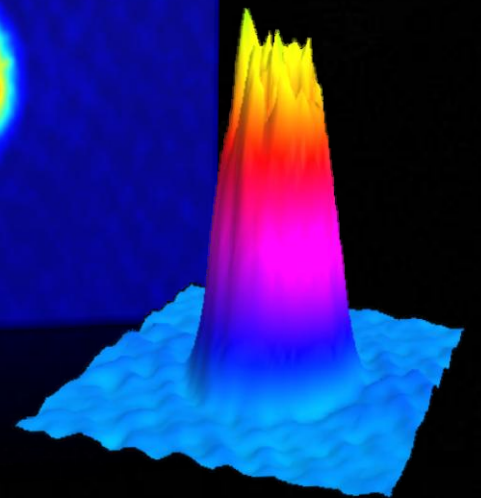
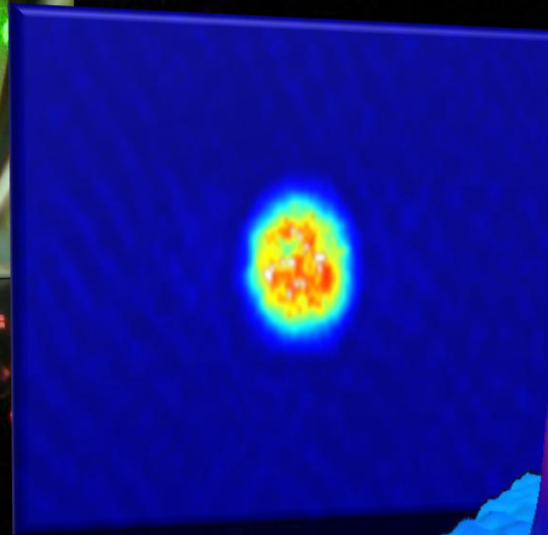
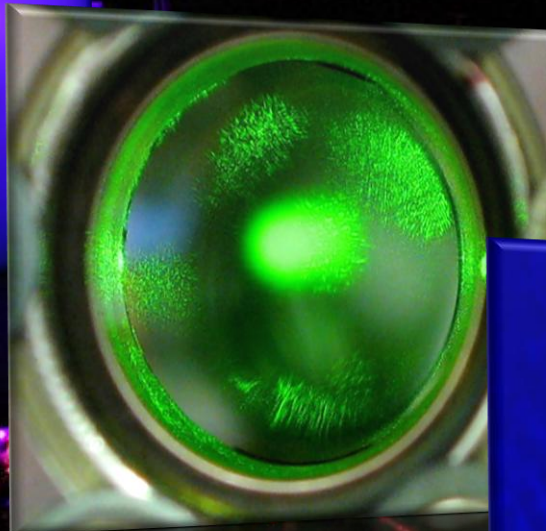
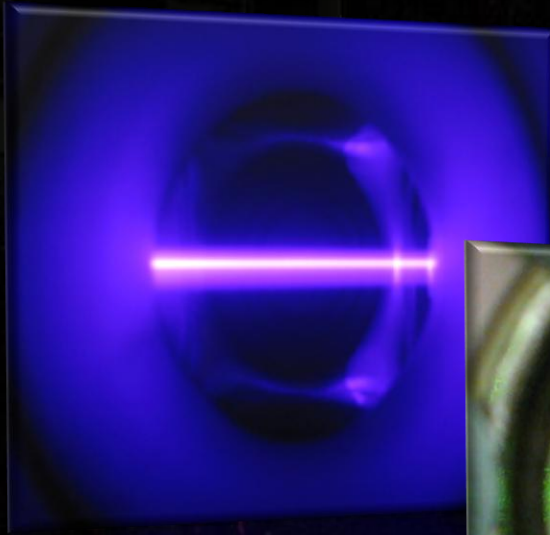


Ytterbium quantum gases in Florence



Leonardo Fallani

University of Florence & LENS

Marco Mancini



Giacomo Cappellini



Guido Pagano



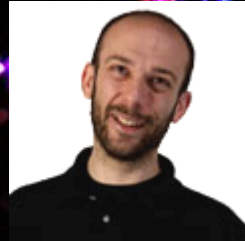
Florian Schäfer



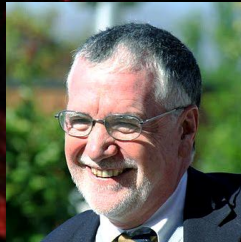
Jacopo Catani



Leonardo Fallani



Massimo Inguscio



and Jonathan T. Green
Pablo Cancio Pastor

Funding from EU FP7 Projects AQUTE, NAMEQUAM
and IIT Istituto Italiano di Tecnologia

Introduction

Bose-Einstein condensation of Ytterbium

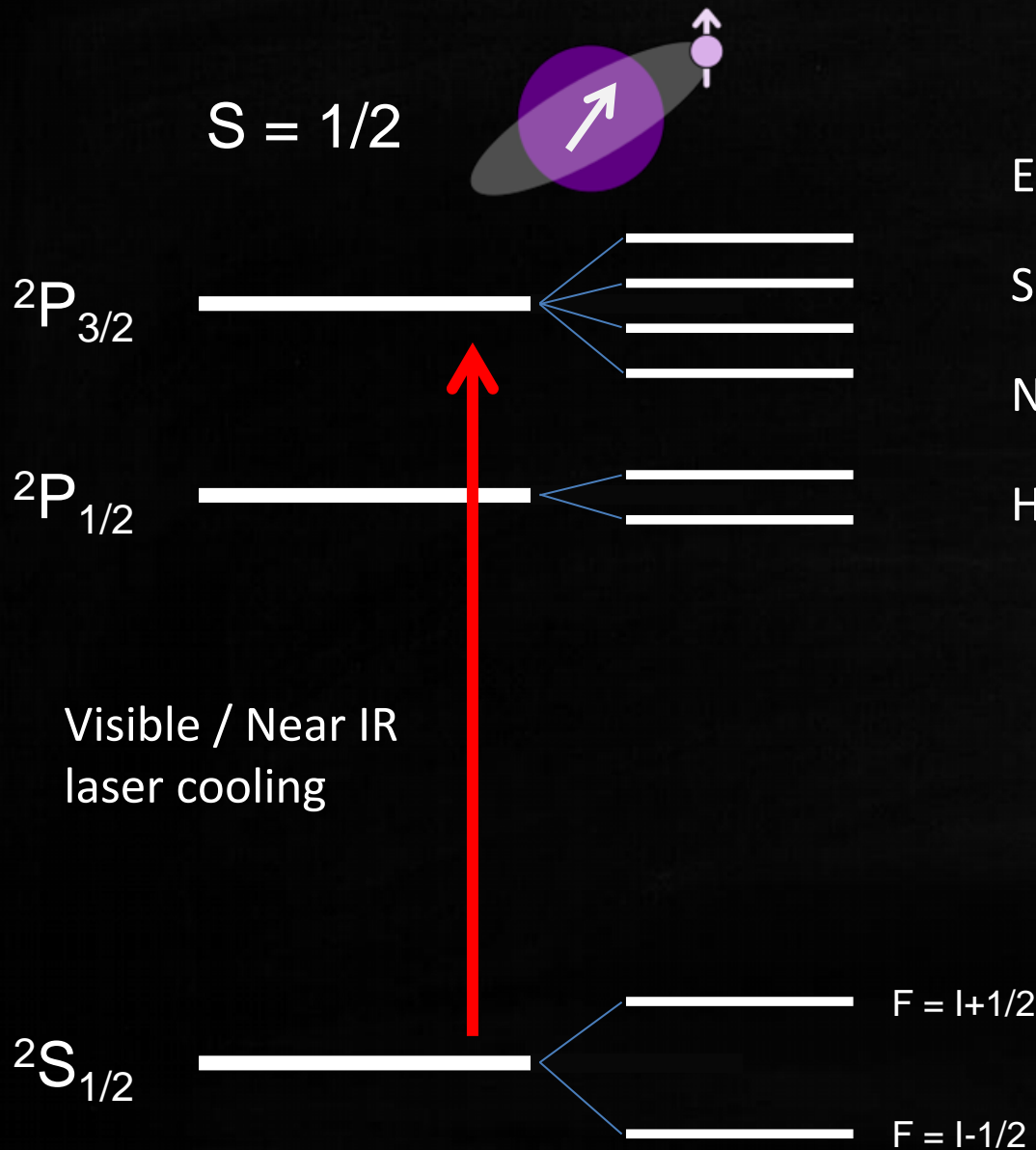
Current and future work

Introduction

Bose-Einstein condensation of Ytterbium

Current and future work

Alkaline atoms



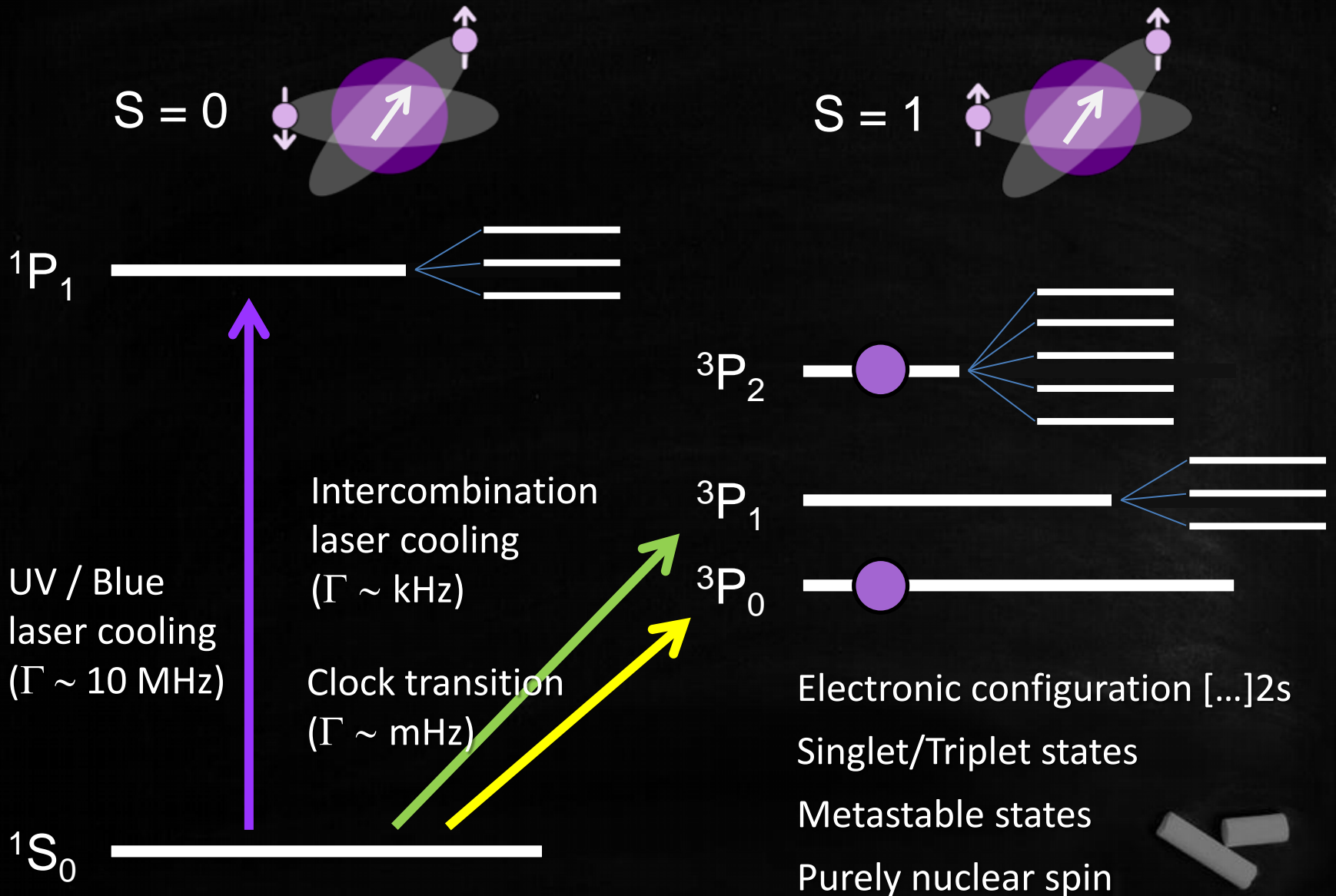
Electronic configuration [...] $1s$

Single-electron structure

Non-zero nuclear spin I

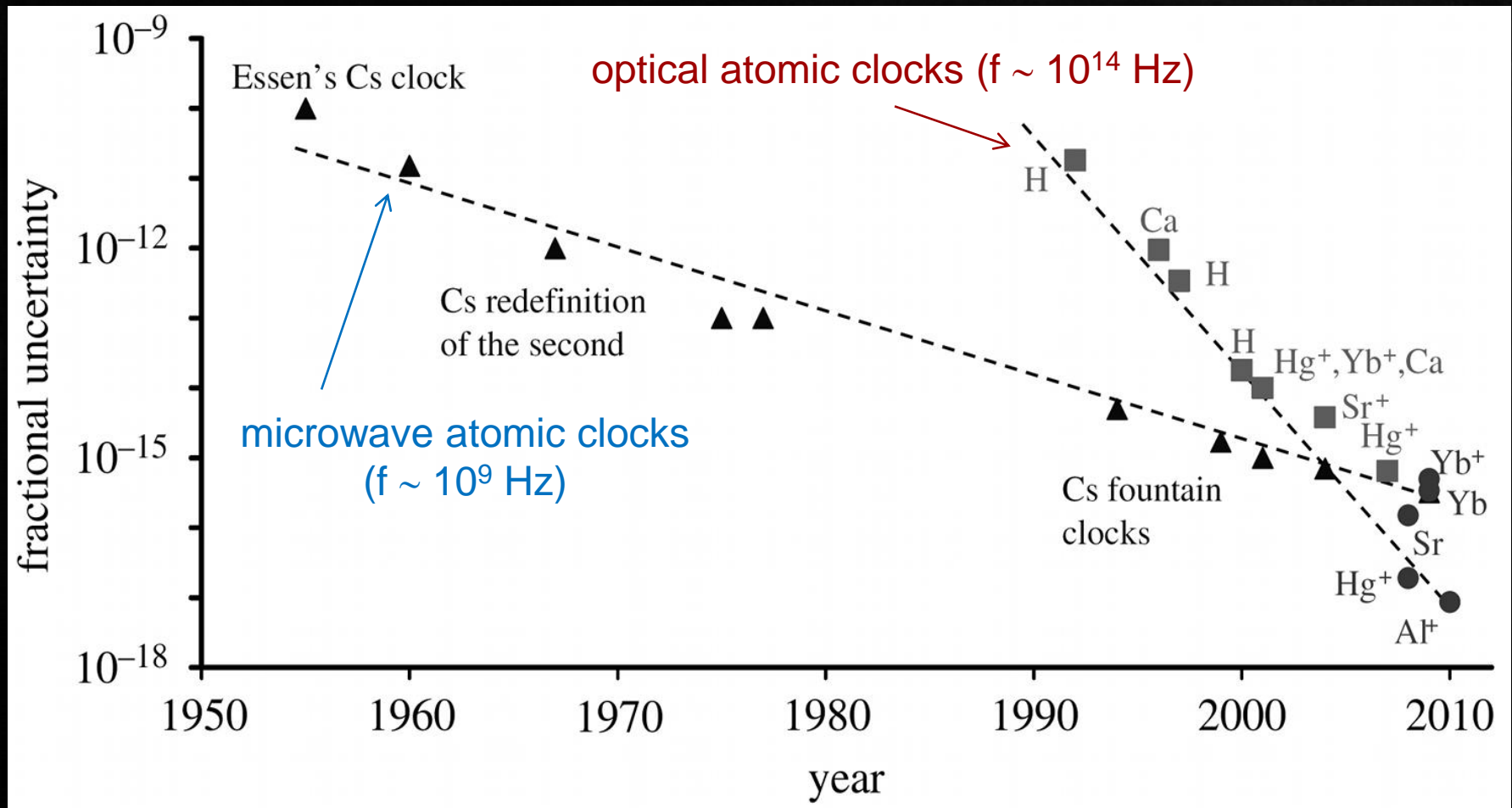
Hyperfine interaction $I \cdot J \neq 0$

Alkaline-earth atoms



Optical clocks

Optical clocks based on $^1S_0 - ^3P_0$ transition in alkaline-earth atoms (and ions)



The Ytterbium family

Yb

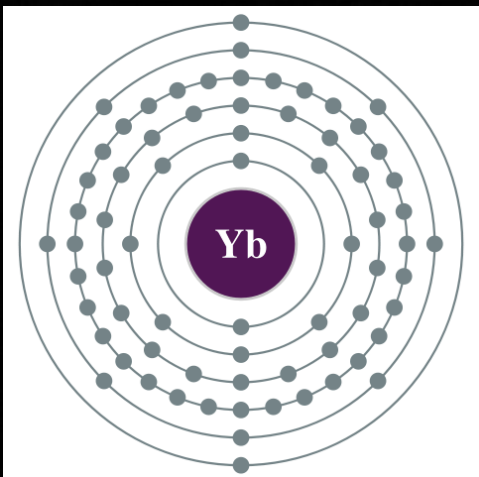
70

173.04



Ytterbium

<http://periodictable.com>

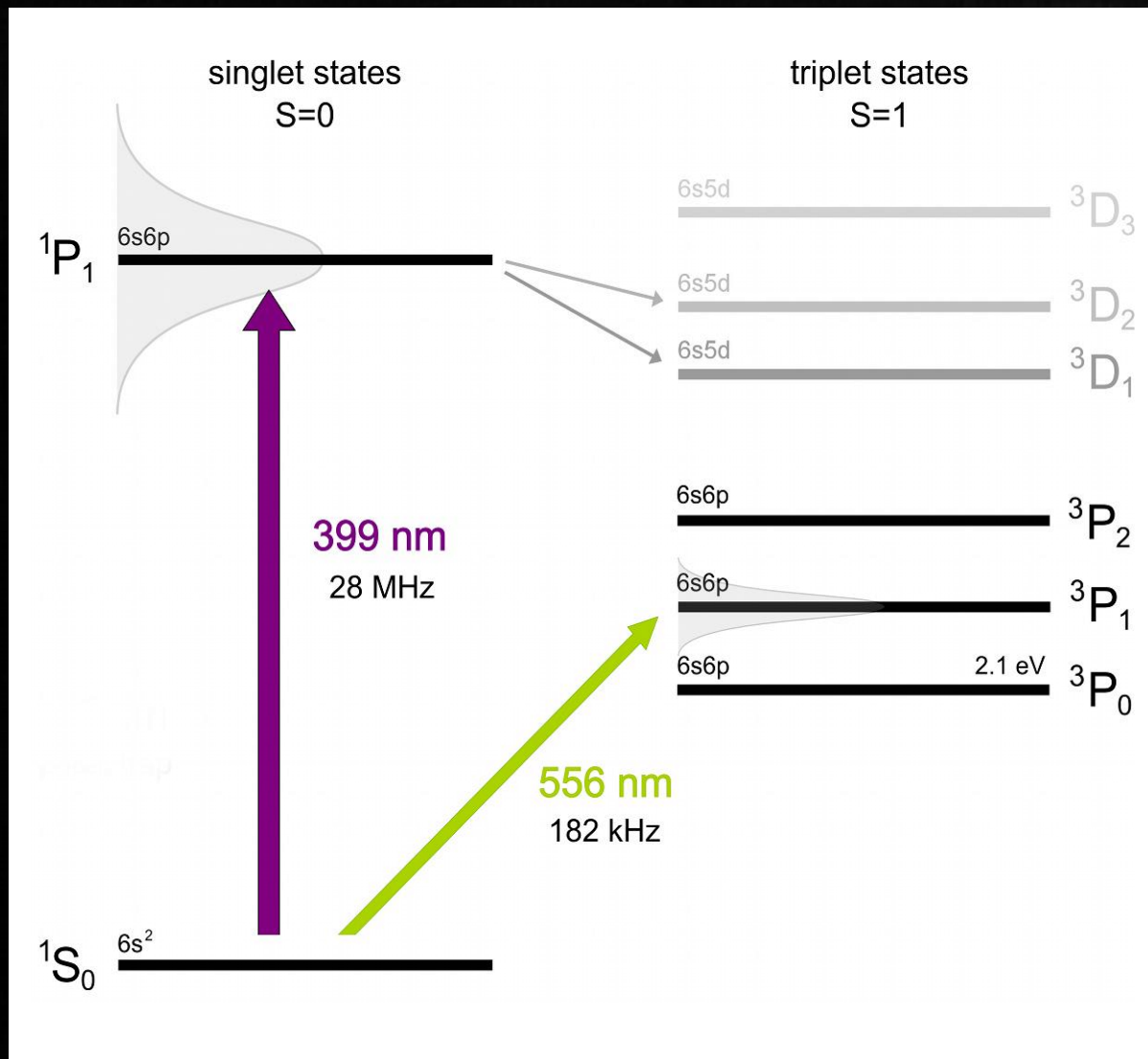


Natural Ytterbium comes in seven stable isotopes:

^{168}Yb	0.13%	$I=0$	boson
^{170}Yb	3.04%	$I=0$	boson
^{171}Yb	14.28%	$I=1/2$	fermion
^{172}Yb	21.83%	$I=0$	boson
^{173}Yb	16.13%	$I=5/2$	fermion
^{174}Yb	31.83%	$I=0$	boson
^{176}Yb	12.76%	$I=0$	boson



Ytterbium levels



Ytterbium interactions

At ultralow temperatures short-range interactions between neutral atoms are completely described by s-wave scattering

s-wave scattering lengths (in a_0 units)

	^{168}Yb	^{170}Yb	^{171}Yb	^{172}Yb	^{173}Yb	^{174}Yb	^{176}Yb
^{168}Yb	252	117	89	65	39	2	-359
^{170}Yb		64	36	-2	-81	-518	209
^{171}Yb			-3	-84	-578	429	142
^{172}Yb				-600	418	200	106
^{173}Yb					200	139	80
^{174}Yb						105	54
^{176}Yb							-24

Kitagawa *et al.*, *PRA* 77, 012719 (2008)

Isotope tuning of the interactions

Introduction

Bose-Einstein condensation of Ytterbium

Current and future work

The experimental setup

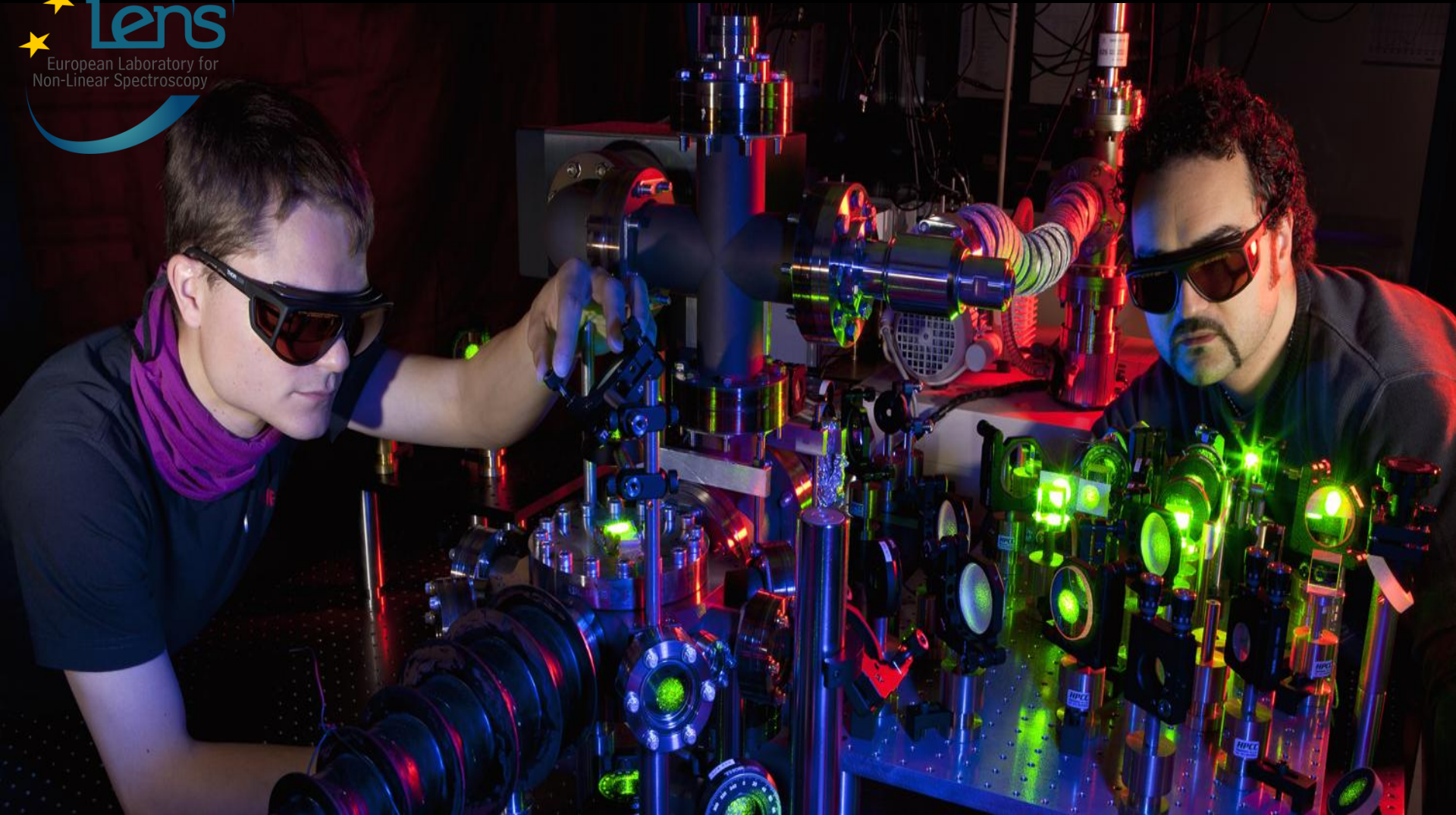


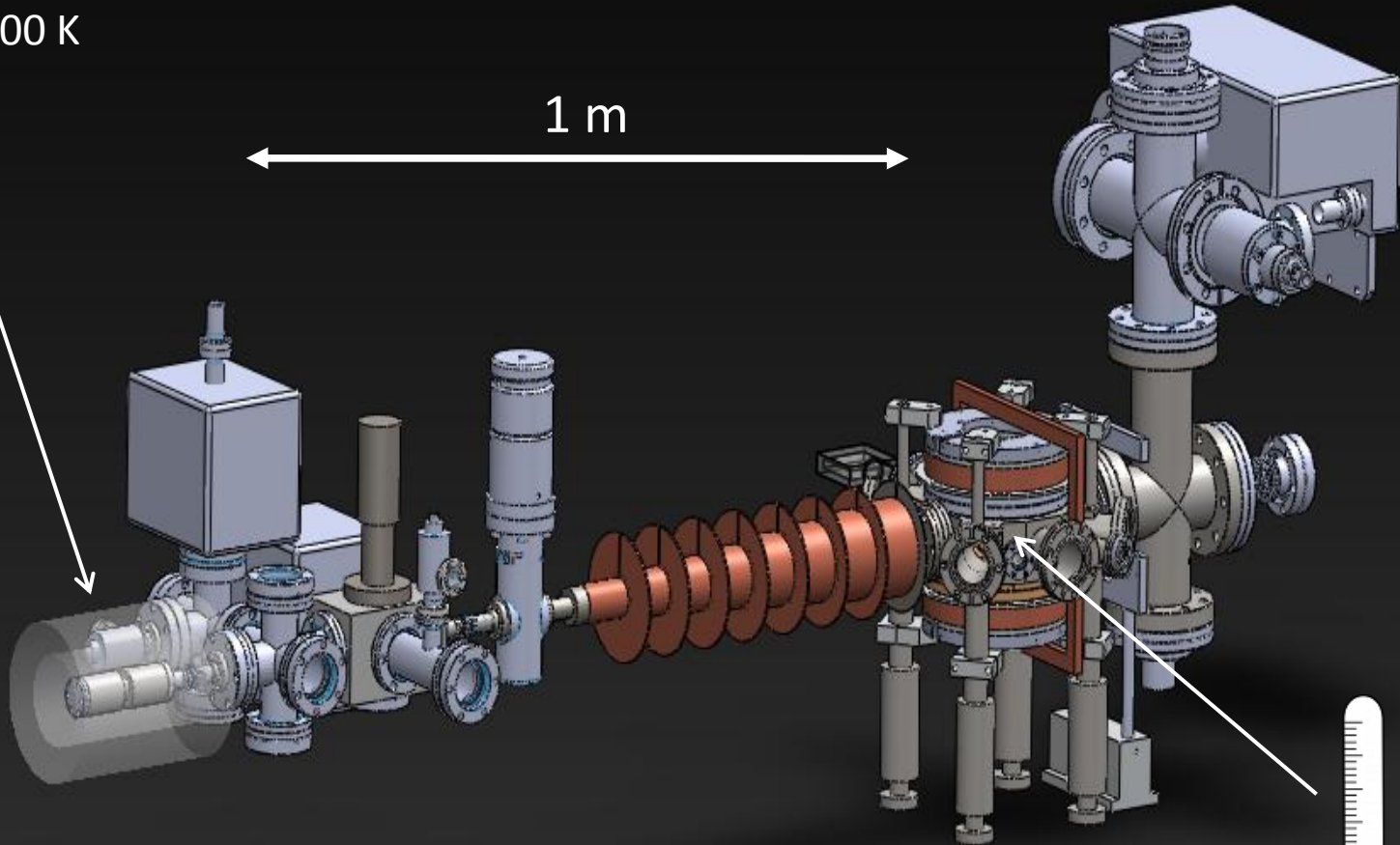
Photo by Marco De Pas

The experimental setup

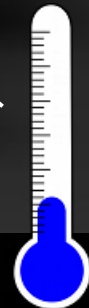


800 K

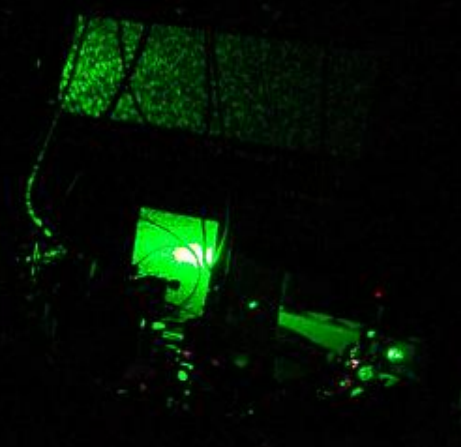
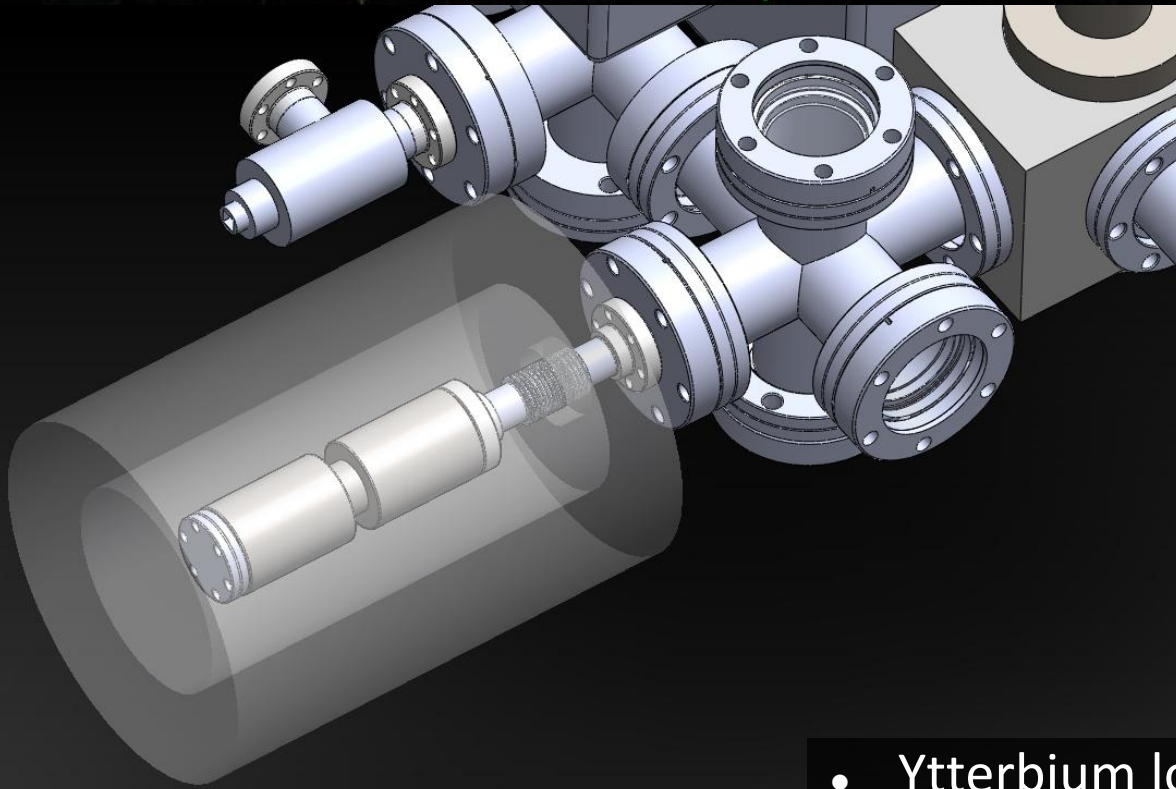
1 m



0.1 μ K

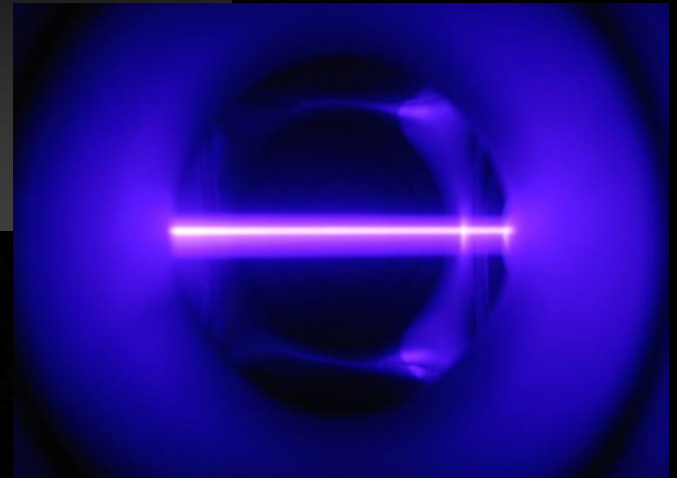
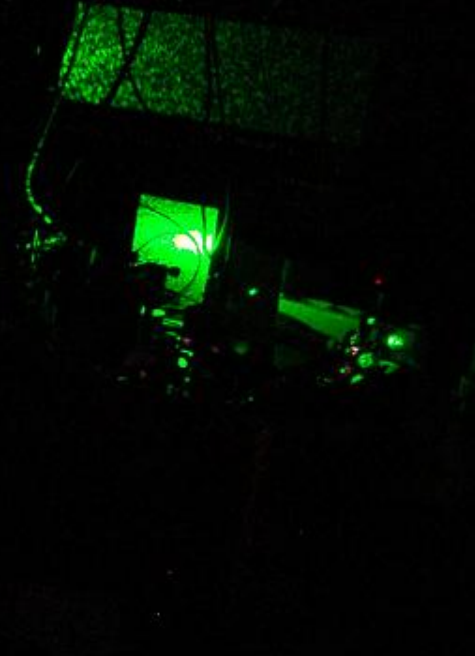
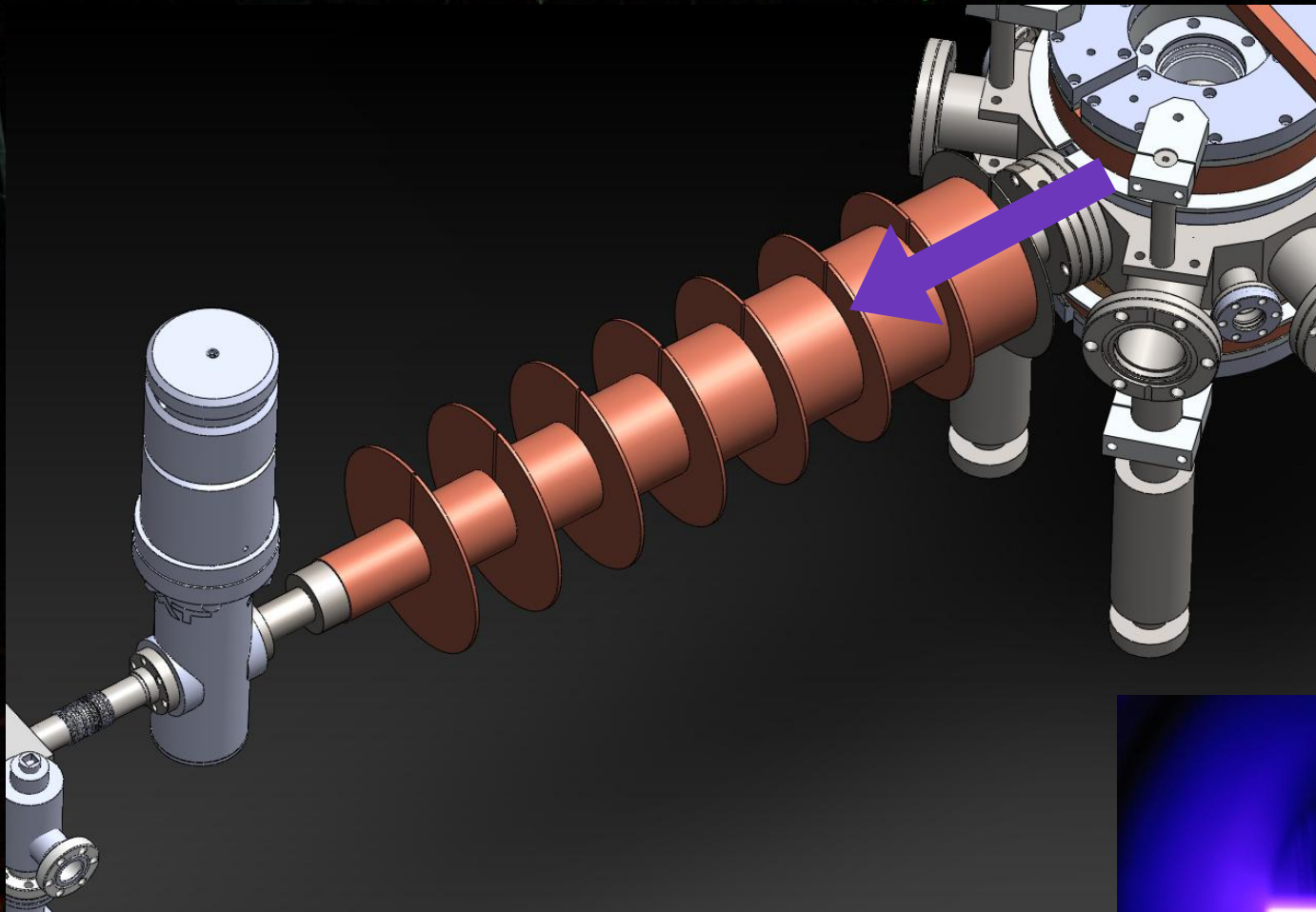


The experimental setup



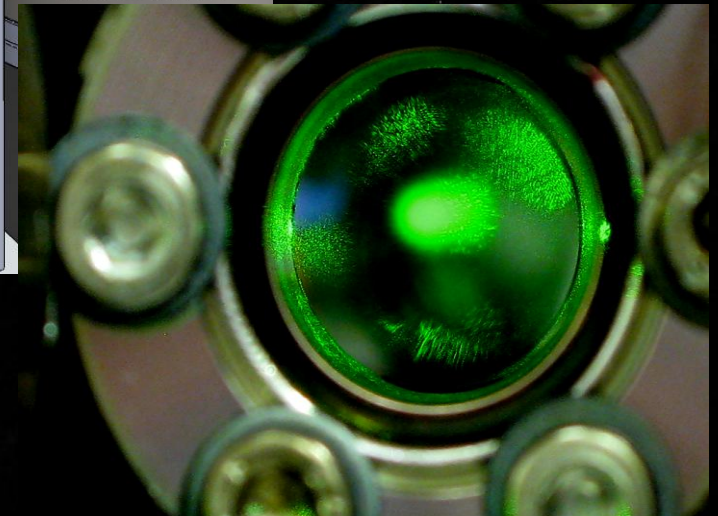
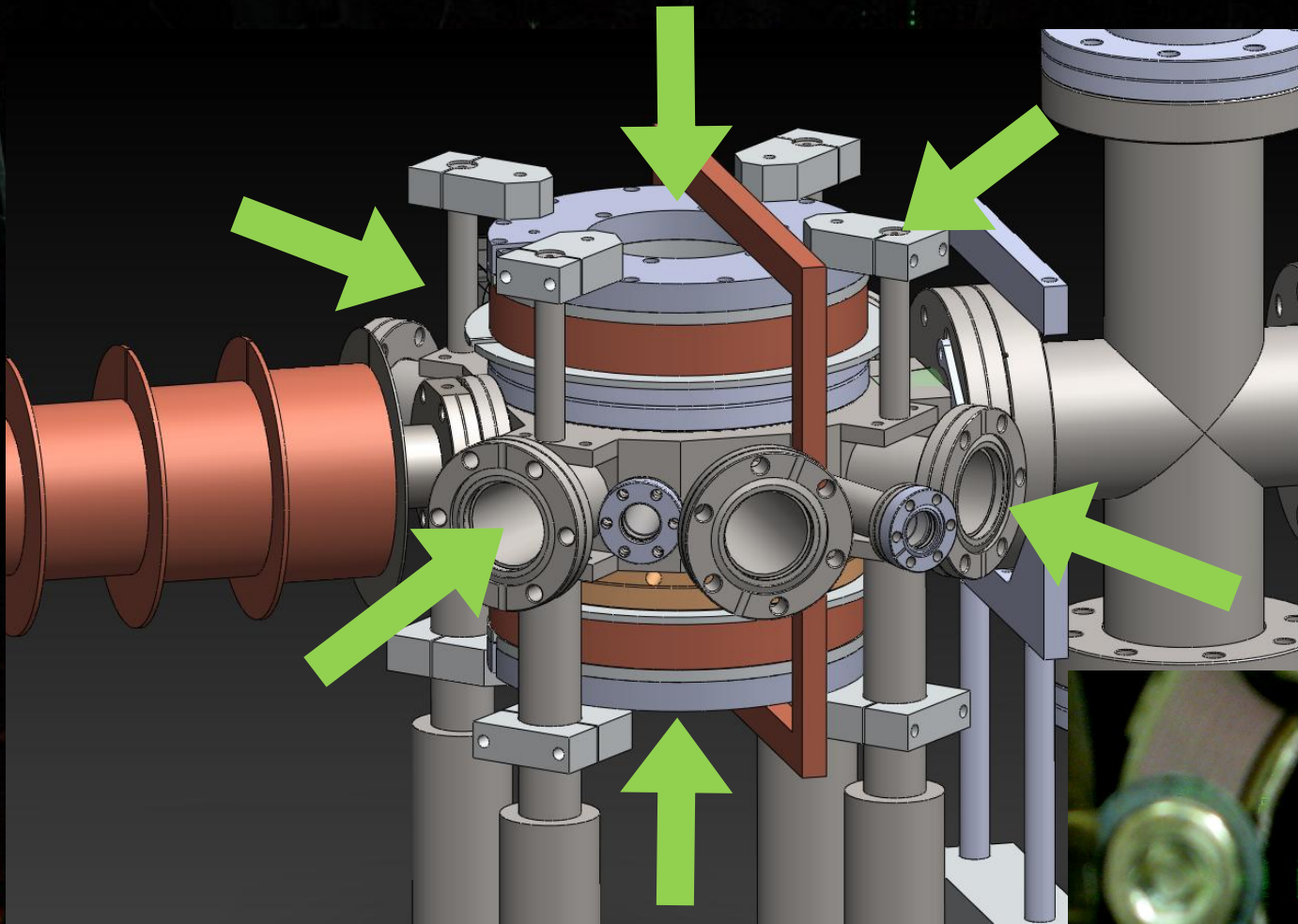
- Ytterbium loaded: 7 g
- Temperature: 800 K
- Atom velocity: ≈ 330 m/s
- Beam diameter: 5 mm

Slowing the atomic beam



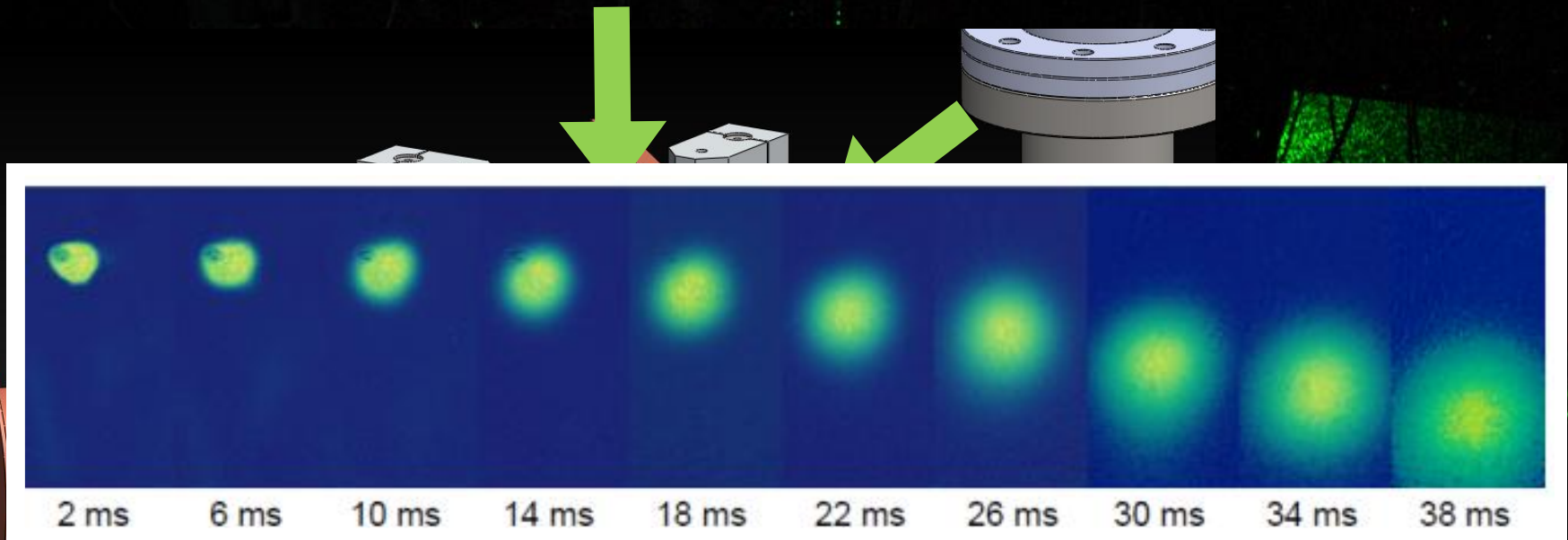
- Strong $^1S_0 \rightarrow ^1P_1$ transition (399 nm)
- Final atom velocity: ≈ 10 m/s

The green MOT

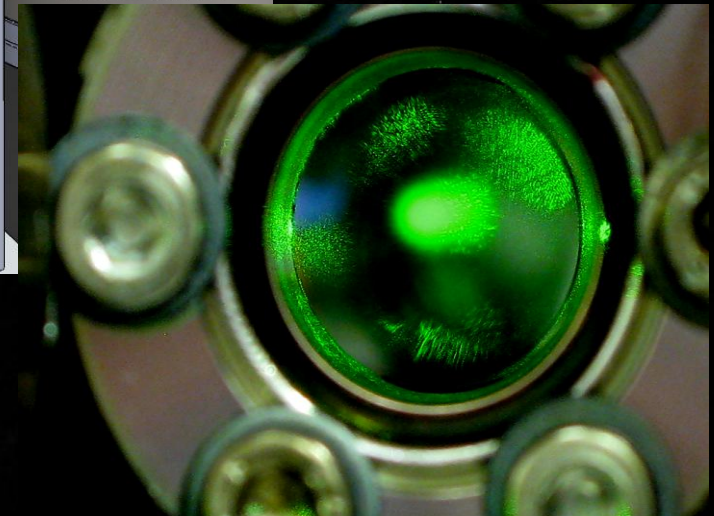


- Narrow $^1S_0 \rightarrow ^3P_1$ transition (556 nm)
- Temperature: $\approx 30 \mu\text{K}$
- Number of atoms: $\approx 2 \cdot 10^9$

The green MOT



- Narrow $^1S_0 \rightarrow ^3P_1$ transition (556 nm)
- Temperature: $\approx 30 \mu\text{K}$
- Number of atoms: $\approx 2 \cdot 10^9$




Optical trapping

Diamagnetic ground state: no magnetic trapping

Optical trap: spatially-dependent ac-Stark shift induced by off-resonant light


1P_1




3P_2




3P_1



3P_0



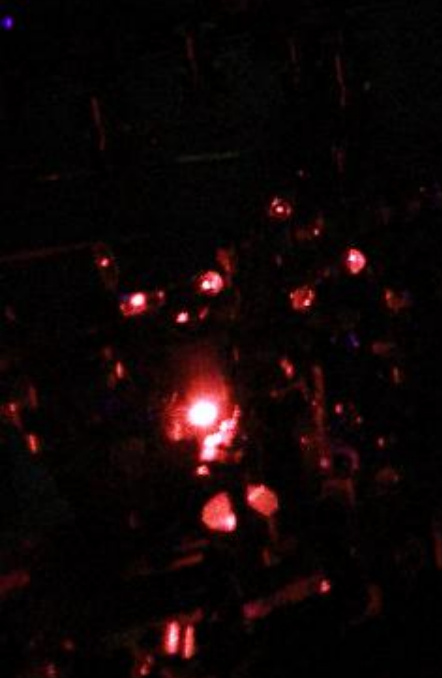
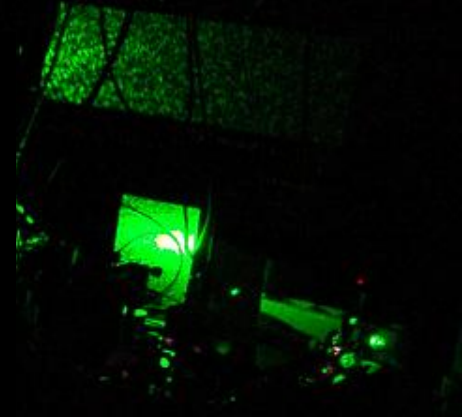
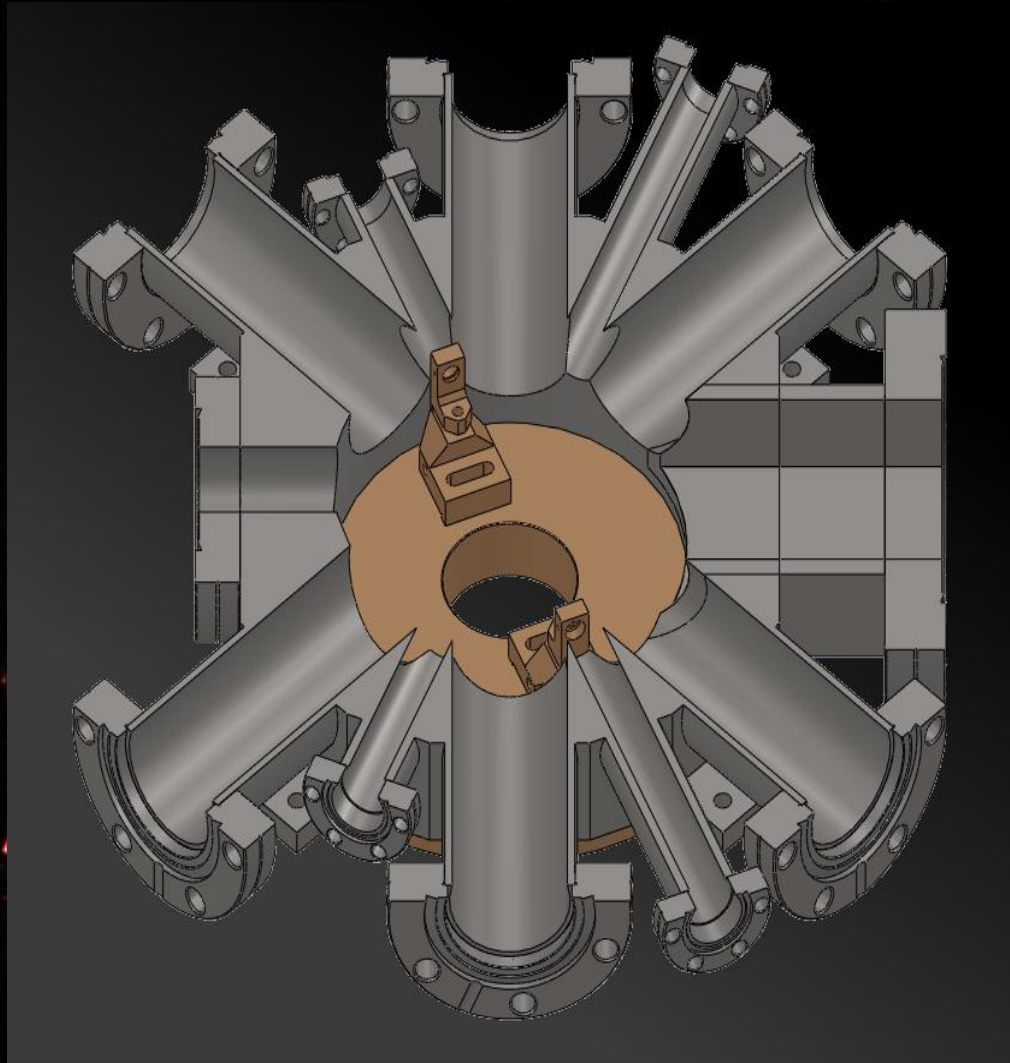
1S_0



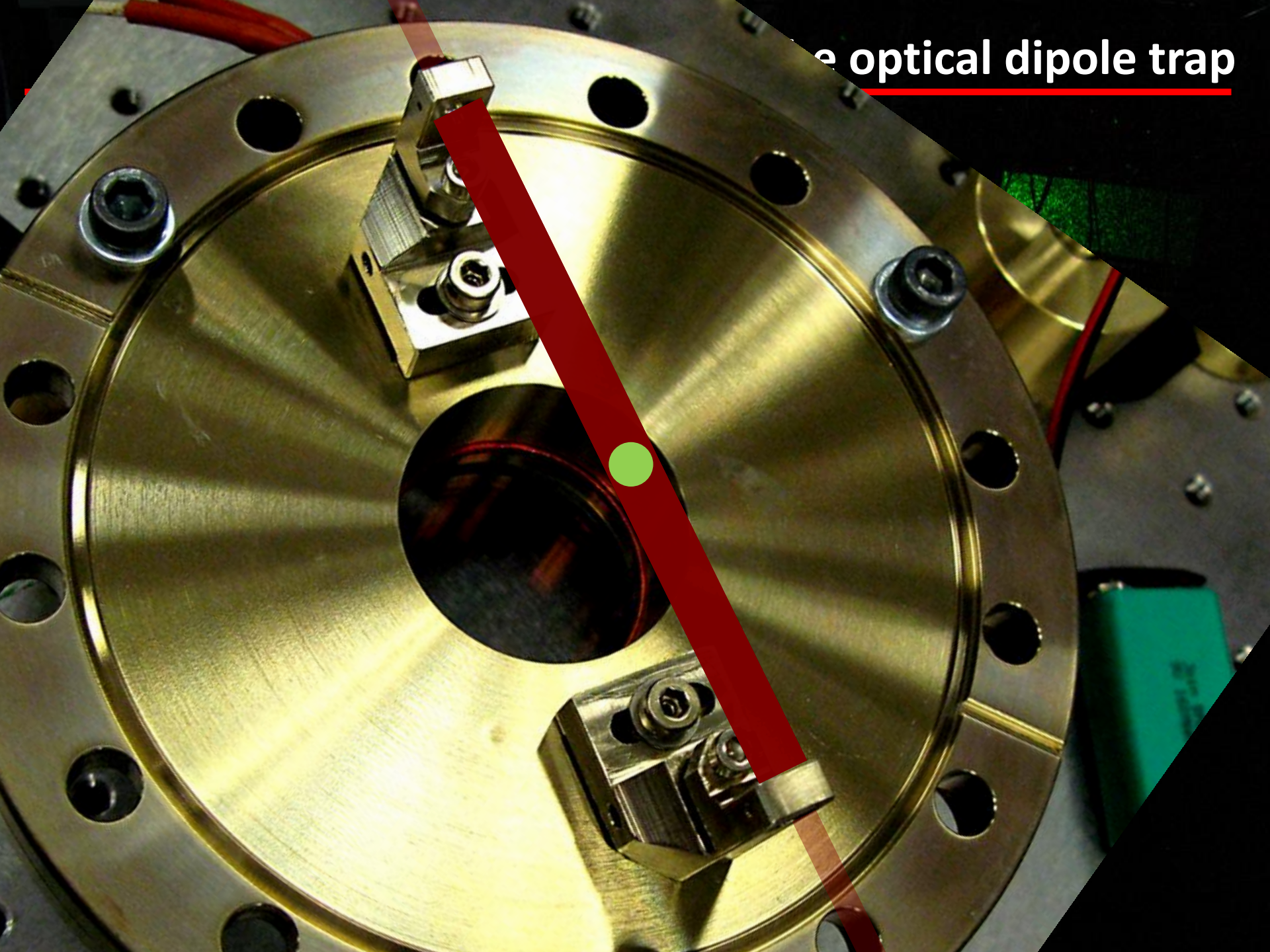
x



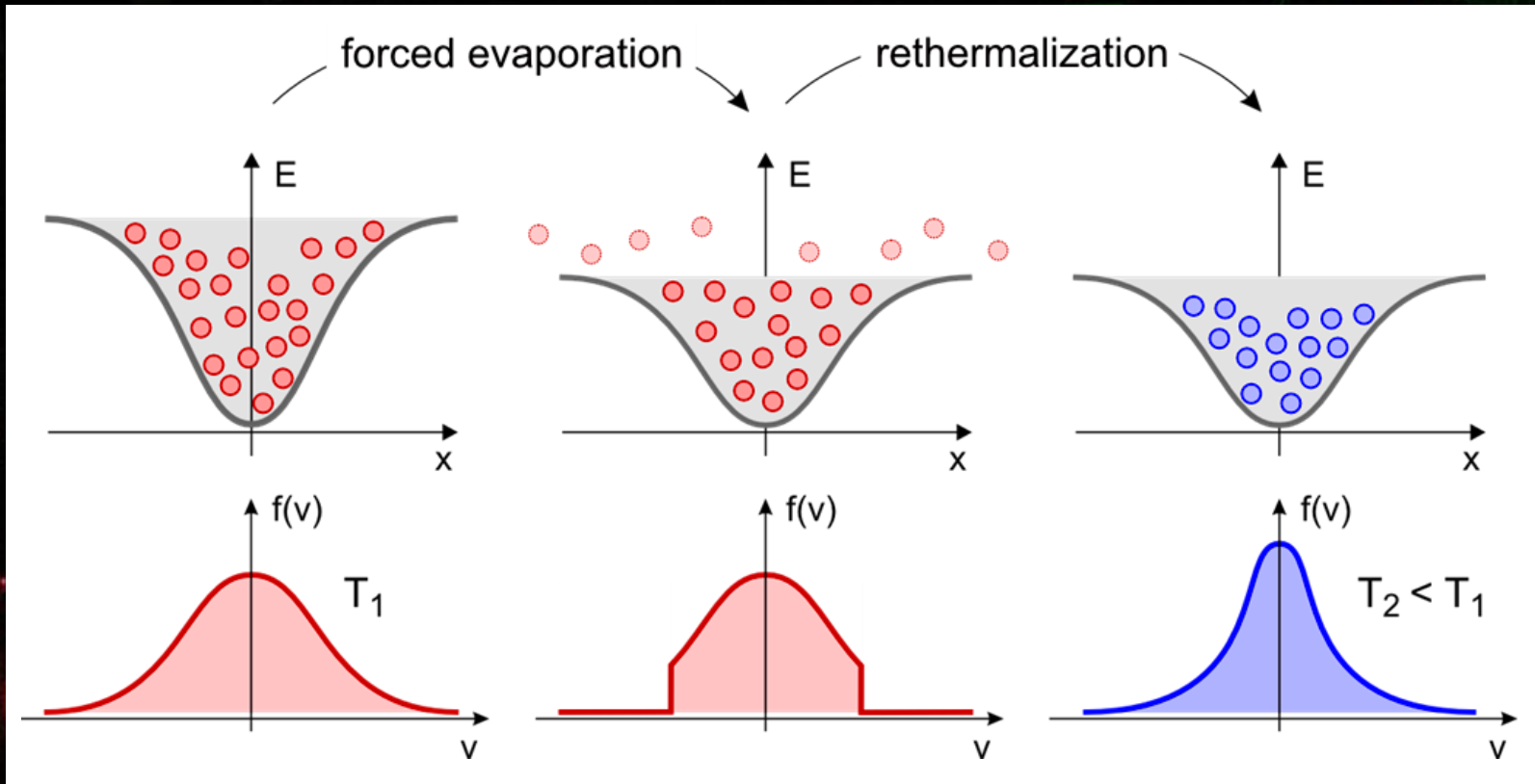
The optical dipole trap



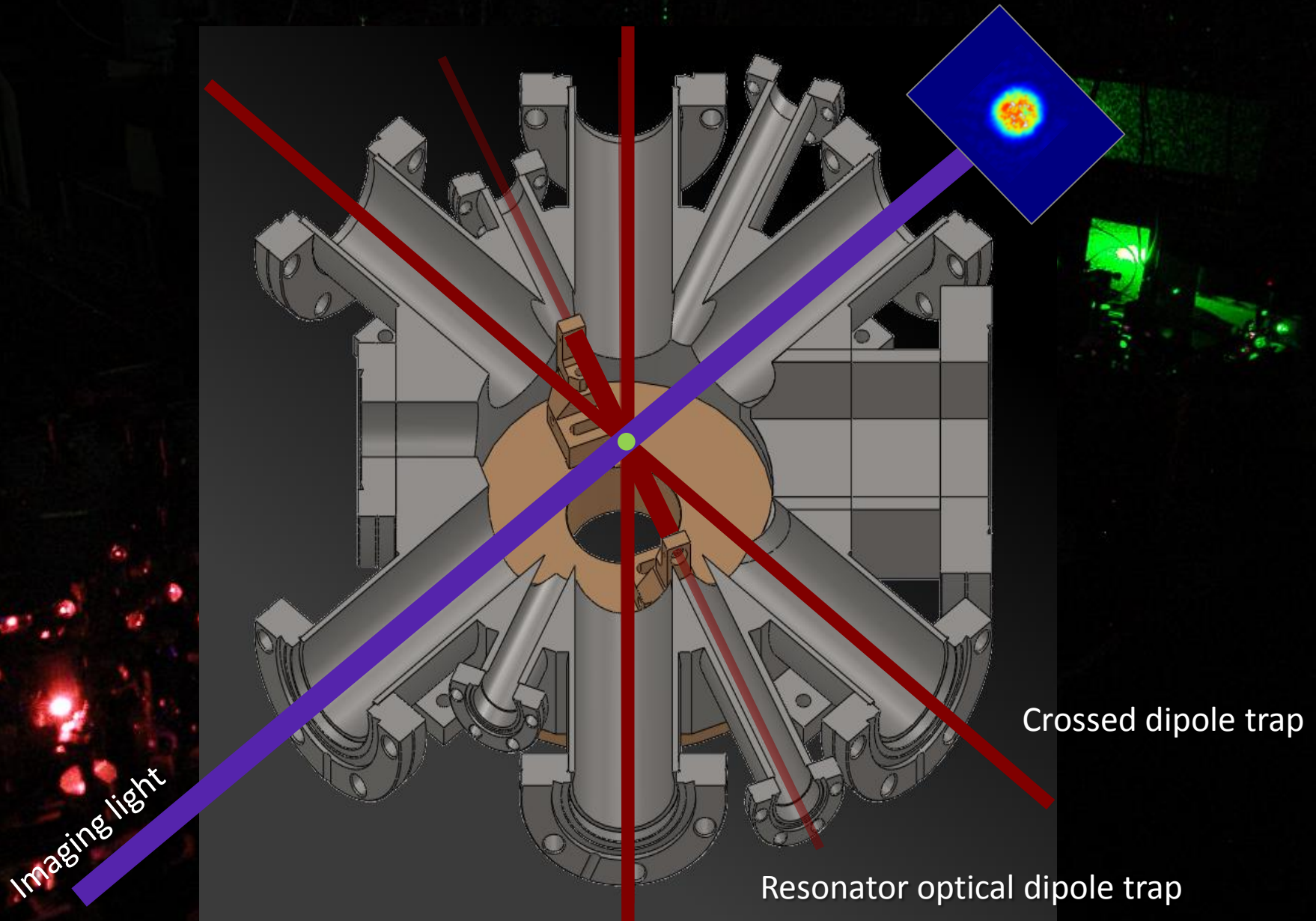
the optical dipole trap



Evaporative cooling



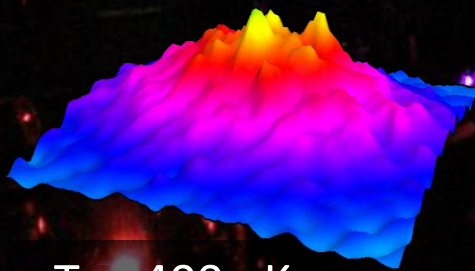
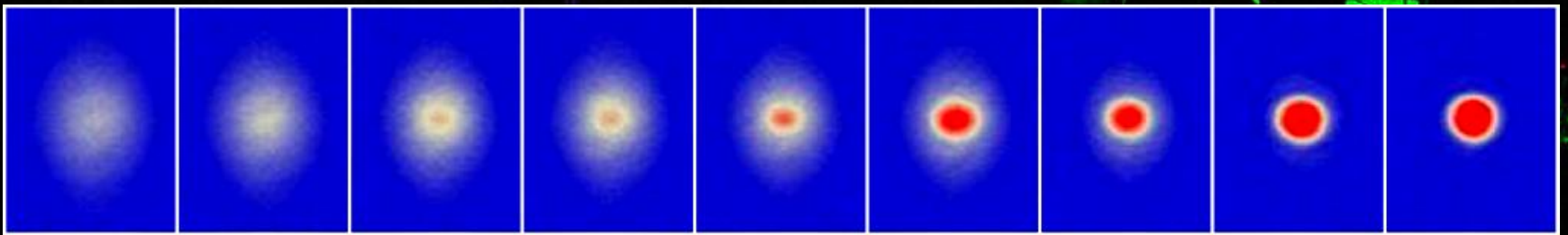
The optical dipole trap



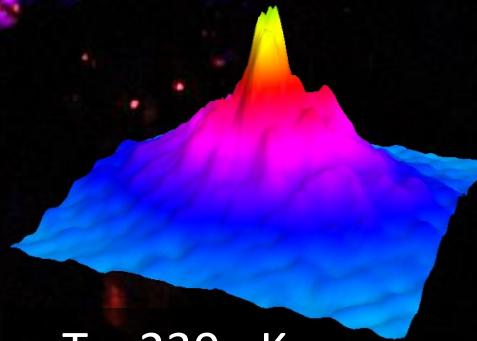
First ^{174}Yb BEC in Florence

Time-of-flight images: momentum distribution

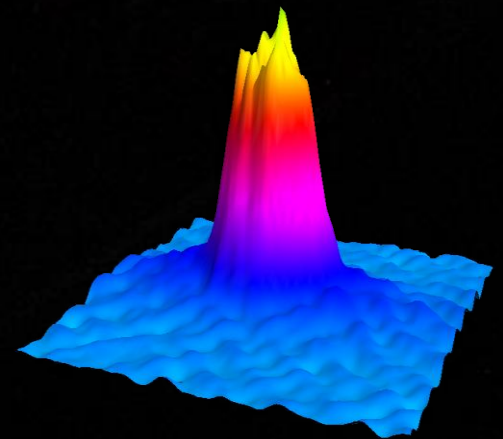
lower temperature \longrightarrow



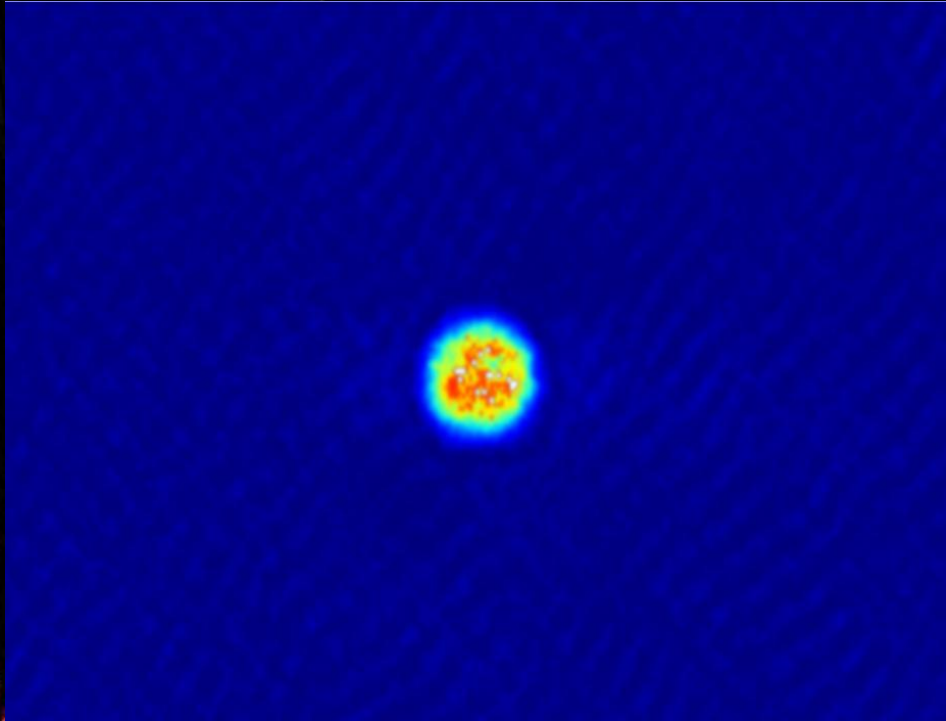
$T \approx 400$ nK



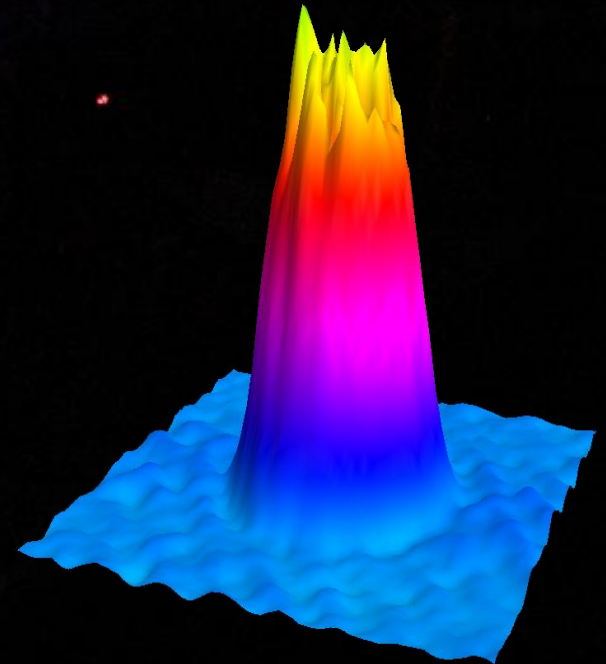
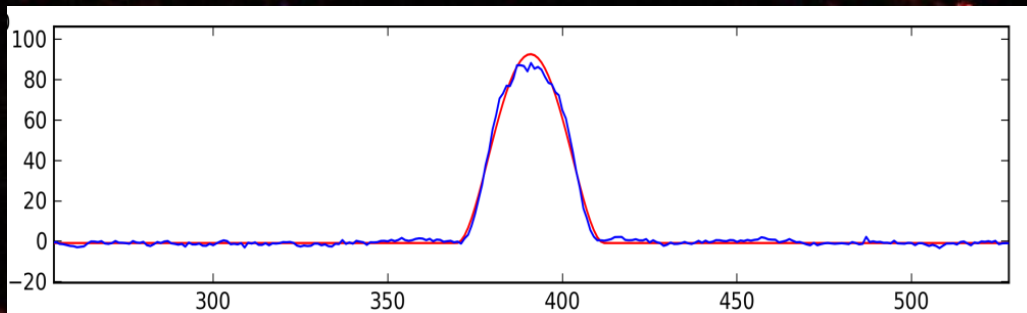
$T \approx 230$ nK



First ^{174}Yb BEC in Florence

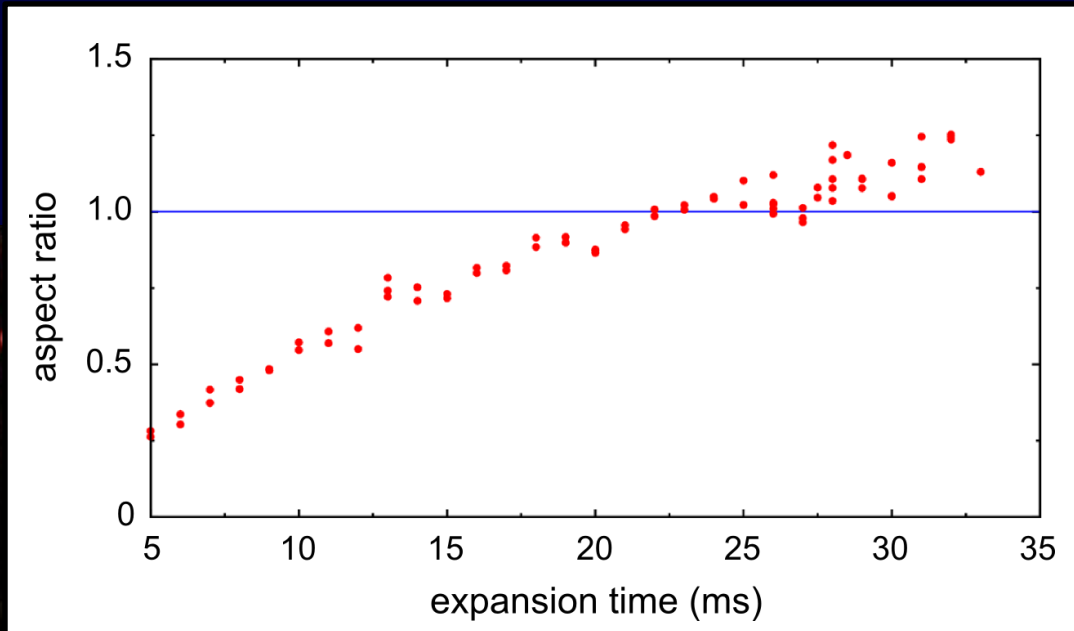
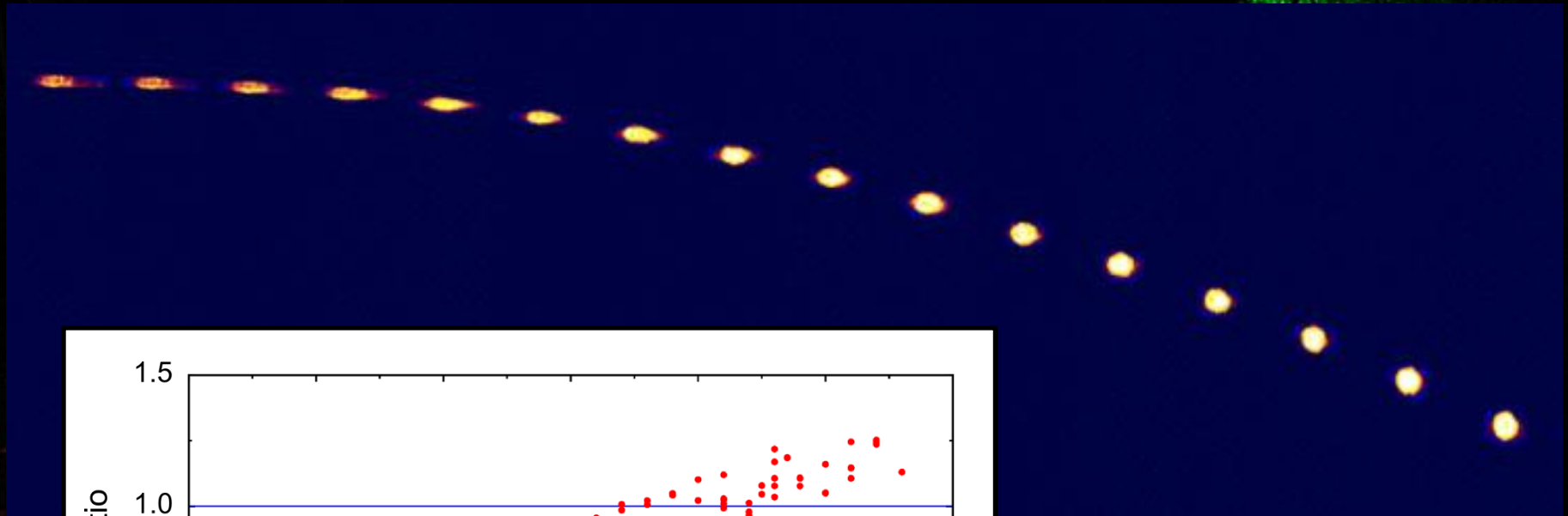


almost pure ^{174}Yb BEC
with $N = 4 \times 10^5$ atoms



First ^{174}Yb BEC in Florence

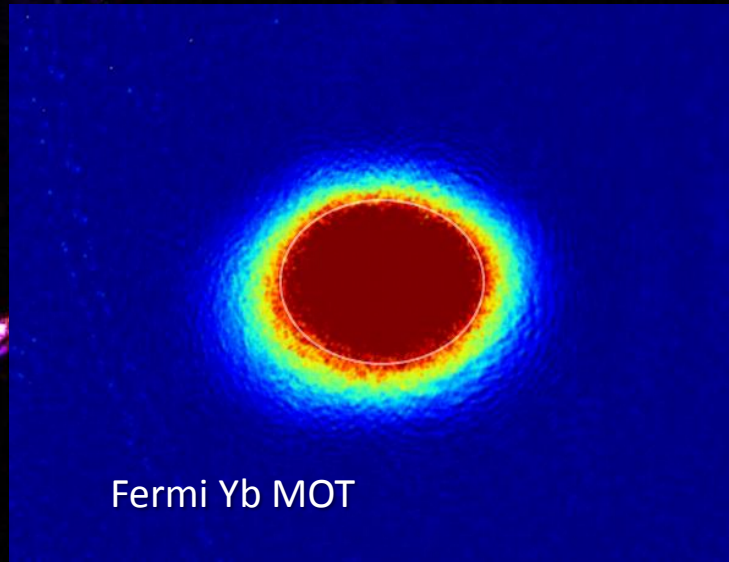
Time-of-flight measurement of anisotropic BEC expansion



Fermionic ^{173}Yb under cooling

Laser cooling and trapping of fermionic ^{173}Yb demonstrated.

Evaporative cooling in progress.



Introduction

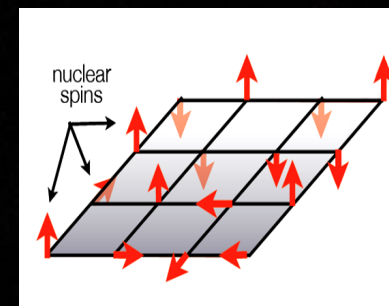
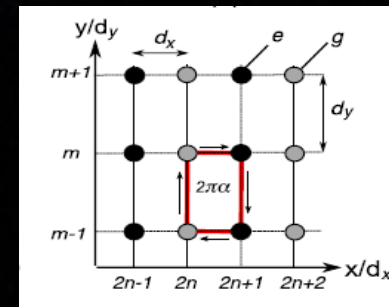
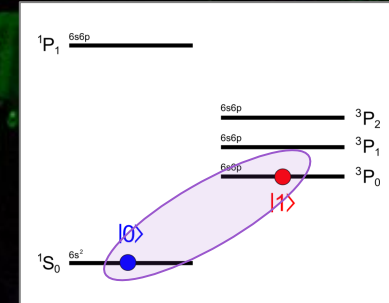
Bose-Einstein condensation of Ytterbium

Current and future work

Why Ytterbium?

Three examples:

- Quantum information
- Synthetic gauge potentials
- SU(N) physics

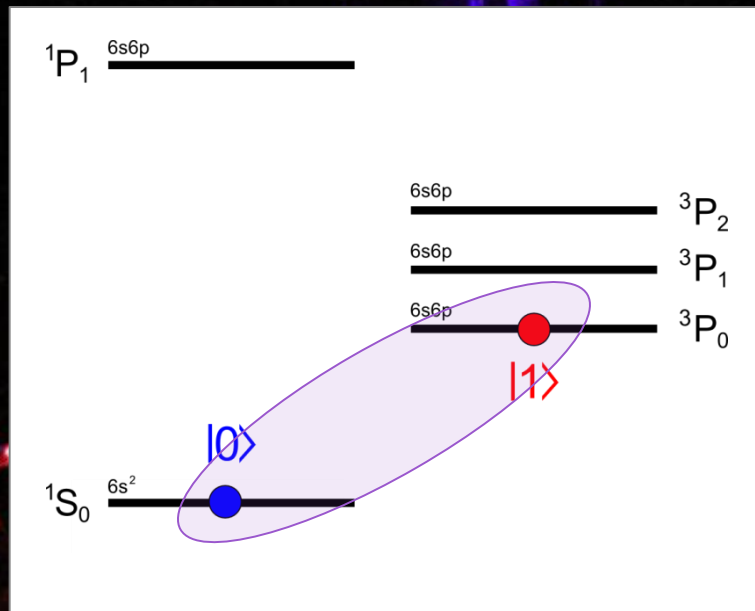


Quantum information with long-lived qubits

Two-electron atoms offer possibilities of encoding quantum information with long coherence times

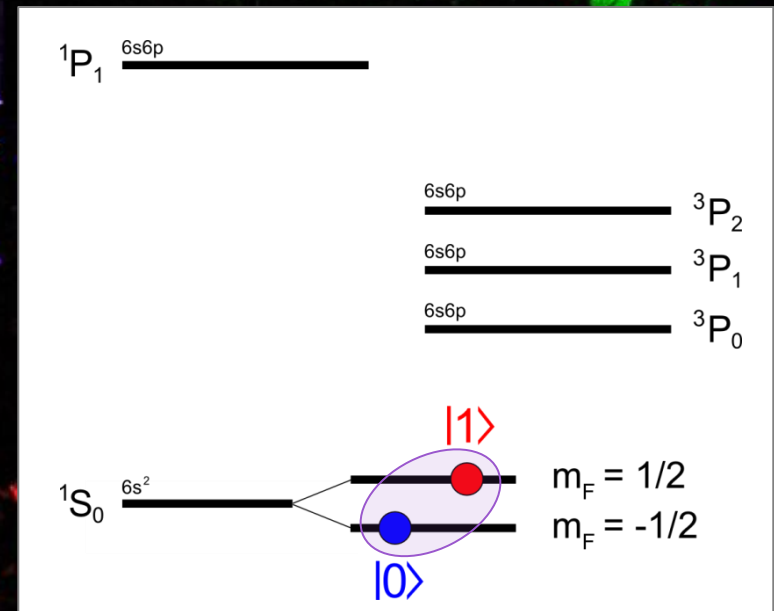
Review paper: A. Daley, arXiv:1106.5712

electronic qubits



- ultra-narrow clock transition
- long coherence times

nuclear qubits

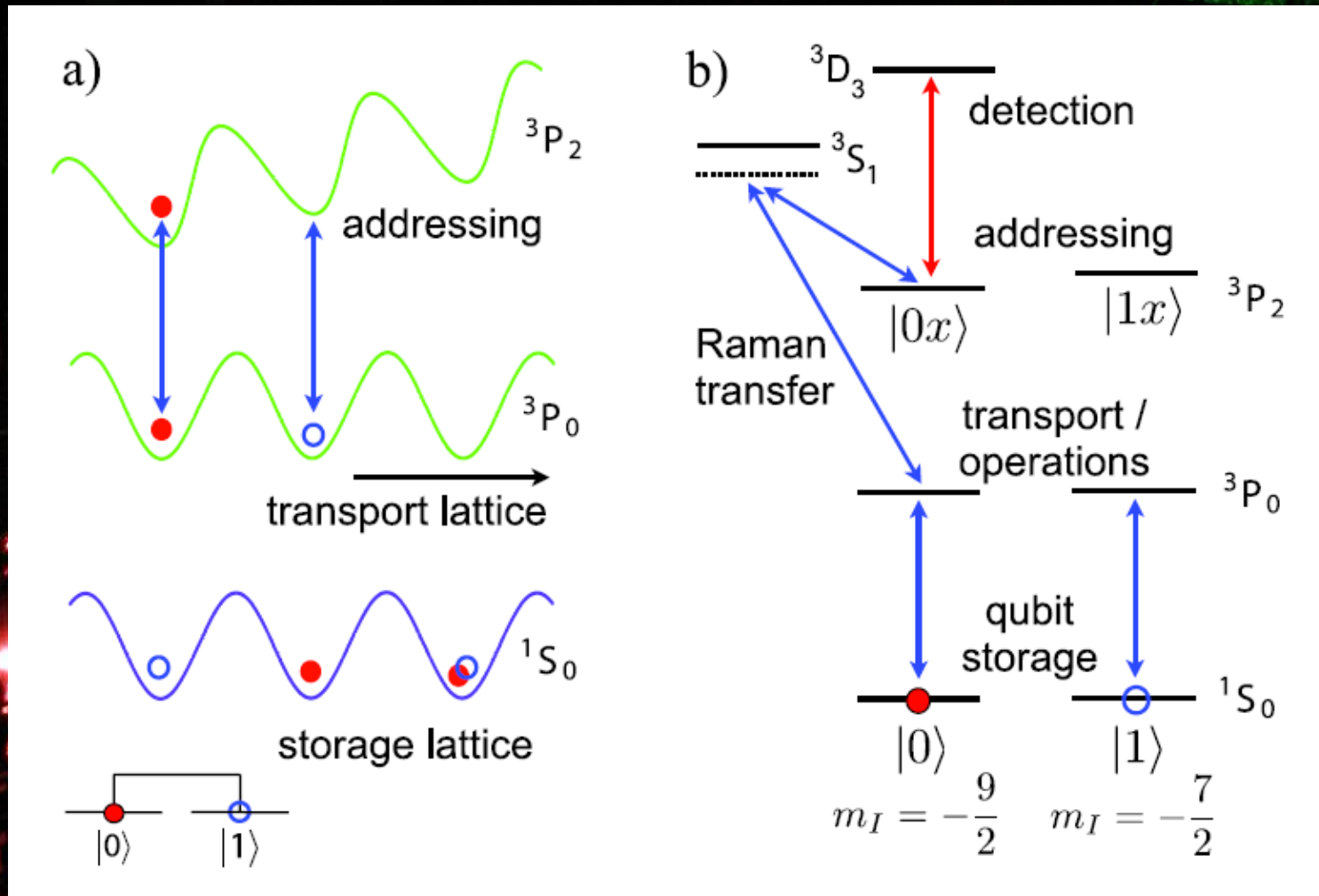


- no hyperfine interaction
- low coupling to magnetic fields

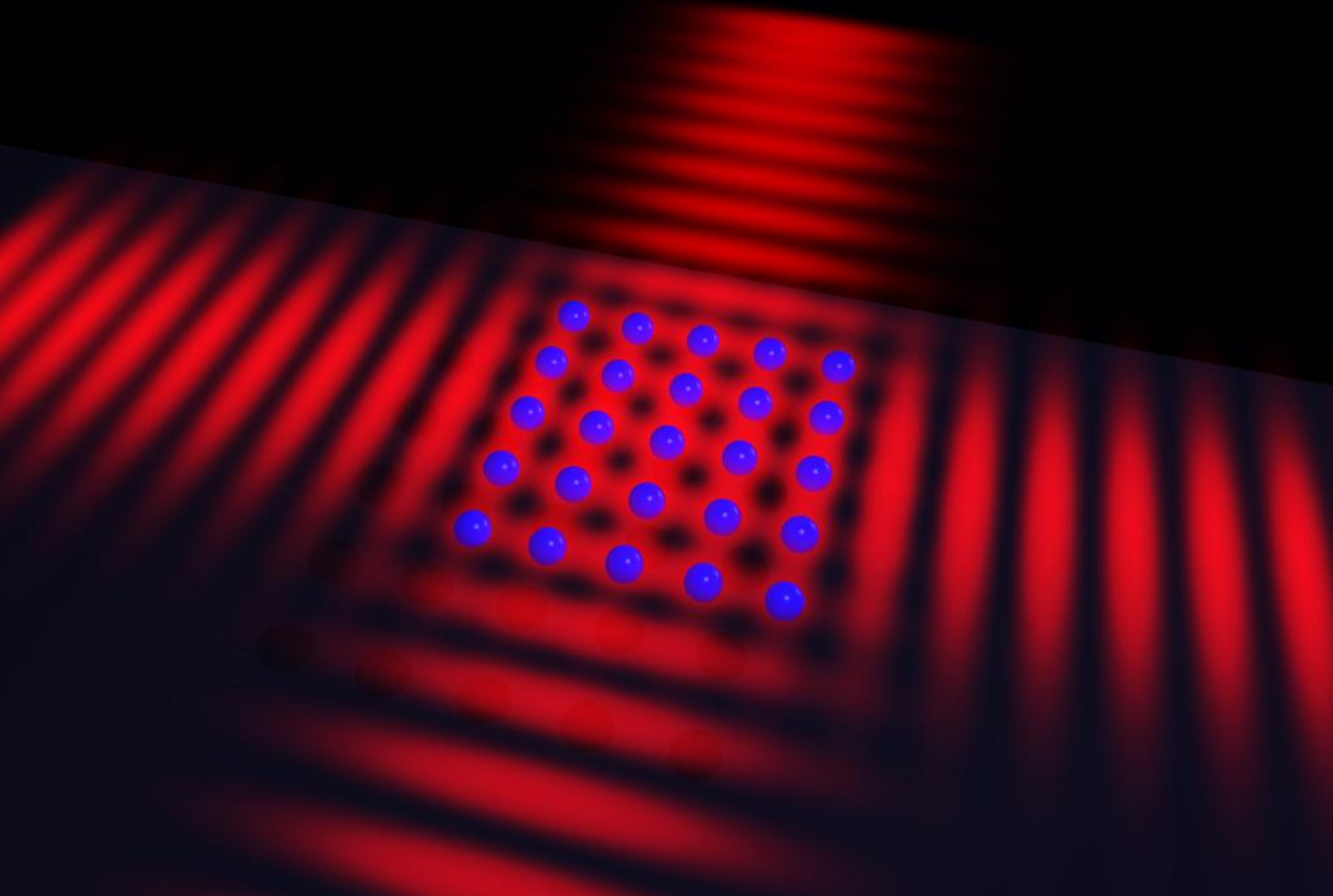
Quantum information with long-lived qubits

Quantum computing with alkaline-earth-metal atoms

A. Daley, M. M. Boyd, J. Ye, P. Zoller, PRL **101**, 170504 (2008)

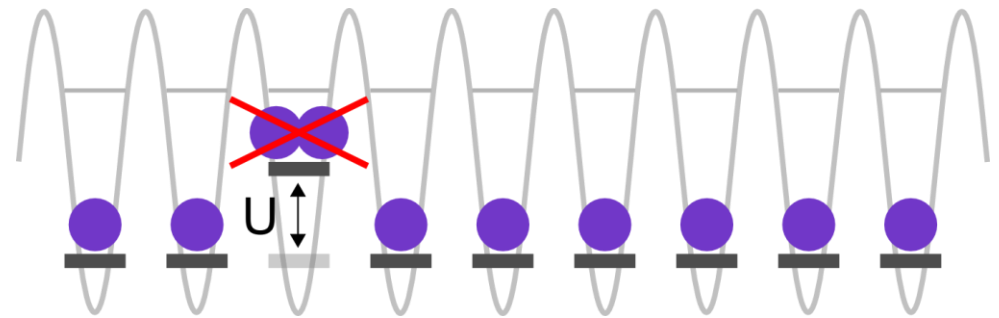


Optical lattices



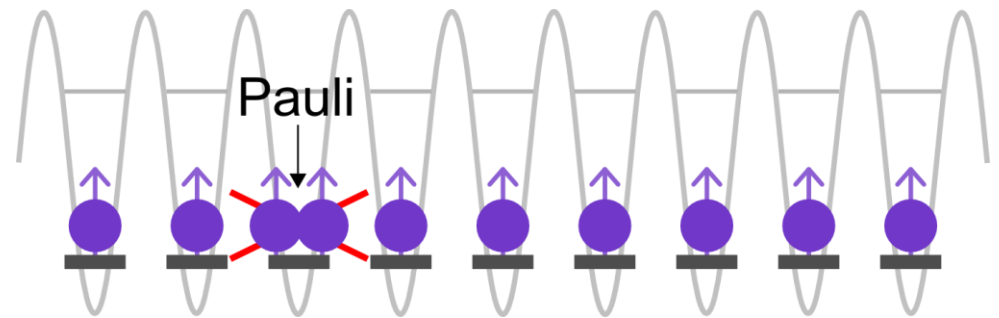
strong repulsive interactions
between bosons

MOTT INSULATOR



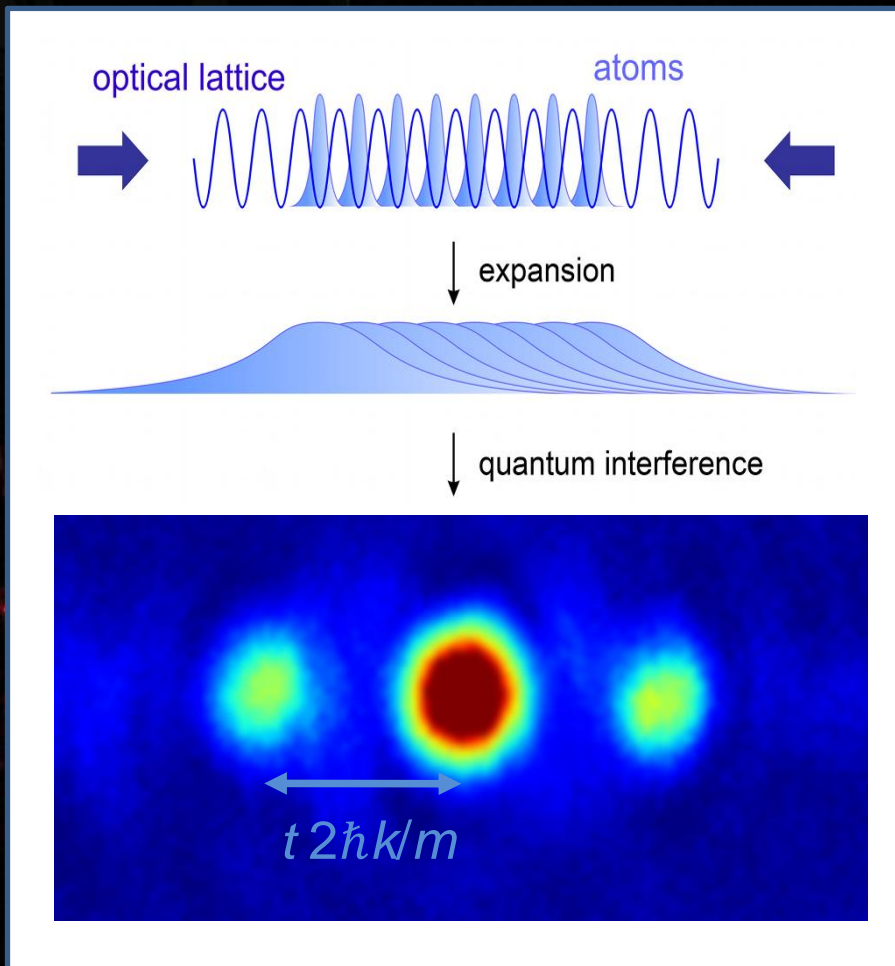
spin polarized fermions

BAND INSULATOR

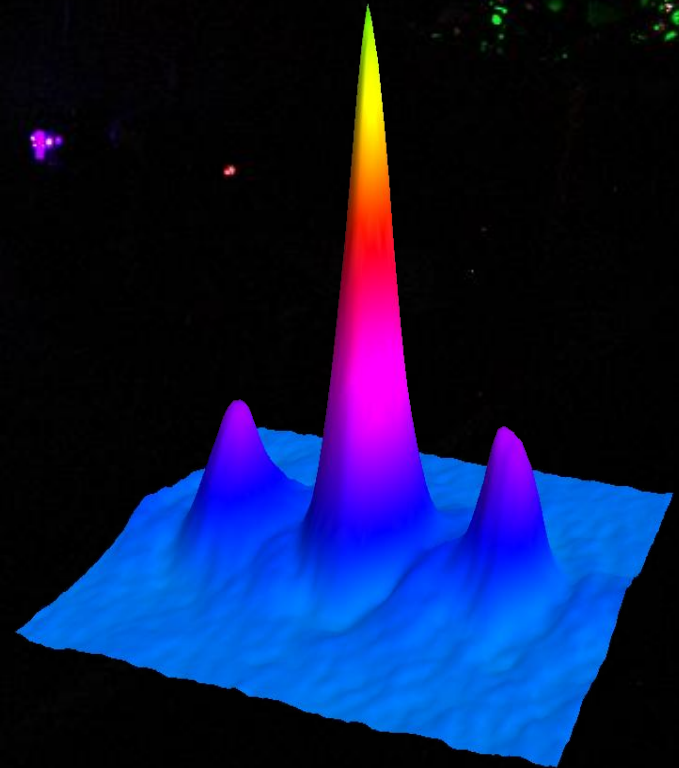


^{174}Yb BEC in optical lattice

1D optical lattice



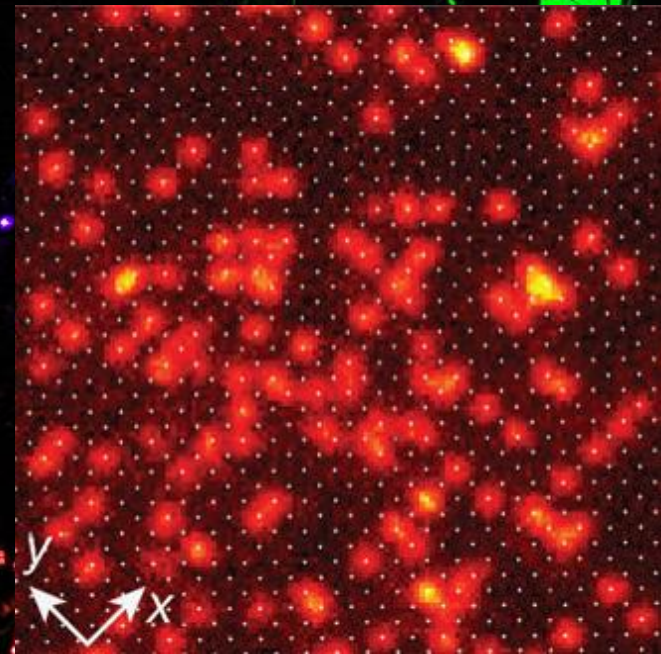
Imaging of momentum distribution
after 30 ms of free expansion



Single-site high-resolution imaging

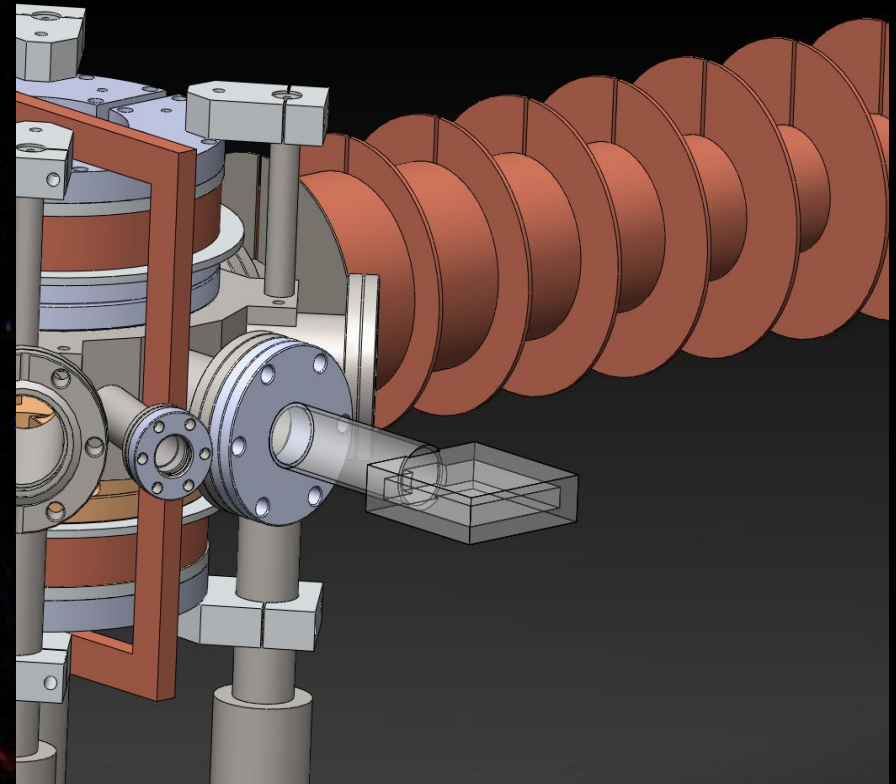
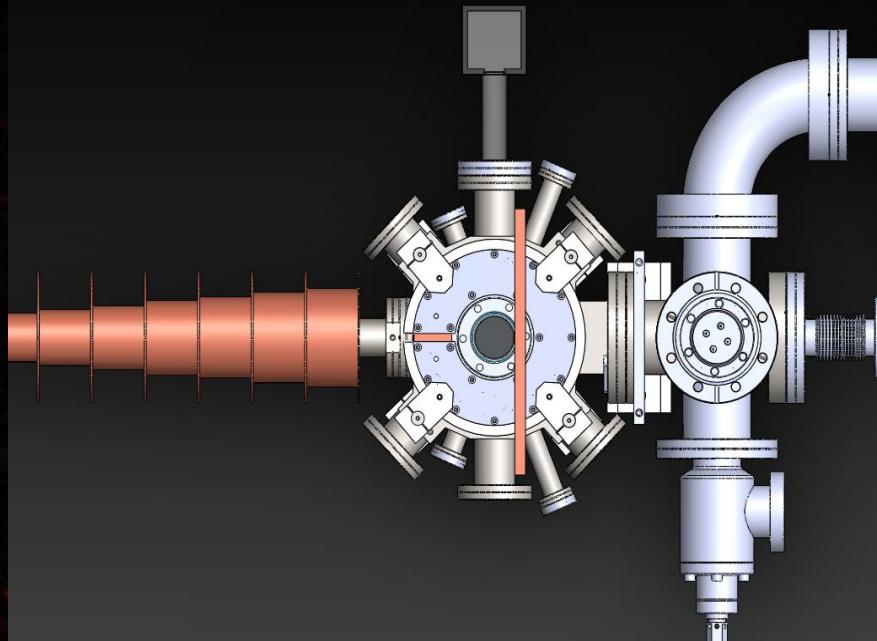


*W. S. Bakr et al.,
Science 329, 547 (2010).*



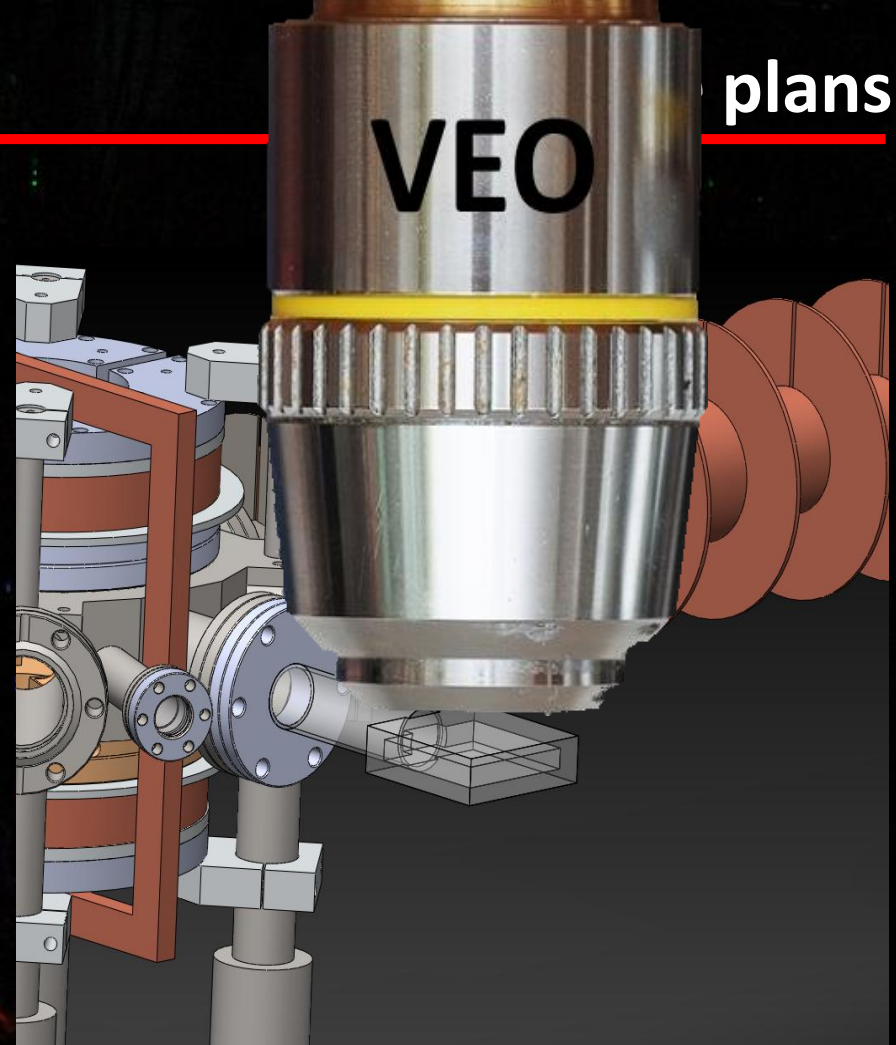
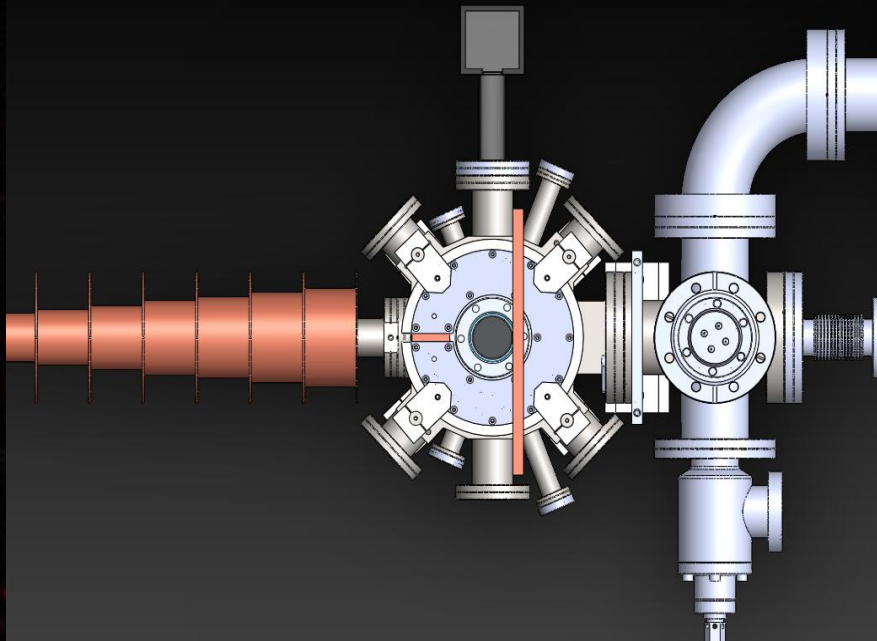
*J. F. Sherson et al.,
Nature 467, 68 (2010).*

Glass cell with large optical access
for high-resolution imaging



plans

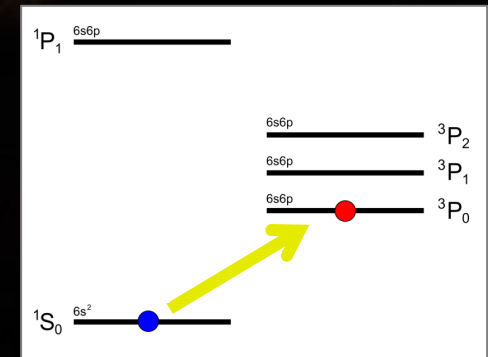
Glass cell with large optical access
for high-resolution imaging



VEO = Very Expensive Objective!

Excitation of the 3P_0 state

Yellow laser @ 578nm for the clock transition $^1S_0 - ^3P_0$



Quantum dot laser 190 mW @ 1156 nm
SHG in bowtie cavity with a PPMgO:CLN crystal (≈ 50 mW)

Narrowing & stabilization by locking to ULE cavity in progress

Synthetic gauge potentials

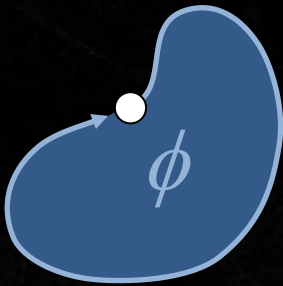
Abelian gauge potentials

→ Artificial magnetic field
QHE (integer and fractional)

$$\hat{H} = \frac{1}{2m} (\mathbf{p} - q\mathbf{A})^2$$

Aharonov-Bohm geometric phase for
the closed loop of an electron in a magnetic field

$$\psi \rightarrow e^{i\phi} \psi \quad \phi = 2\pi \frac{\Phi}{\Phi_0} \quad \Phi_0 = \frac{h}{e}$$



Non-Abelian gauge potentials

→ Non-Abelian anyons
Fractional statistics
Topological insulators

$$|\psi\rangle \rightarrow \hat{U} |\psi\rangle$$

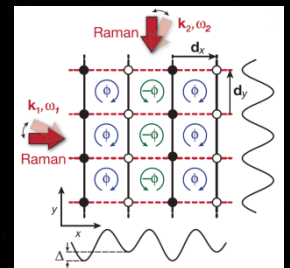
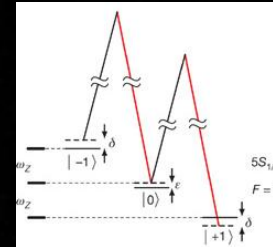
U unitary transformation
of a multi-component wavefunction

Synthetic gauge potentials

Different ways to produce artificial (Abelian) gauge potentials

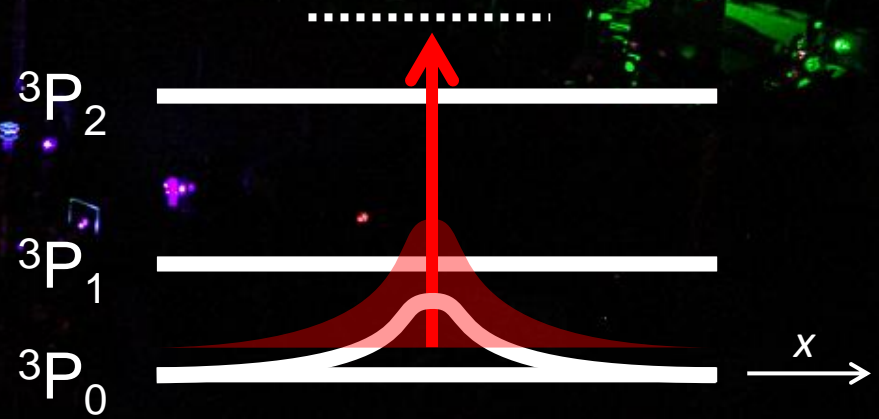
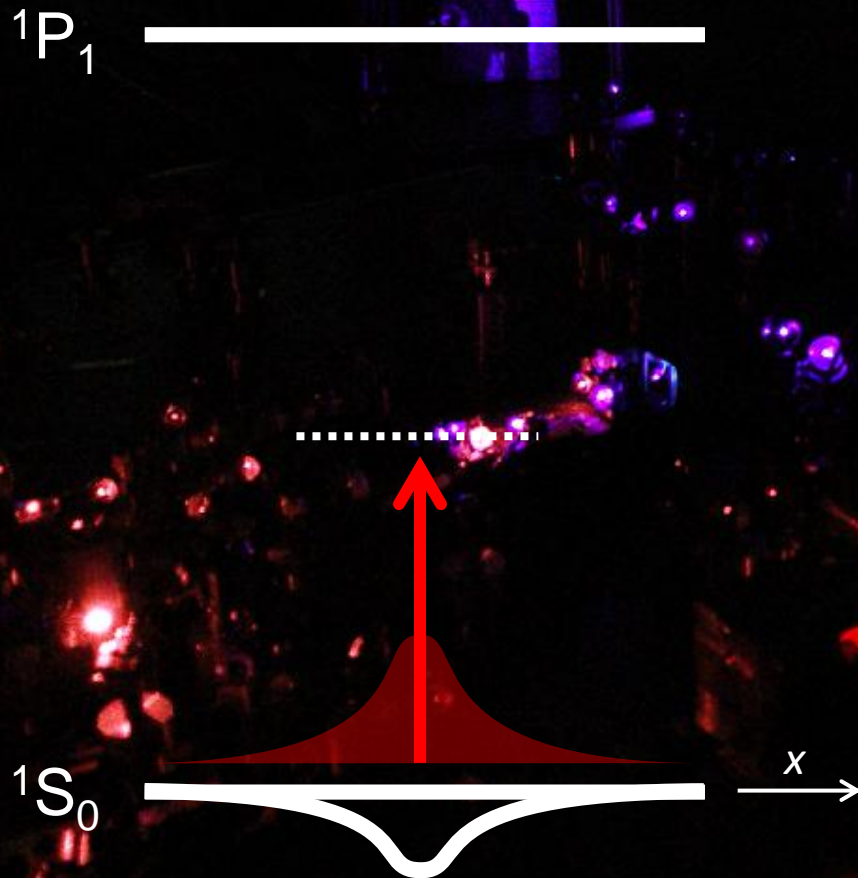
J. Dalibard et al., Rev. Mod. Phys. **83**, 1523 (2011)

- Rotating traps
- Optical dressing in multilevel atoms
Y.-J. Lin et al., Nature **462**, 628 (2009).
- Laser-assisted tunnelling in state-dependent lattices
M. Aidelsburger et al., Phys. Rev. Lett. **107**, 255301 (2011).



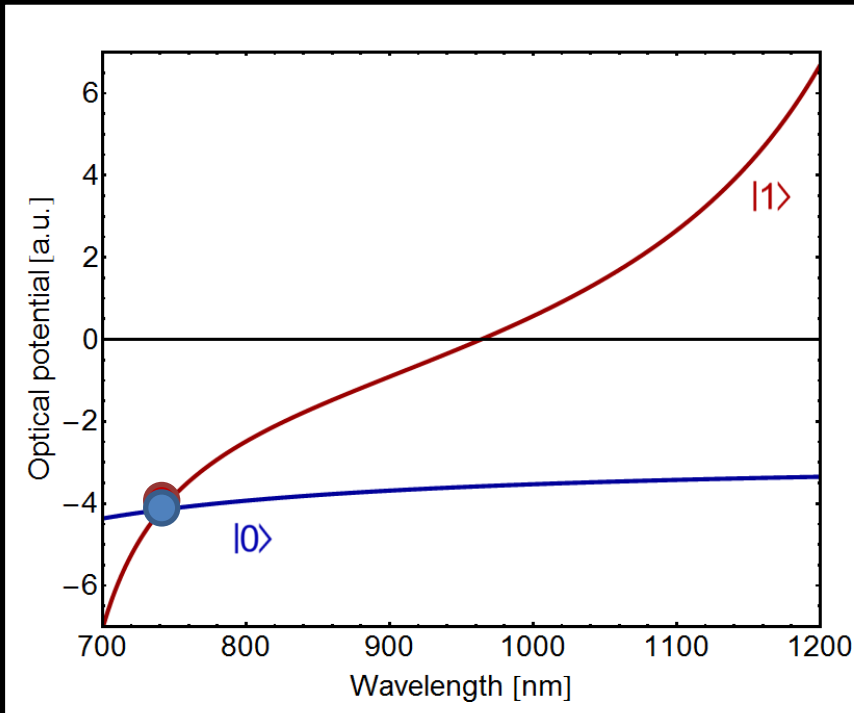
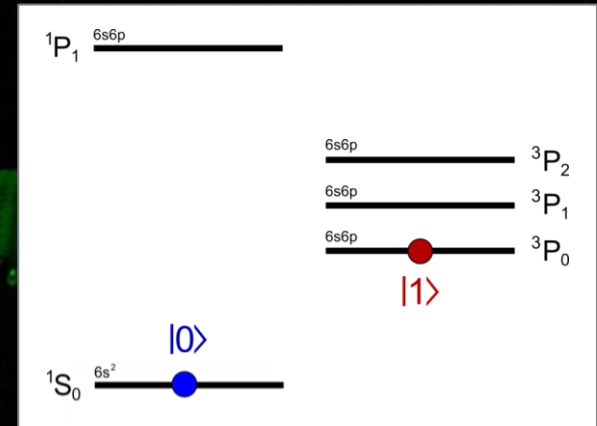
State-dependent optical trapping

Spatially-dependent ac-Stark shift induced by off-resonant light

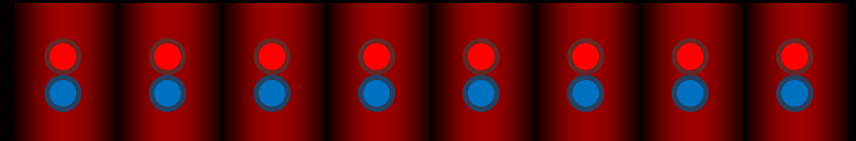


Synthetic gauge potentials

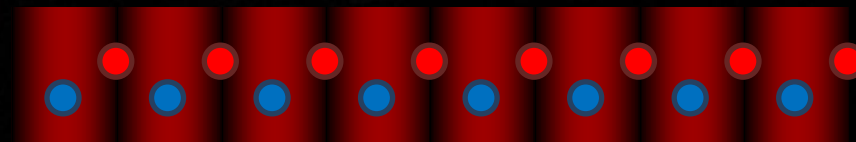
State-dependent potentials for Ytterbium



Magic wavelength 760nm



Antimagic wavelength 1120nm



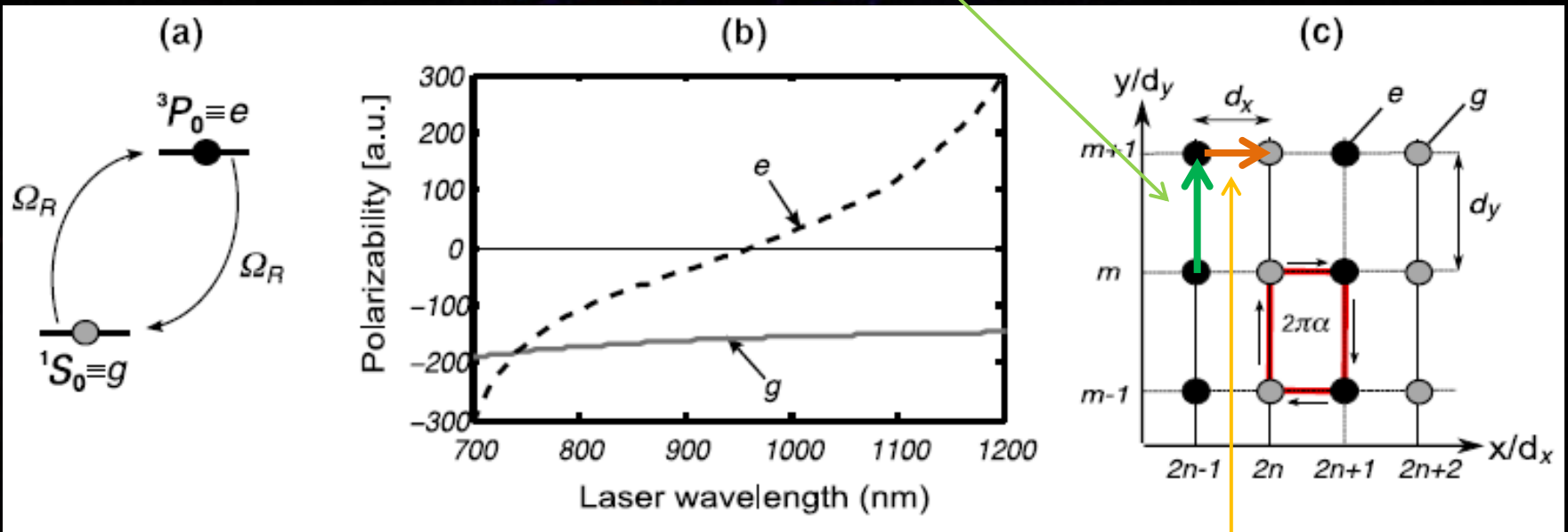
Synthetic gauge potentials

Laser-assisted tunnelling in state-dependent potentials

D. Jaksch and P. Zoller, *New J. Phys.* **5**, 56 (2003)

F. Gerbier and J. Dalibard, *New J. Phys.* **12**, 033007 (2010)

Ordinary tunnelling: J

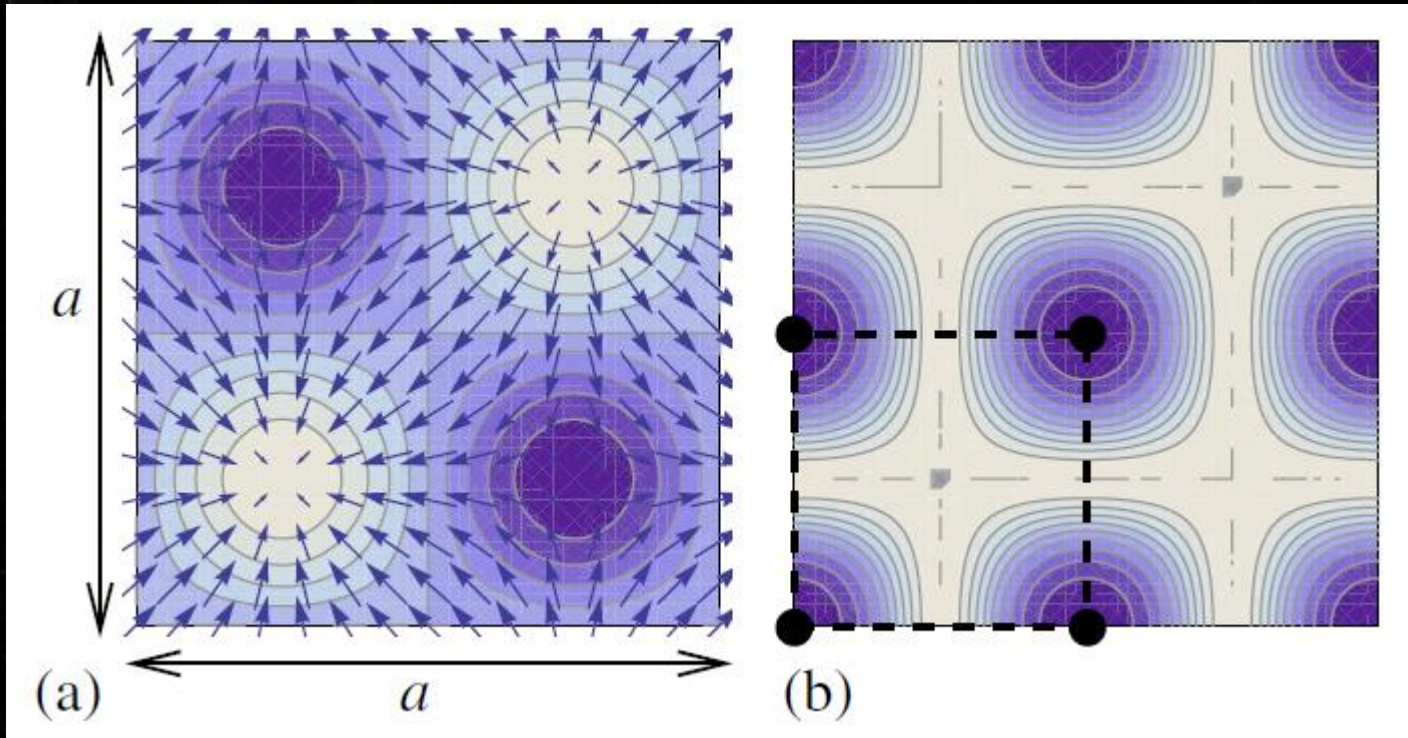


Laser-assisted tunnelling: $J \exp(ikx)$

Synthetic gauge potentials

Optical flux lattices

N. Cooper, PRL **106**, 175301 (2011)



Absence of hyperfine interaction



Interaction strength between different nuclear spin states are the same!

SU(2I+1) symmetry

SU(2) for ^{171}Yb

$I=1/2$



SU(6) for ^{173}Yb

$I=5/2$



SU(N) magnetism

Example: interacting fermions (repulsive) on a square lattice

Fermi-Hubbard model

$$\hat{H} = t \sum_{\langle ij \rangle, \alpha=1, \dots, N} \left(\hat{f}_{i, \alpha}^\dagger \hat{f}_{j, \alpha} + h.c. \right) + U \sum_{i, \alpha, \beta} \hat{f}_{i, \alpha}^\dagger \hat{f}_{i, \alpha} \hat{f}_{i, \beta}^\dagger \hat{f}_{i, \beta}$$

↓ $U \gg t$

SU(N) symmetric Heisenberg model

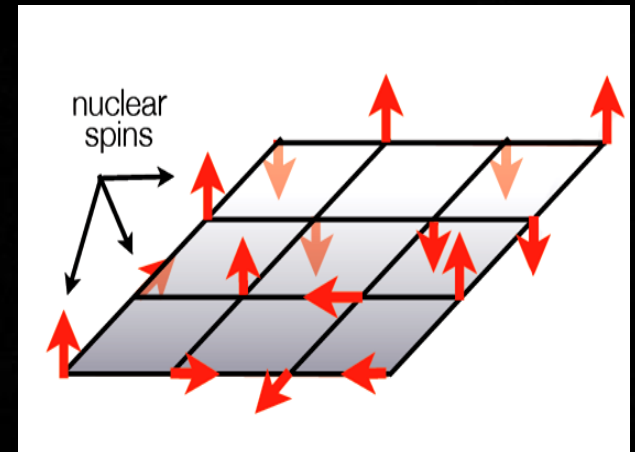
$$S_{\alpha\beta}(\mathbf{r}) = f_{r\alpha}^\dagger f_{r\beta} \quad \text{SU(N) spin}$$

$$\mathcal{H} = J \sum_{\langle \mathbf{r}\mathbf{r}' \rangle} S_{\alpha\beta}(\mathbf{r}) S_{\beta\alpha}(\mathbf{r}')$$

$$J = \frac{2t^2}{U}$$

superexchange interaction

Independent of spin projection α !

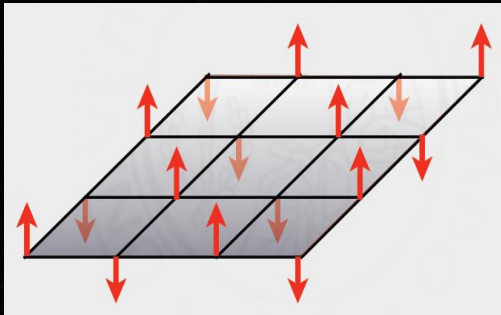


SU(N) magnetism

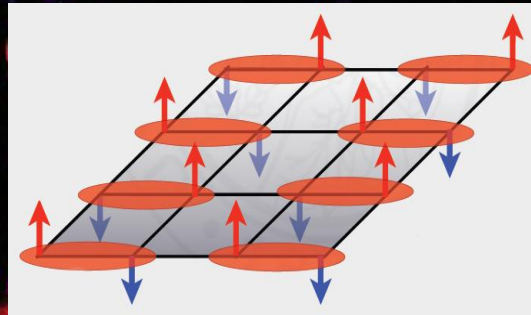
Increased symmetry \longrightarrow Exotic ground states, Topological excitations

Possible ground states (phase diagram largely unknown):

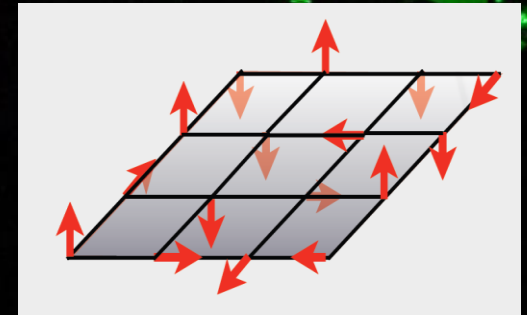
Neel state



Valence Bond Solids



Chiral Spin Liquids



Figures from V. Gurarie

KITP "Beyond Standard Optical Lattices" (2010) online talk

Non-Abelian excitations
Fractional statistics

Some references to SU(N):

M. A. Cazalilla et al., *New J. Phys.* **11**, 103033 (2009).

M. Hermele et al., *Phys. Rev. Lett.* **103**, 135301 (2009).

A. V. Gorshko et al., *Nature Physics* **6**, 289 (2010).

Conclusions

Key properties of ytterbium:

Many isotopes
Metastable states
Ultra-narrow transitions
Purely nuclear spin
State-selective optical potentials

Experiment at Lens:

^{174}Yb Bose-Einstein condensation
 ^{173}Yb Fermi gas under cooling

What can be studied:

Long coherence times for Q.I.
Synthetic gauge potentials
SU(N) physics