



PDFs in the LHC Era

... a new perspective

Earth, as viewed from GGI

Fred Olness

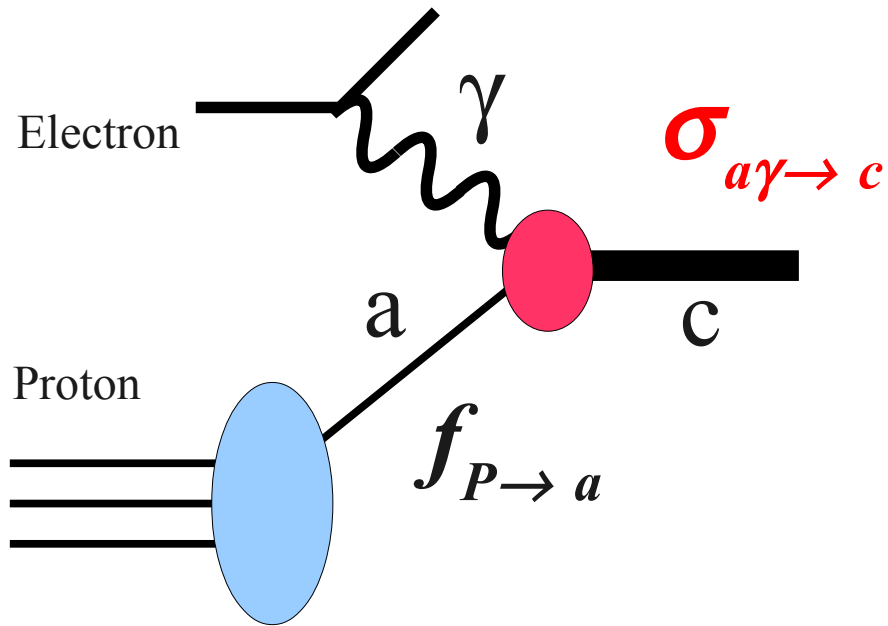
GGI

SMU

16 September 2011



The Parton Model and Factorization



Parton Distribution Functions

(PDFs) $f_{P \rightarrow a}$

are the key to calculations
involving hadrons!!!

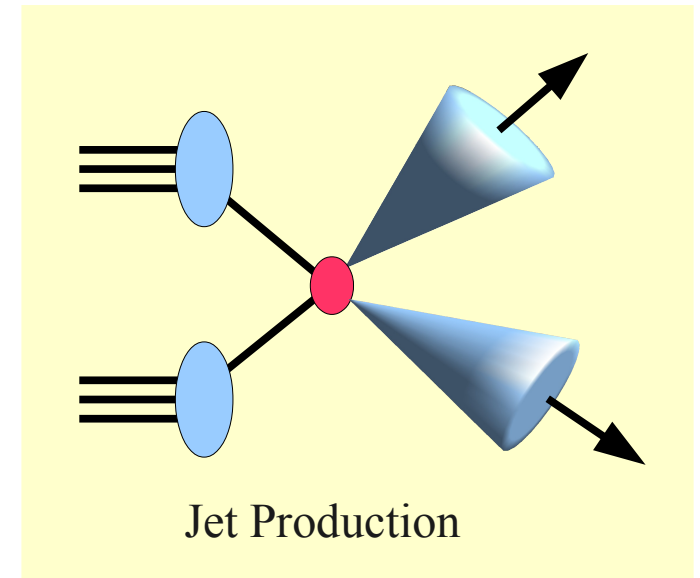
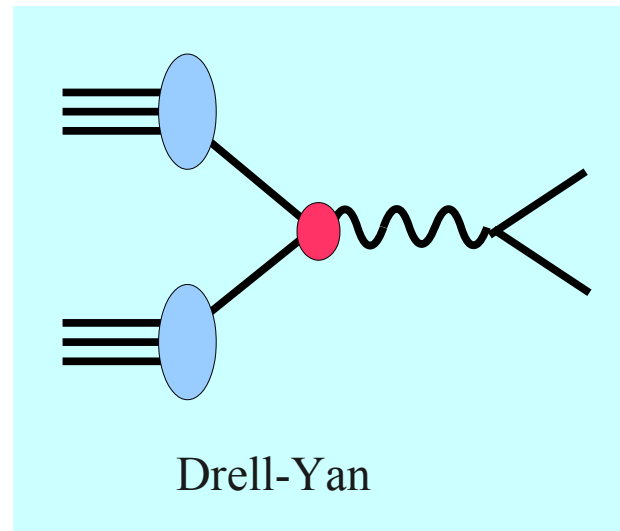
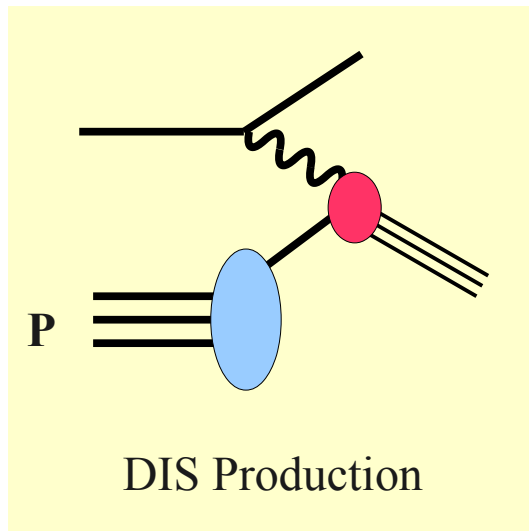
$$\sigma_{P \gamma \rightarrow c} = f_{P \rightarrow a} \otimes \hat{\sigma}_{a \gamma \rightarrow c}$$

must extract from
experiment

calculable from
theoretical model

Corrections of
order (Λ^2/Q^2)

Cross section is product of independent probabilities!!! (Homework Assignment)



$$F_2^\nu \sim [d + s + \bar{u} + \bar{c}]$$

$$F_2^{\bar{\nu}} \sim [\bar{d} + \bar{s} + u + c]$$

$$F_3^\nu = 2 [d + s - \bar{u} - \bar{c}]$$

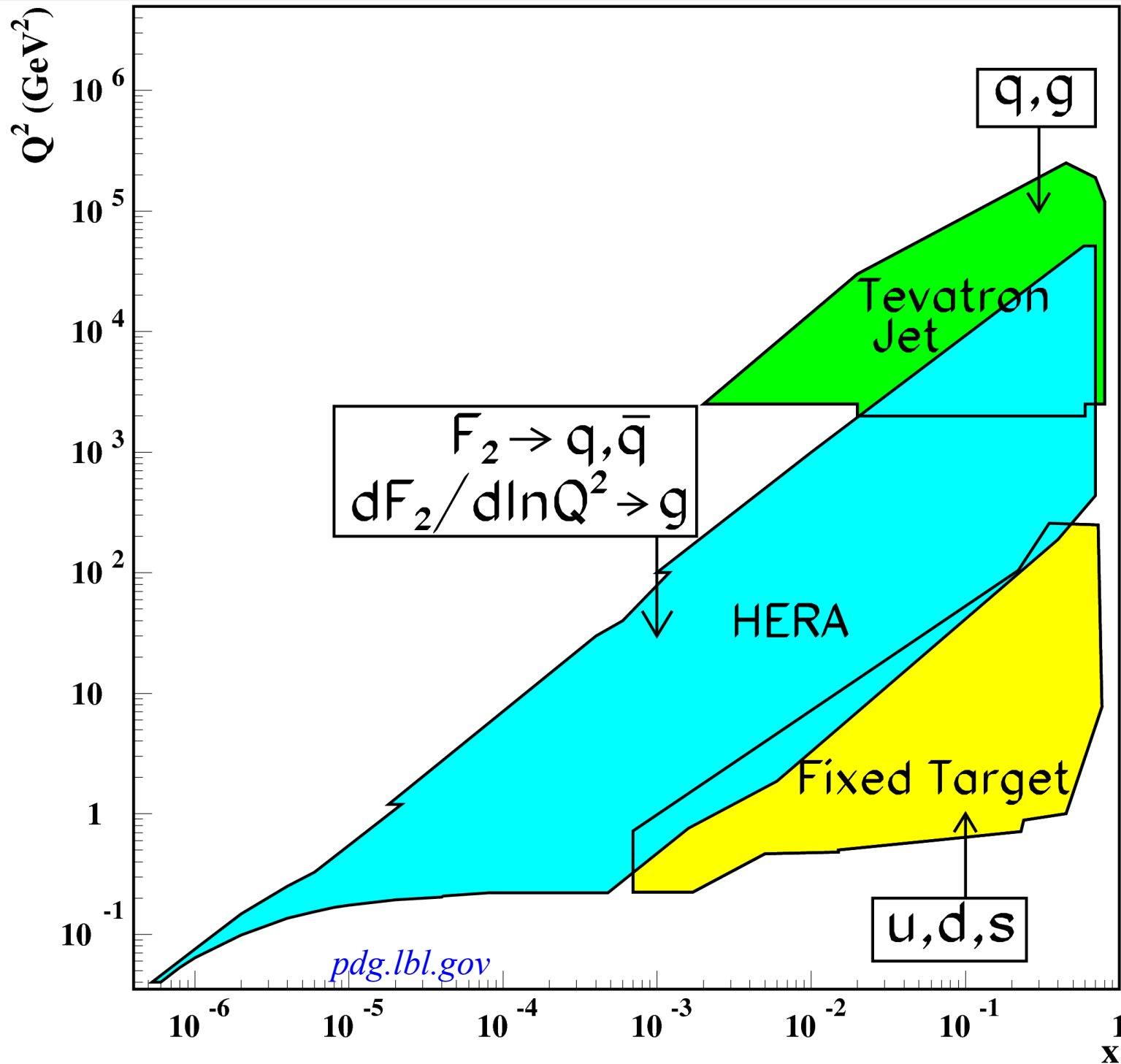
$$F_3^{\bar{\nu}} = 2 [u + c - \bar{d} - \bar{s}]$$

$$F_2^{\ell^\pm} \sim \left(\frac{1}{3}\right)^2 [d + s] + \left(\frac{2}{3}\right)^2 [u + c]$$

In particular, the DIS combinations have historically been particularly useful

Different linear combinations – key for flavor differentiation

The ν -DIS data typically use heavy targets, and this requires the application of nuclear corrections



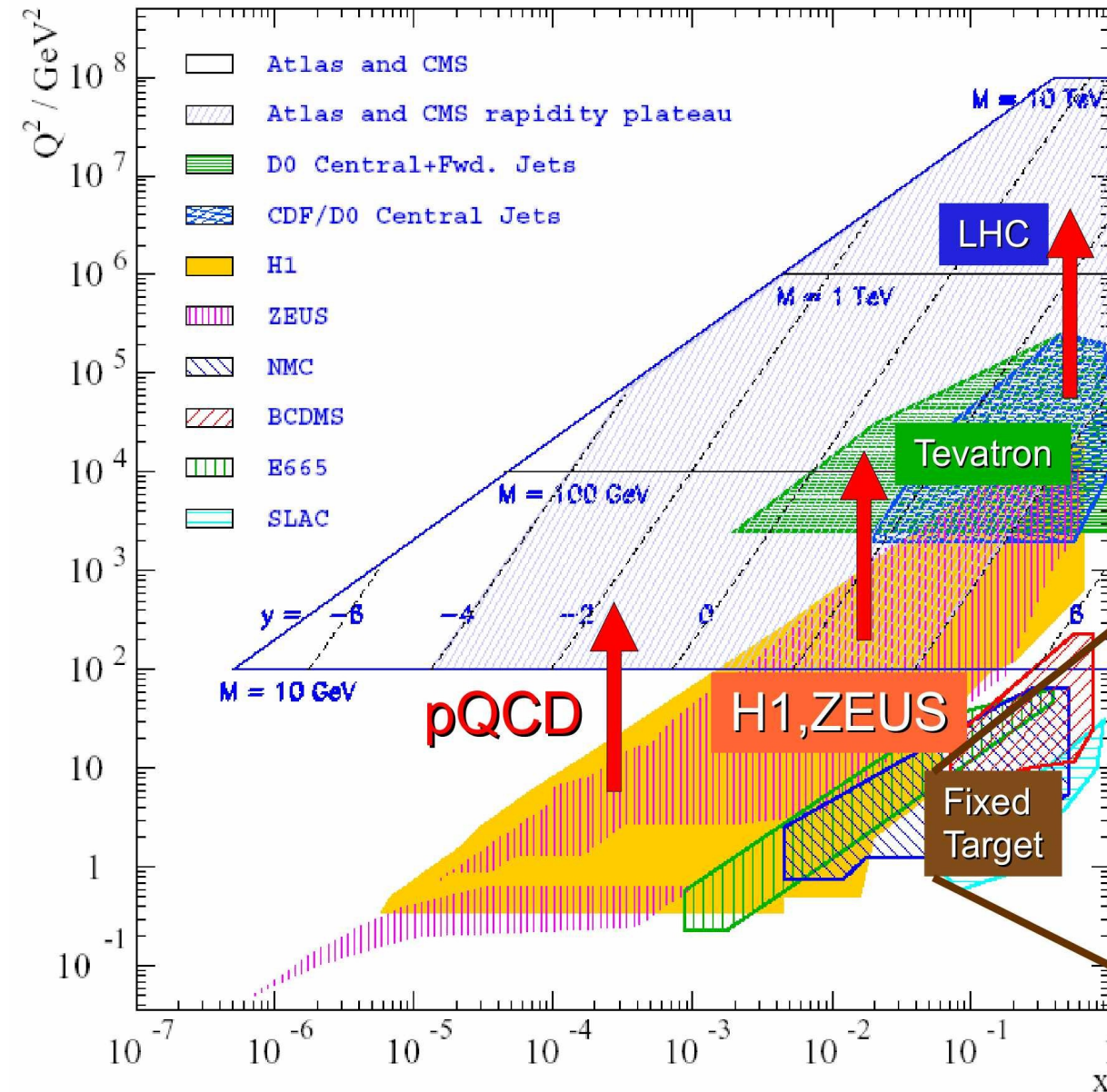
Fixed Target:
 DIS charged lepton
 DIS neutrino
 DY
 Direct Photon

HERA (ep)

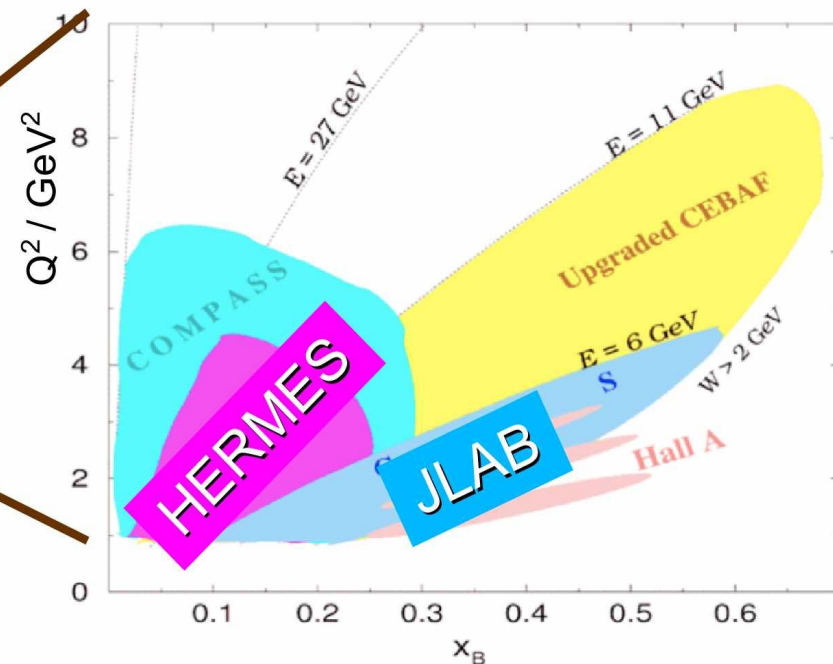
Tevatron:
 Jets
 W Asymmetry

LHC DATA
just now being included

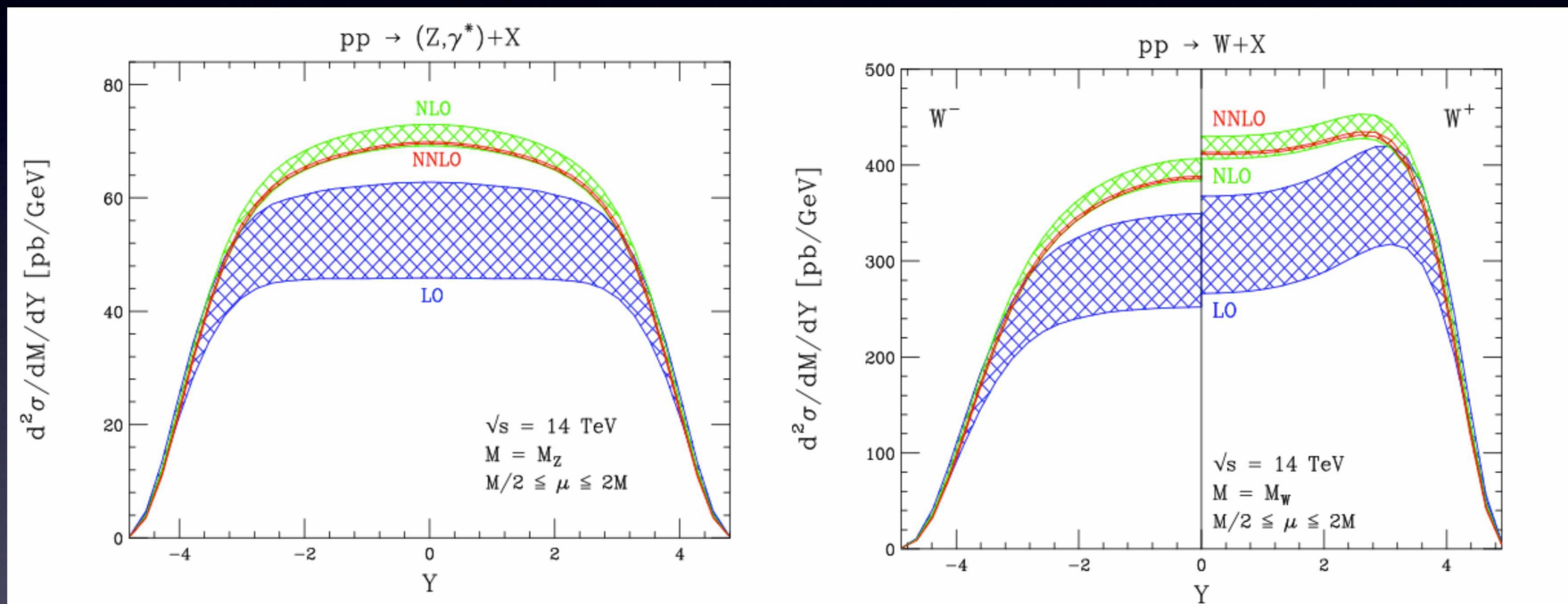
Kinematics



- H1, ZEUS: large coverage in x and Q^2
- HERMES, JLAB: access to large x and low Q^2
- pQCD (DGLAP) predicts Q^2 dependence



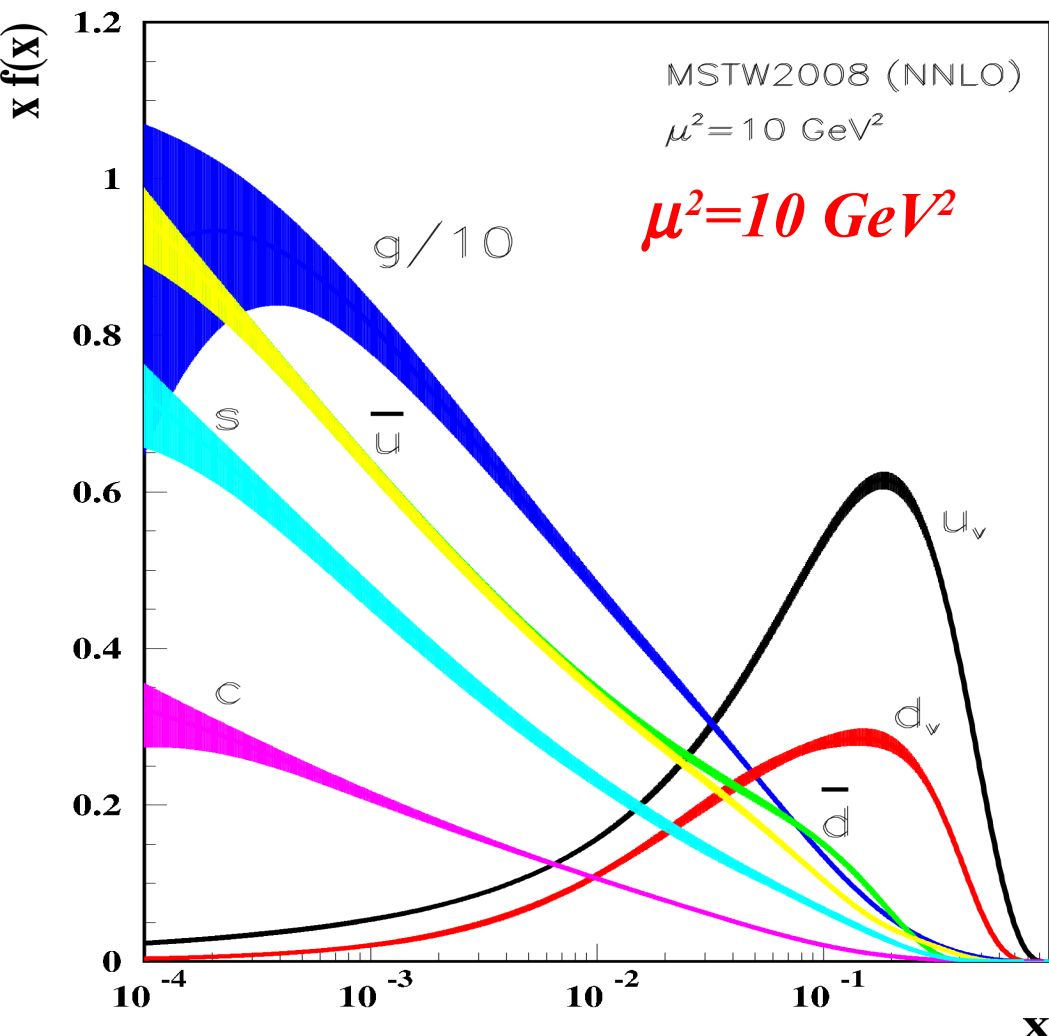
Rapidity distributions



Anastasiou, Dixon, Melnikov, Petriello '03, '05; Melnikov, Petriello '06

👉 LHC: perturbative accuracy of the order of 1%. This is absolutely unique.

PDFs are certainly one of the foundations
that our search for “new physics”
is built upon



LHC results

LHC is ideal for certain measurements

For full picture, we need to combine with other measurements

**These measurements are interdependent
and part of the foundation we use
to calibrate the search for “new physics”**

I will look at some select examples

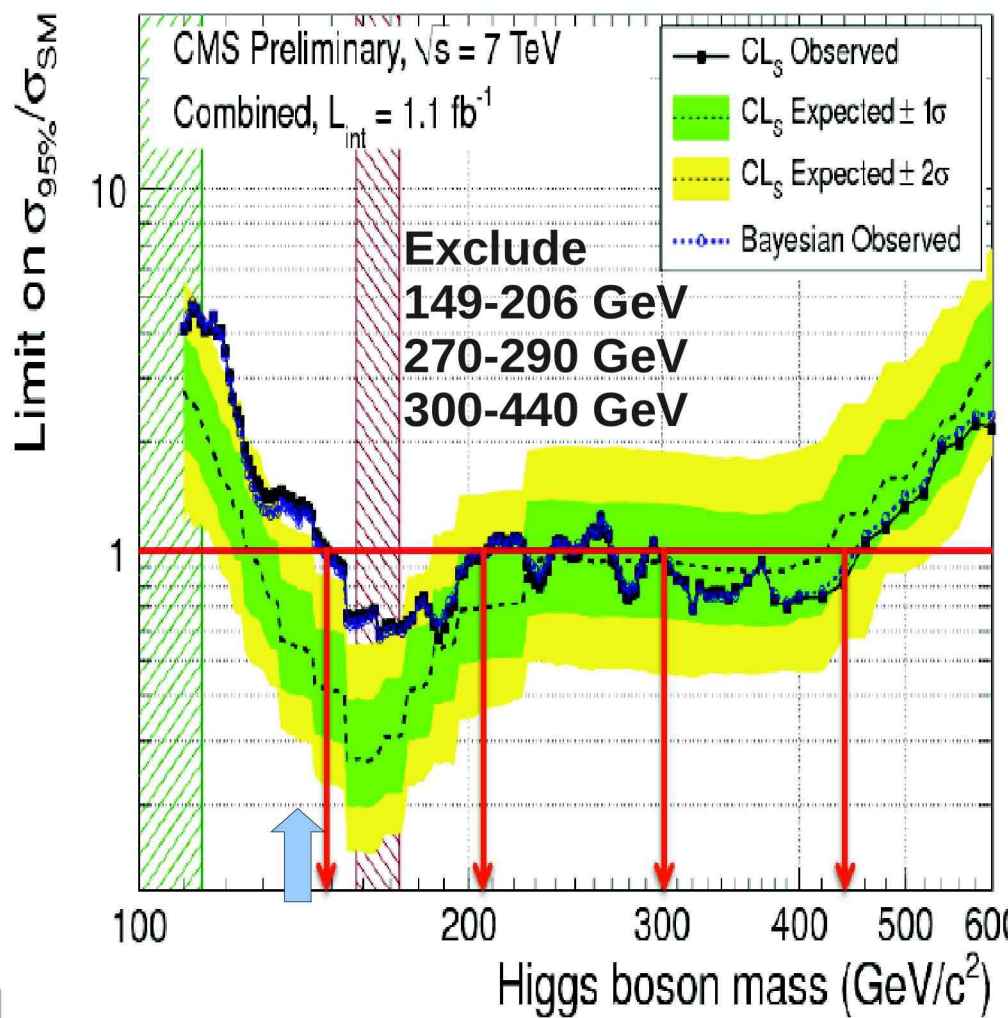
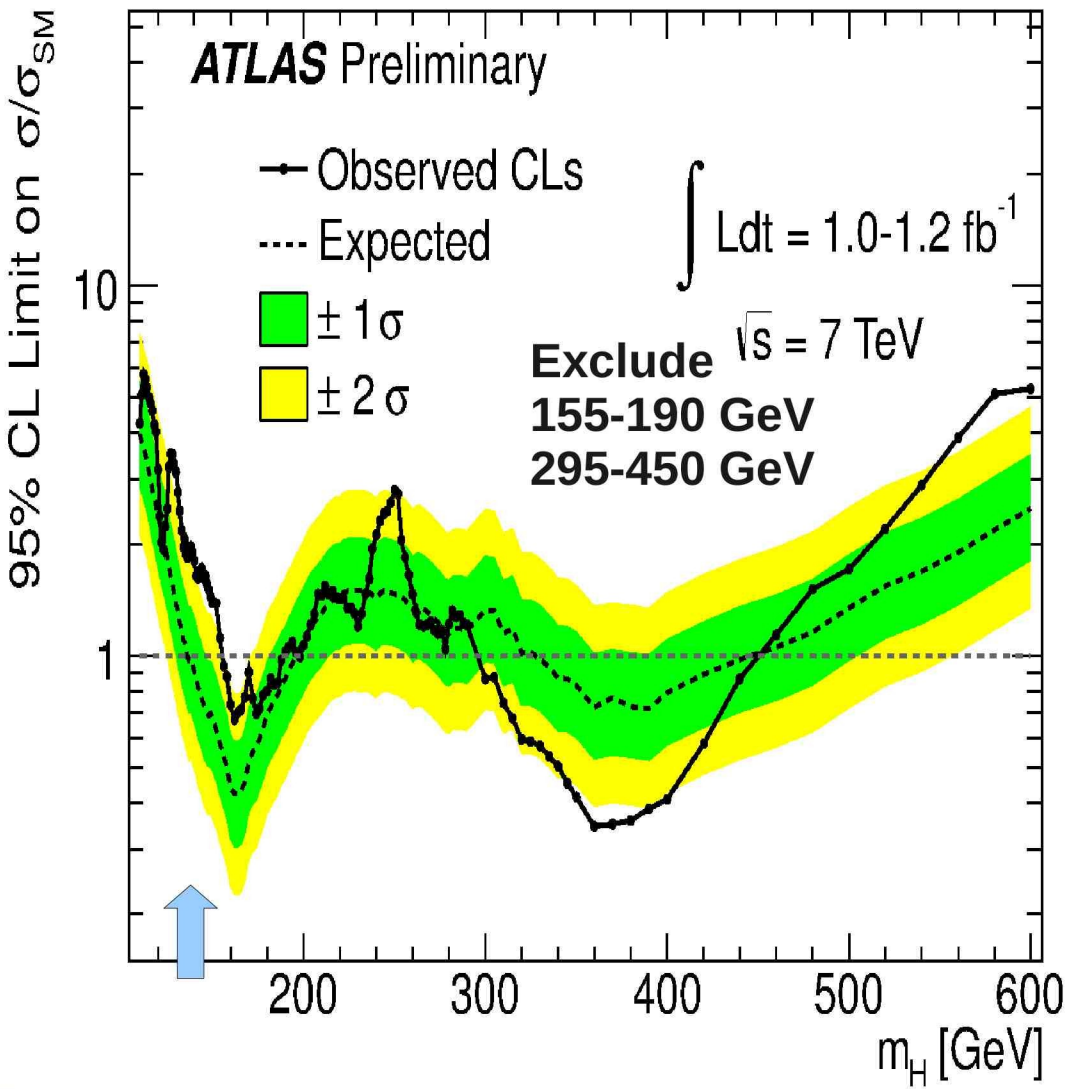
Higgs Production

...



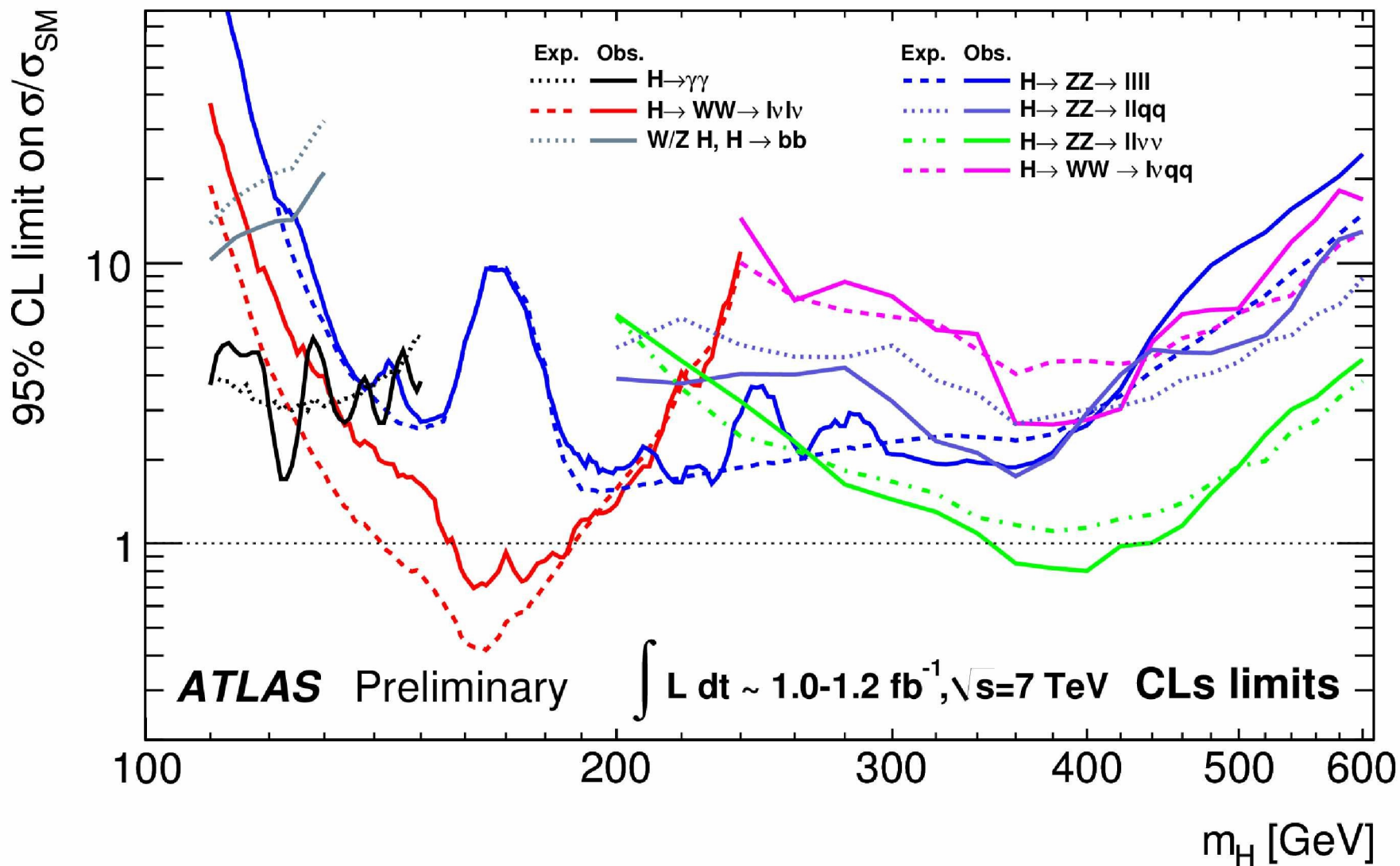
ATLAS & CMS limits

Sensitivities differ in detail
But on average similar





Channels reviewed (ATLAS)



W/Z

Production

“Benchmark Calculations”

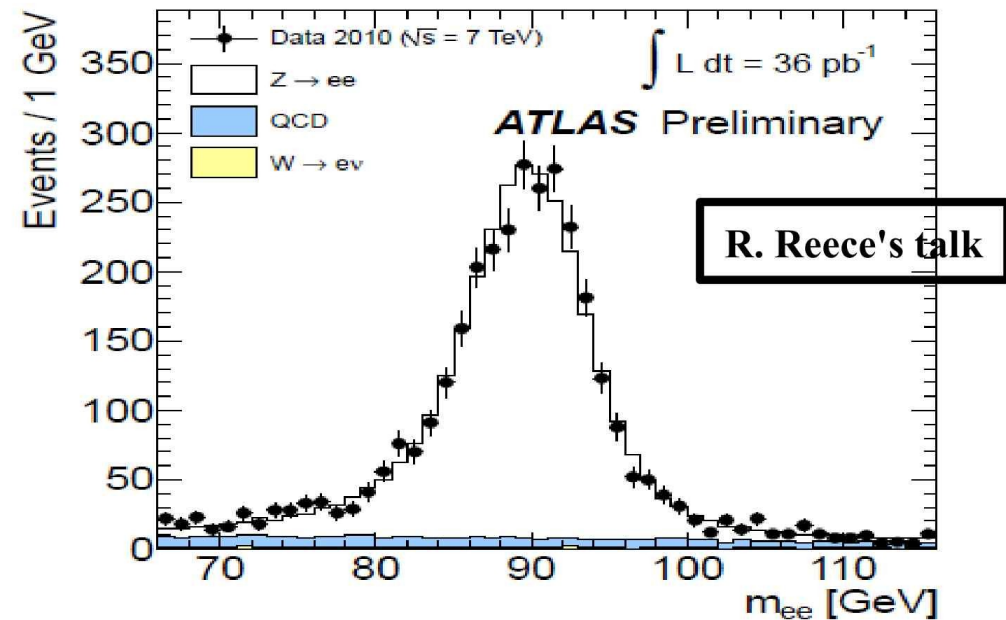
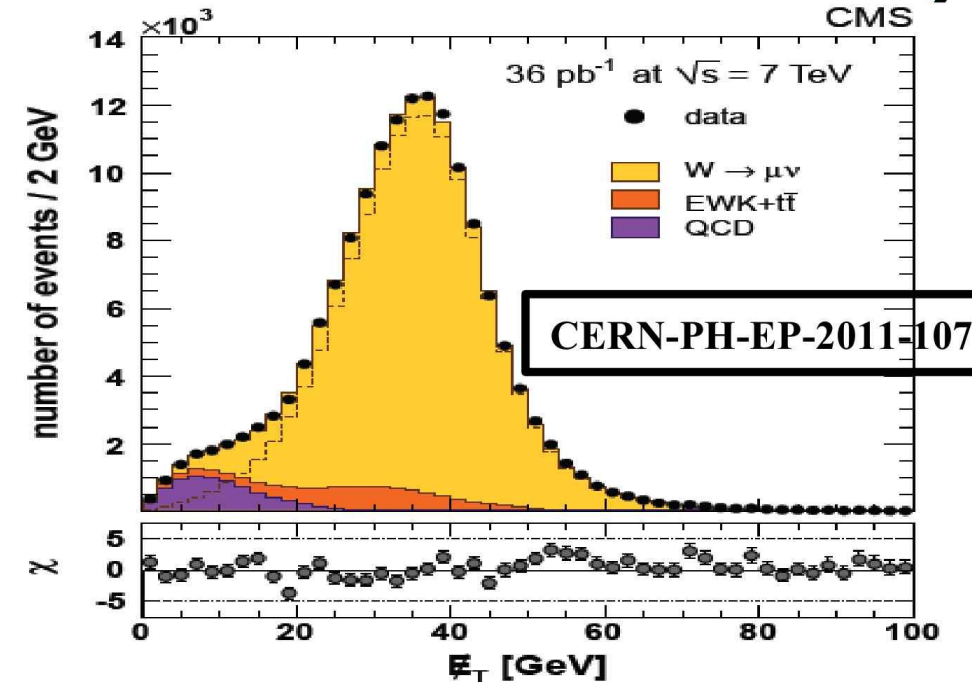
W and Z selection (muons/electrons)

W selection:

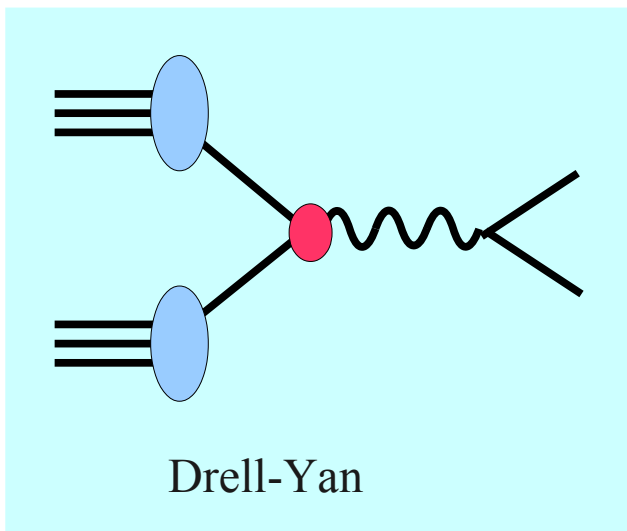
- High- p_T lepton ($p_T > 20-25$ GeV)
- Well Isolated from hadronic activity
- Loose cut on missing transverse energy (or not cut at all (CMS))
- Efficiencies, resolutions, signal and background shapes studied / extracted from data.

Z selection:

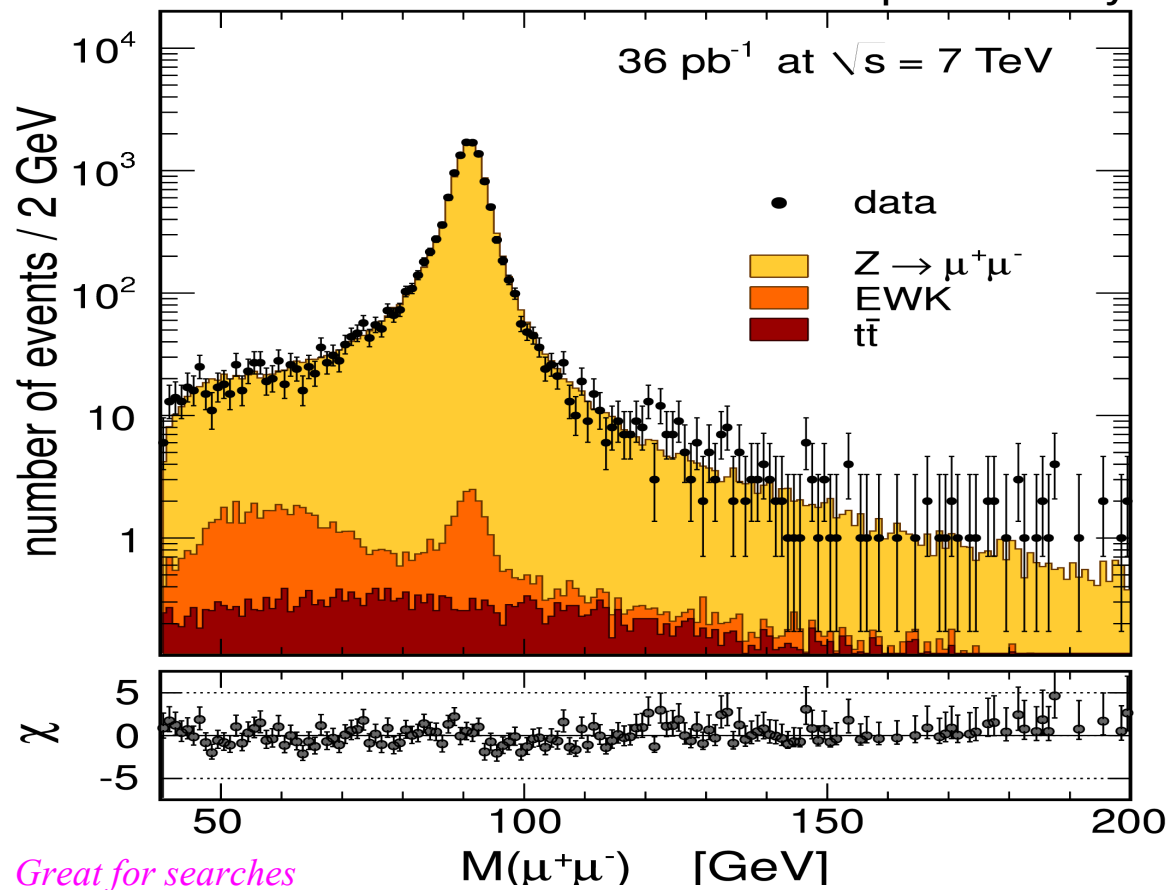
- Two high- p_T leptons ($p_T > 20-25$ GeV), also isolated
- Dilepton mass consistent with a Z
- Efficiencies, resolutions studied / extracted from data.
- Almost background free



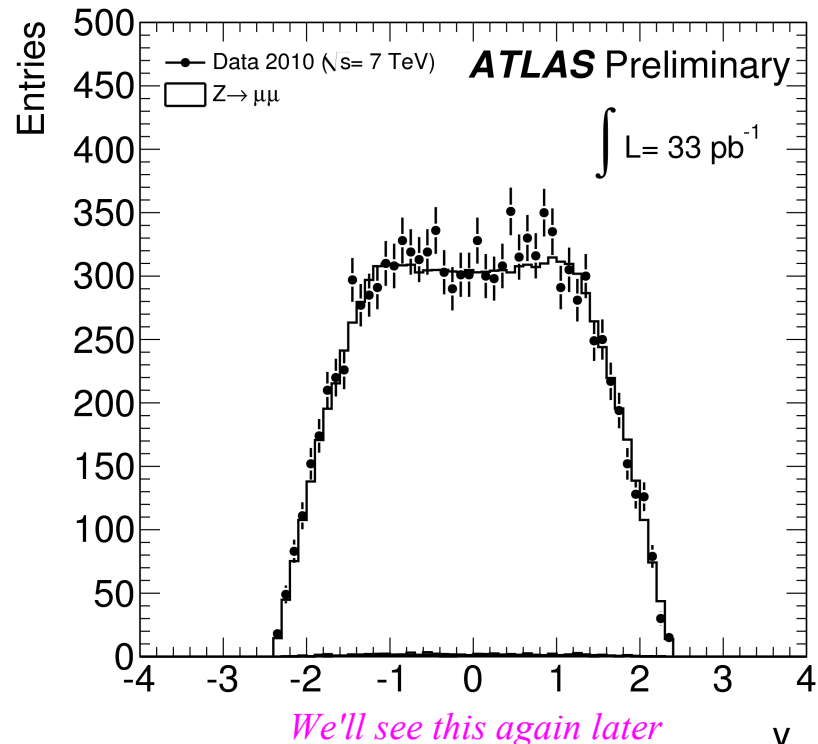
W/Z PRODUCTION



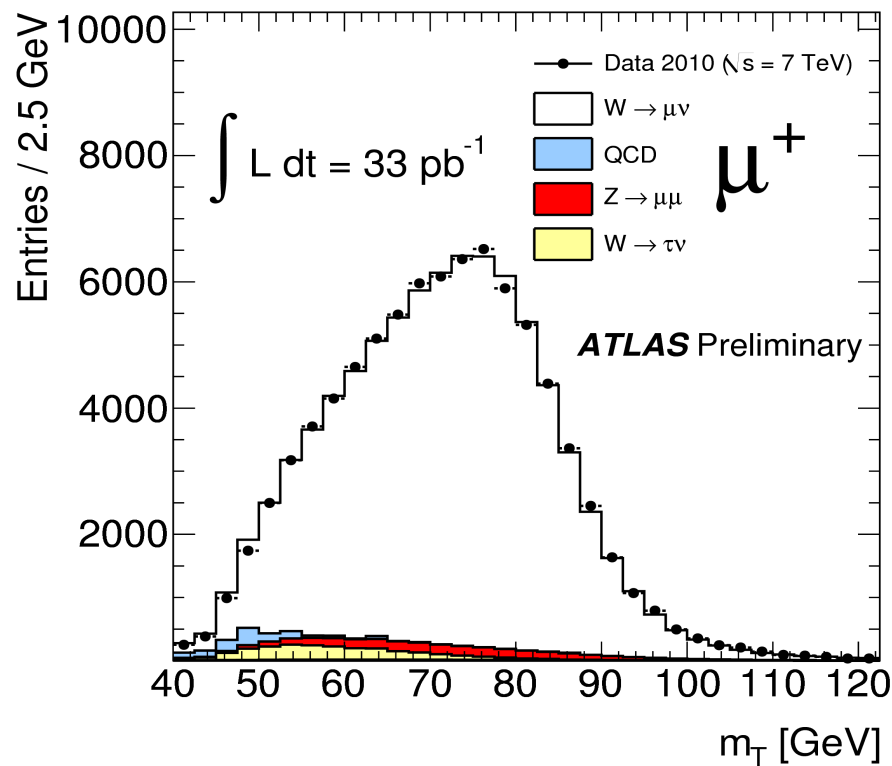
CMS preliminary



Great for searches



We'll see this again later



Systematic uncertainties (CMS)

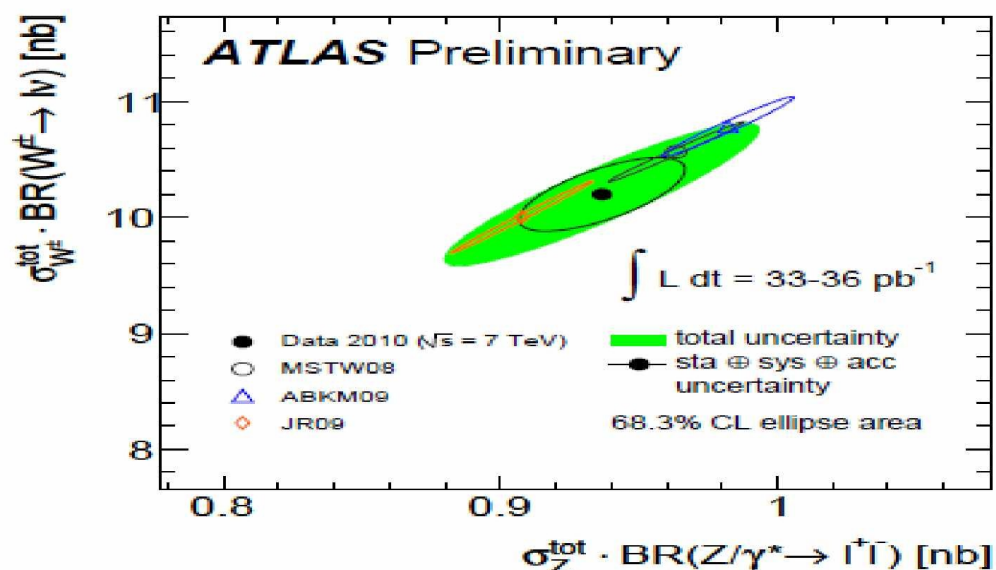
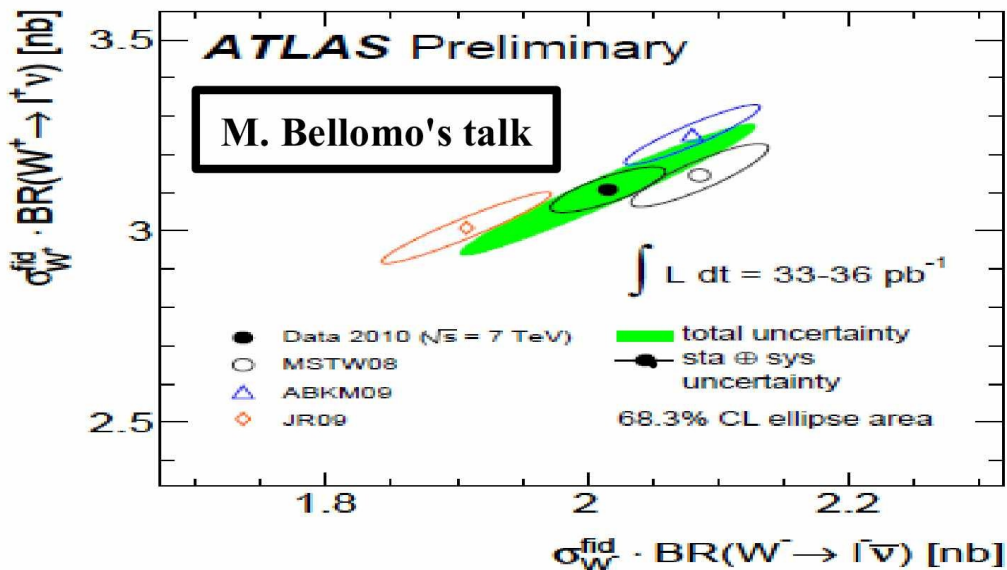
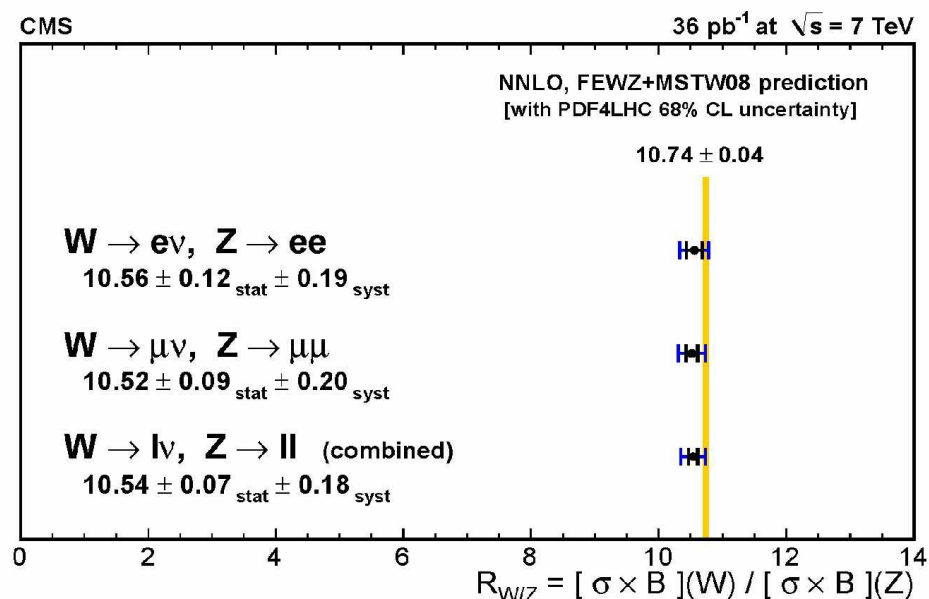
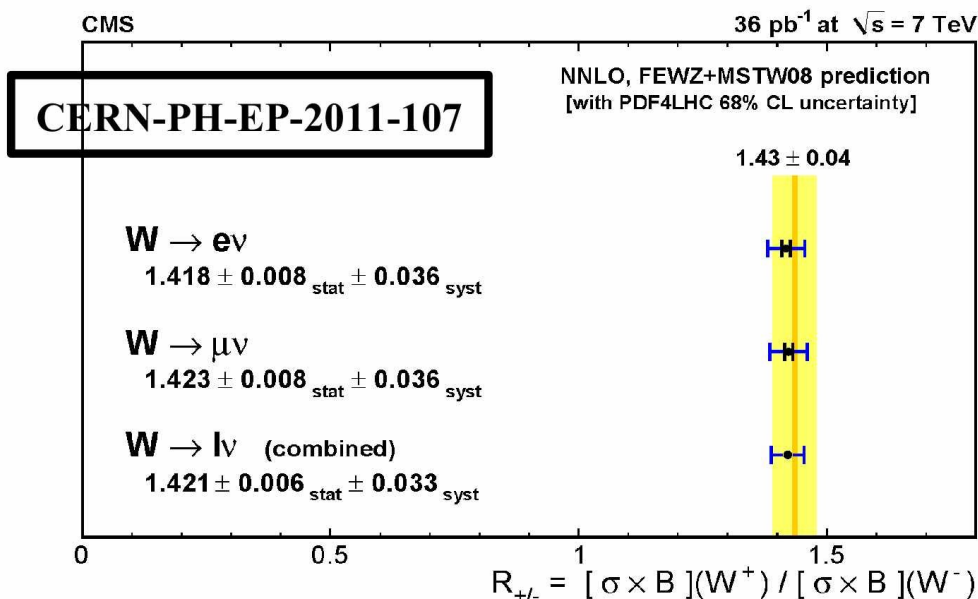
Source (%)	$W \rightarrow e\nu$	$W \rightarrow \mu\nu$	$Z \rightarrow e^+e^-$	$Z \rightarrow \mu^+\mu^-$
Lepton reconstruction & identification	1.4	0.9	1.8	n/a
Trigger prefiring	n/a	0.5	n/a	0.5
Energy/momentum scale & resolution	0.5	0.22	0.12	0.35
E_T scale & resolution	0.3	0.2	n/a	n/a
Background subtraction / modeling	0.35	0.4	0.14	0.28
Trigger changes throughout 2010	n/a	n/a	n/a	0.1
Total experimental	1.6	1.1	1.8	0.7
PDF uncertainty for acceptance	0.6	0.8	0.9	1.1
Other theoretical uncertainties	0.7	0.8	1.4	1.6
Total theoretical	0.9	1.1	1.6	1.9
Total (excluding luminosity)	1.8	1.6	2.4	2.0

- These are the final 2010 results from CMS (already submitted for publication)
- Experimental uncertainties are significantly reduced thanks to the extensive use of data-driven methods to control efficiencies, backgrounds and signal shapes
- Theoretical and experimental uncertainties have similar sizes
- Measurements in fiducial volume are also provided (smaller theory uncertainty)

w/z

Ratios

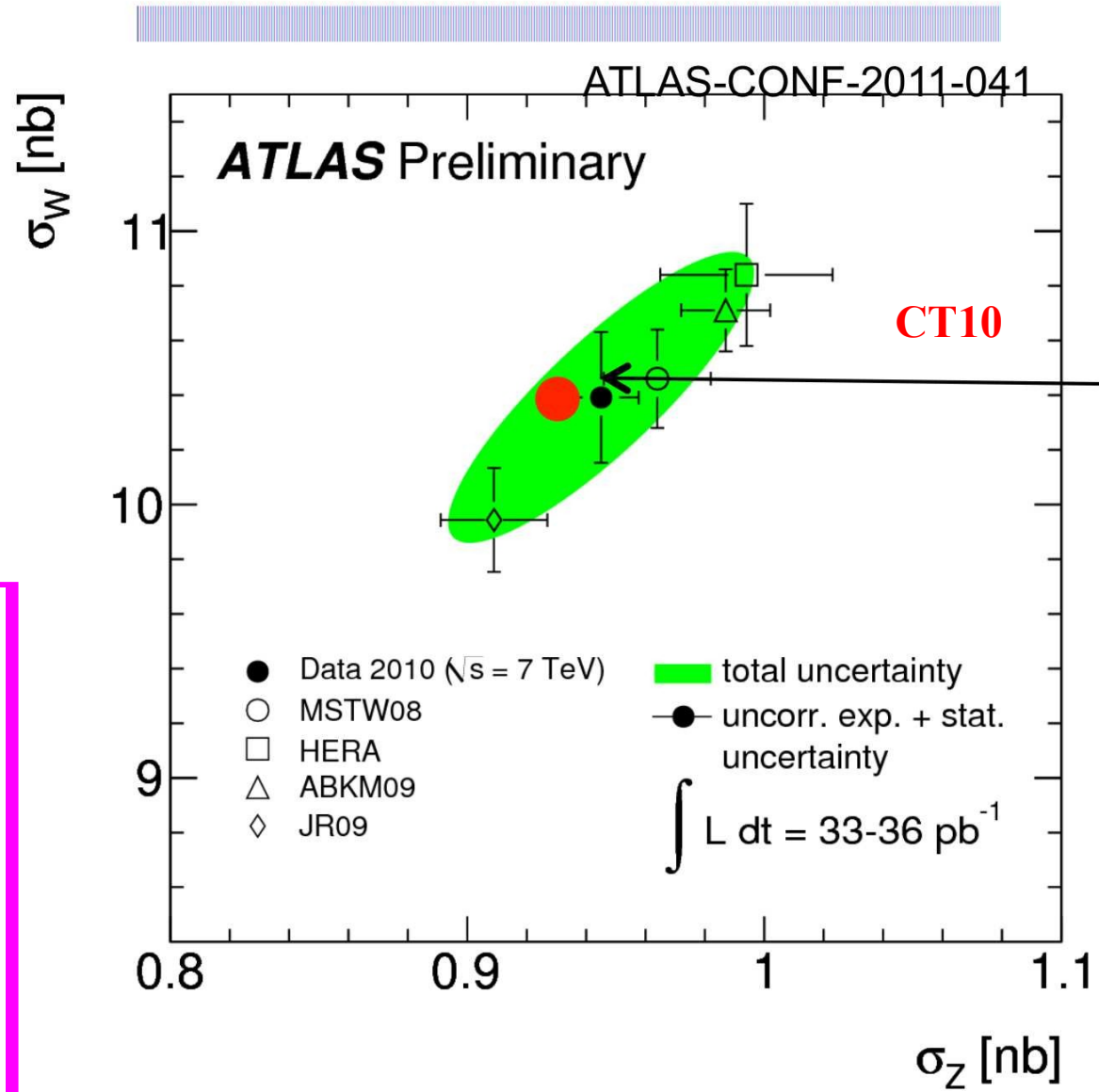
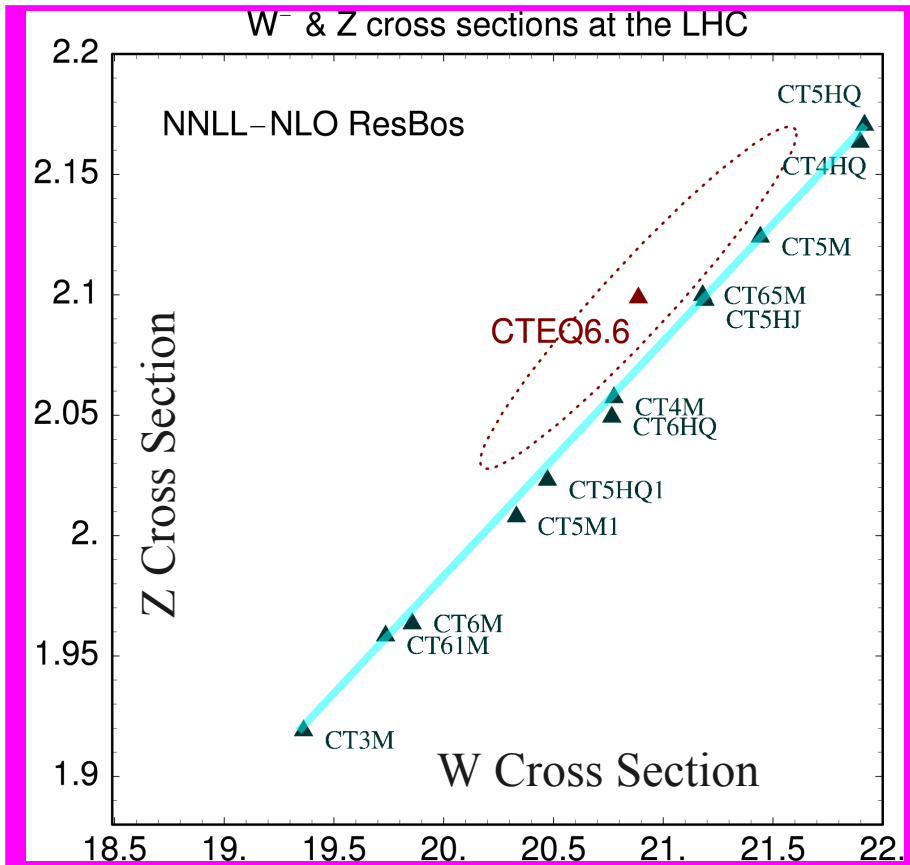
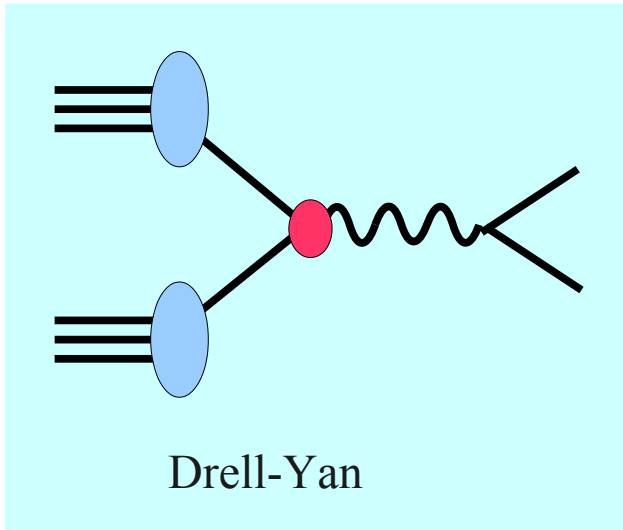
LHC W+/W- and W/Z ratios



W+/W-: potential to constrain PDF uncertainties

W/Z: stringent test of theoretical expectations

W/Z PRODUCTION



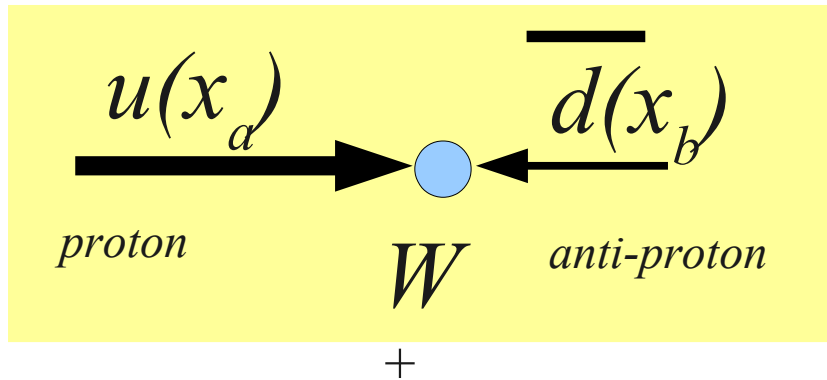
ATLAS W/Z cross section ratio in good agreement with NNLO predictions from the PDF groups shown

W

Asymmetry

@ LHC

A bit of calculation



With the previous approximation,

$$A \approx \frac{u(x_a)d(x_b) - d(x_a)u(x_b)}{u(x_a)d(x_b) + d(x_a)u(x_b)} = \frac{R_{du}(x_b) - R_{du}(x_a)}{R_{du}(x_b) + R_{du}(x_a)}$$

where $R_{du}(x) = \frac{d(x)}{u(x)}$

We can make Taylor expansions:

$$x_{1,2} = x_0 e^{\pm y} \simeq x_0 (1 \pm y)$$

$$R_{du}(x_{1,2}) \approx R_{du}(x_0) \pm y x_0 R'_{du}(\sqrt{\tau})$$

Thus, the asymmetry is:

$$A(y) = -y x_0 \frac{R'_{du}(x_0)}{R_{du}(x_0)}$$

W lepton charge asymmetry

- A first natural extension of the W inclusive studies is the study of the W^+/W^- ratio, R_W , as a function of different kinematic variables.
- Experimentally, a clean way to do this study is to measure the charge asymmetry as a function of the lepton pseudo-rapidity

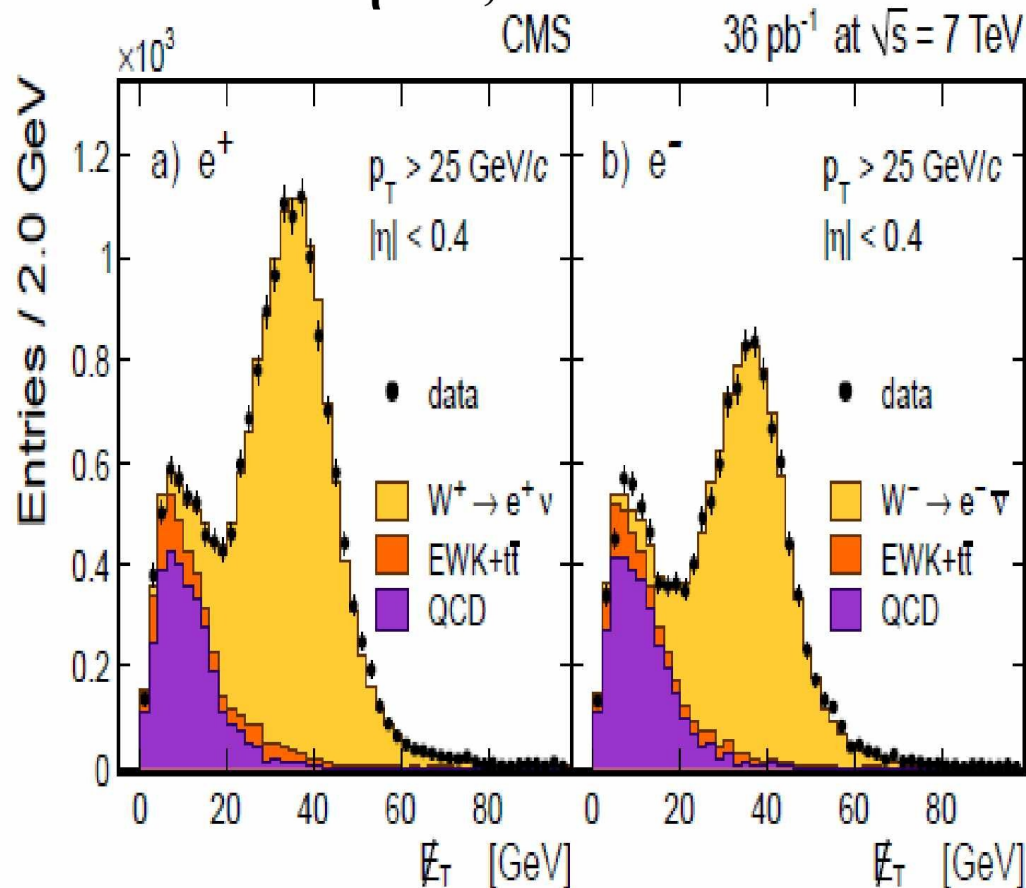
$$A(\eta) = \frac{\frac{d\sigma}{d\eta}(W^+ \rightarrow l^+ \nu) - \frac{d\sigma}{d\eta}(W^- \rightarrow l^- \bar{\nu})}{\frac{d\sigma}{d\eta}(W^+ \rightarrow l^+ \nu) + \frac{d\sigma}{d\eta}(W^- \rightarrow l^- \bar{\nu})}$$

$$\left(A(\eta) \equiv \frac{R_W(\eta) - 1}{R_W(\eta) + 1} \right)$$

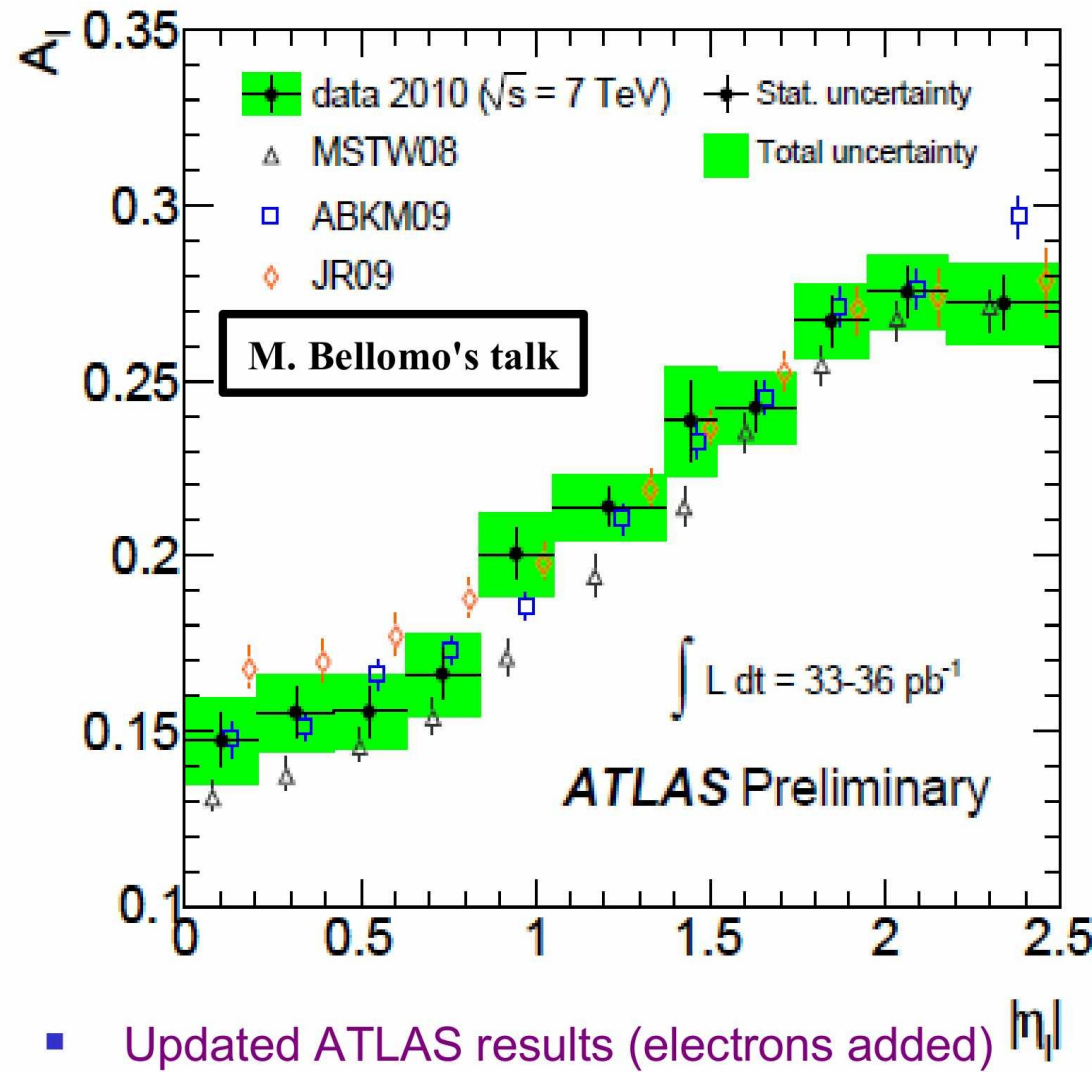
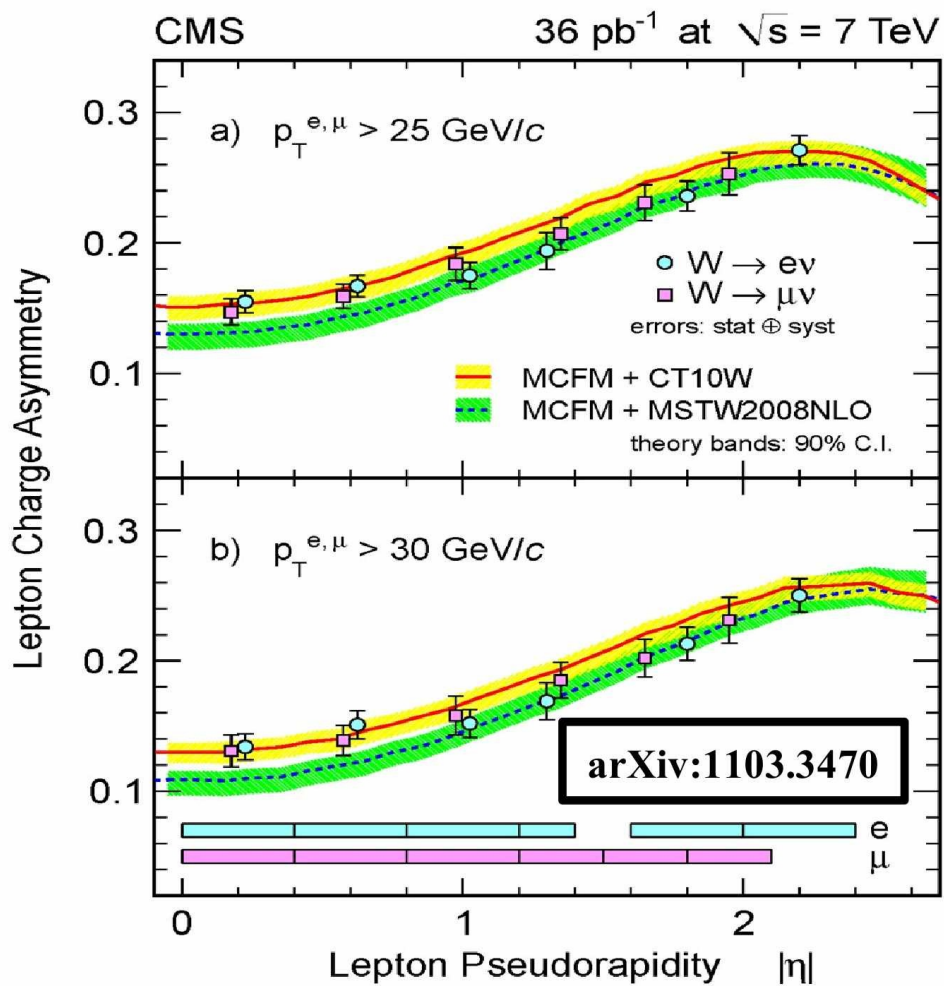
- This measurement is rather sensitive to PDFs because most systematic uncertainties cancel in the ratio
- Selections follow closely the criteria used in inclusive measurements.

arXiv:1103.3470

First η bin, electron channel



W lepton charge asymmetry

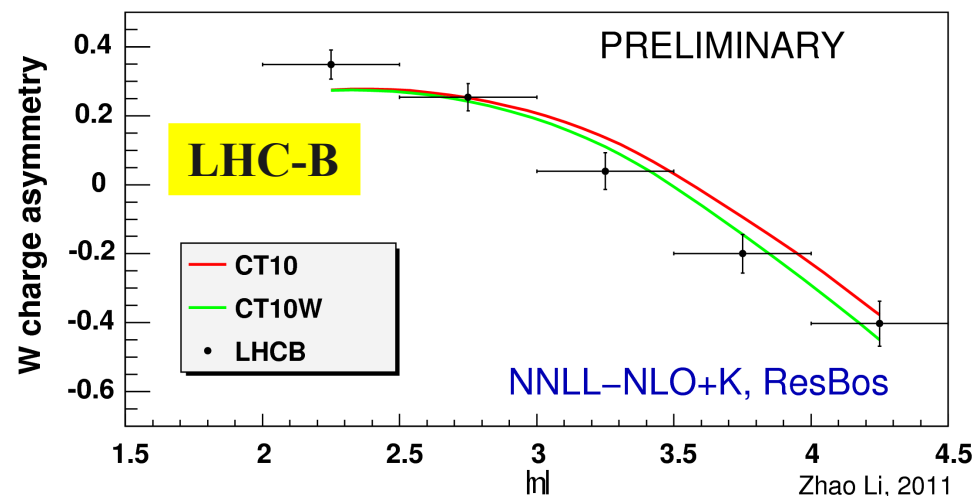
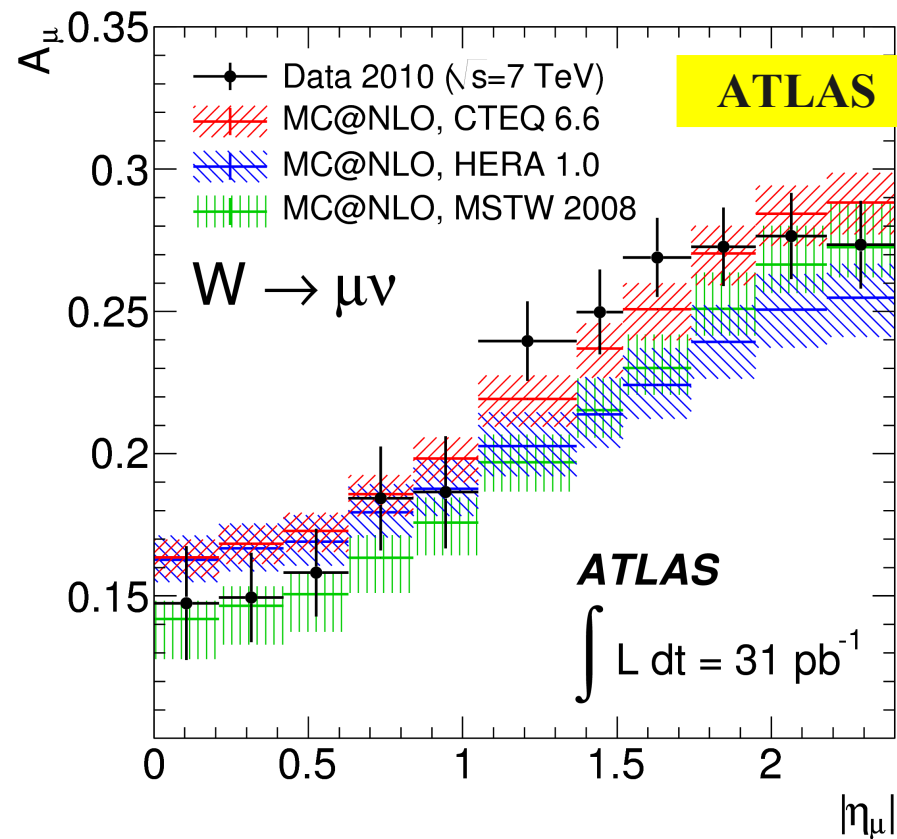
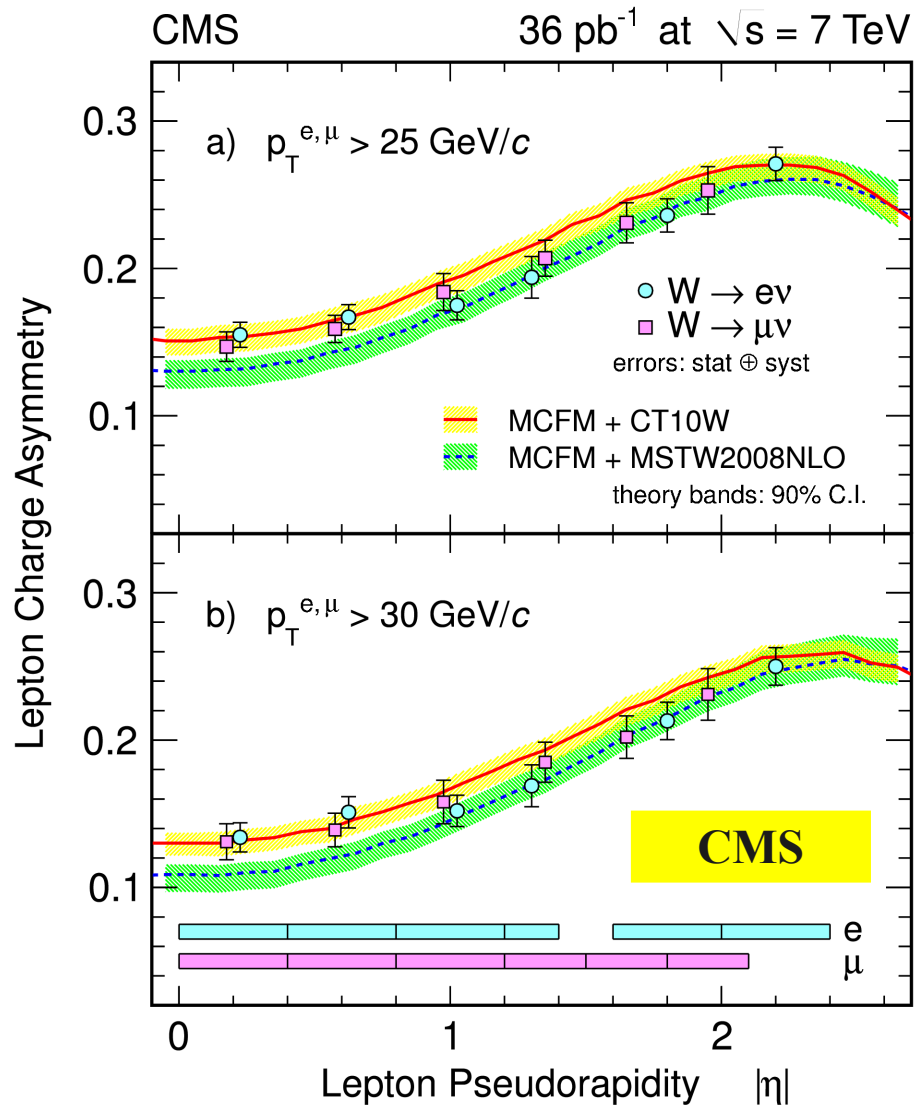


■ CMS published results

■ Updated ATLAS results (electrons added)

In reasonable agreement with different PDF predictions, but extremely sensitive to shape details

$$A_{\mu} = \frac{d\sigma(W_{\mu^+}) - d\sigma(W_{\mu^-})}{d\sigma(W_{\mu^+}) + d\sigma(W_{\mu^-})}$$



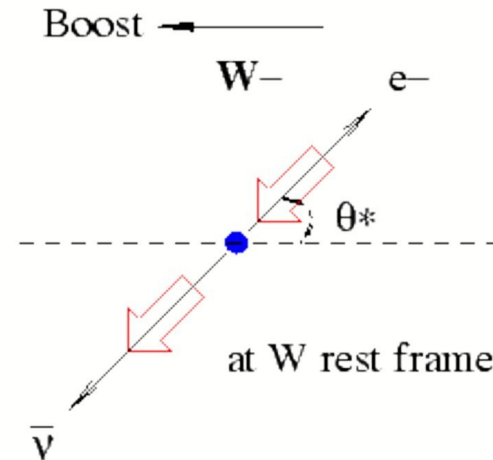
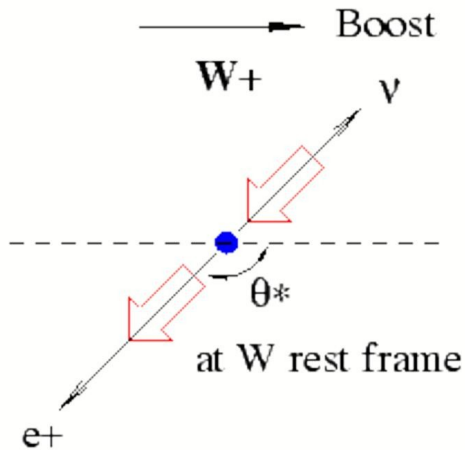
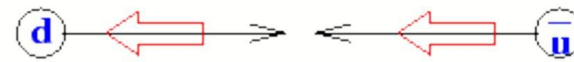
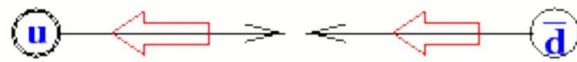
W

Asymmetry

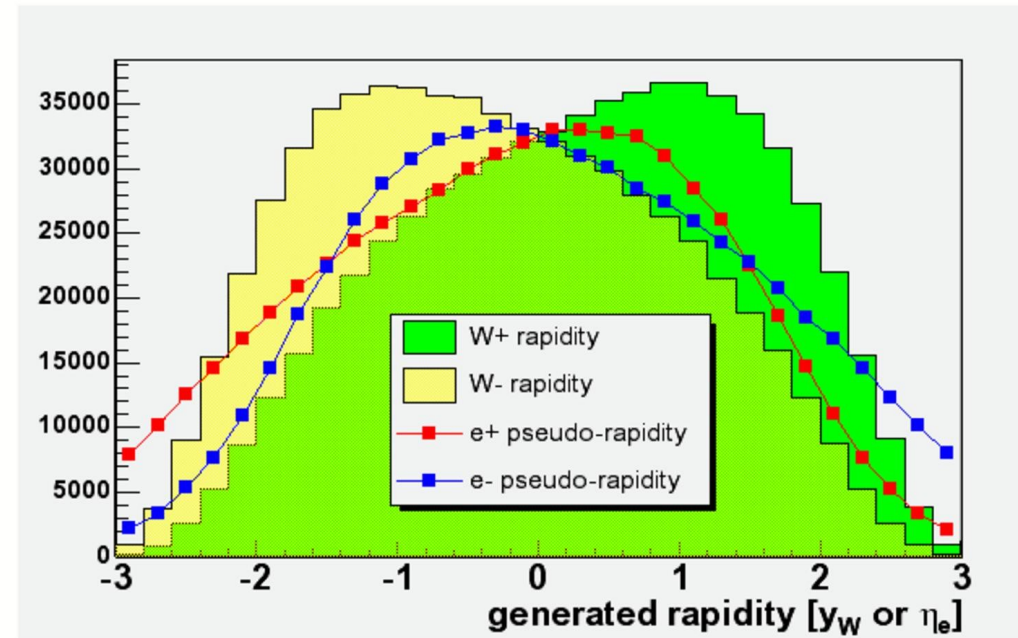
@ Tevatron

$$A_\ell = \frac{d\sigma(W^+ \rightarrow \ell^+) - d\sigma(W^- \rightarrow \ell^-)}{d\sigma(W^+ \rightarrow \ell^+) + d\sigma(W^- \rightarrow \ell^-)}$$

W asymmetry in p-pbar collisions

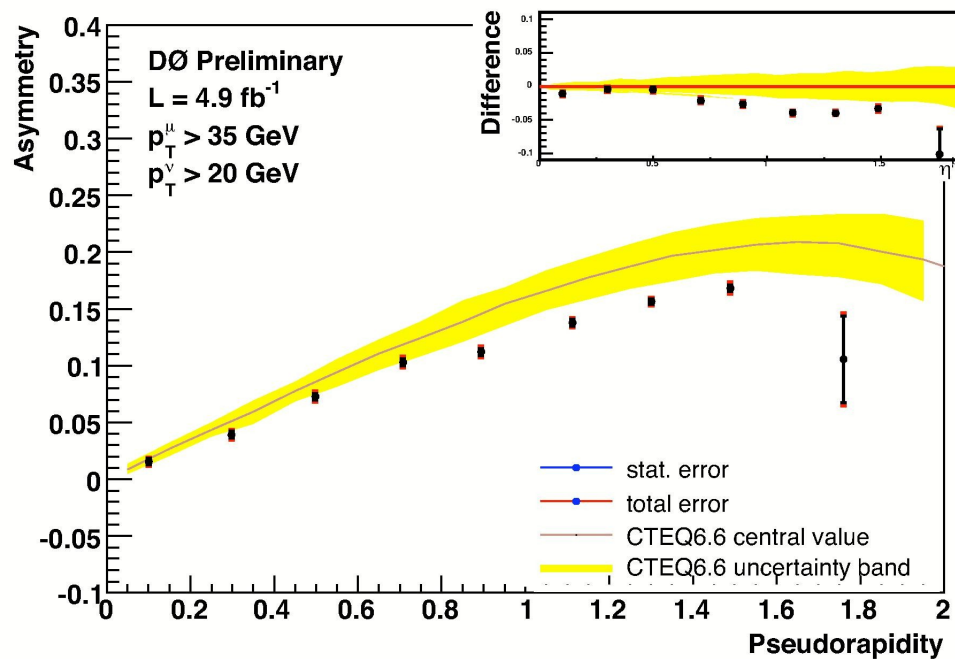
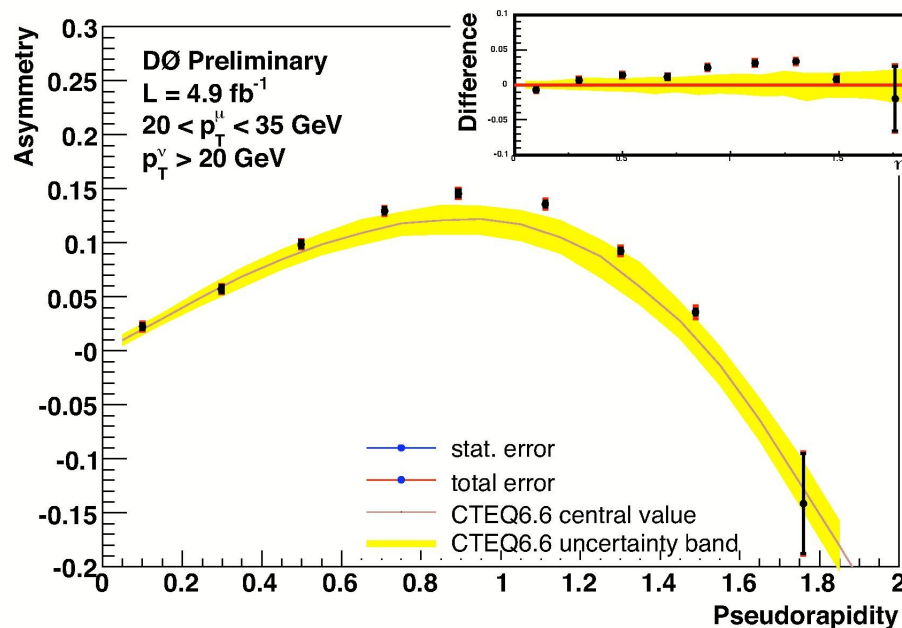
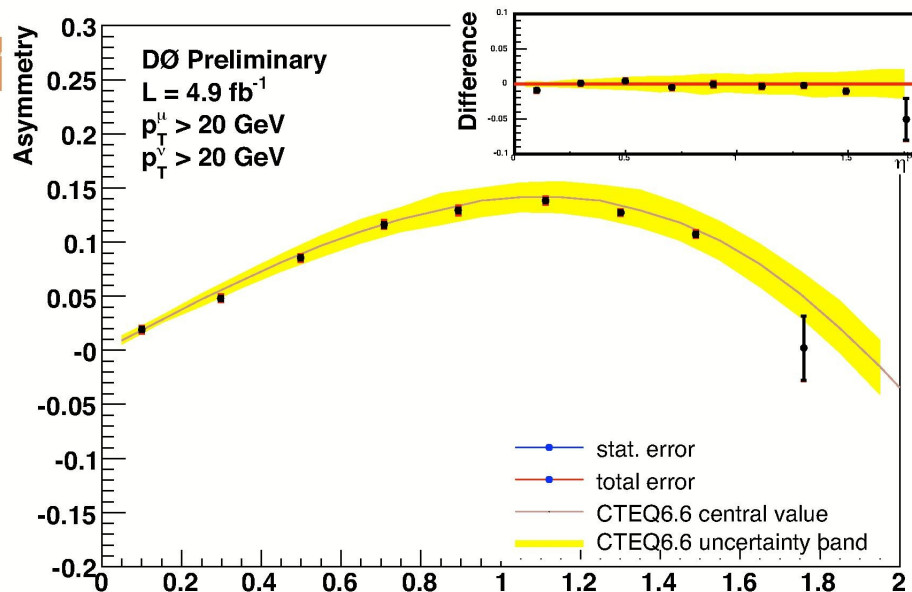


- If u quarks carry more momentum than d quarks, the W^+ will head in the proton direction preferentially.
- Unfortunately, the V-A interaction means that the charged lepton from W decay heads backwards in the W frame



Results compared to RESBOS+CTEQ6.6M

1



3 muon PT bins, $PT(\nu) > 20 \text{ GeV}$

Upper Left – $PT(\mu) > 20 \text{ GeV}$

Upper Right $20 \text{ GeV} < PT(\mu) < 35 \text{ GeV}$

Lower Left $PT(\mu) > 35 \text{ GeV}$

The puzzle of the CDF/D0 W lepton asymmetry

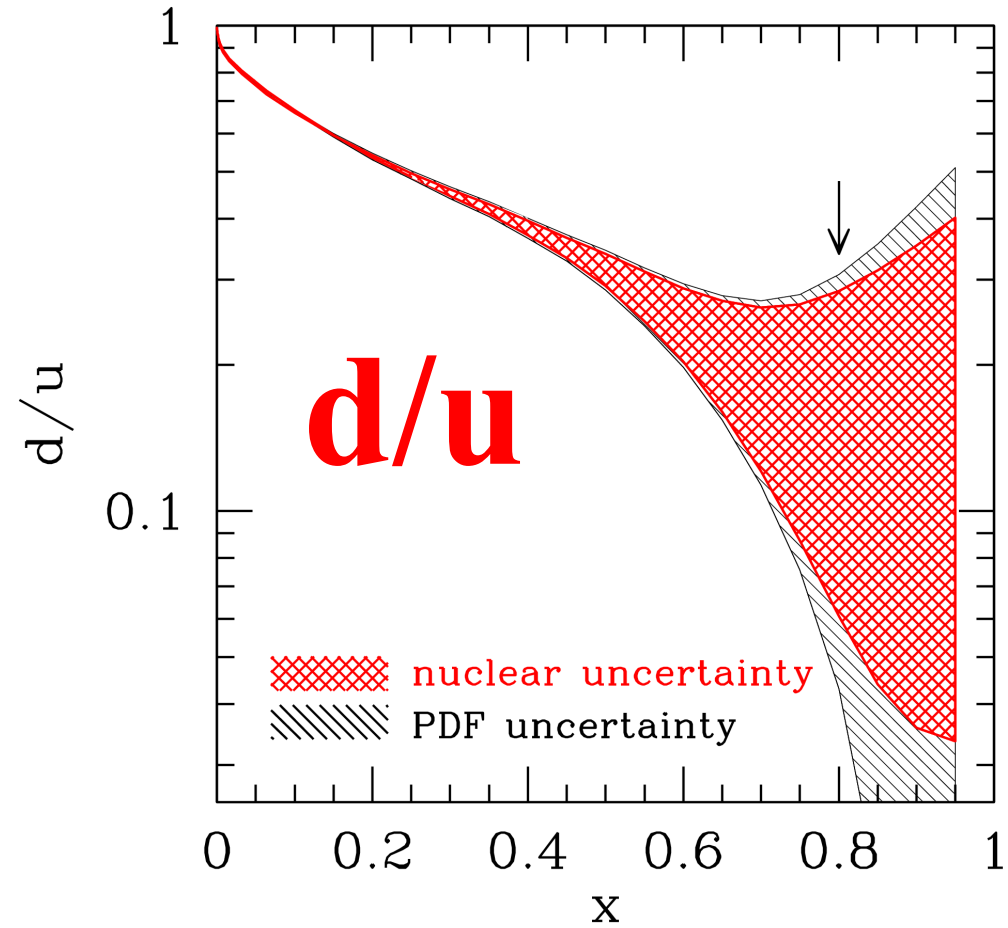
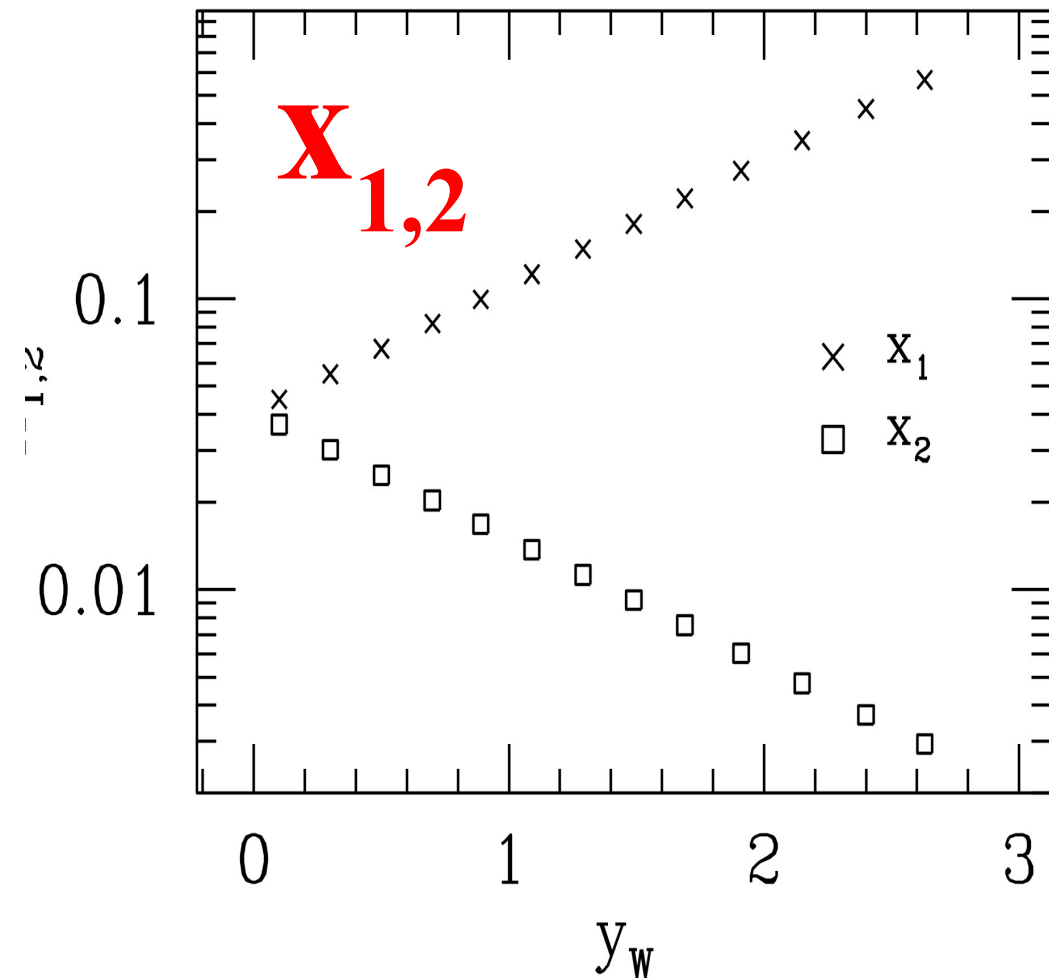
- CT10W set reasonably agrees with 3 $p_{T\ell}$ bins of $A_e(y_e)$ and one bin of $A_\mu(y_\mu)$ from D0 Run-2 (2008).
- NNPDF 2.0 (*arXiv:1012.0836*) agrees with $A_\mu(y_\mu)$, disagrees with two p_{T_e} bins of $A_e(y_e)$.
- CT10, many other PDFs fail.

Agreement of PQCD with D0 $A_e(y_e)$	χ^2/n_{pt}	Source or comments
CTEQ6.6, NLO	191/36=5.5	<i>Our study;</i> <i>Resbos, NNLL-NLO</i>
CT10W, NLO	78/36=2.2 With $A_\mu(y_\mu)$: 88/47=1.9	
ABKM'09, NNLO	540/24=22.5	<i>Catani, Ferrera, Grazzini,</i> <i>JHEP 05, 006 (2010)</i>
MSTW'08, NNLO	205/24=8.6	
JR09VF, NNLO	113/24=4.7	

What is
happening
with d/u

This combination rather unique

$$A_\ell = \frac{d\sigma(W^+ \rightarrow \ell^+) - d\sigma(W^- \rightarrow \ell^-)}{d\sigma(W^+ \rightarrow \ell^+) + d\sigma(W^- \rightarrow \ell^-)}$$



Determined from DIS and DY on p and d

$$x_{1,2} \sim \frac{M}{\sqrt{s}} e^{\pm y}$$

LHC values scaled appropriately

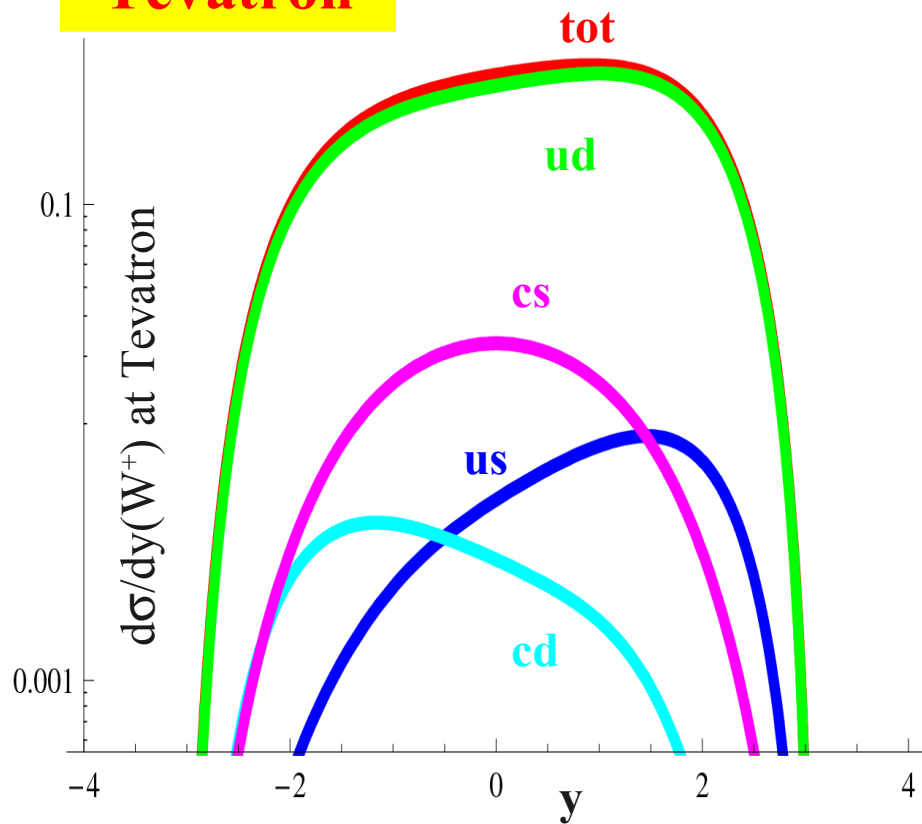
W/Z

Production

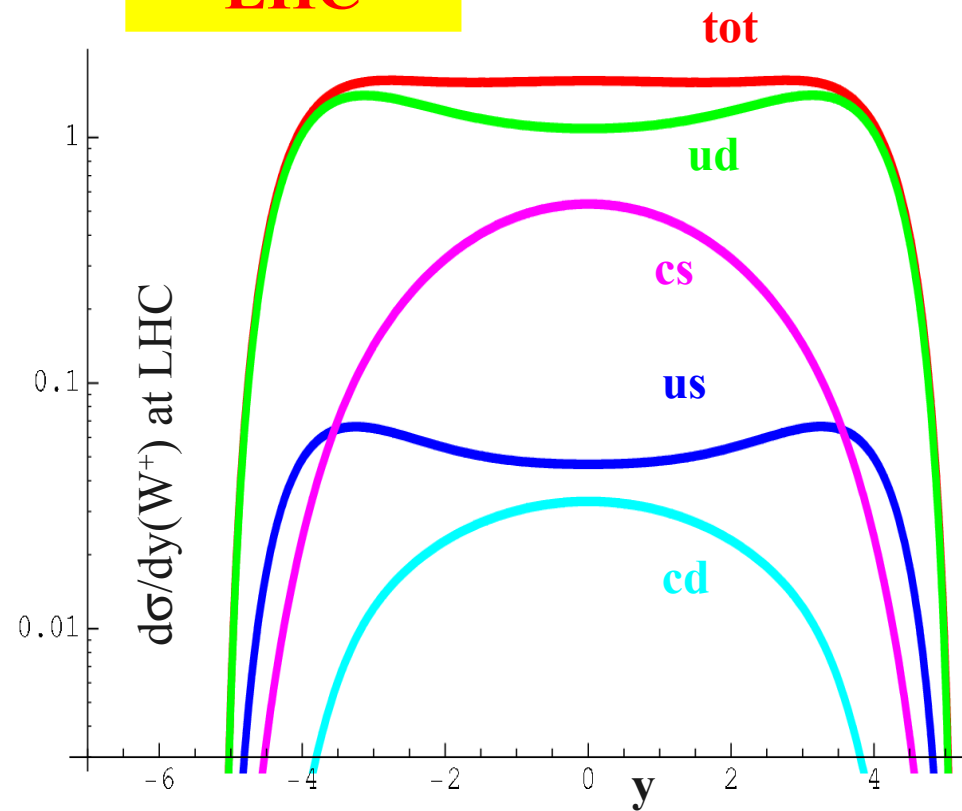
On the theoretical side ...

Heavy quark PDFs are essential ingredient

Tevatron

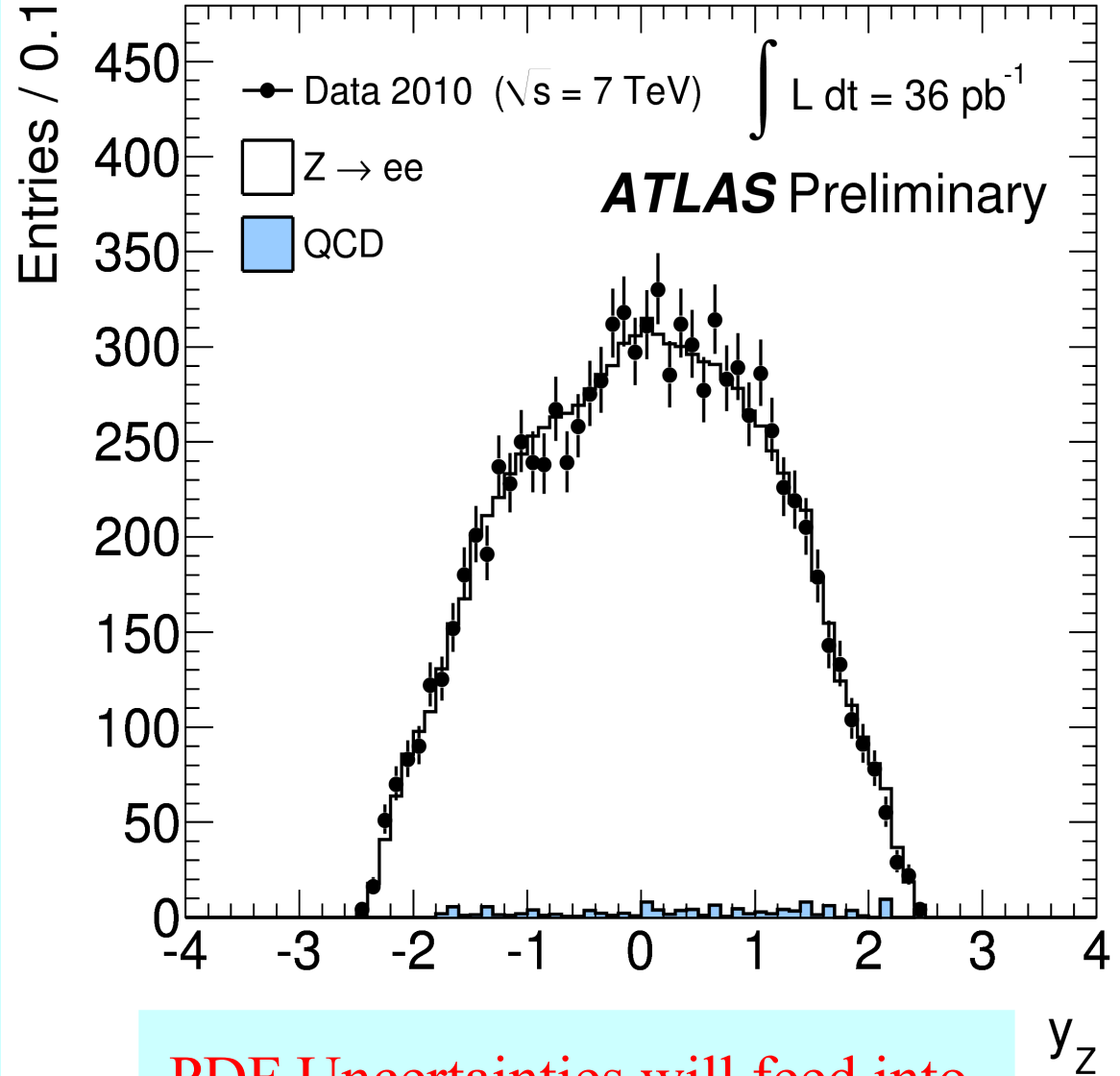
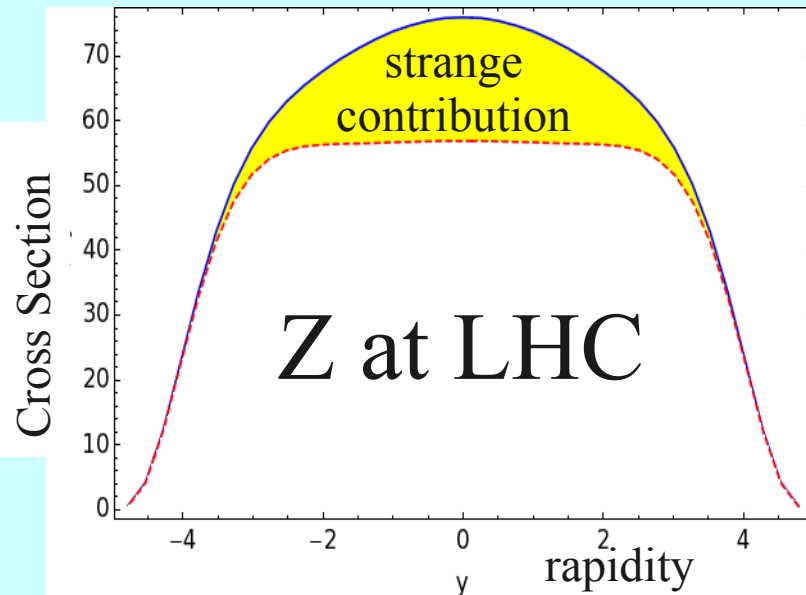
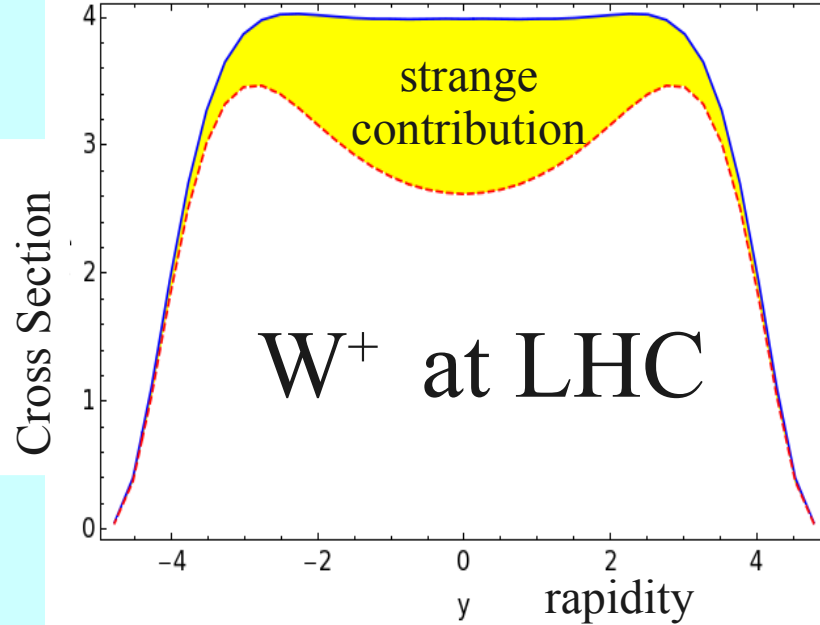


LHC



Heavy Quark components play an increasingly important role at the LHC

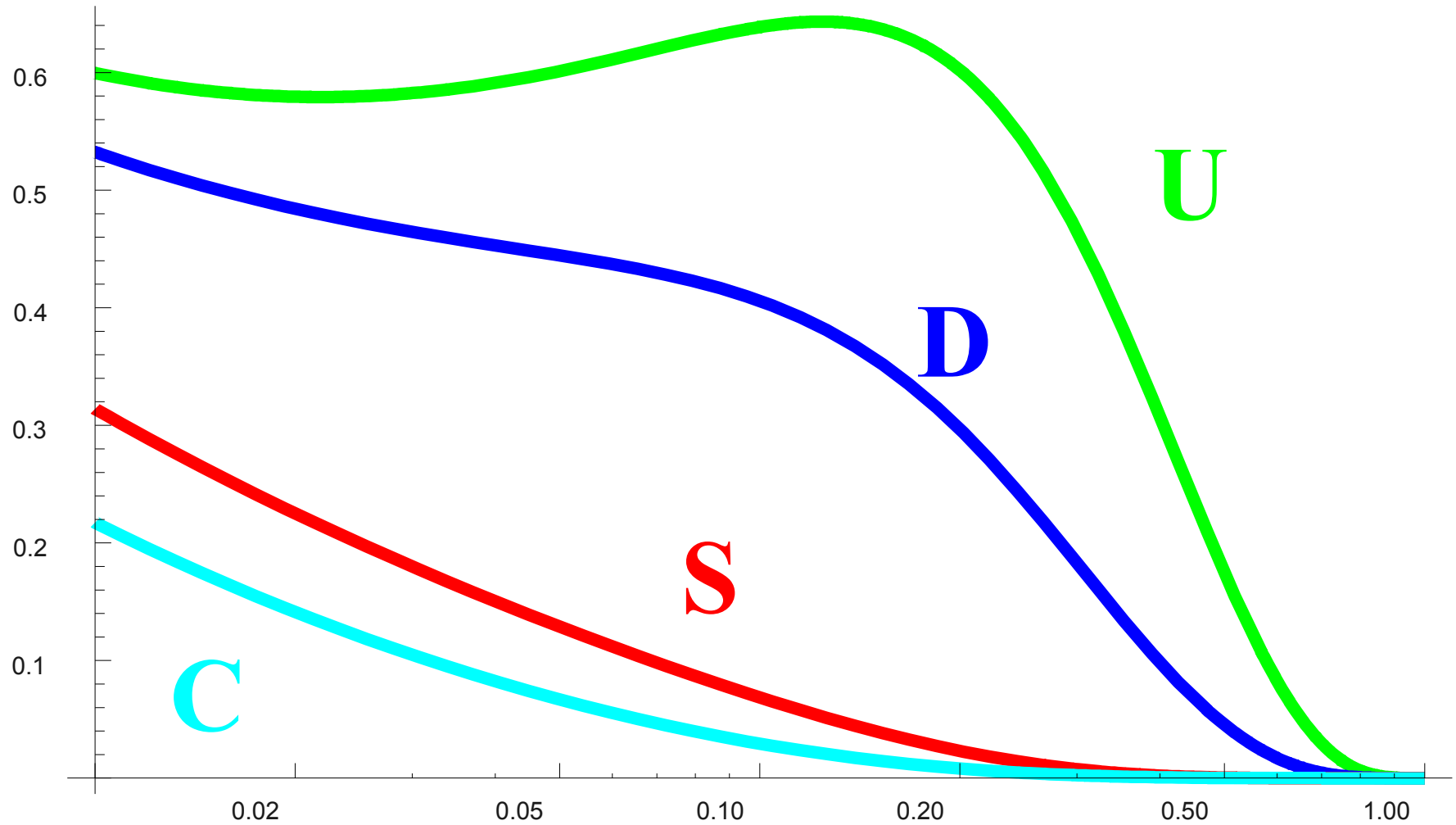
PDF Uncertainties \Rightarrow $S(x)$ PDF \Rightarrow W/Z at LHC

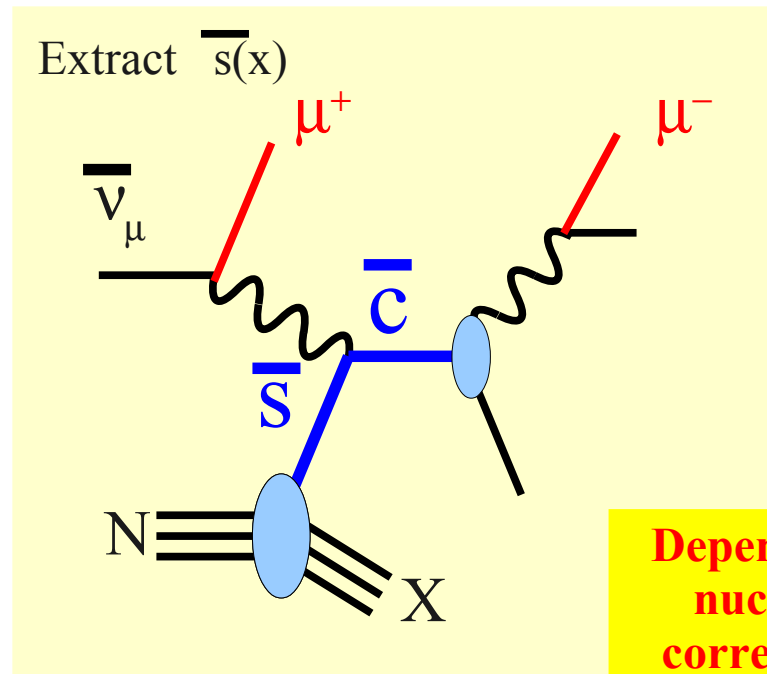
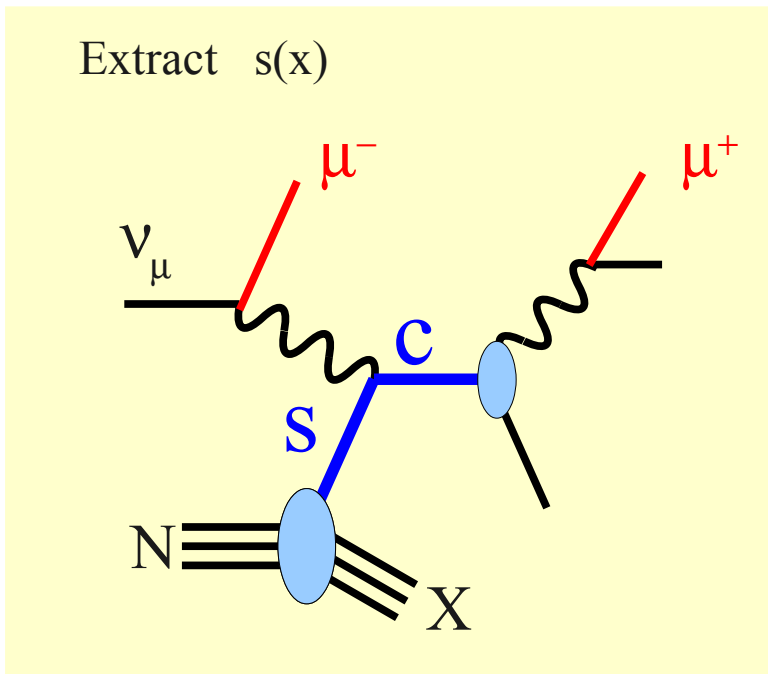


PDF Uncertainties will feed into
LHC "Benchmark" processes

Comparison with new NNPDF sets: Les Houches 2009

What constrains the Strange???



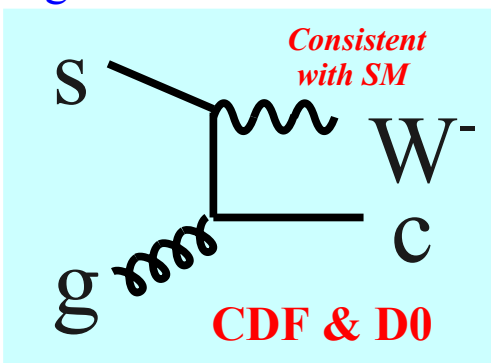


Depends on nuclear corrections

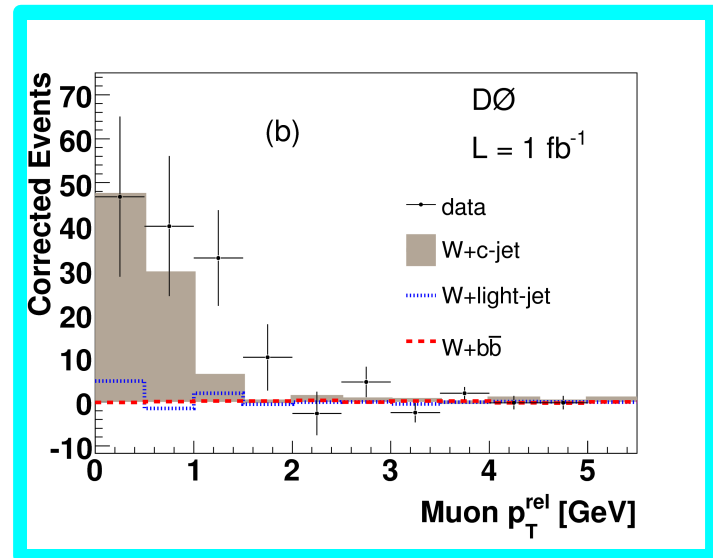
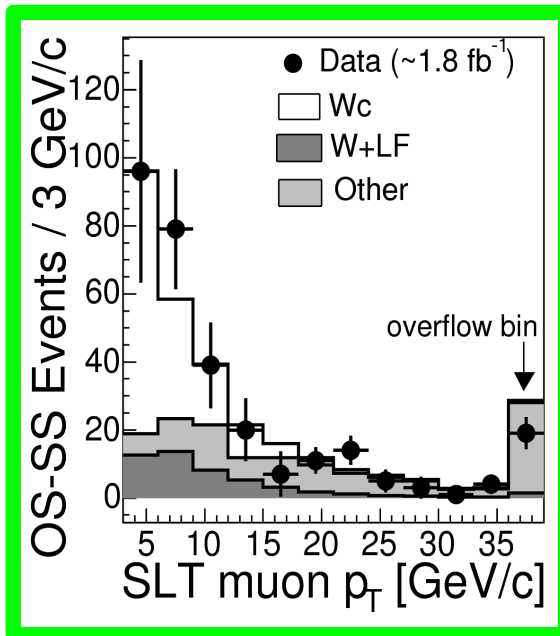
$s(x)$ and $\bar{s}(x)$ are essential in extraction of $\text{Sin}\theta_W$

Used in CTEQ6 Fits

$s g \rightarrow Wc$ at the Tevatron



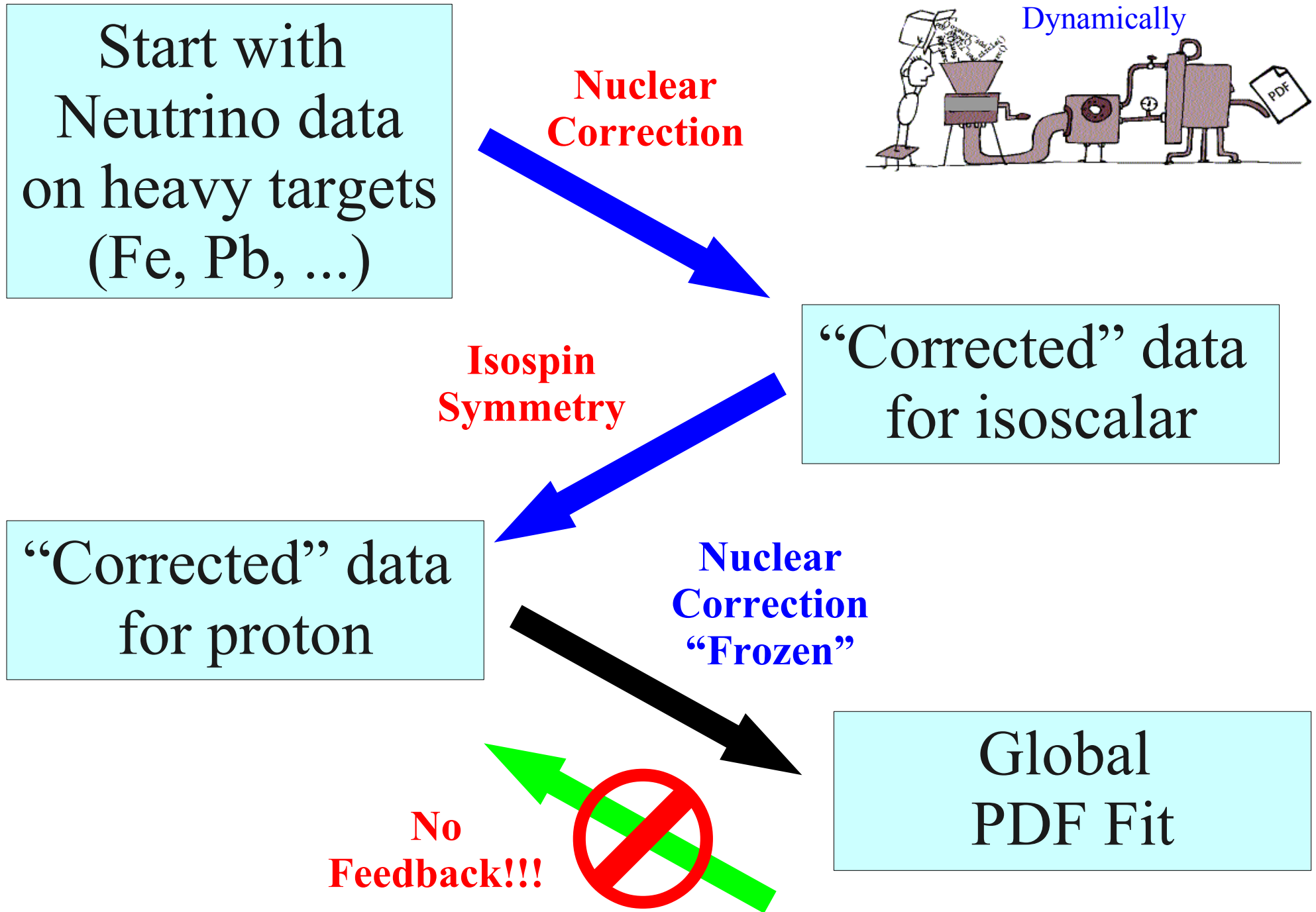
CDF: PRL 100:091803,2008.
D0: PLB666:23,2008.



Also a challenge at LHC

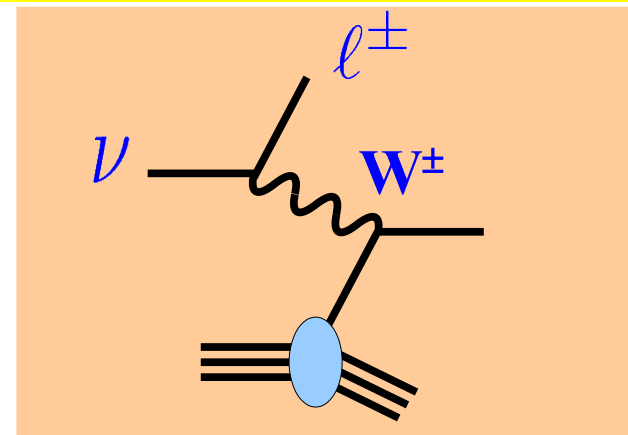
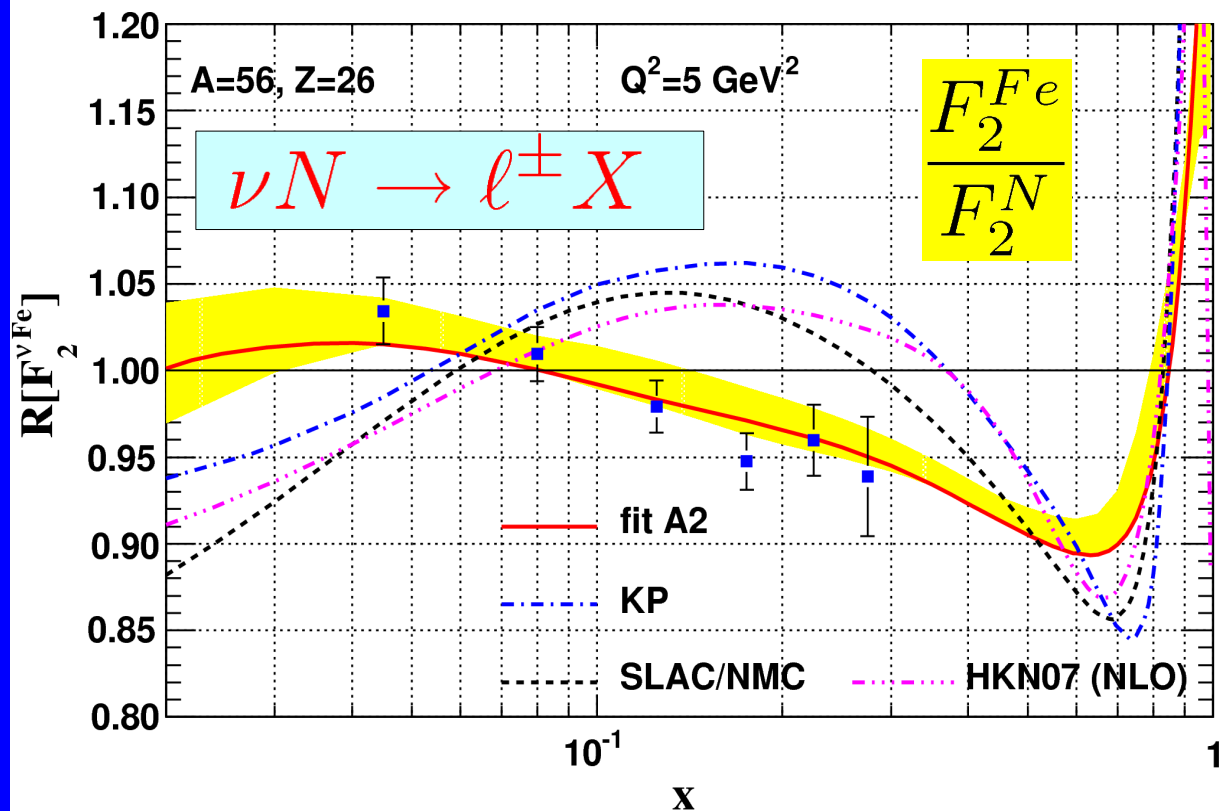
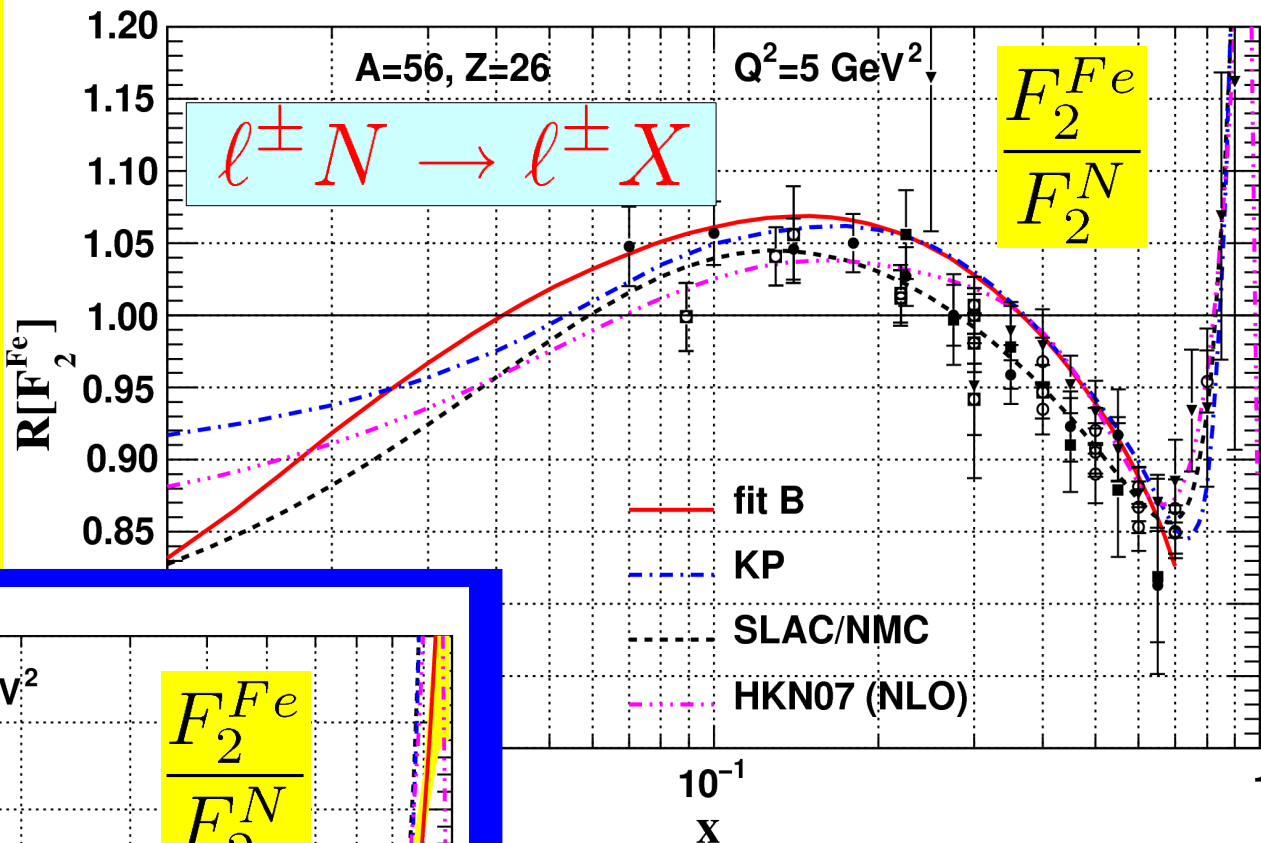
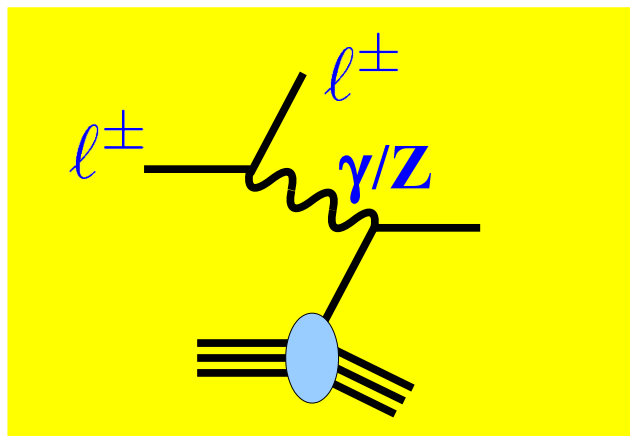
Nuclear Corrections

???



Oooooops!

Charged Lepton DIS \Rightarrow



\Leftarrow **Neutrino DIS**

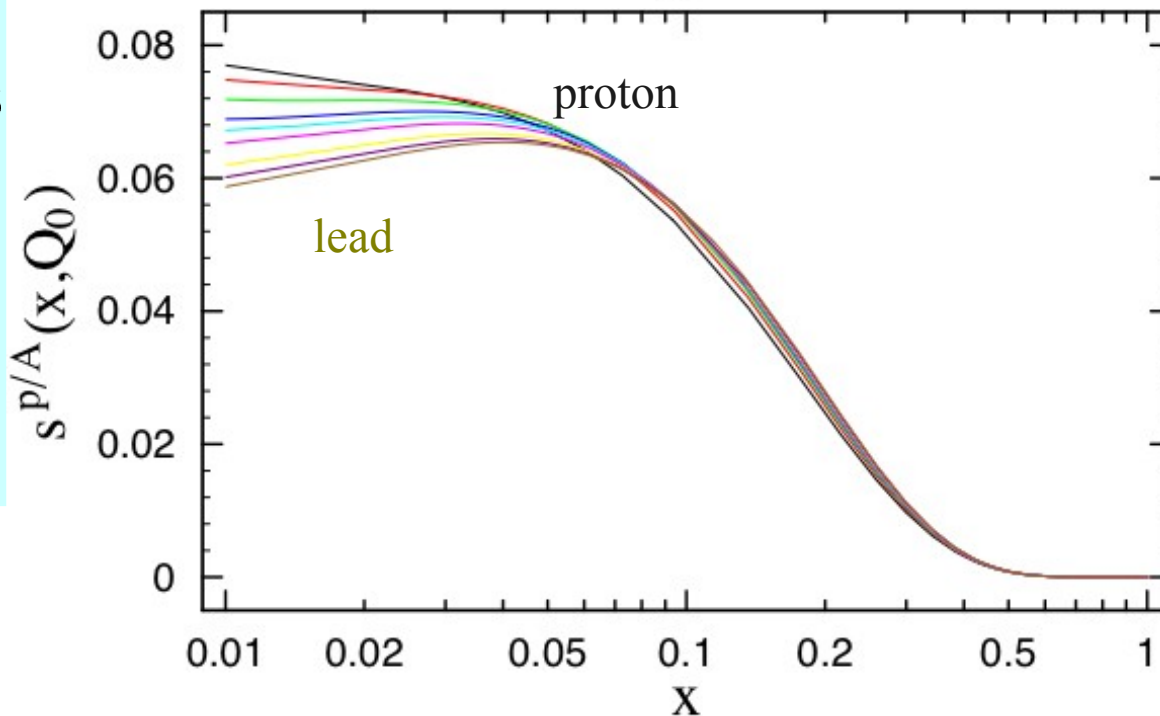
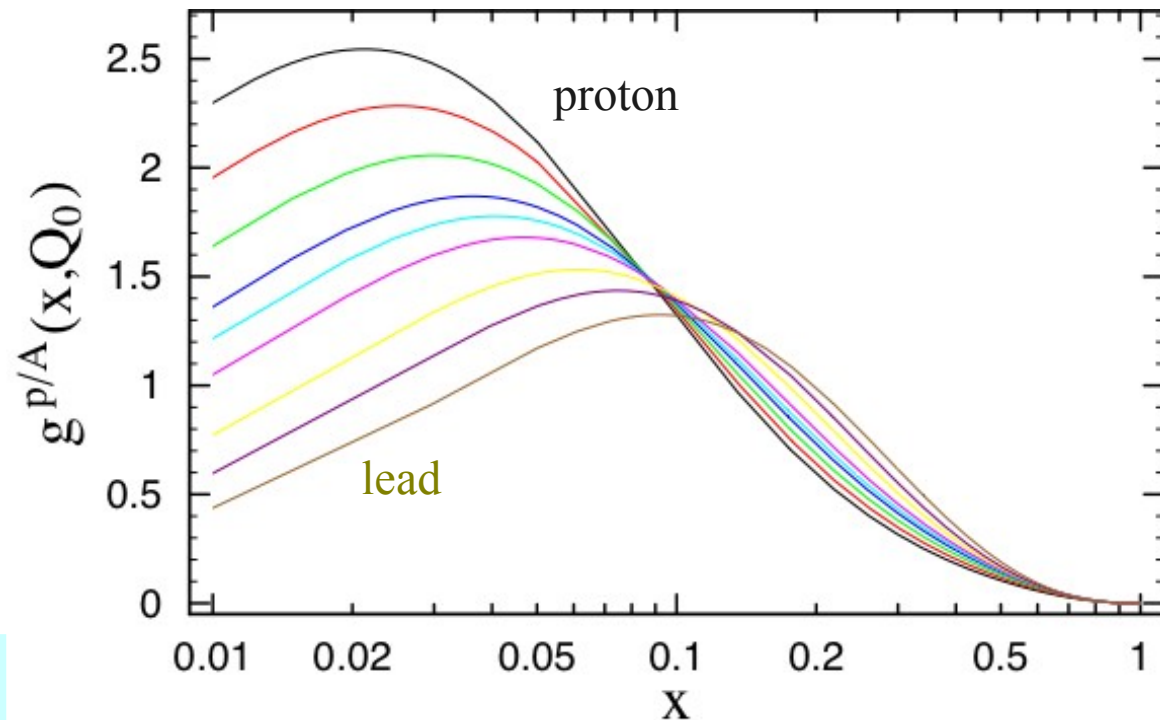
- ✓ CTEQ style global fit extended
handle various nuclear targets
- ✓ CTEQ Data + nuclear DIS & DY
[~15 targets; ~2000+ data]
- ✓ A-dependence modeled;
NLO fits work well

A-Dependent PDFs

$$xf(x) = x^{a_1} (1-x)^{a_2} e^{a_3 x} (1 + e^{a_4 x})^{a_5}$$

$$a_i \rightarrow a_i(A)$$

$$a_k = a_{k,0} + a_{k,1} (1 - A^{-a_{k,2}})$$



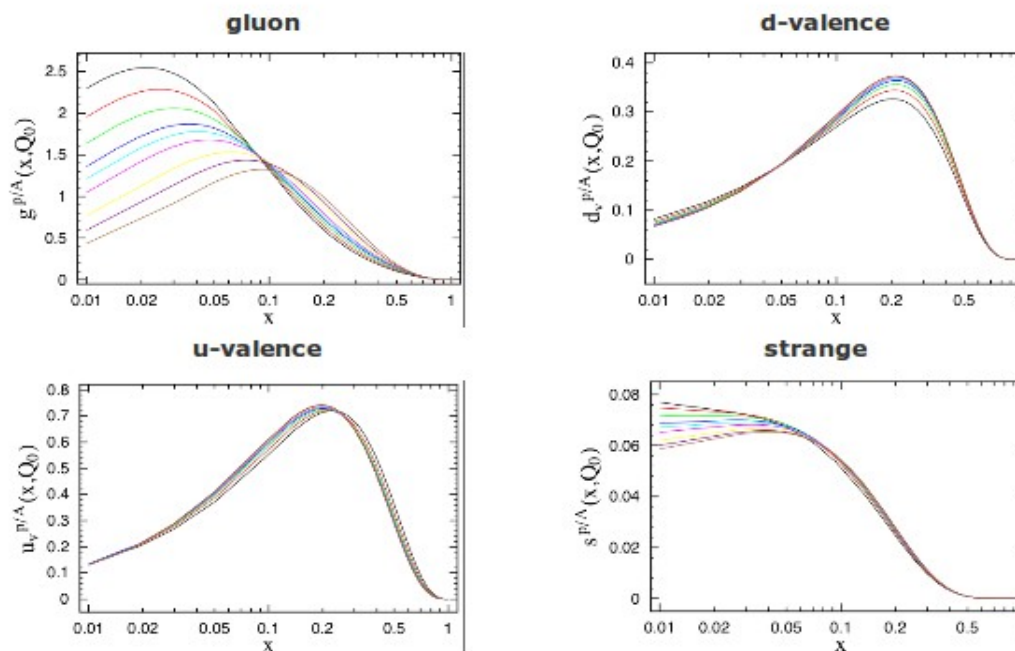
nCTEQ

nuclear parton distribution functions

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- Papers & Talks
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- Tracker
- Wiki

nCTEQ project is an extension of the CTEQ collaborative effort to determine parton distribution functions inside of a free proton. It generalizes the free-proton PDF framework to determine densities of partons in bound protons (hence nCTEQ which stands for nuclear CTEQ). More details on the framework and the first results can be found in [arXiv:09072357 \[hep-ph\]](https://arxiv.org/abs/09072357).

The effects of the nuclear environment on the parton densities can be shown as modified parton densities



where all black curves stand for free proton PDF and red, green, blue, cyan, pink, yellow, magenta and brown curves show PDF in protons bound in nuclei - from deuterium (red) to lead (brown).

K Kovarik,
I. Schienbein,
J.Y. Yu,
T. Stavreva,
T Jezo,
C. Keppel,
J.G. Morfin,
F. Olness,
J.F. Owens.

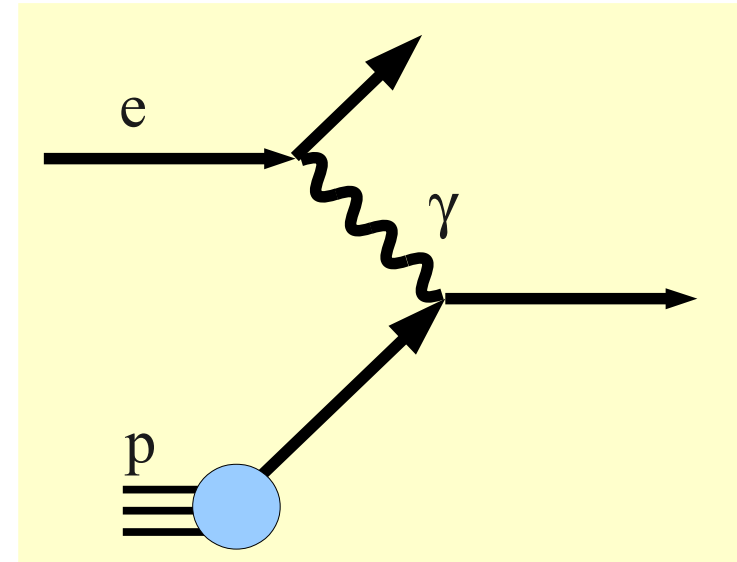
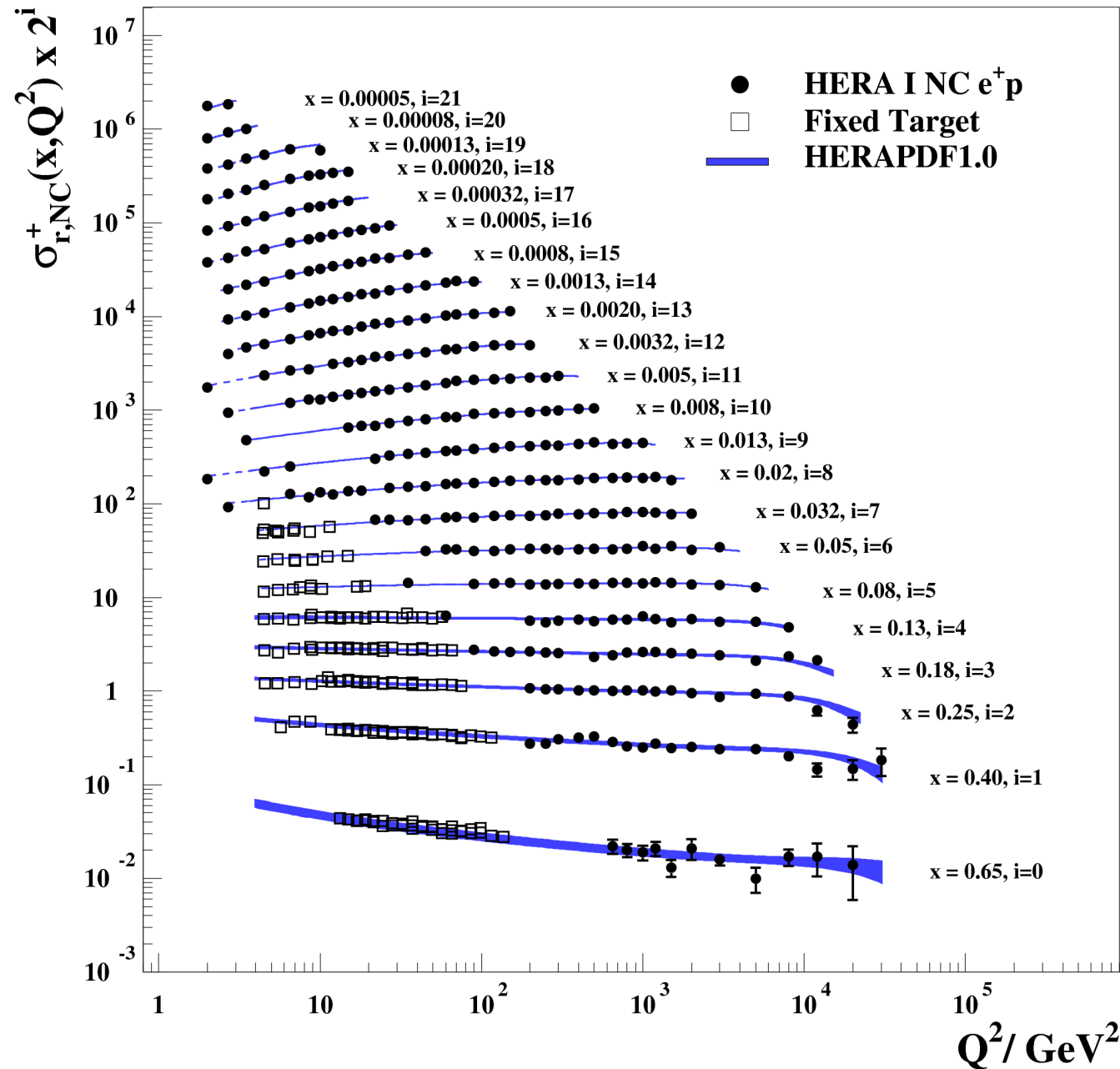
... what about the

Heavy Quarks

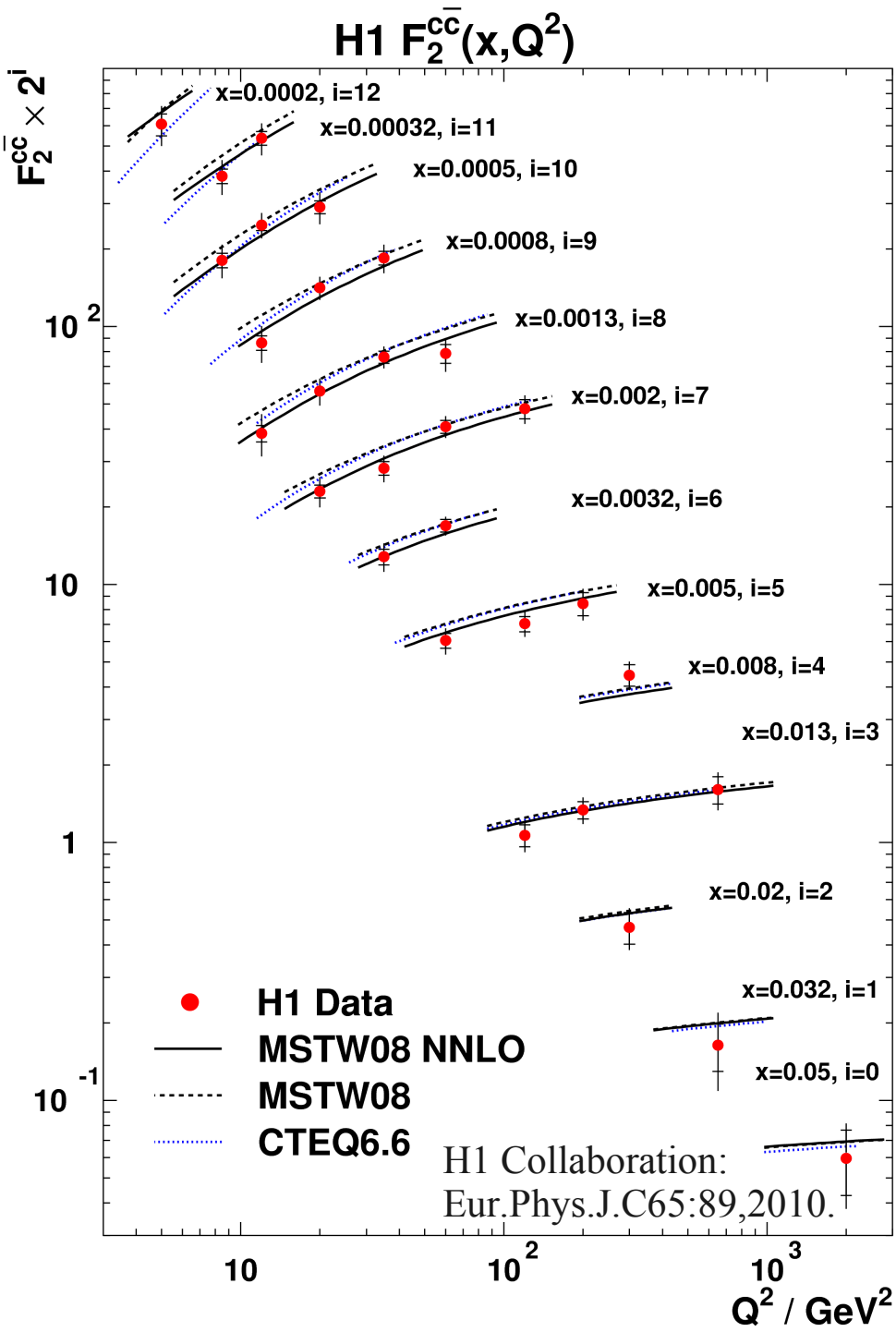
c & b

Extrinsic & Intrinsic

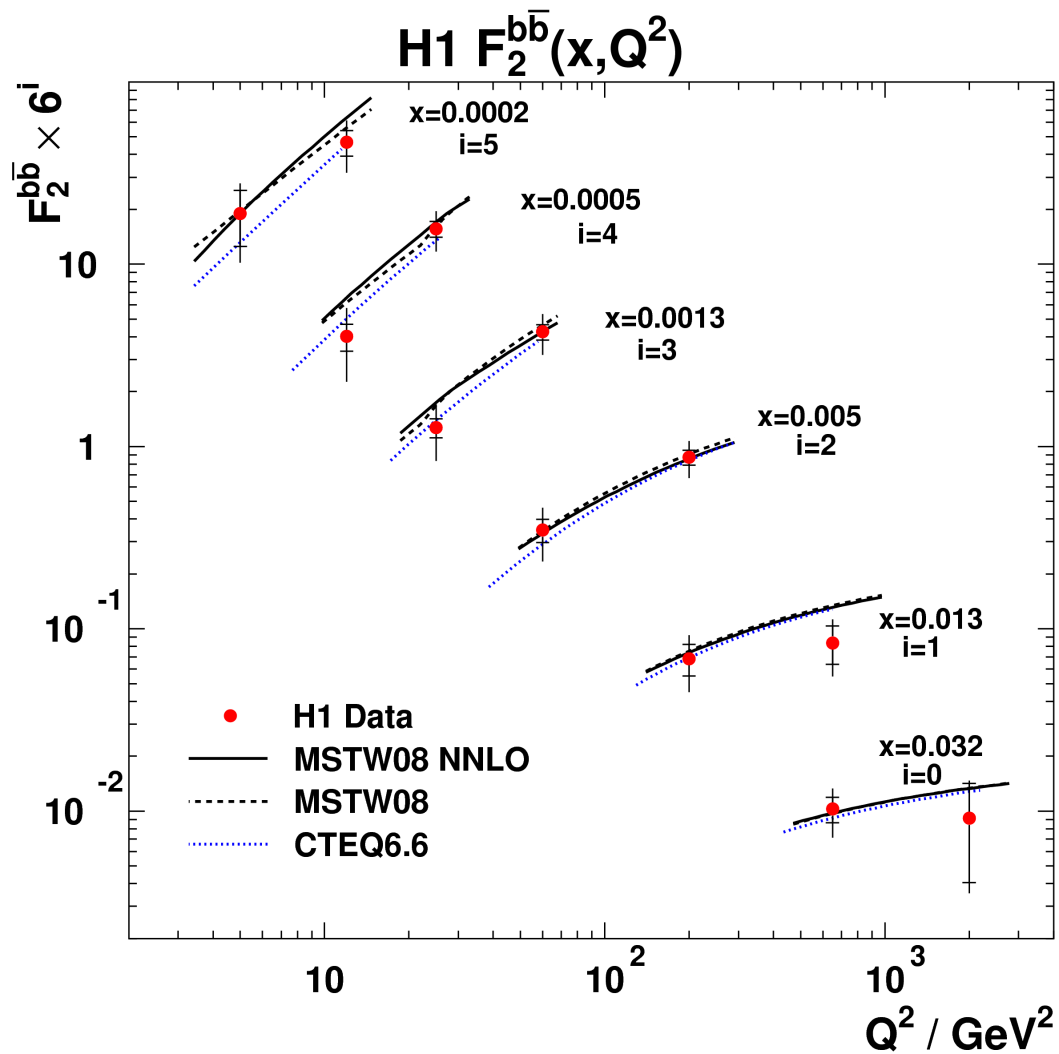
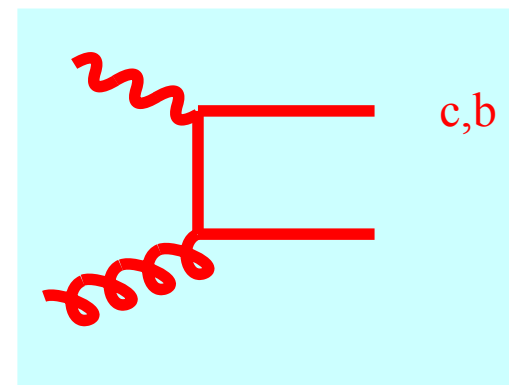
H1 and ZEUS

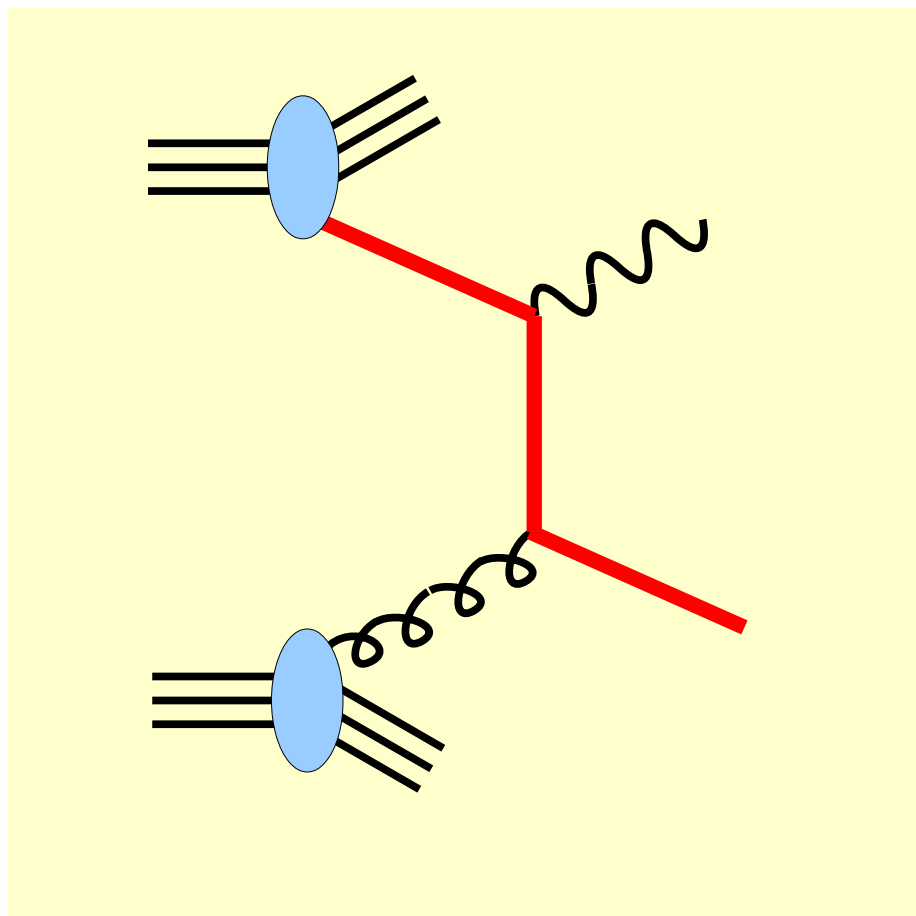


$$F_2^{ep} \sim \frac{4}{9} (u + \bar{u} + c + \bar{c}) + \frac{1}{9} (d + \bar{d} + s + \bar{s})$$



**c & b
tied to
gluon PDFs**





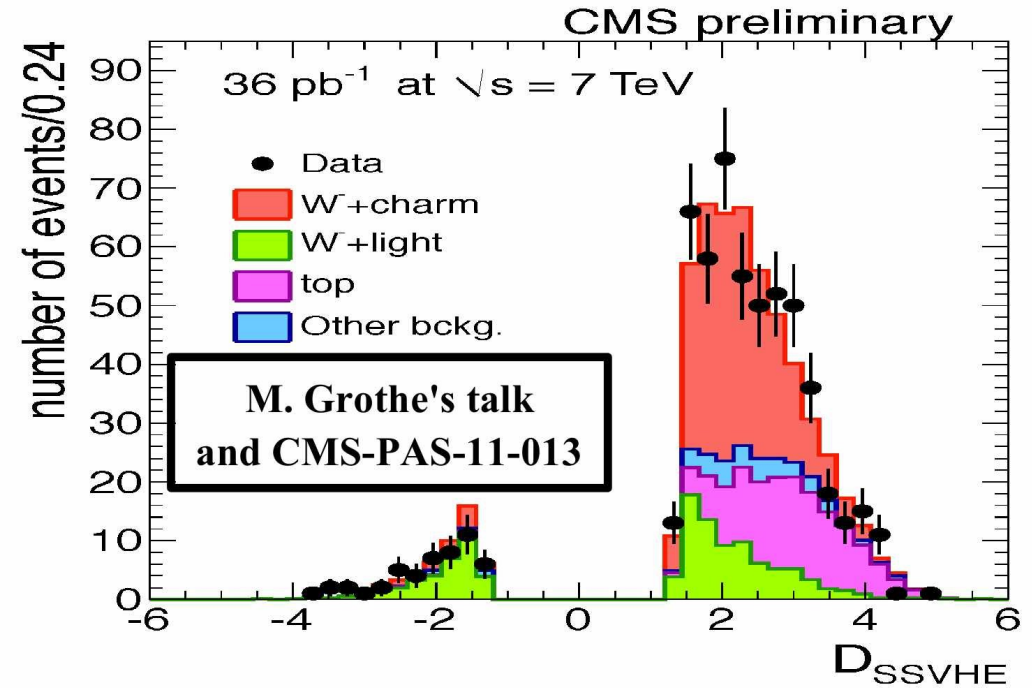
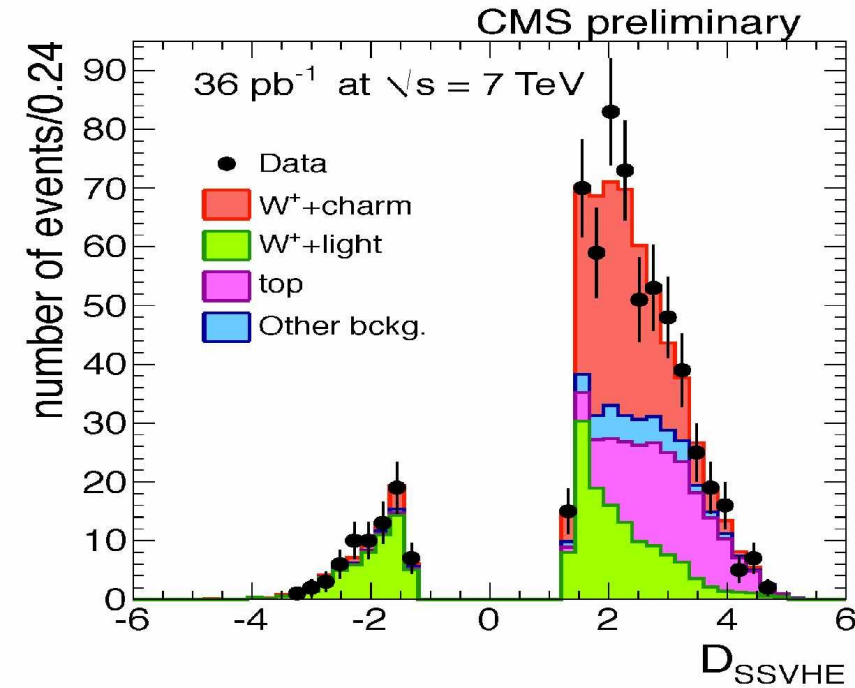
$$c \quad g \rightarrow c \quad \gamma$$

$$b \quad g \rightarrow b \quad \gamma$$

$$s \quad g \rightarrow c \quad W$$

$$c \quad g \rightarrow b \quad W$$

First LHC results on W+charm (CMS)



- Sensitive to strange quark PDFs (process dominated by $s+g \rightarrow W + \text{charm}$):

- PDF uncertainties from the second quark generation are a potential source of uncertainty for the W mass measurement at the LHC
- Data-driven control of light-quark and top backgrounds
- Enormous margin for improvement (only 2010 statistics used), new method (secondary vertex tagging), complementary to the one employed until now at Tevatron (semileptonic charm decay tagging):

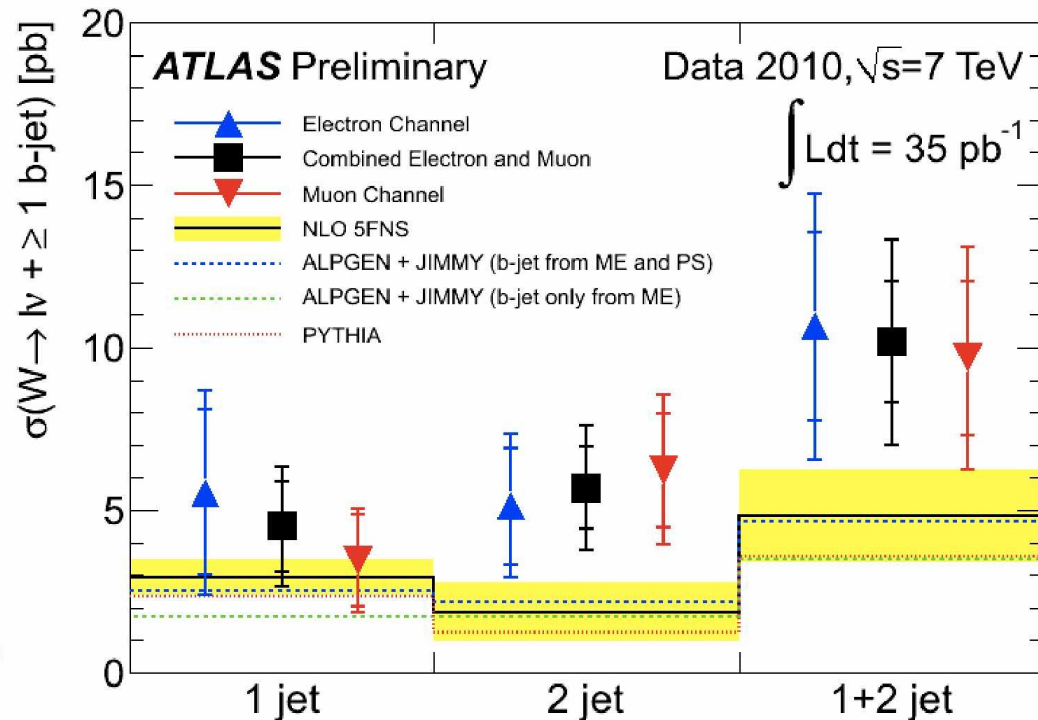
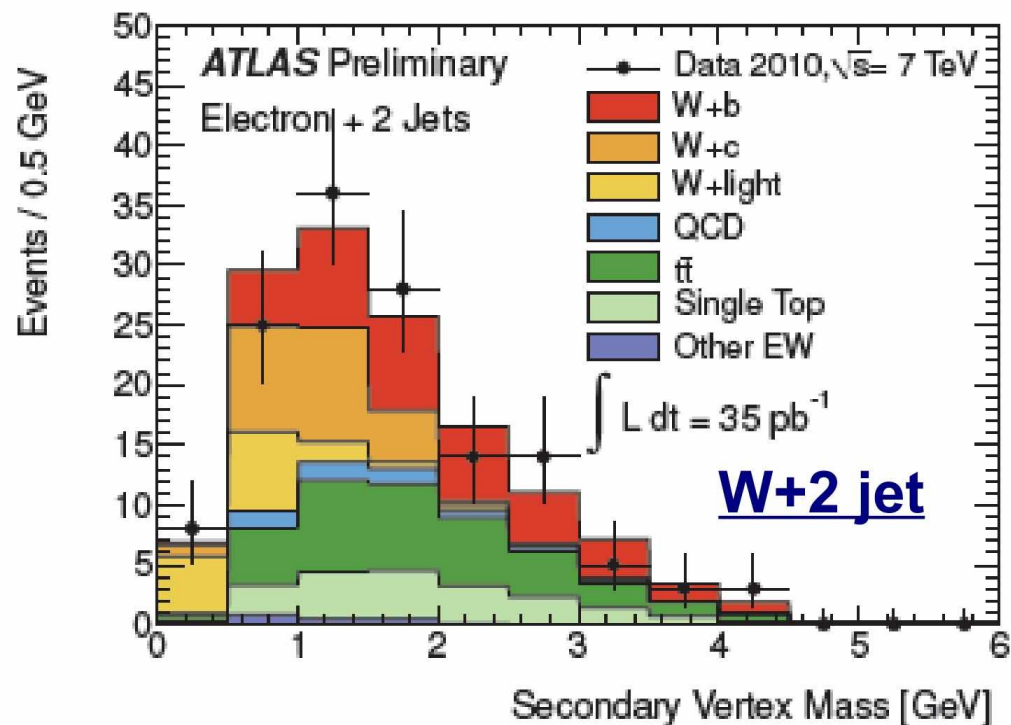
For $p_T^{jet} > 20$ GeV, $|\eta^{jet}| < 2.1$:

$$\frac{\sigma(W^+ + \text{charm})}{\sigma(W^- + \text{charm})} = 0.92 \pm 0.19(\text{stat.}) \pm 0.04(\text{syst.}); \quad \frac{\sigma(W + \text{charm})}{\sigma(W + \text{jets})} = 0.142 \pm 0.015(\text{stat.}) \pm 0.024(\text{syst.})$$

First LHC W+b results (ATLAS)

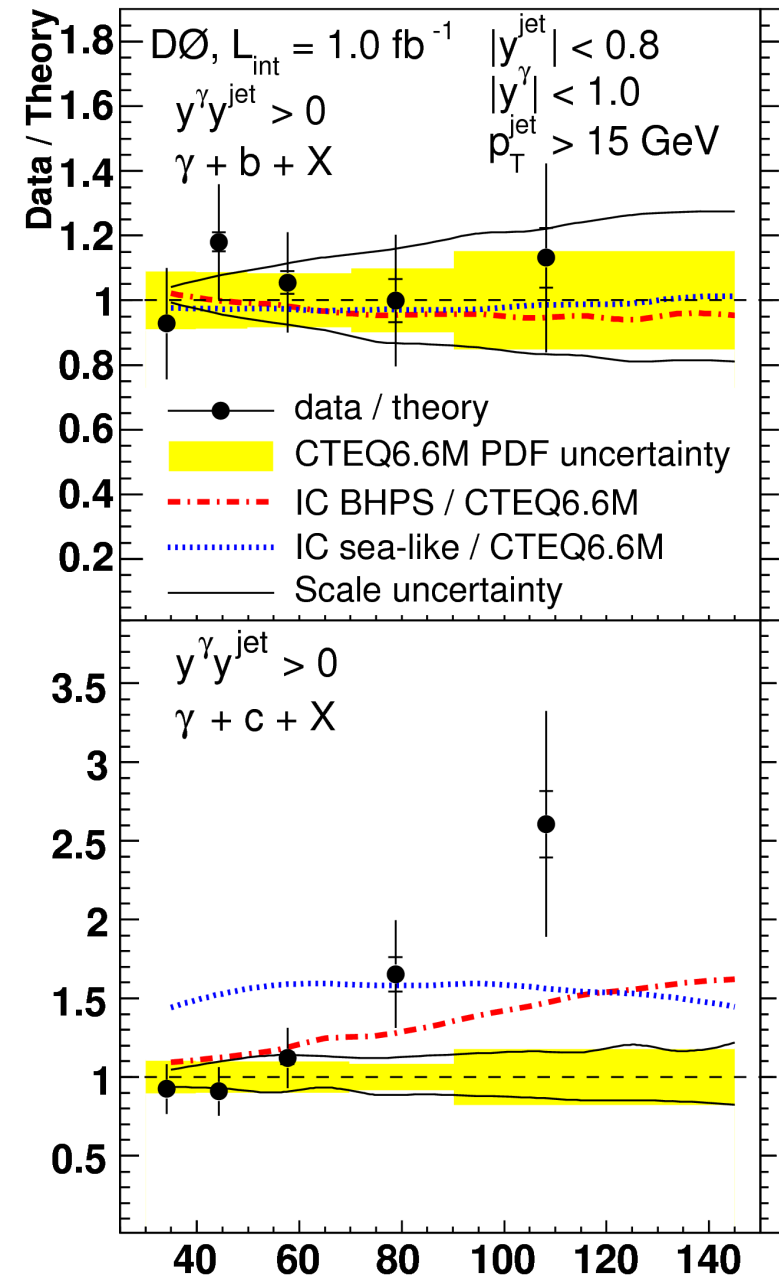
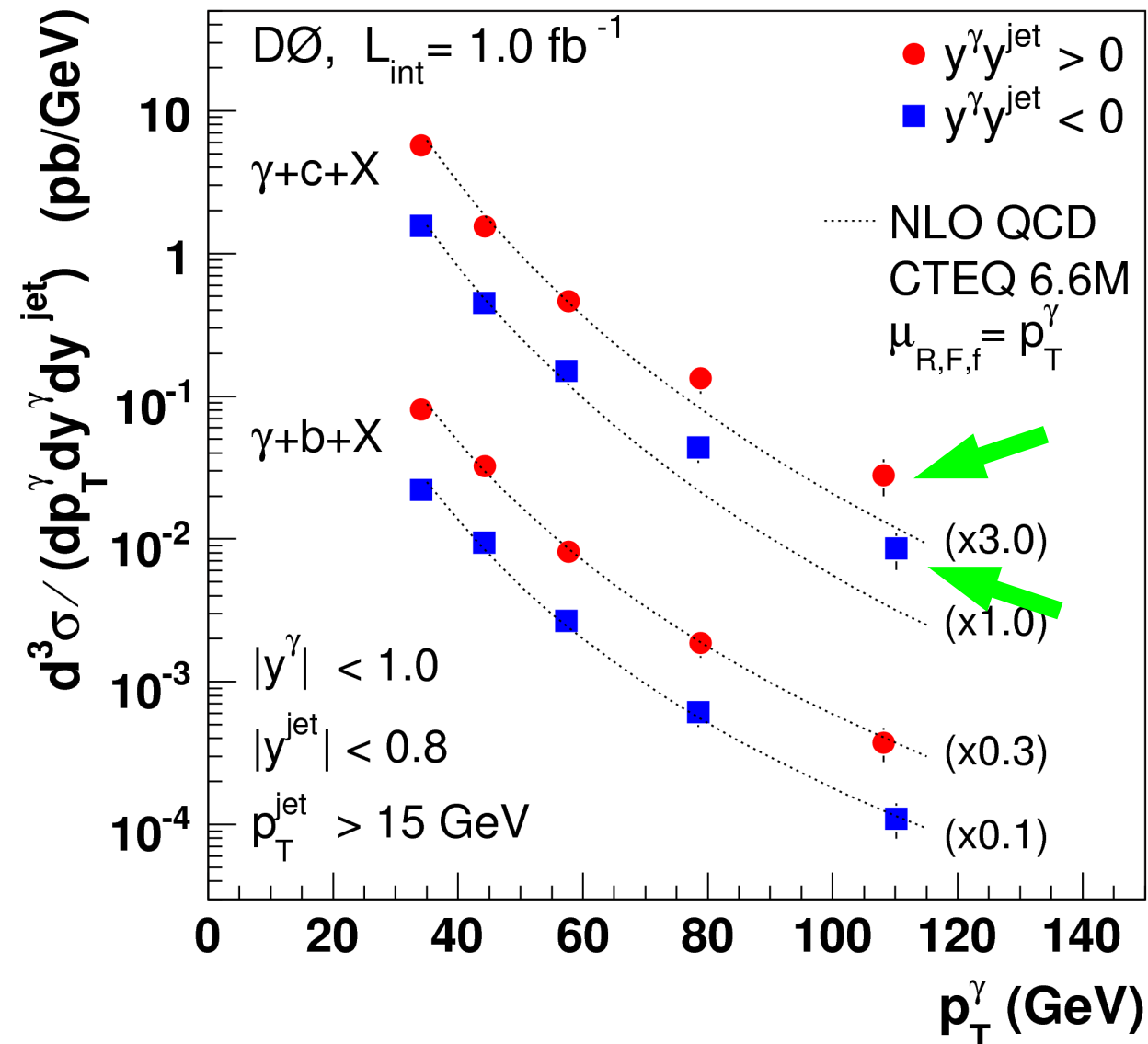
- Important background for Higgs searches: $W+H$ ($H \rightarrow b\bar{b}$) at low Higgs masses. Also a background for $t\bar{t}$ and single-top measurements
- W+b excess over expectations published by CDF

A. Messina's talk

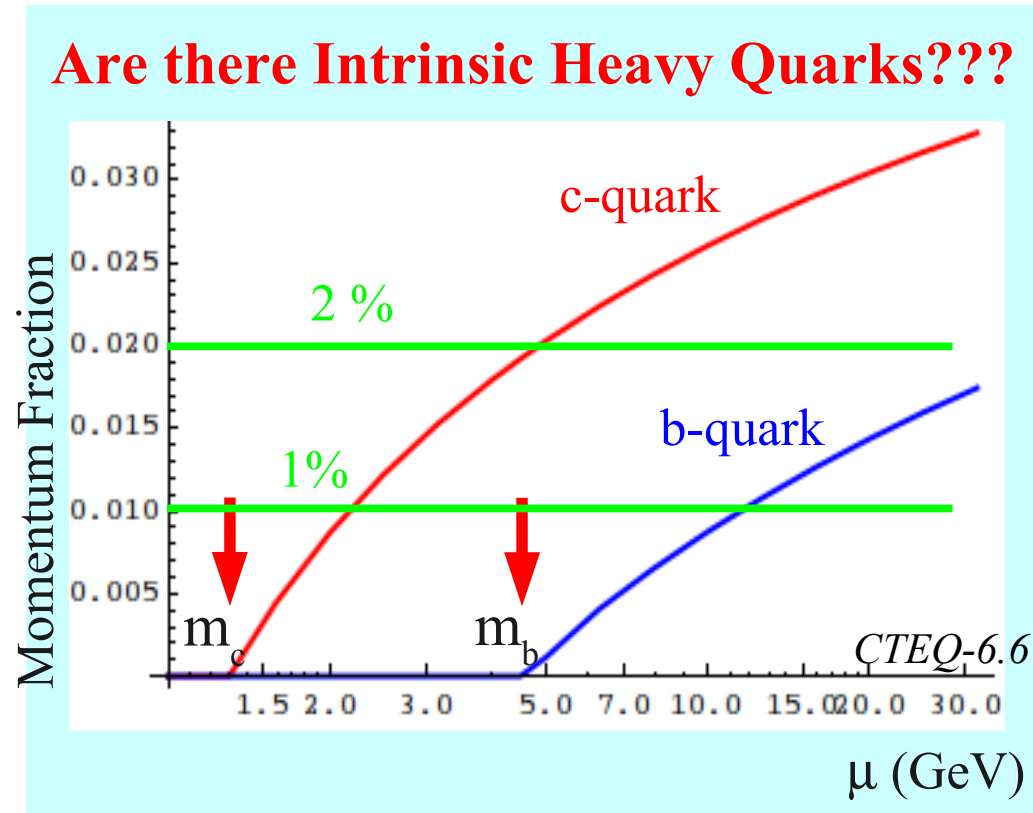


- Significant decay length ($>5.85 \sigma$), fit to the reconstructed mass at secondary vertex
- Challenging analysis: it requires significant reduction and control of top backgrounds and W+charm. Analysis performed independently for 1 and 2 b-tags in the event

Agreement with theoretical predictions at the 1.5σ level



Are there Intrinsic Heavy Quarks??? Do they matter???

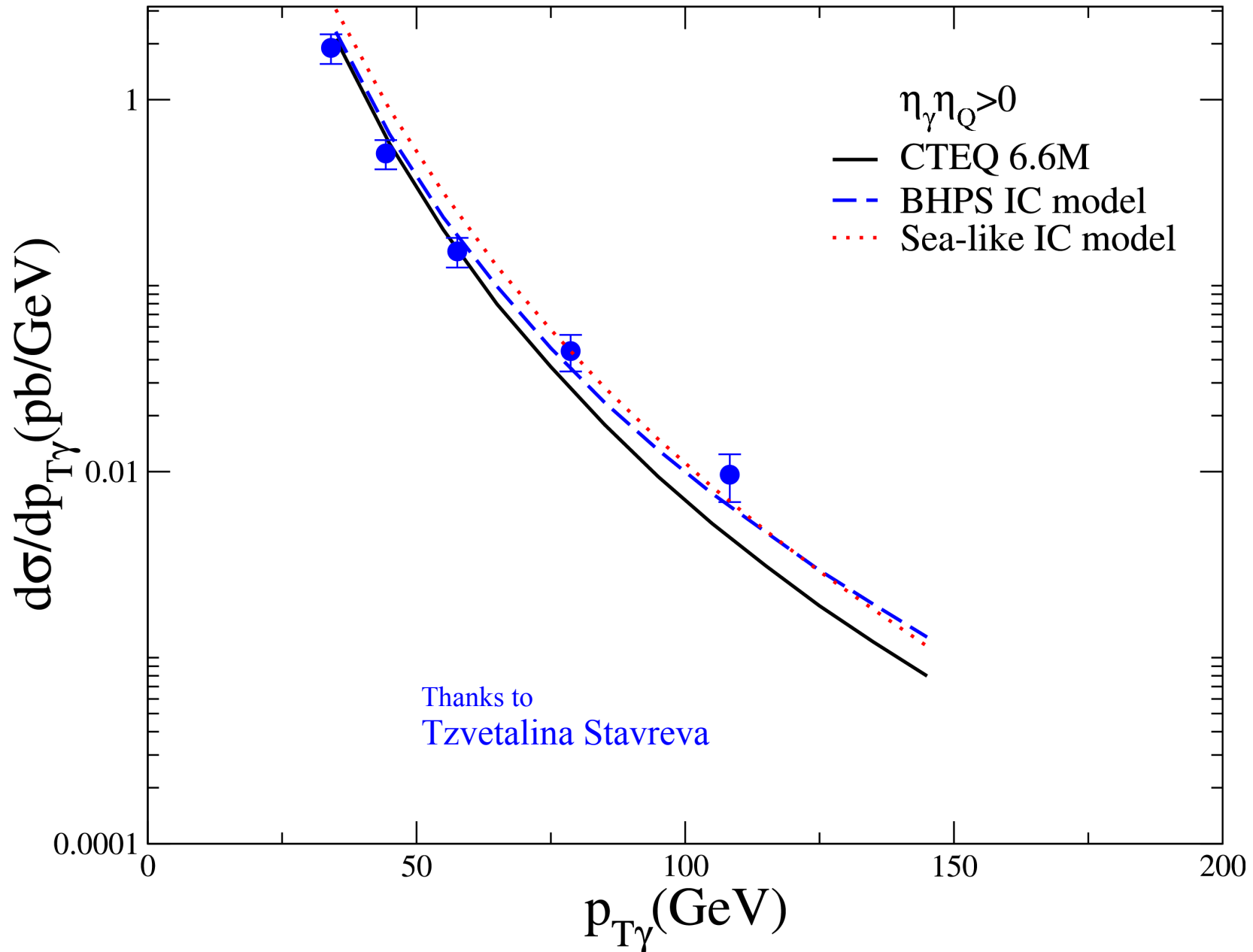


- * Most sensitive near threshold
- * What happens if we allow the evolution to determine charm?

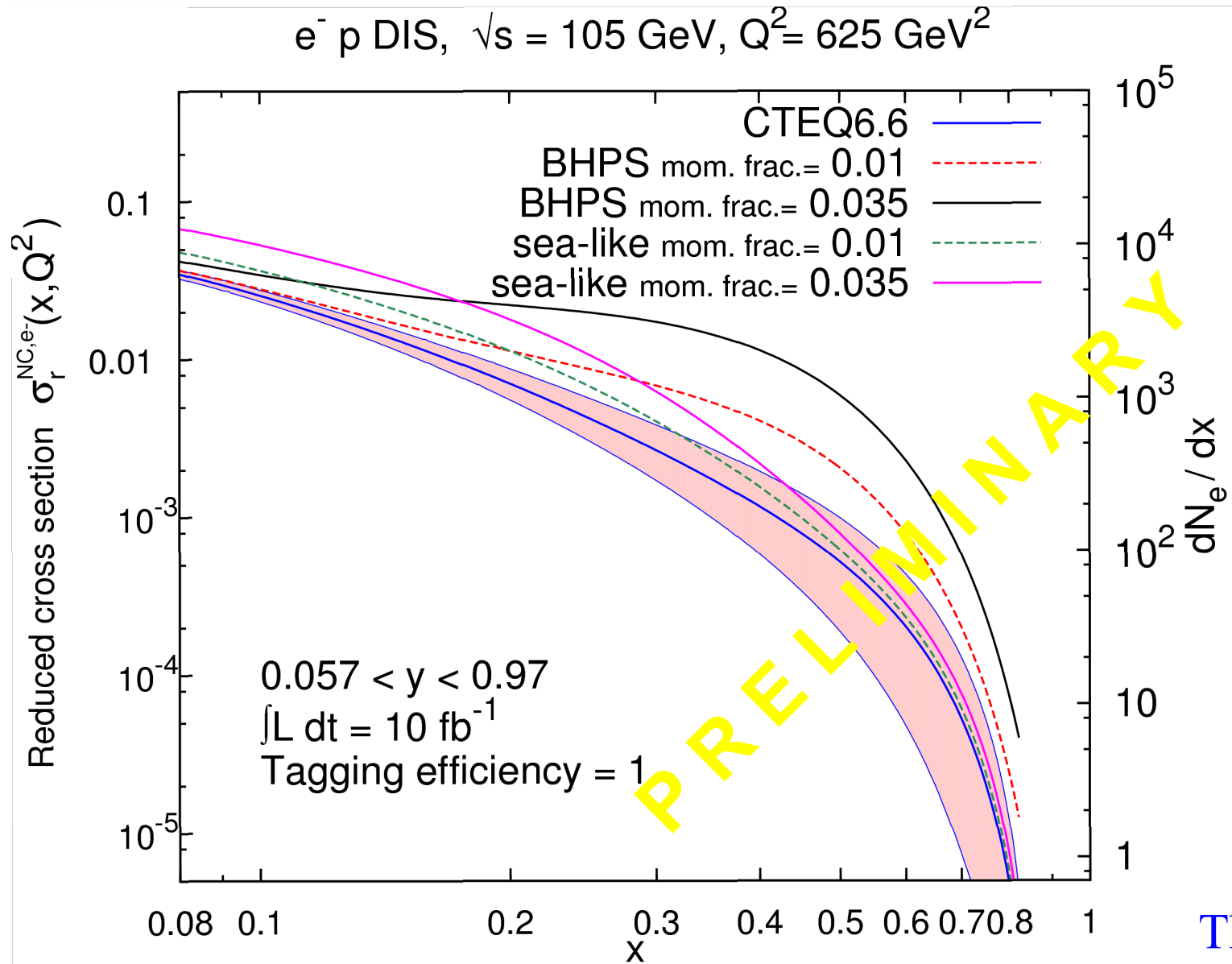
Zero: No intrinsic charm
 Positive: Intrinsic charm
 Negative: Inconsistent

$$p+p \rightarrow \gamma+c+X$$

$$\sqrt{S} = 1.96 \text{ TeV}$$



Sample Cross Section for an Electron Ion Collider



Thanks to
Marco Guzzi
for this
calculation

What is the
proper treatment
of masses???

2009 Les Houches Comparative Studies

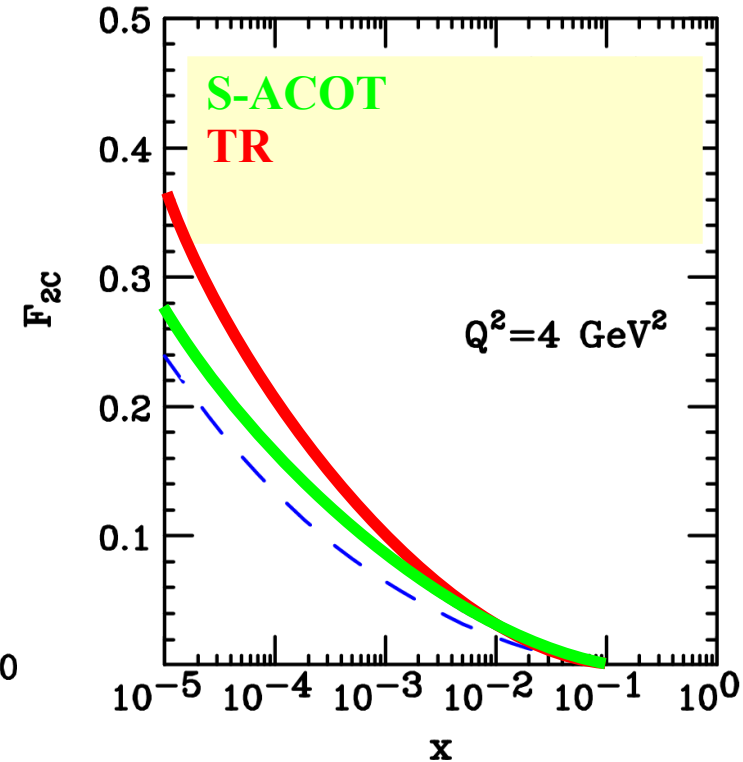
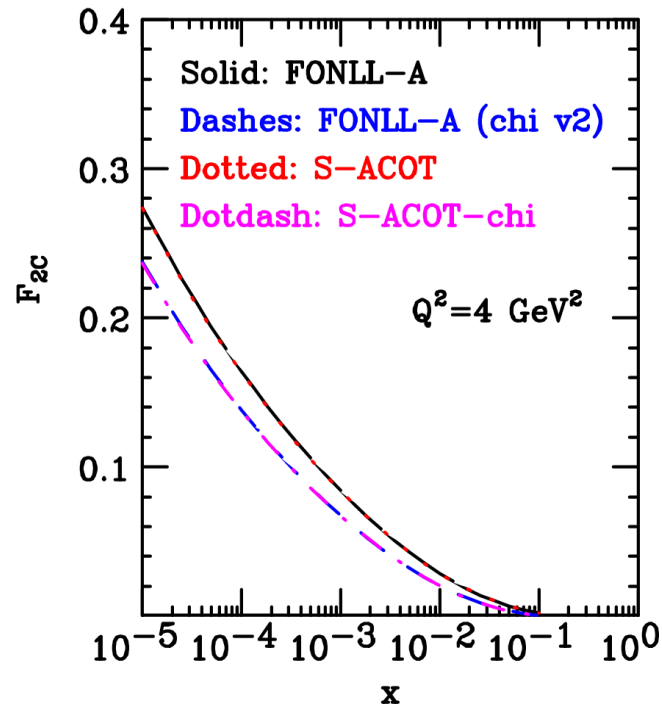
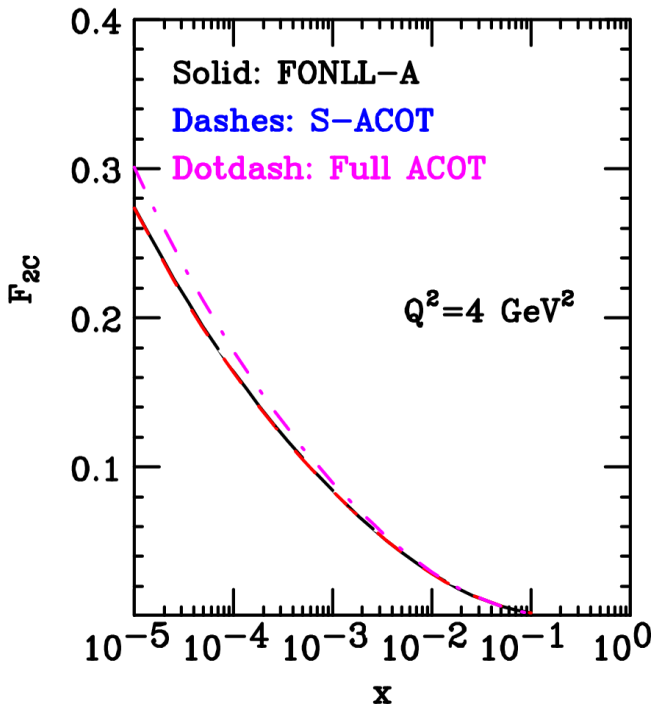
The SM and NLO Multileg Working Group: Summary report.

e-Print: arXiv:1003.1241 [hep-ph]



Physics at TeV Colliders
Les Houches 8-26 June 2009





ACOT & S-ACOT
 essentially
 identical
 ... scheme
 differences are
 higher order

FONNL & S-ACOT
 Numerically similar

MSTW09

We can quantify
theoretical scheme
 differences

TR type schemes

ACOT type schemes

$Q < m_H$

$Q > m_H$

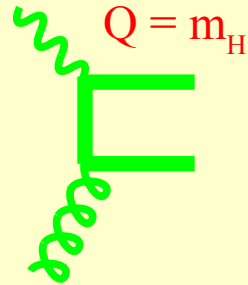
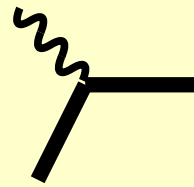
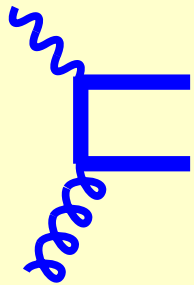
constant
term

$Q < m_H$

$Q > m_H$

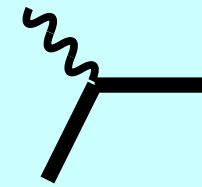
constant
term

LO



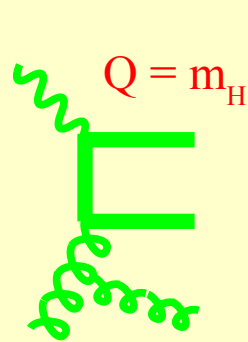
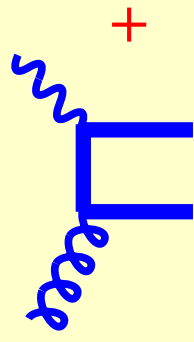
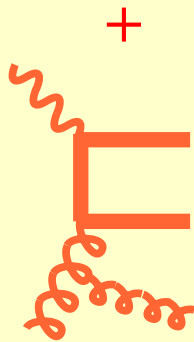
LO

\emptyset



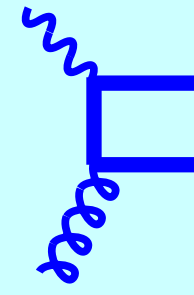
+
 \emptyset

NLO

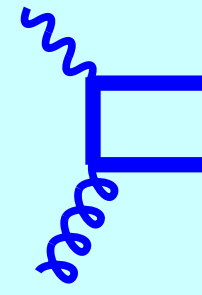


NLO

+

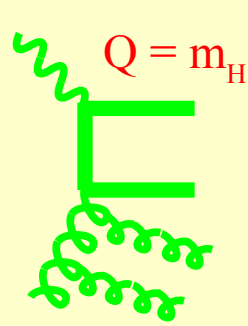
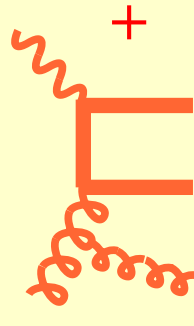
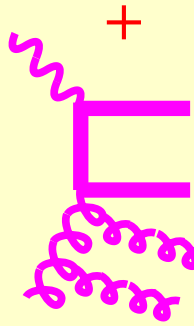


+



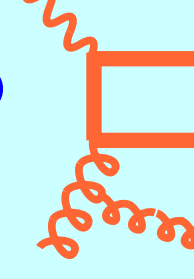
+
 \emptyset

NNLO

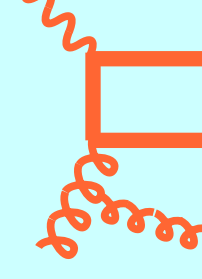


NNLO

+

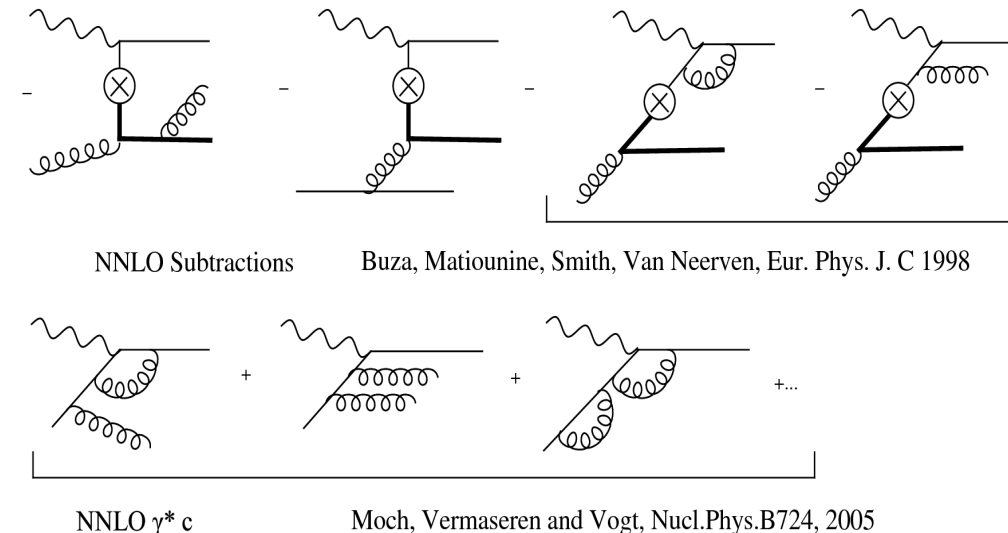
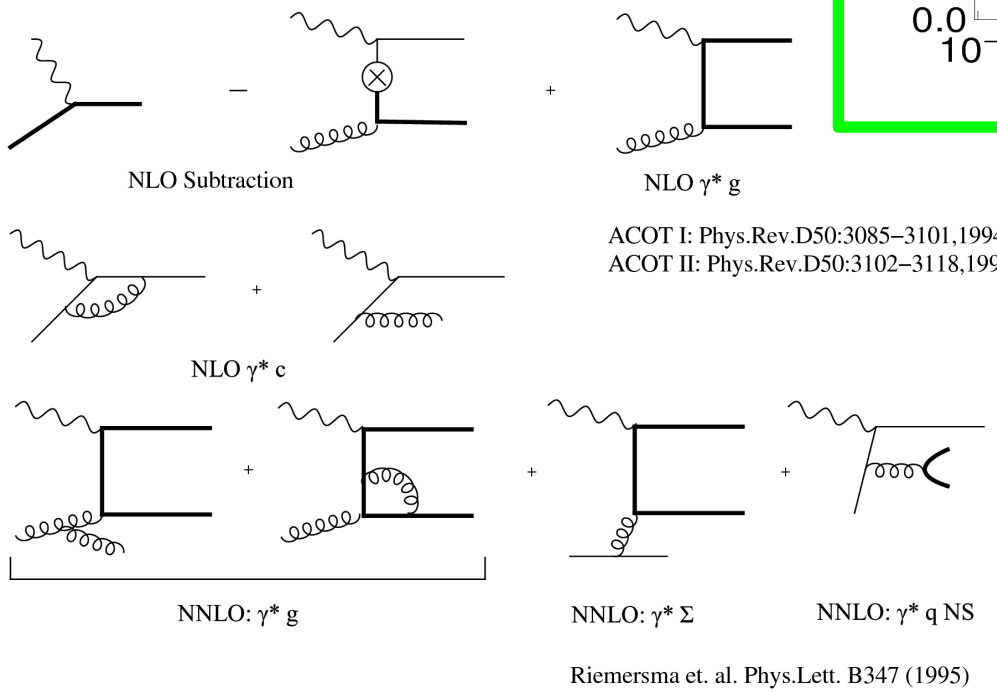
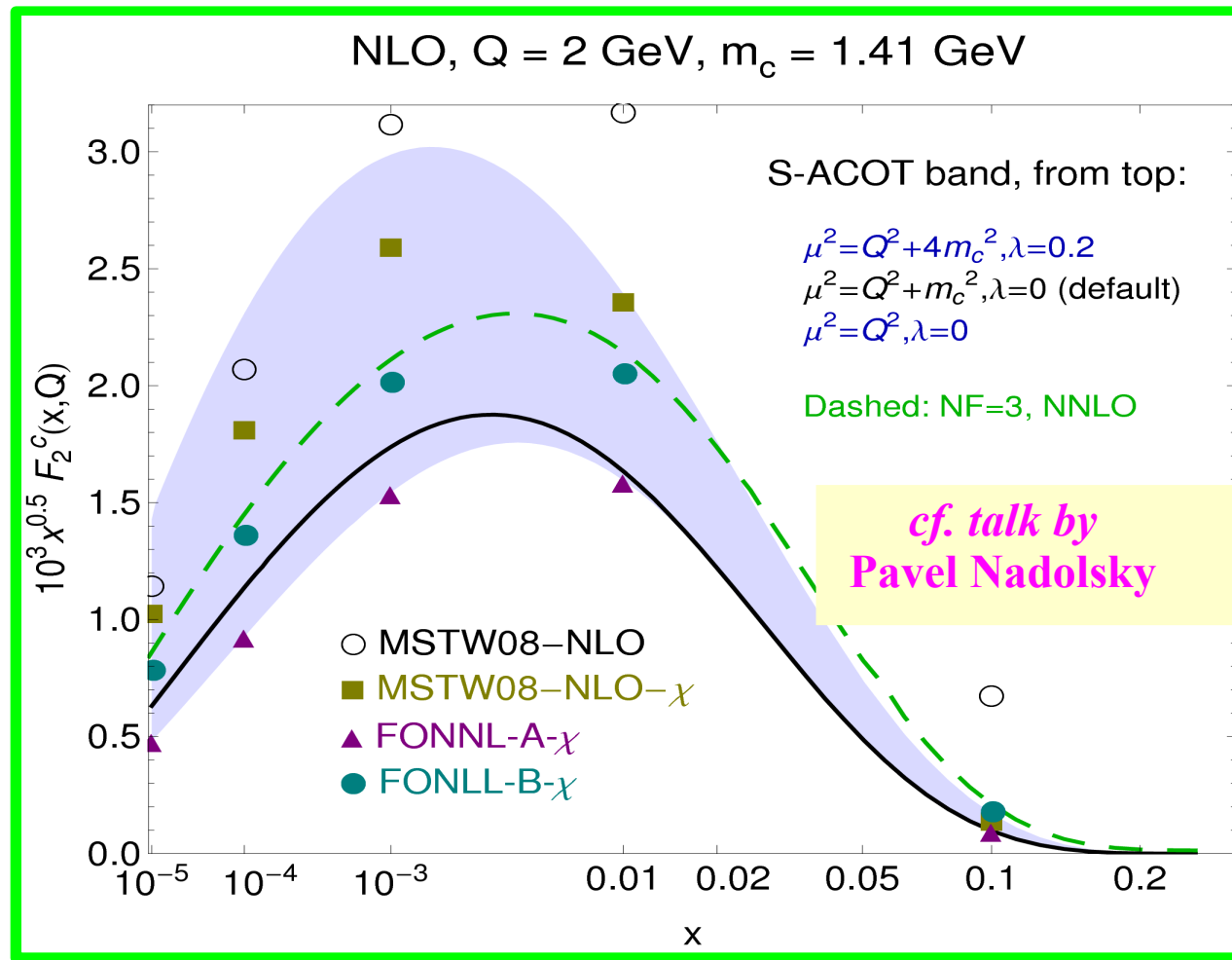
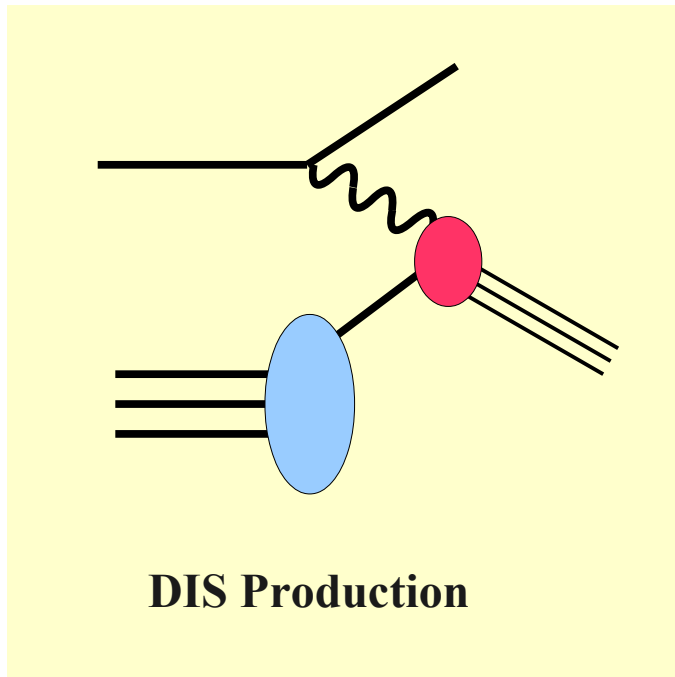


+



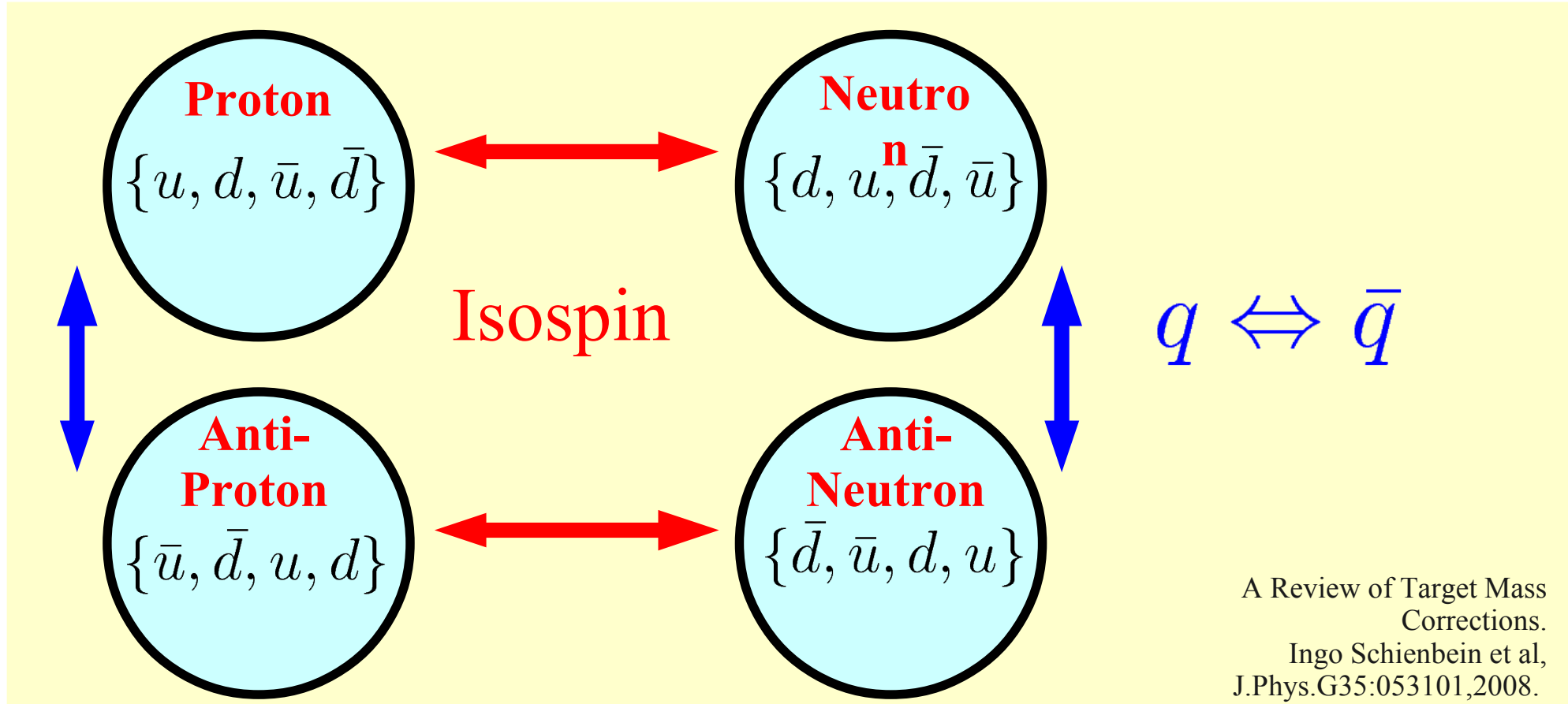
+
 \emptyset

Step toward NNLO CTEQ



Isospin Symmetry

... taken for granted



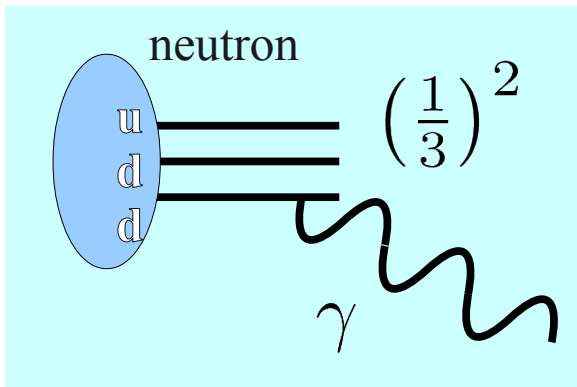
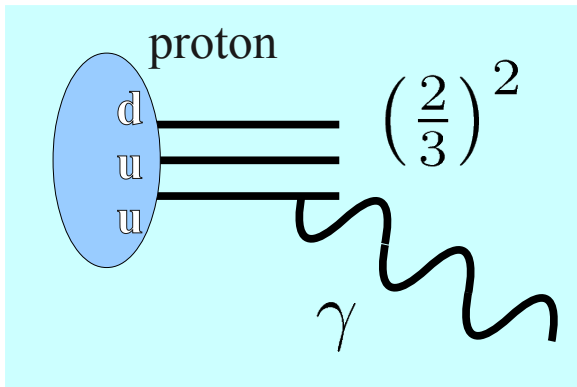
$$\Delta x F_3^A = x F_3^{\nu A} - x F_3^{\bar{\nu} A} = +2x s_A^+ - 2x c_A^+ + x \delta I_A$$

$$\Delta F_2^A = \frac{5}{18} F_2^{CC} - F_2^{NC} = +\frac{1}{6} x s_A^+(x) - \frac{1}{6} x c_A^+(x) + \frac{1}{6} x \frac{N}{A} \delta I_A$$

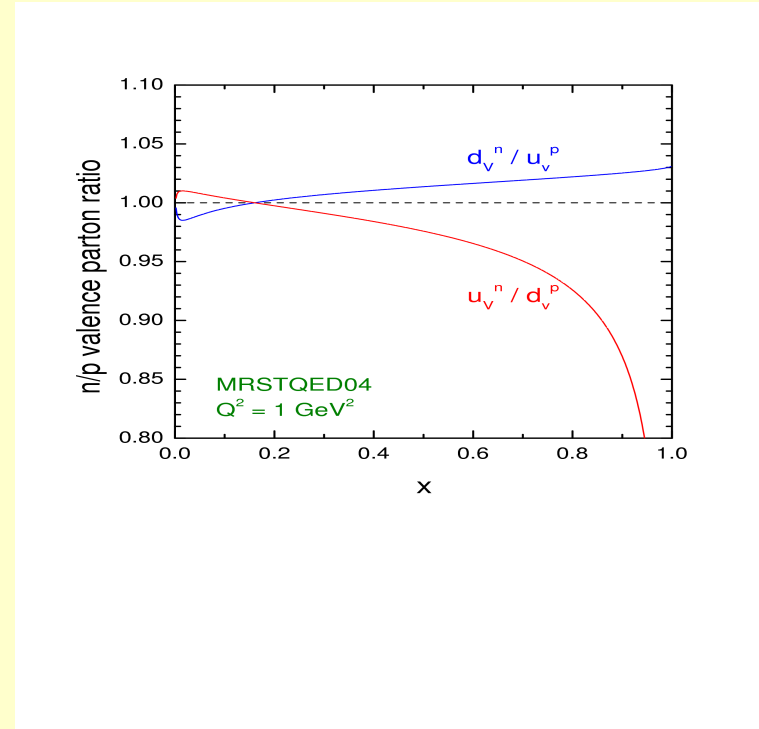
$$\delta I_A = (\delta d + \delta \bar{d}) - (\delta u + \delta \bar{u})$$

$$\delta u = \delta u_p - \delta d_n$$

Photon is not flavor blind!!!



MRST-QED 04



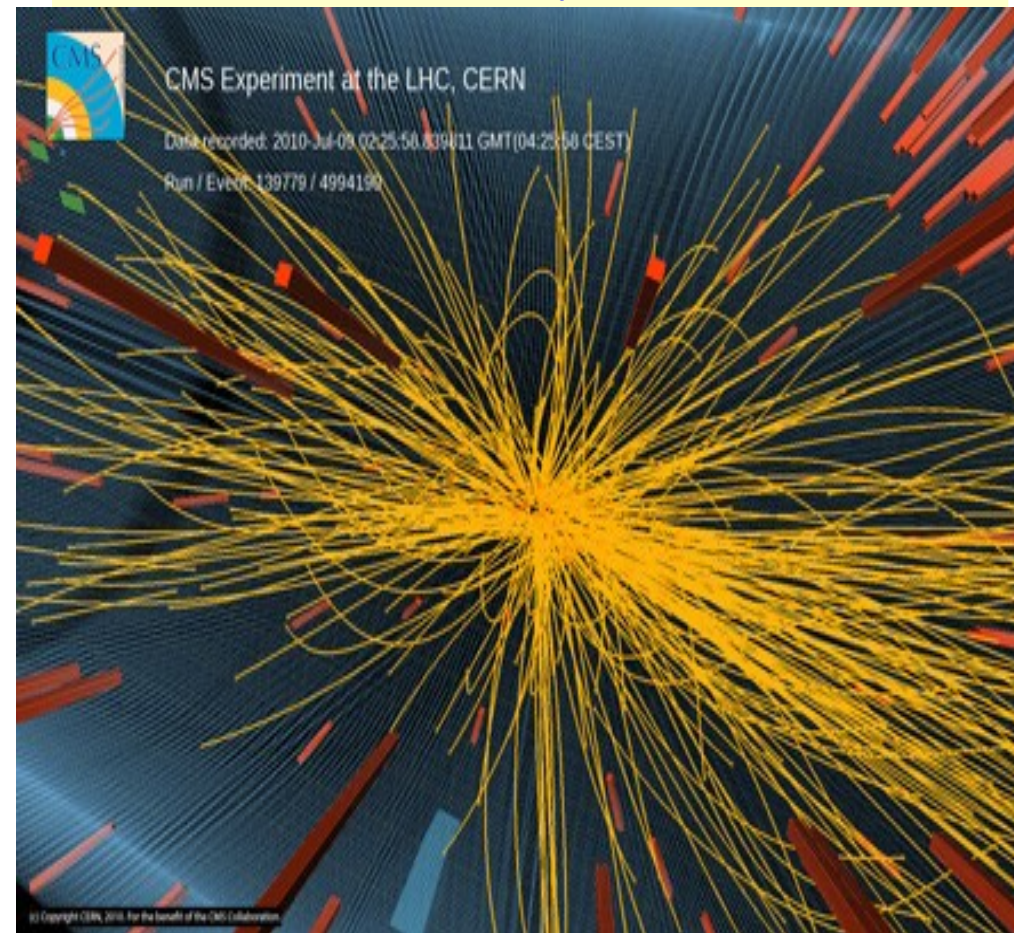
MRST, Eur.Phys.J.C39:155-161,2005.

Isospin terms are comparable to NNLO QCD

Could Isospin terms affect Tevatron W-Asymmetry???

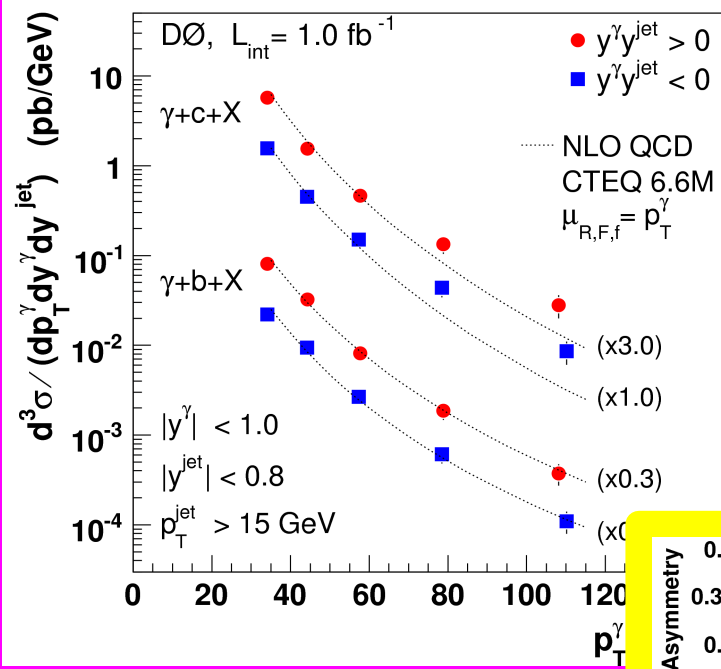


Ability to discover
“New Physics”
is dependent on distinguishing
“Old Physics”



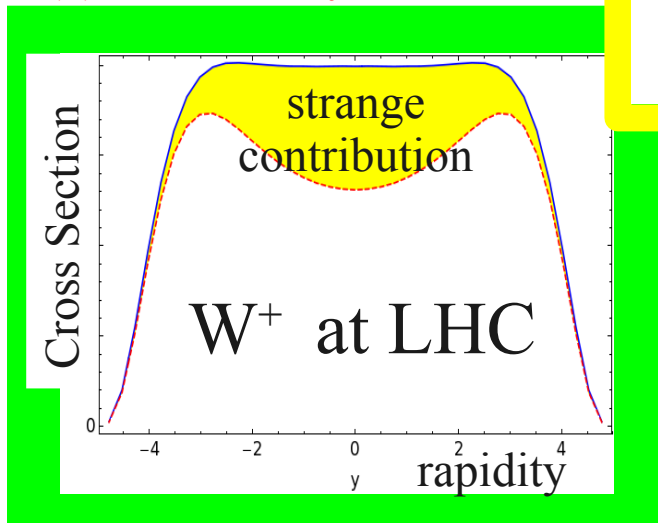
As experimental precision has
increased, we need to be concerned
about the details

The Puzzles

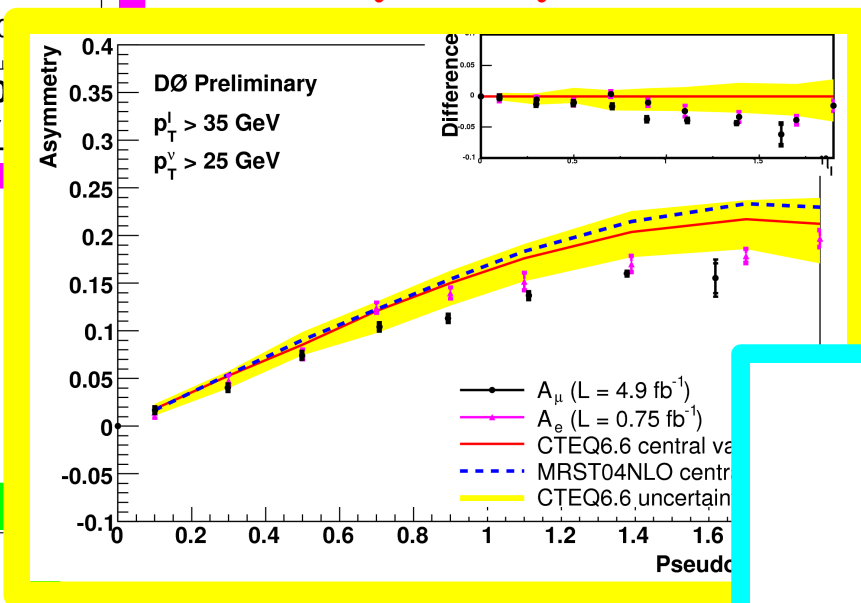


Charm+ γ at Tevatron

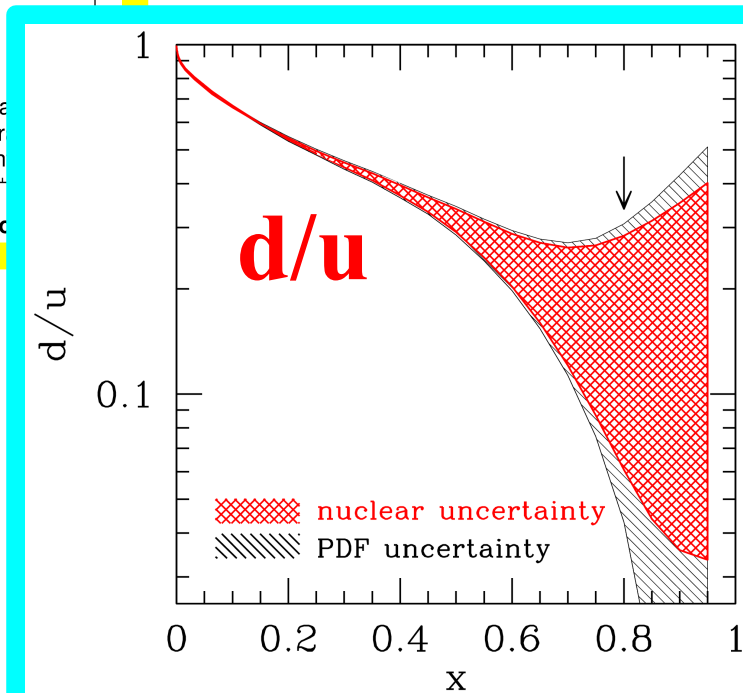
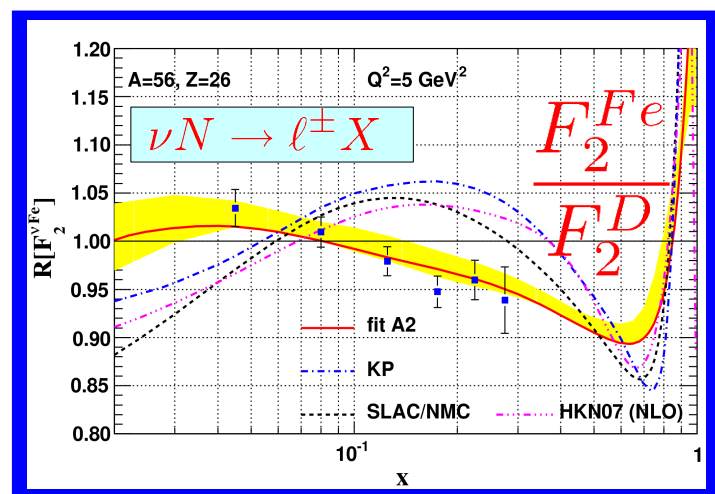
$s(x)$ uncertainty



W Asymmetry



Nuclear Modifications
limit PDFs



Backup

F_L

Why is F_L so special ???

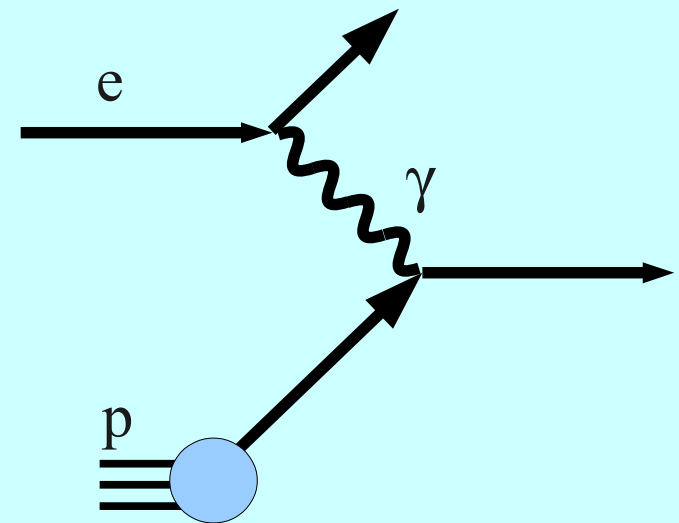
$$\frac{d\sigma^{\nu DIS}}{dx dy} = (1-y)^2 \bar{q}(x) + (1-y) \phi(x) + q(x)$$

$$\frac{d\sigma^{\nu DIS}}{dx dy} = (1-y)^2 F_+(x) + (1-y) F_0(x) + F_-(x)$$

$$F_0 = \frac{F_2}{2x} - F_1$$

$$F_0 = 0 \implies F_2 = 2x F_1$$

Callan-Gross



$$F_L \sim \frac{m^2}{Q^2} q(x) + \alpha_S \{c_g \otimes g(x) + c_q \otimes q(x)\}$$

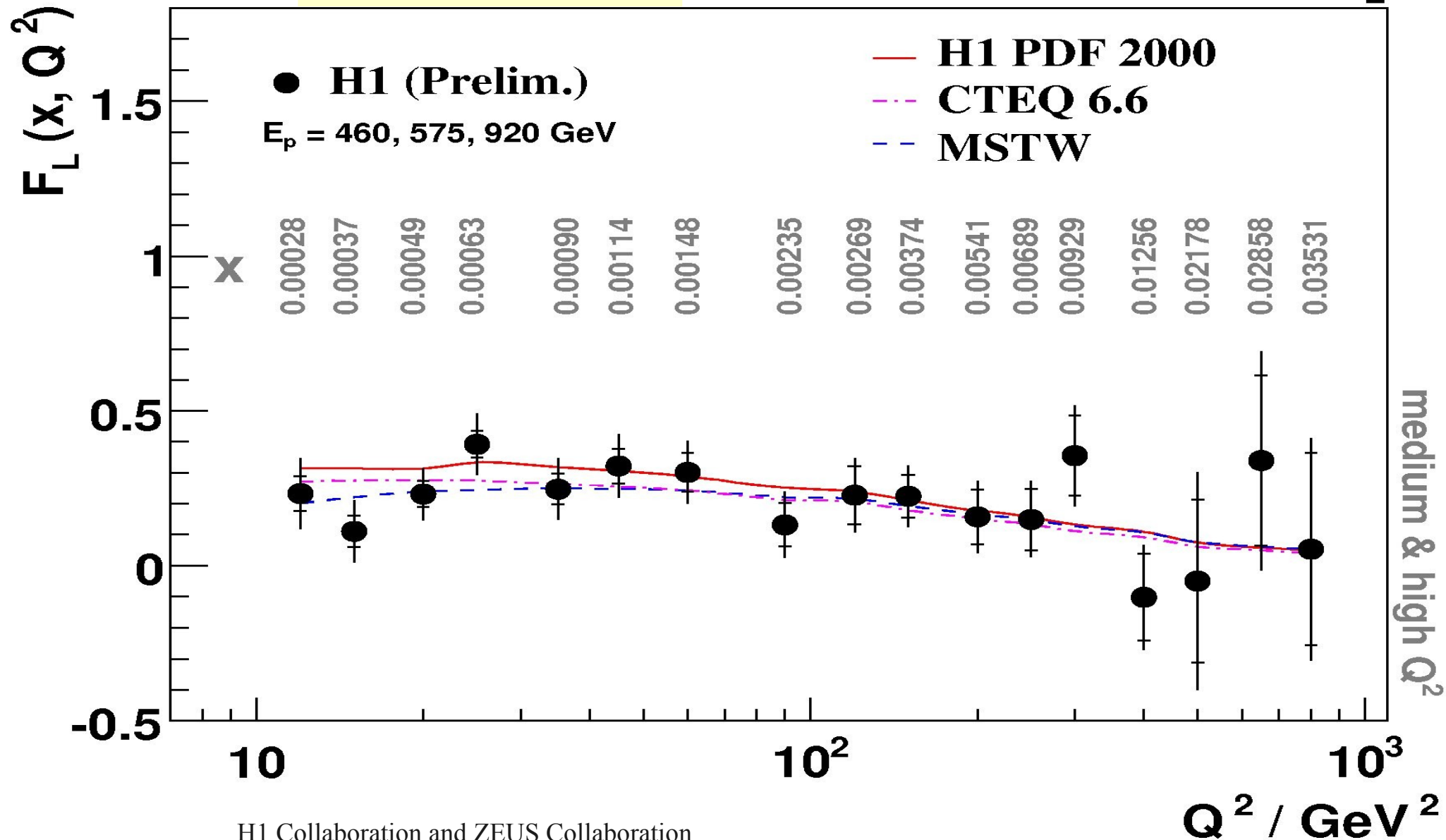
Masses
are
important

Higher Orders
are important

New F_L Measurements: New Perspective

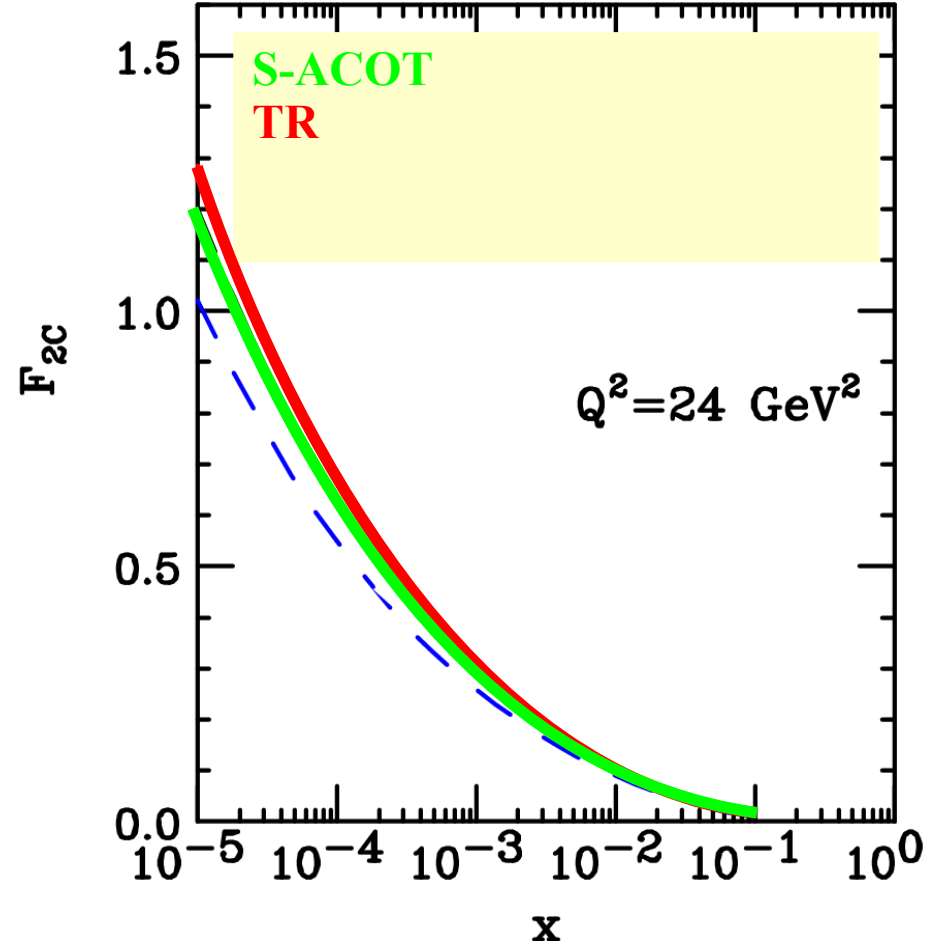
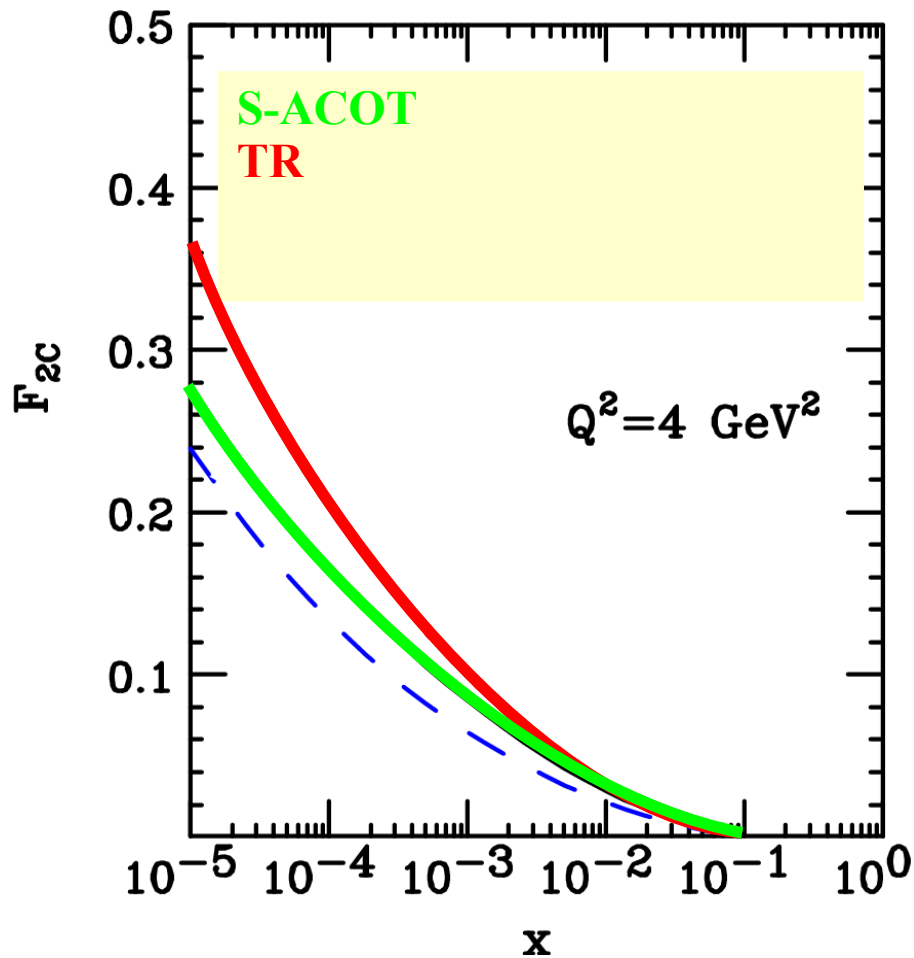
Updated results soon ...

H1 Preliminary F_L



H1 Collaboration and ZEUS Collaboration
 (S. Glazov for the collaboration).
 Nucl.Phys.Proc.Suppl.191:16-24, 2009.

Comparison of ACOT & TR Schemes



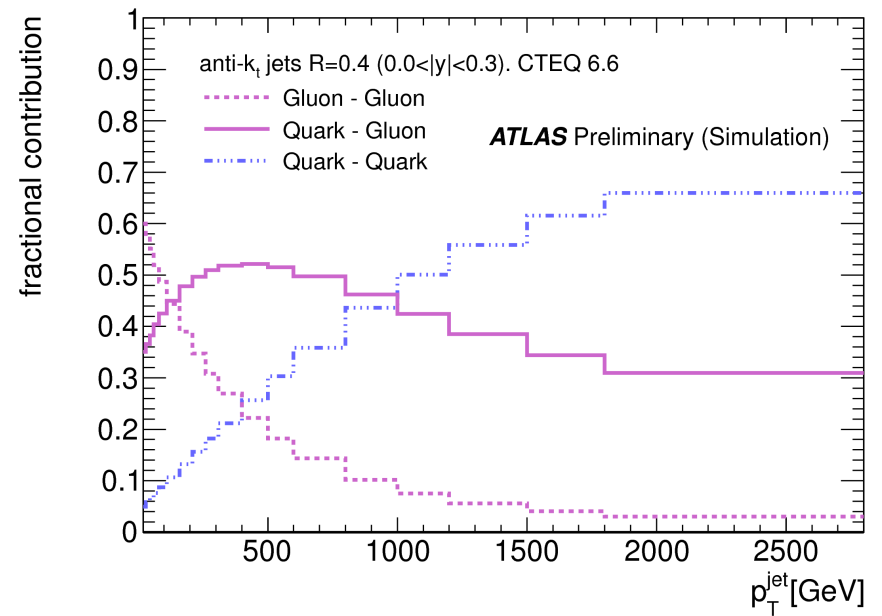
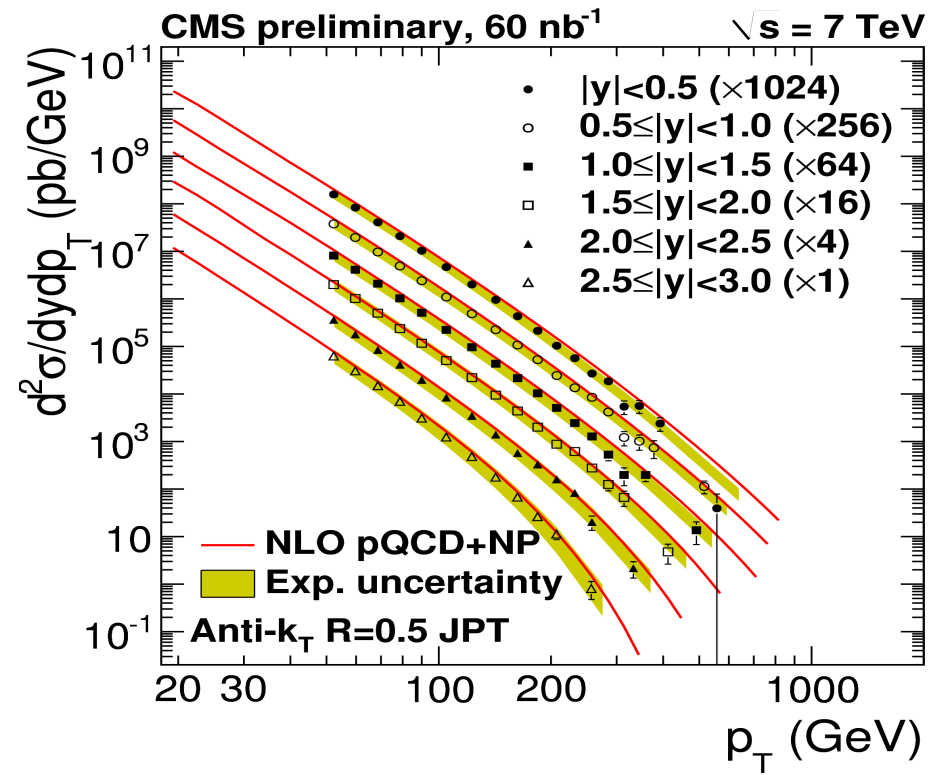
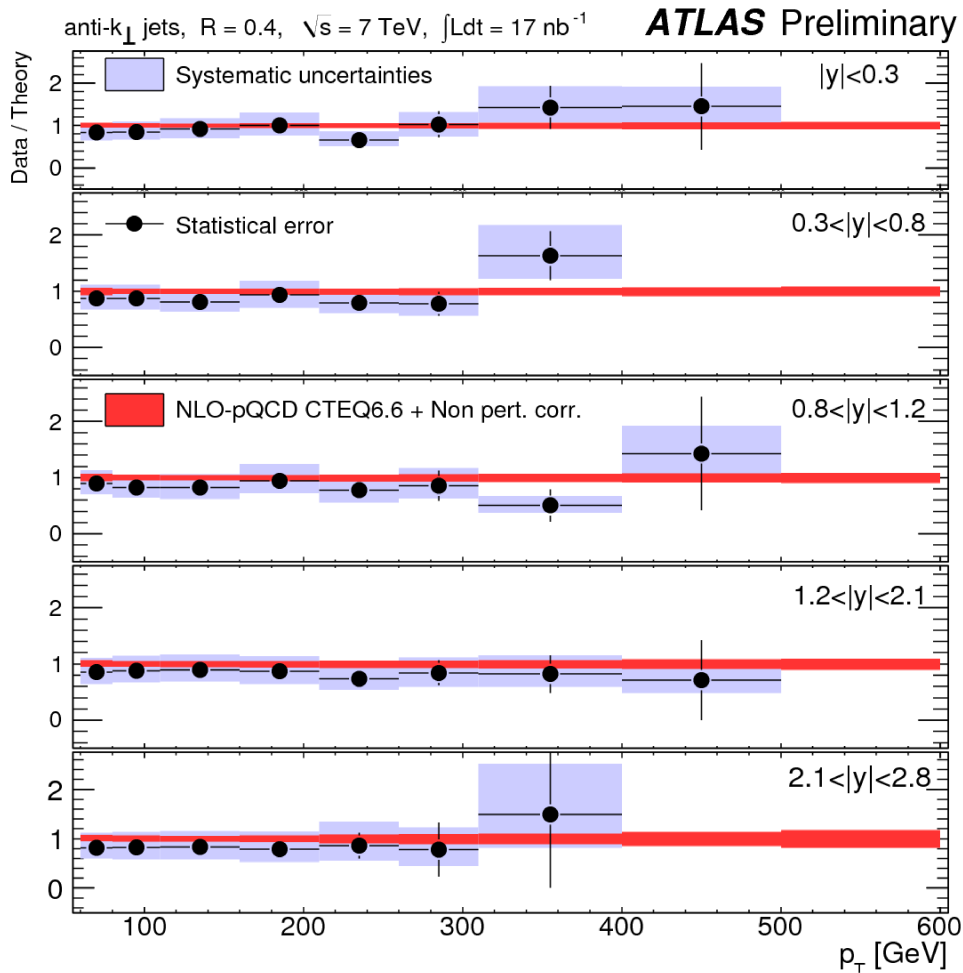
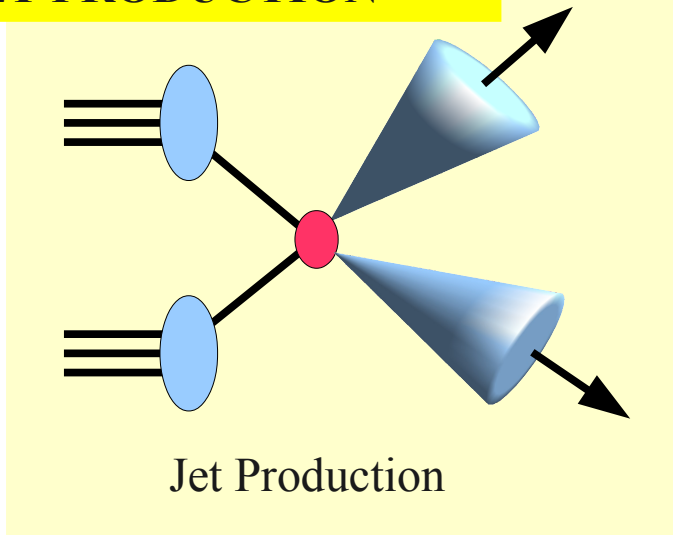
Different schemes \Rightarrow Different PDFs \Rightarrow yet consistent σ

Differences reduce at:

- 1) higher Q ,
- 2) higher order

If experiments are sensitive, time to compute to higher order

JET PRODUCTION



Parton discrimination