log^k N}{N} \Rightarrow$ subleading

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- the difference in the included subleading terms is useful to estimate the importance of these terms

Impact in phenomenology: rapidity distributions (1)

$$\frac{1}{\tau} \frac{d\sigma}{dQ^2 dY} = \int_{\sqrt{\tau}e^Y}^1 \frac{dx_1}{x_1} \int_{\sqrt{\tau}e^{-Y}}^1 \frac{dx_2}{x_2} f_1(x_1) f_2(x_2) C\left(\frac{\tau}{x_1 x_2}, Y - \frac{1}{2} \log \frac{x_1}{x_2}\right)$$

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or, back to y space,

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After changing variables we get the compact expression

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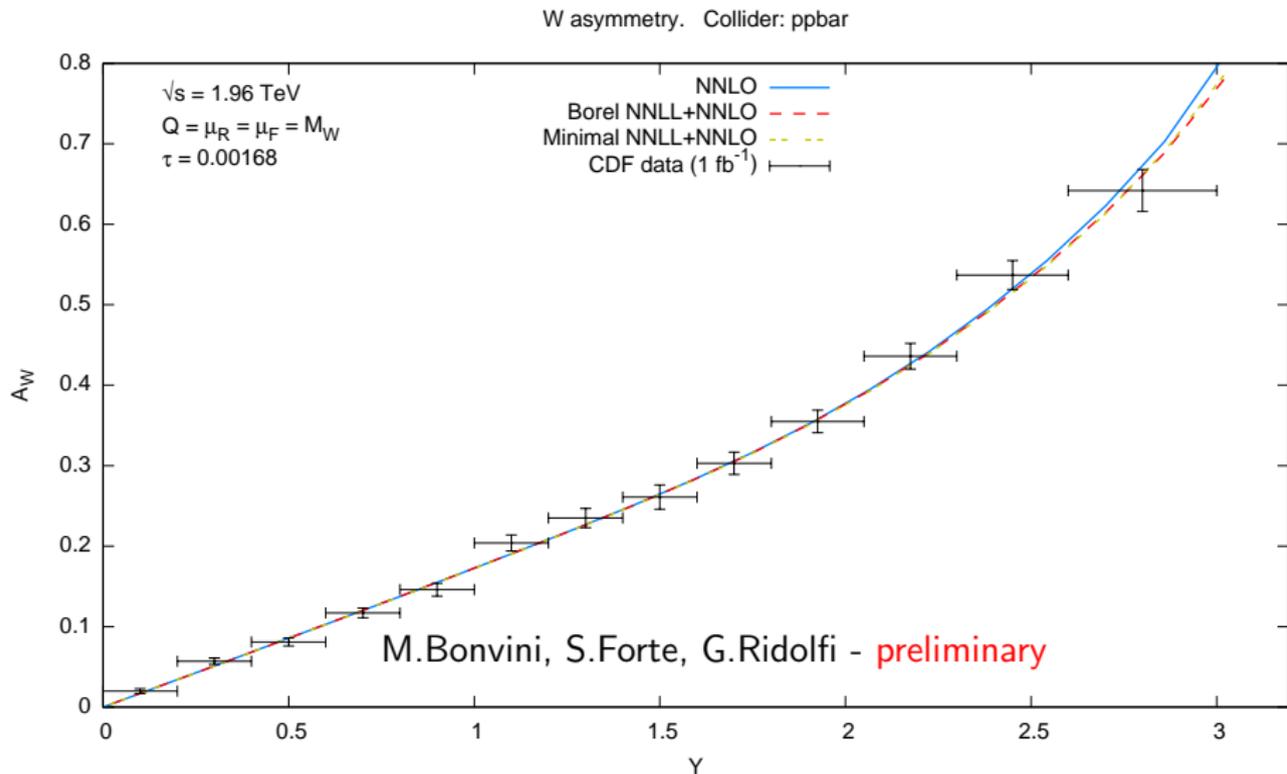
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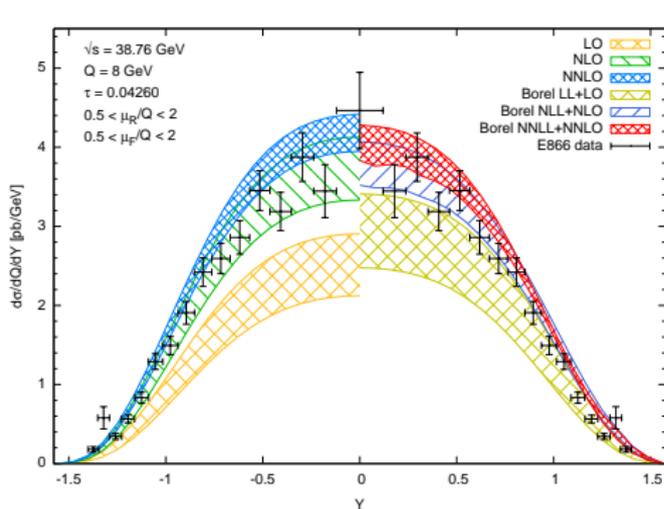
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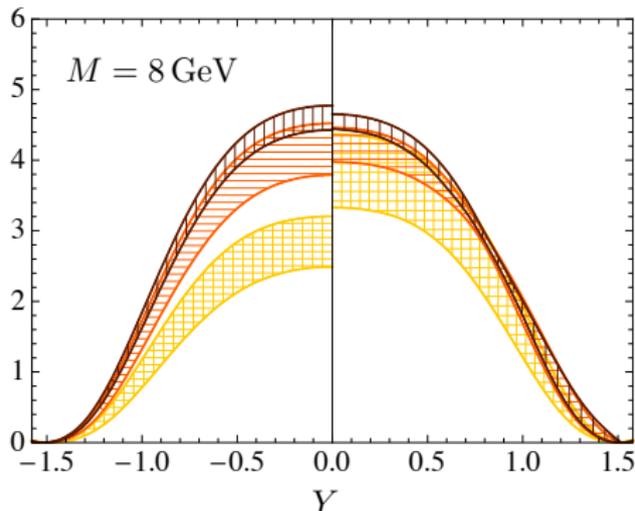
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M. Bonvini, S. Forte, G. Ridolfi
preliminary

NNPDF2.0

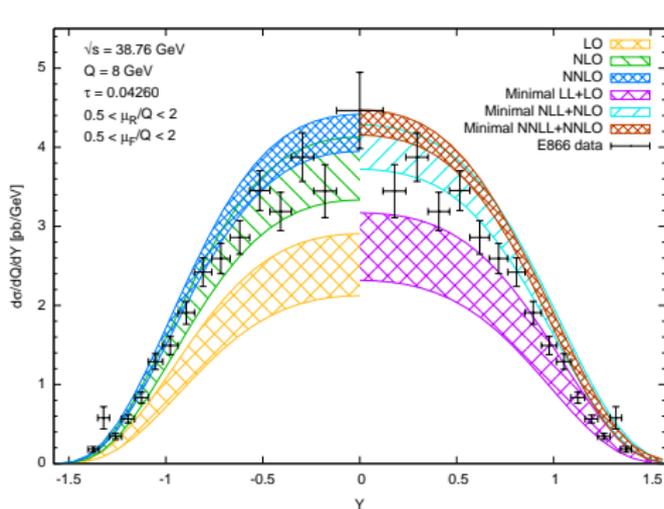


T. Becher, M. Neubert, G. Xu
(hep-ph/0710.0680)

MRST04NNLO

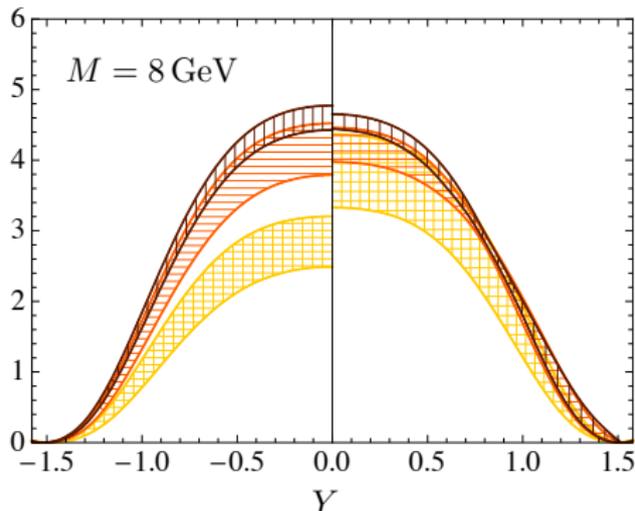
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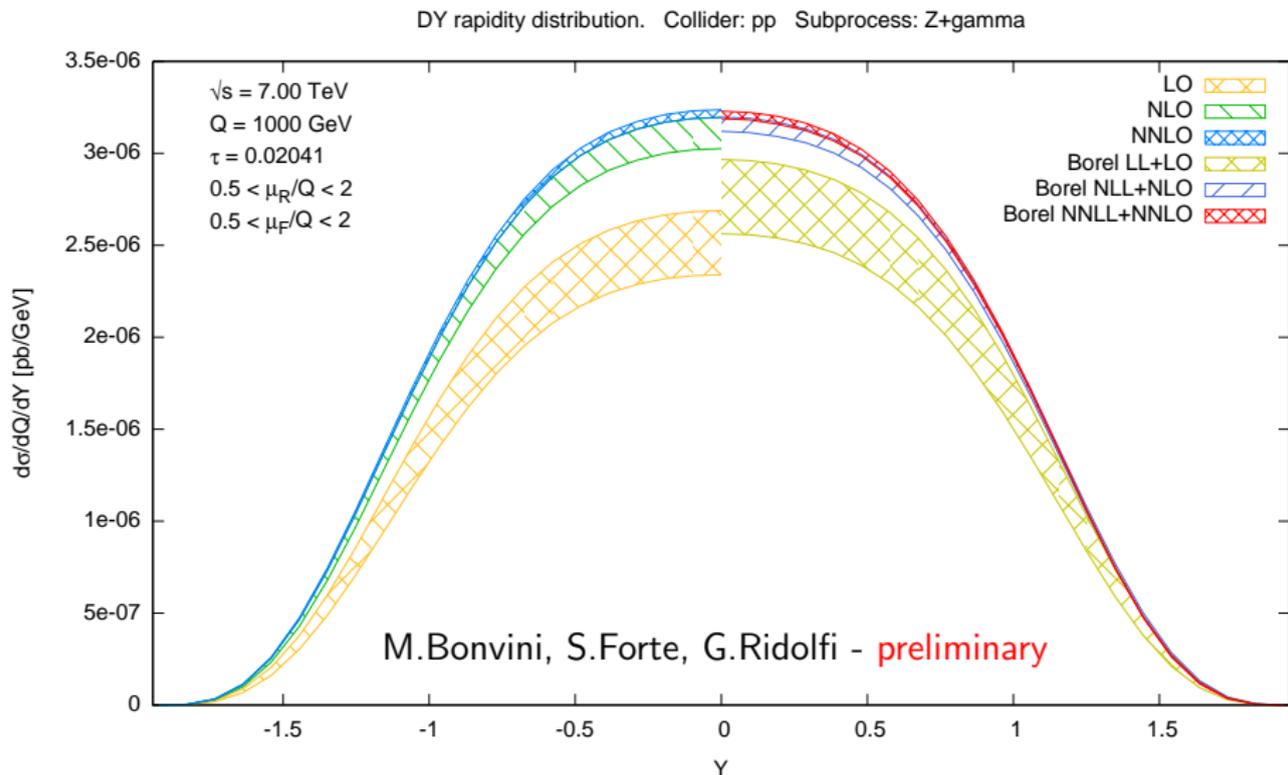
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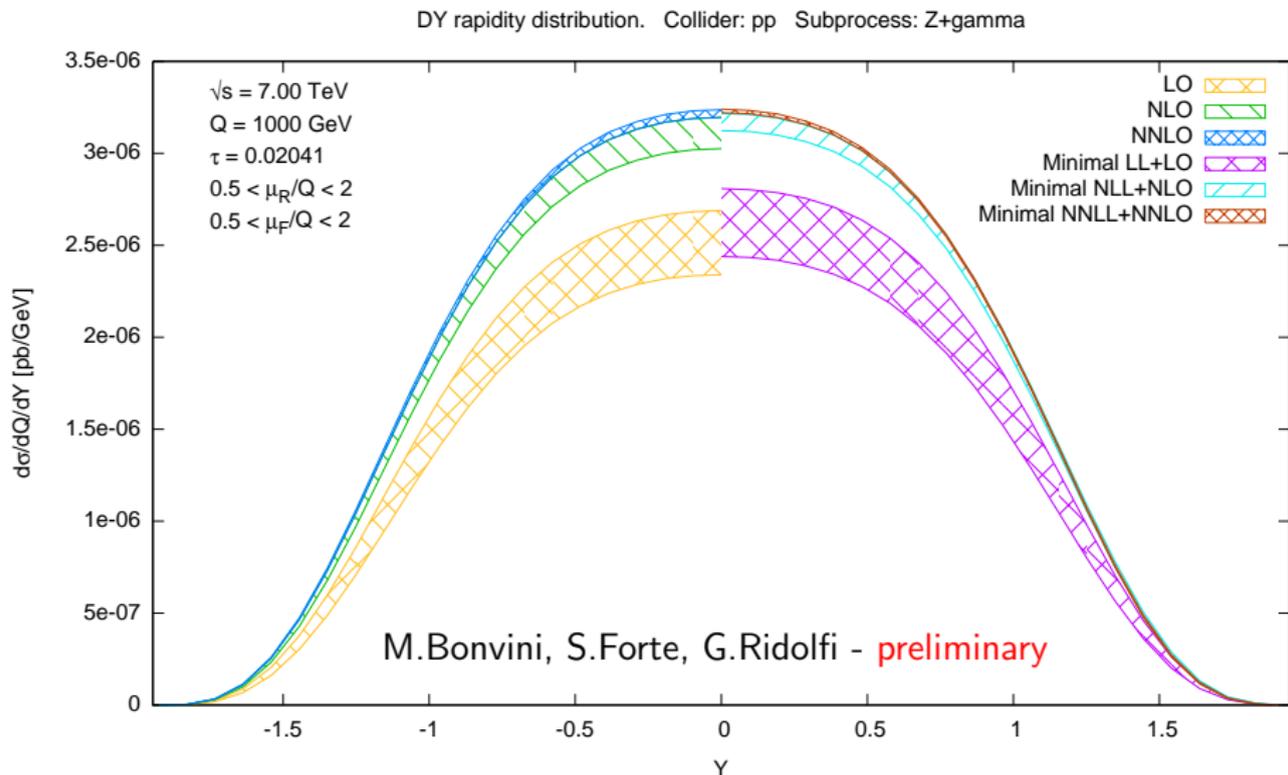
NNPDF2.0

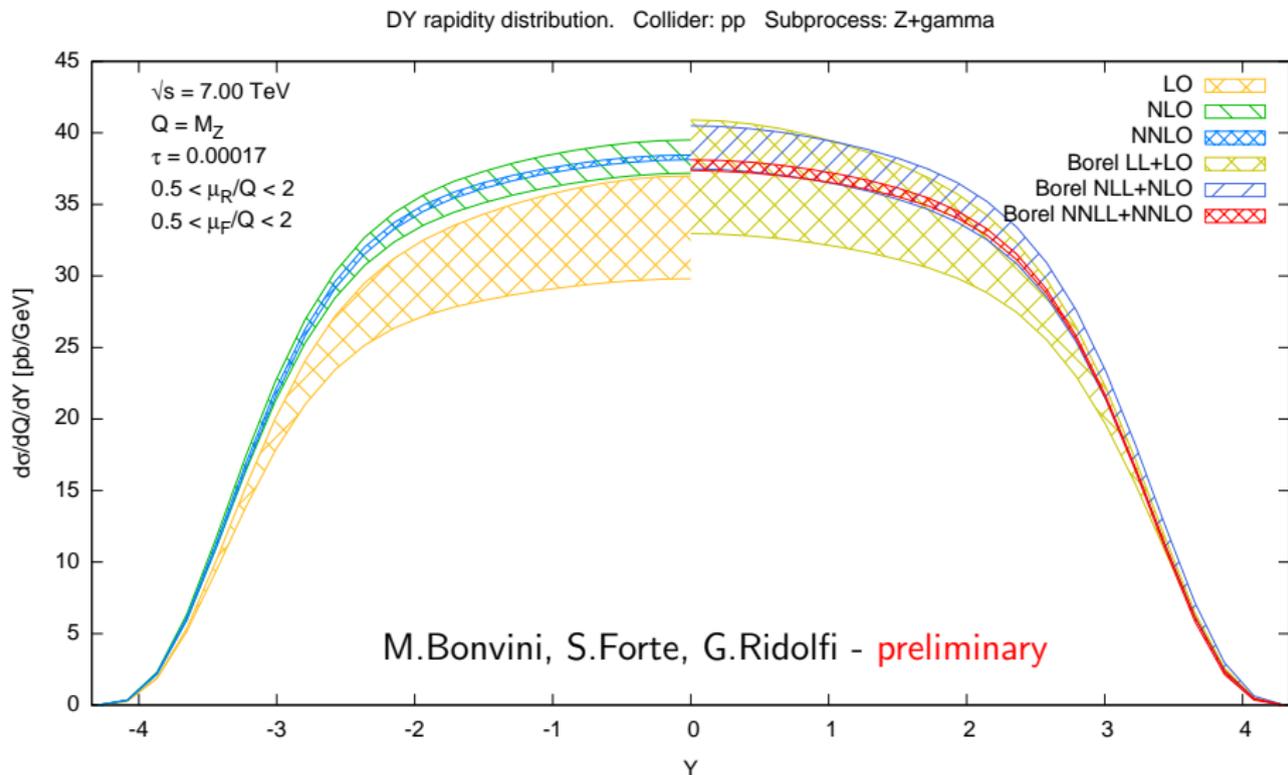


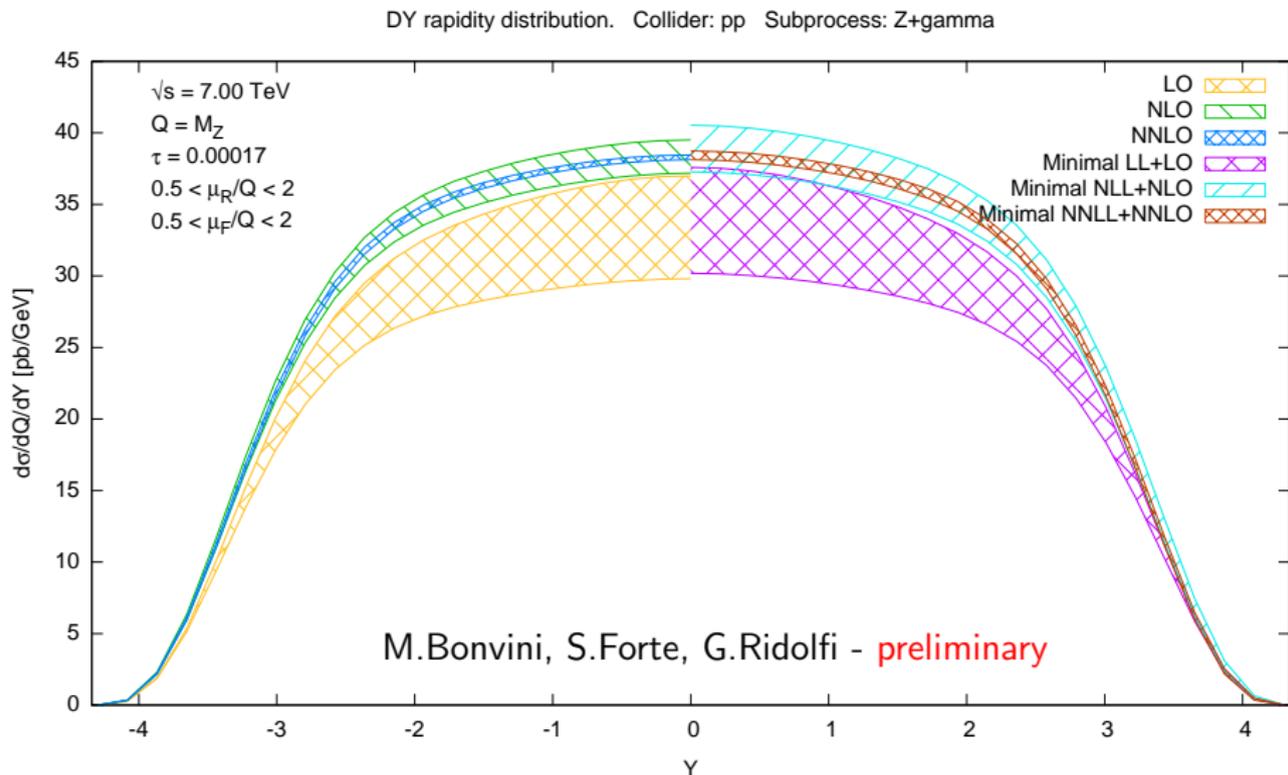
T. Becher, M. Neubert, G. Xu
(hep-ph/0710.0680)

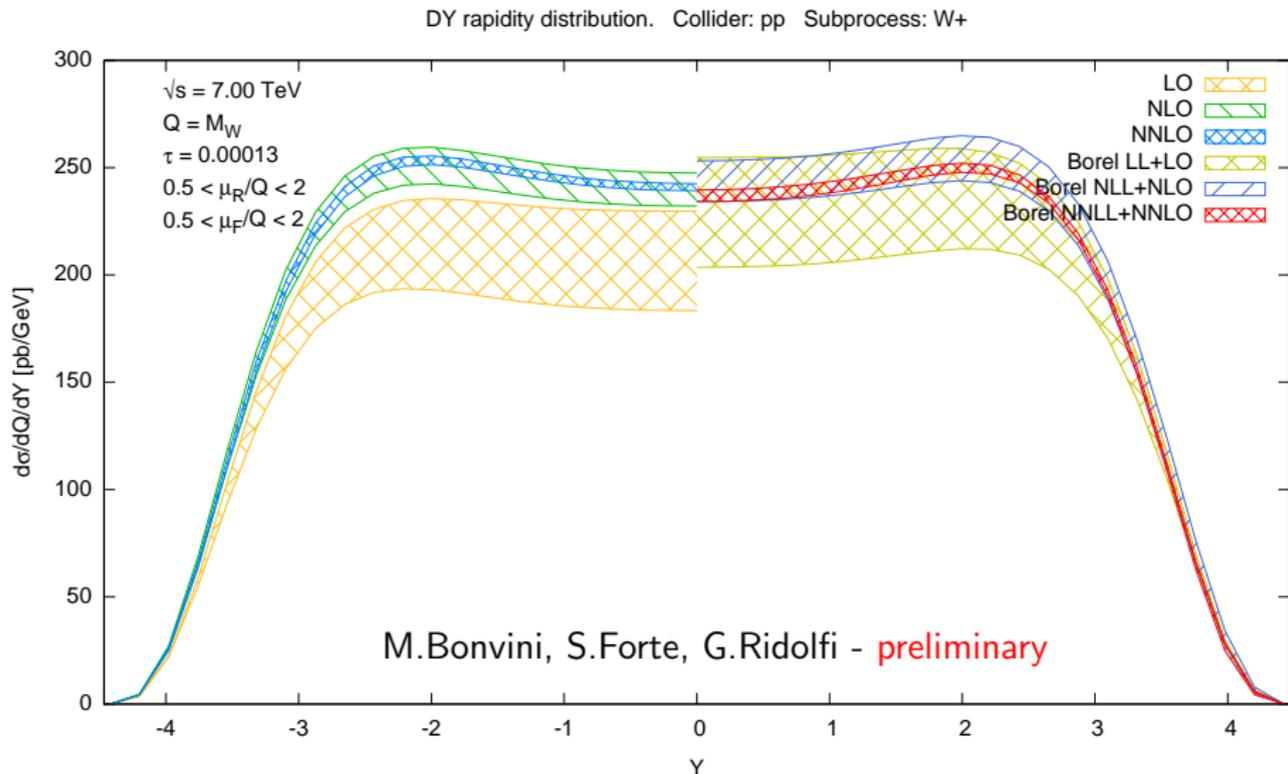
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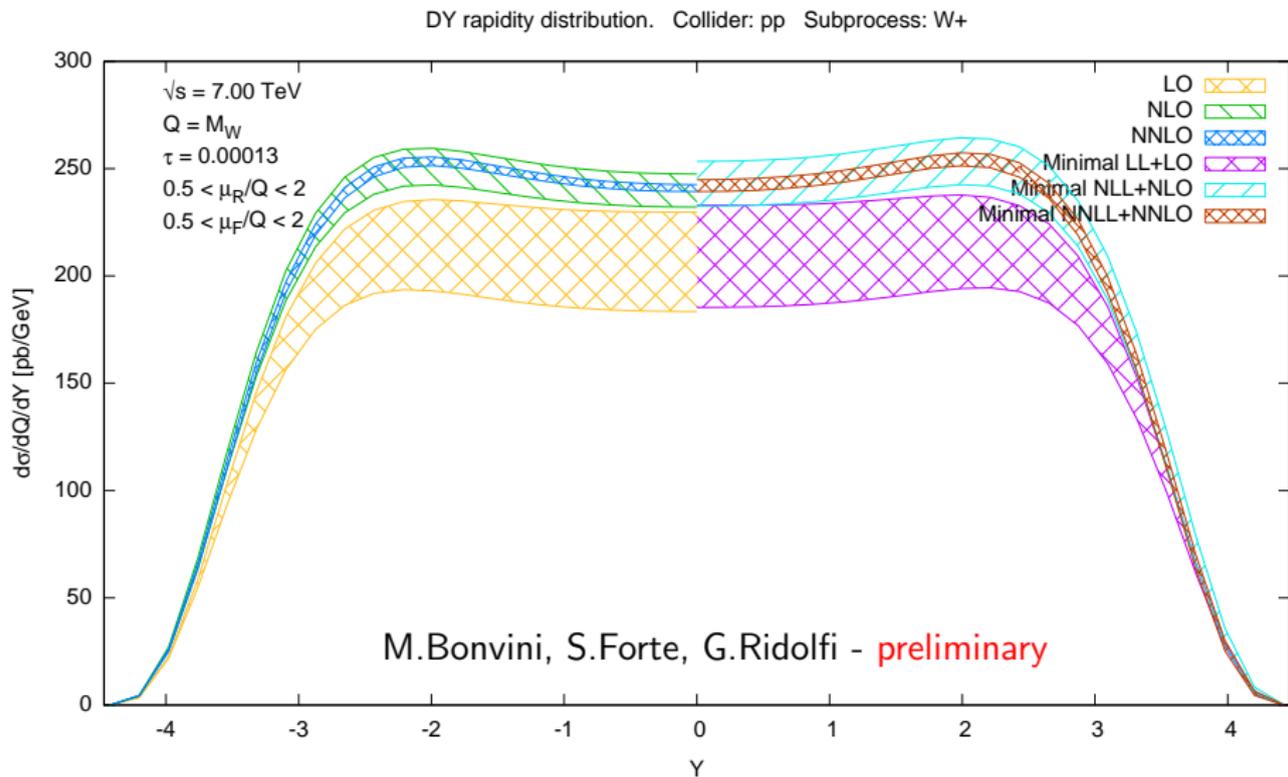


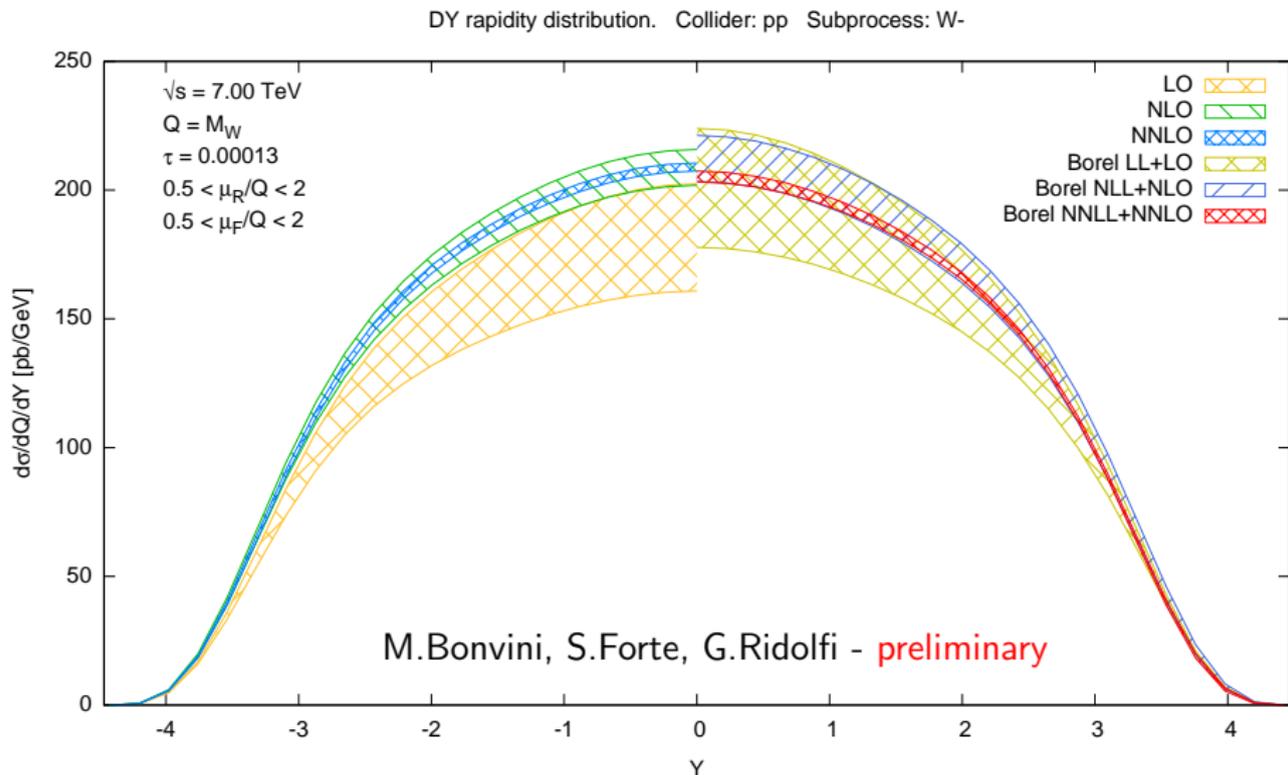


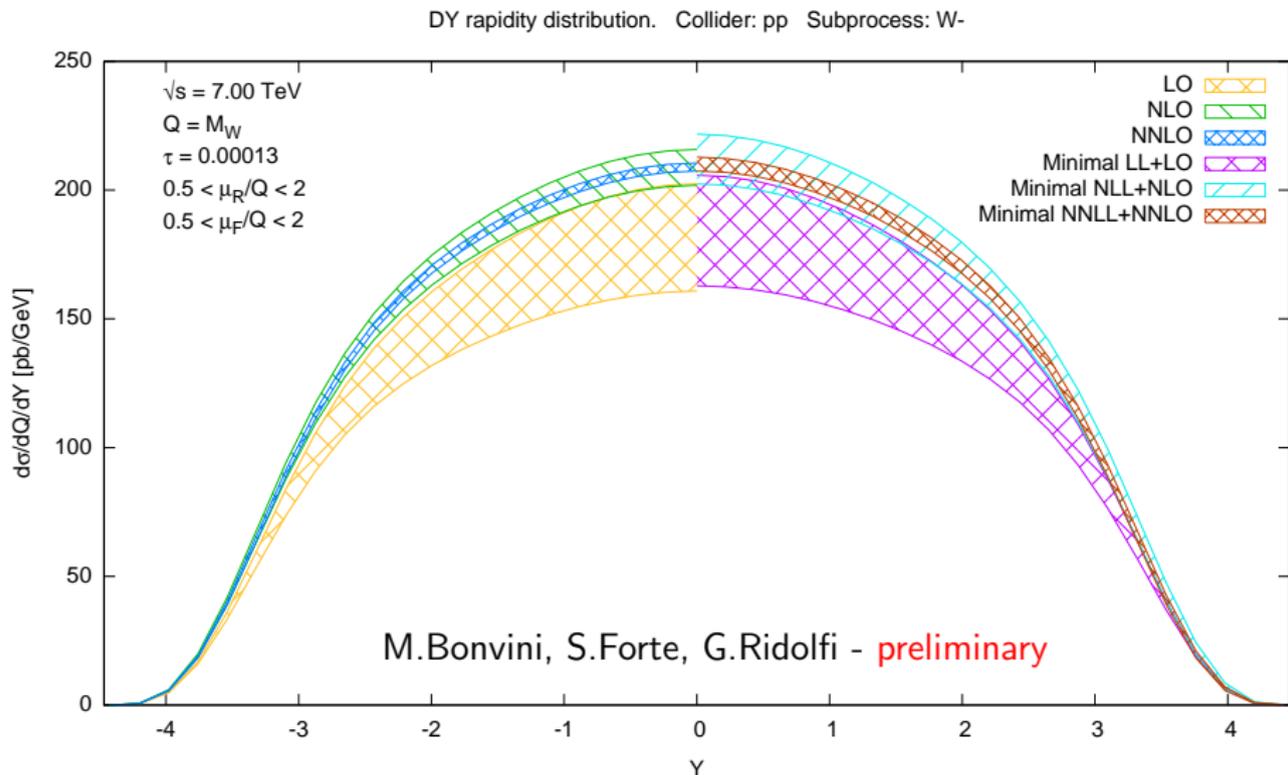












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- Include subdominant $1/N$ contributions
(S.Moch, A.Vogt: hep-ph/0909.2124 and today talk)

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- Apply to other processes such as Higgs production

Backup slides

- Expand the function

$$\frac{z^\alpha}{(1-z)^\beta} \mathcal{L}(z)$$

on a polynomial basis (with suitable $\alpha, \beta > 0$)

- Compute the Mellin transform of $\mathcal{L}(z)$ analytically
- Compute the complex Mellin inversion integral numerically

- Compute the convolution integral

$$\int_{\tau}^1 \frac{dz}{z} \mathcal{L}\left(\frac{\tau}{z}\right) [(1-z)^{\xi-1}]_+$$

It is convenient to expand on a polynomial basis the function

$$\frac{1}{1-z} \left[\frac{1}{z} \mathcal{L}\left(\frac{\tau}{z}\right) - \mathcal{L}(\tau) \right]$$

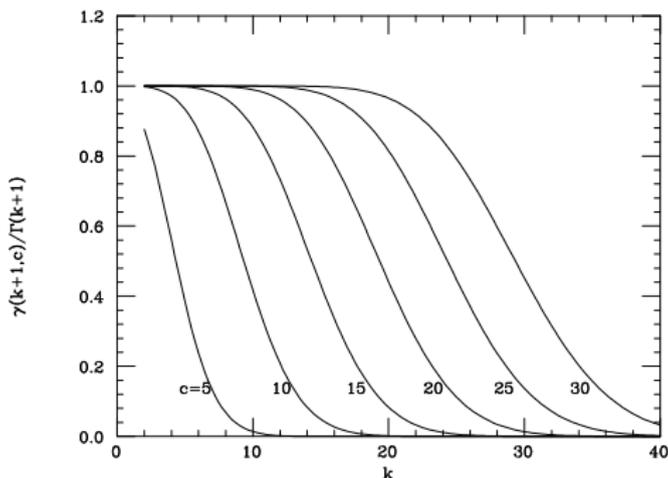
and compute the integral analytically

- Compute the complex ξ integral numerically

How BP works

Apply the BP to a power of $\log \frac{1}{N}$

$$\mathcal{M}^{-1} \left(\log^k \frac{1}{N} \right) \Big|_{\text{BP}} = \frac{\gamma(k+1, C/\bar{\alpha})}{\Gamma(k+1)} \mathcal{M}^{-1} \left(\log^k \frac{1}{N} \right)$$



The BP essentially truncates the divergent sum