

Dynamic correlations,
fluctuation-dissipation relations and
effective temperatures
after a quantum quench (of the Ising chain)



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In collaboration with:

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(LPTHE Paris)

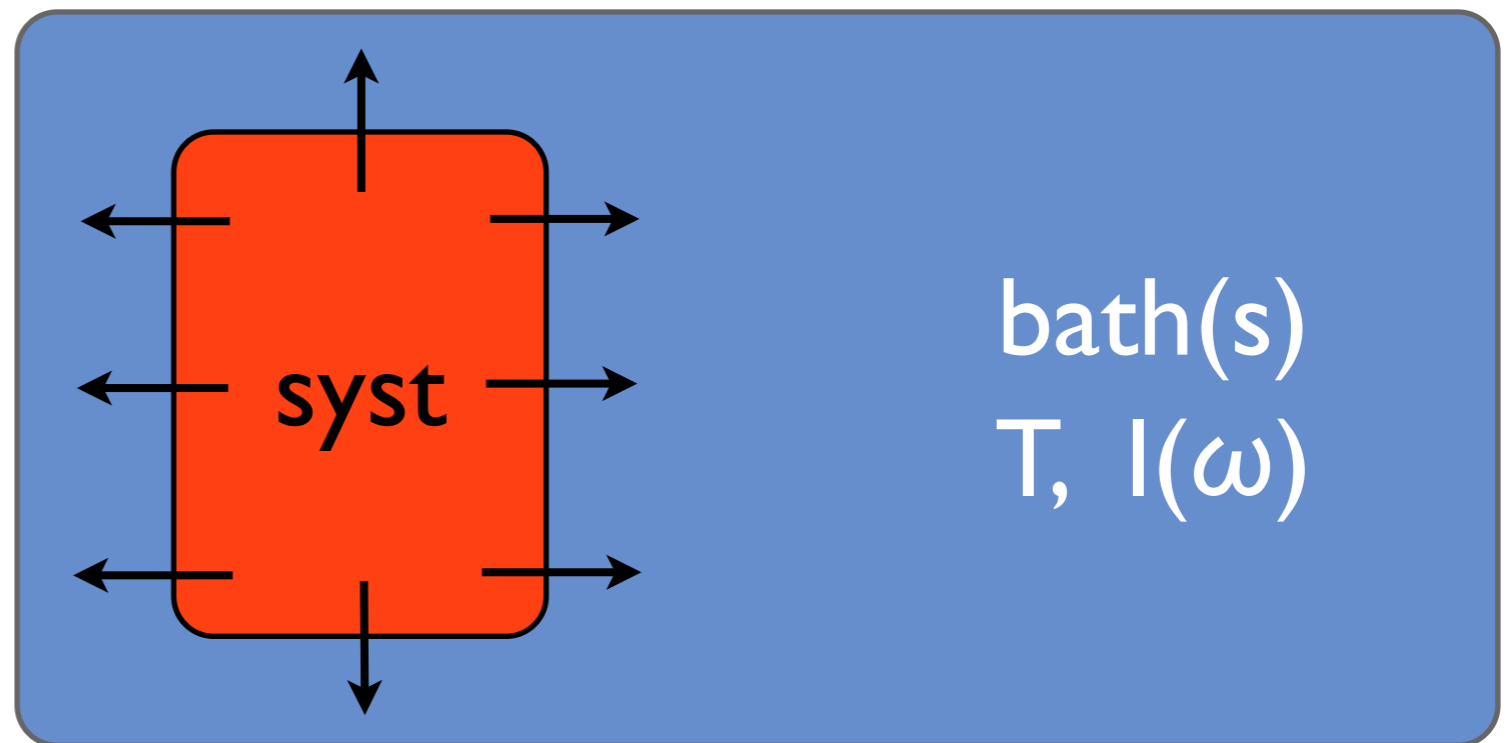
GGI, Florence 18.05.12

Non-equilibrium dynamics

Classical/quantum open systems

$$H = H_{syst} + H_{bath} + H_{int}$$

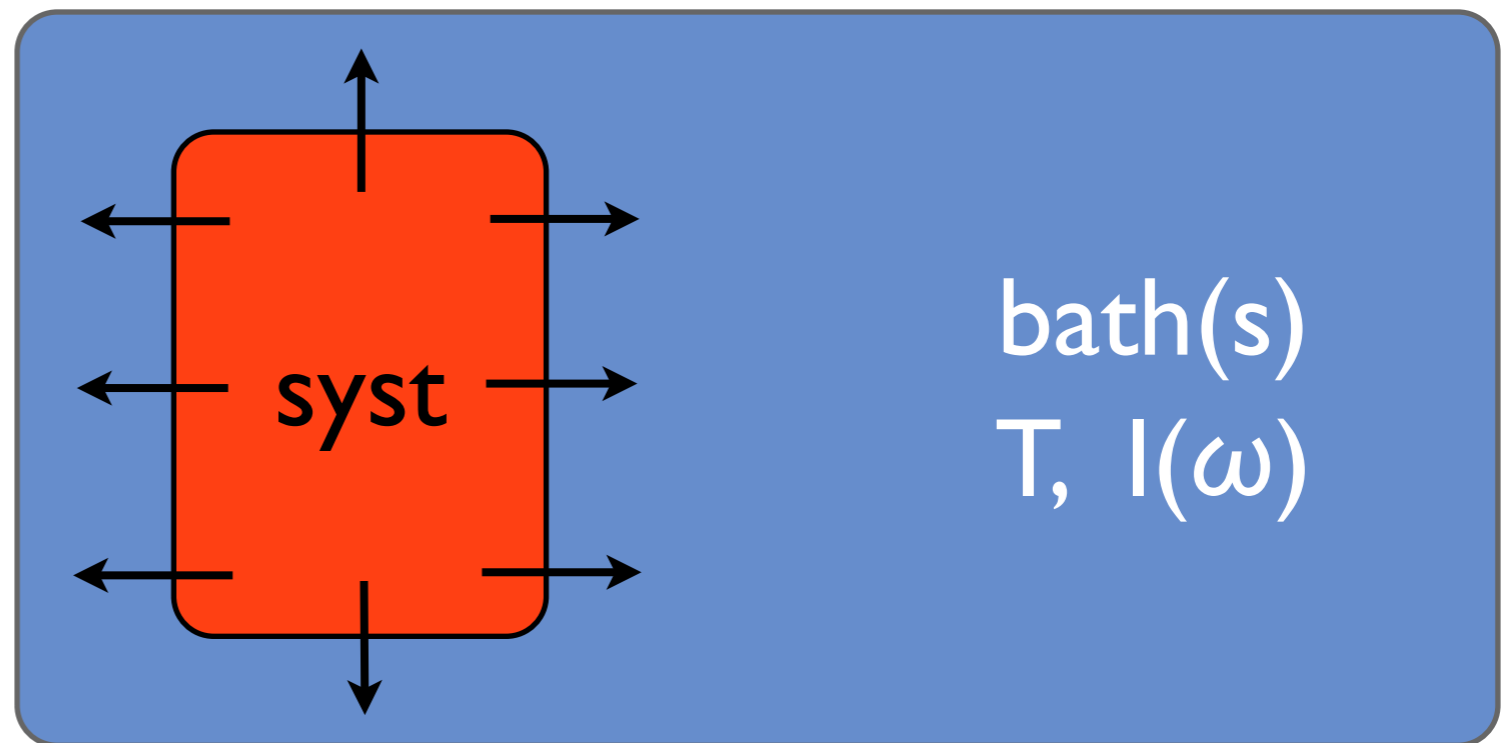
- Transport
 - Pumping
 - Thermal quenches
- phase transitions, glasses



Non-equilibrium dynamics

Classical/quantum open systems

$$H = H_{\text{sys}} + \cancel{H_{\text{bath}}} + \cancel{H_{\text{int}}}$$



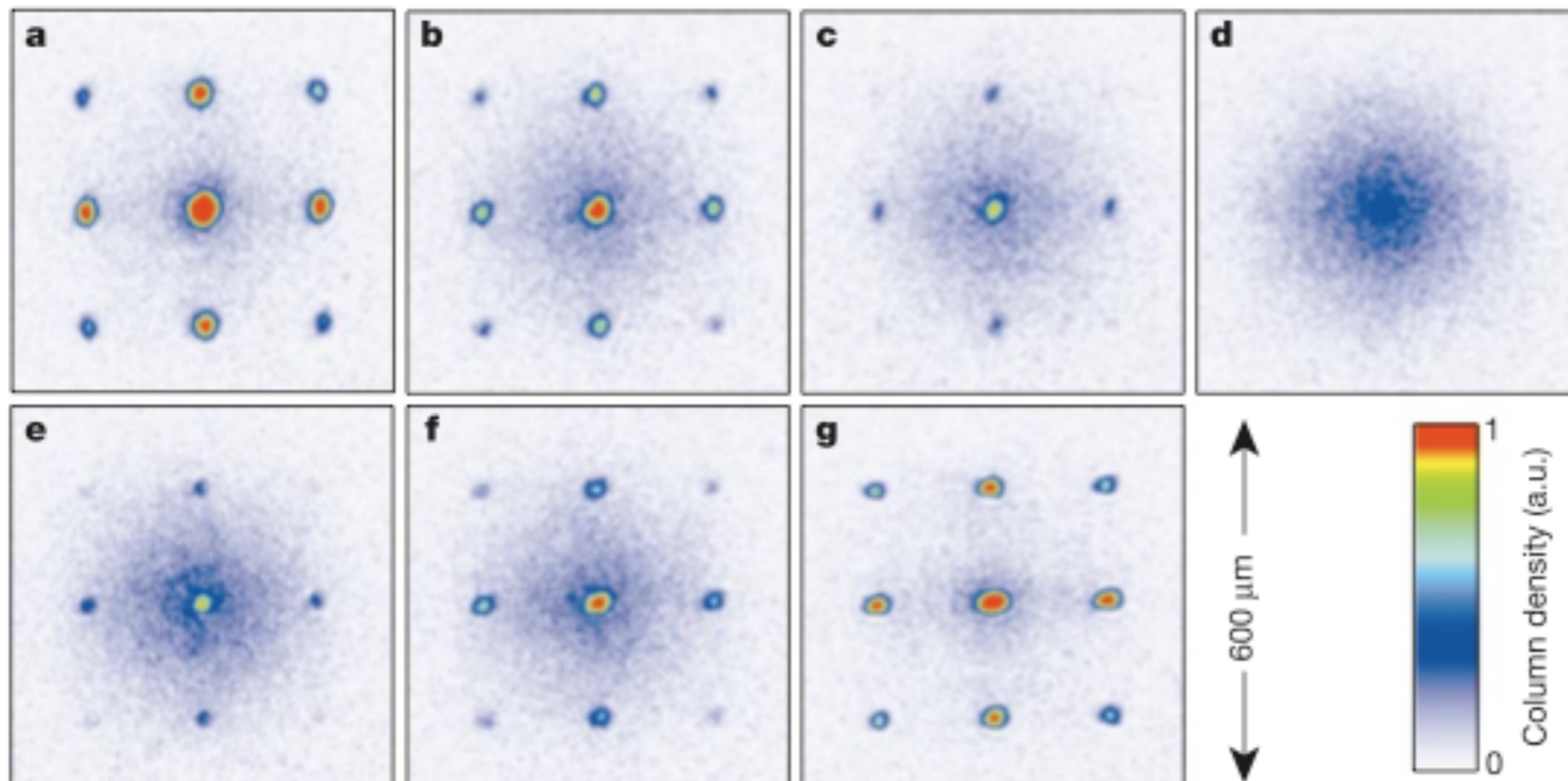
- Transport
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Non-equilibrium dynamics

Quantum isolated systems

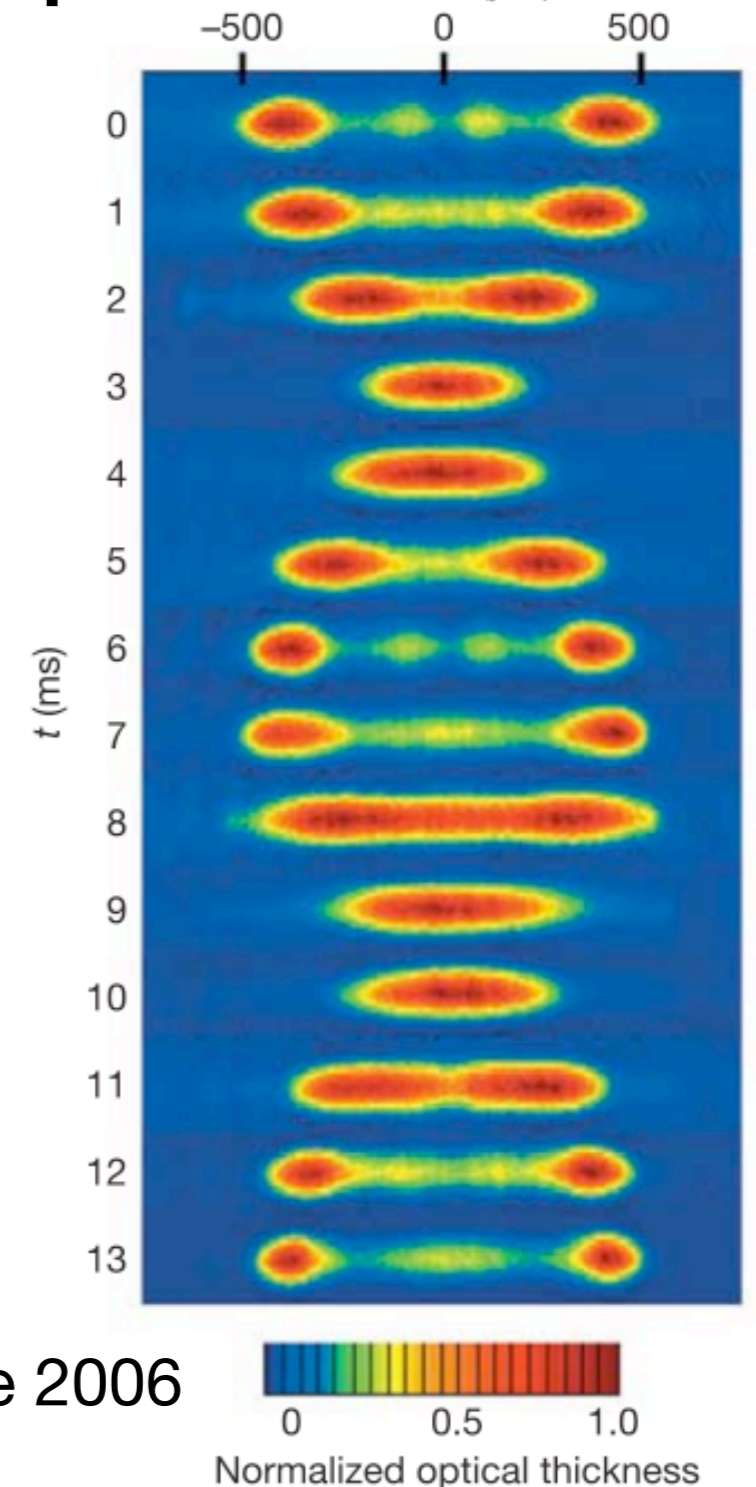
Revived interest → ultra cold atoms

- Control of H
- Coherent dynamics



- Approach to equilibrium \rightarrow Foundation, cosmology
- Quantum annealing: quantum computation, Kibble-Zurek mechanism
- Lack of thermal behavior \rightarrow localization, integrability ...
- Exploration of highly excited states

[Polkovnikov, Sengupta, Silva, Vengalattore, RMP'11]



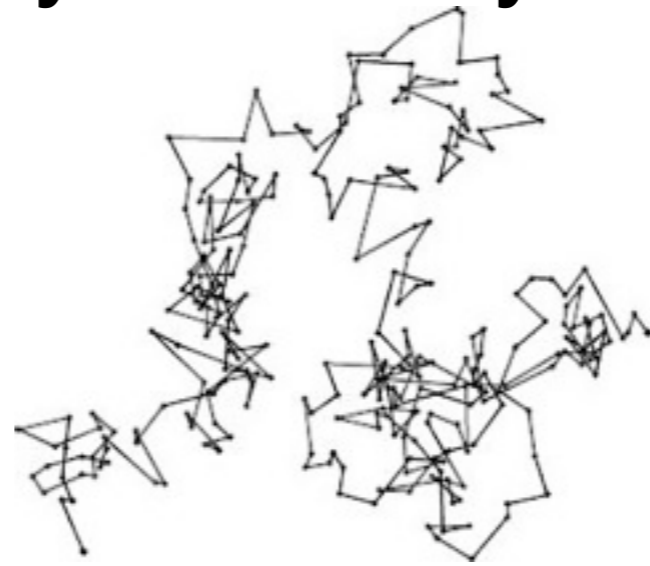
Kinoshita et al., Nature 2006

Fluctuation-dissipation relations:

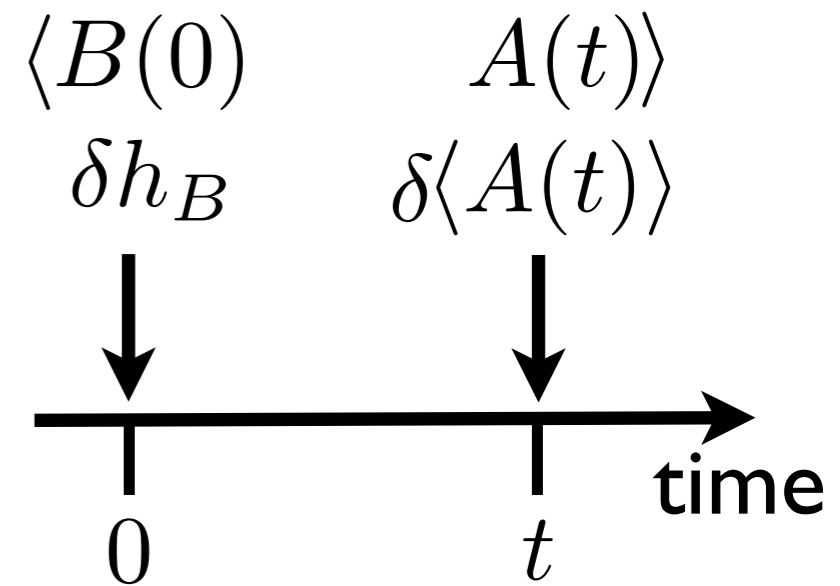
statistical averages \Leftrightarrow dynamical responses

(diffusivity/viscosity ...)

$$D \propto \frac{T}{\eta}$$



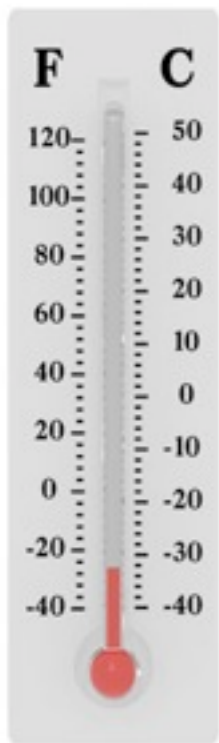
$$-\frac{d}{dt} C_{AB}(t) \underset{\forall A, B}{\propto} T R_{AB}(t)$$



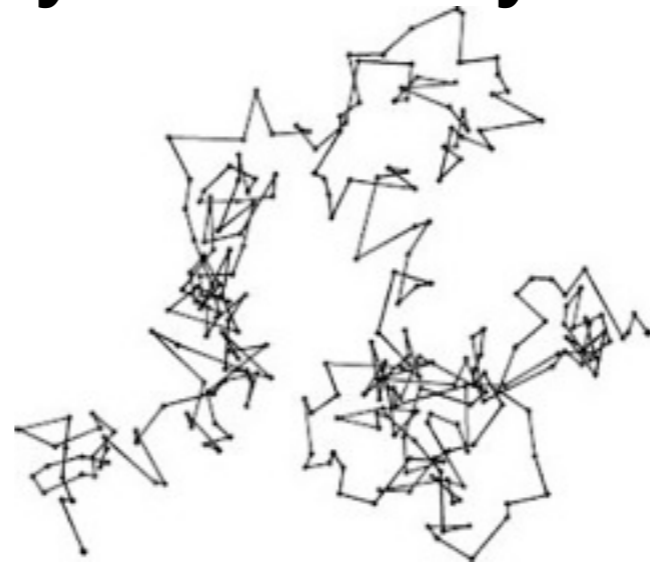
Fluctuation-dissipation relations:

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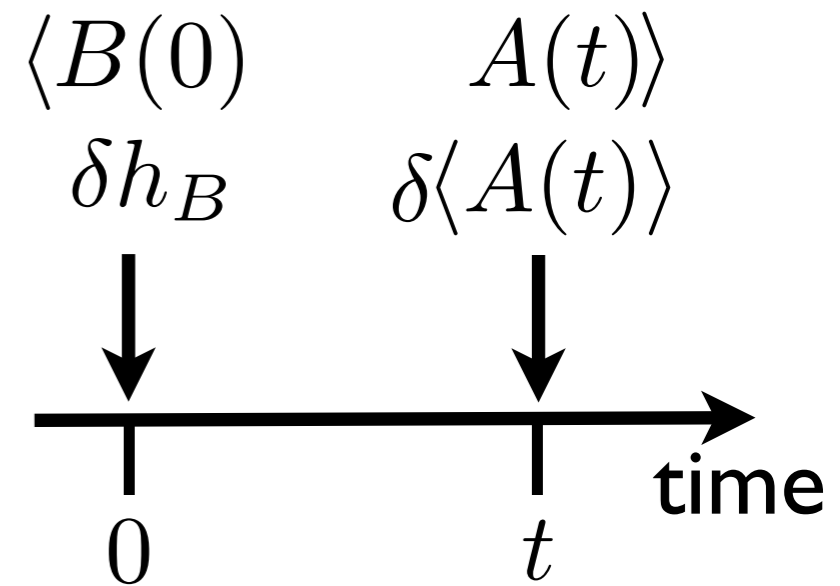
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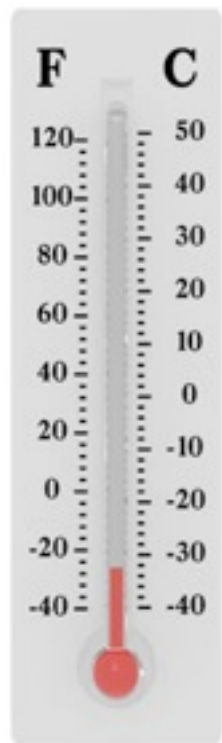
$$-\frac{d}{dt} C_{AB}(t) \propto T R_{AB}(t) \quad \forall A, B$$



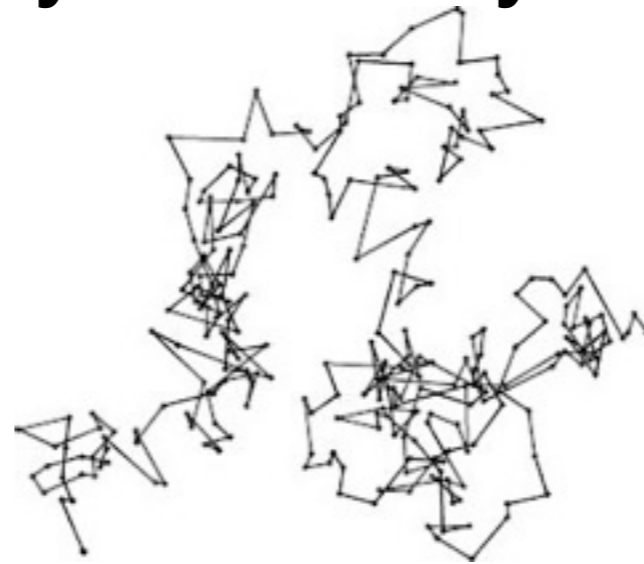
Fluctuation-dissipation relations:

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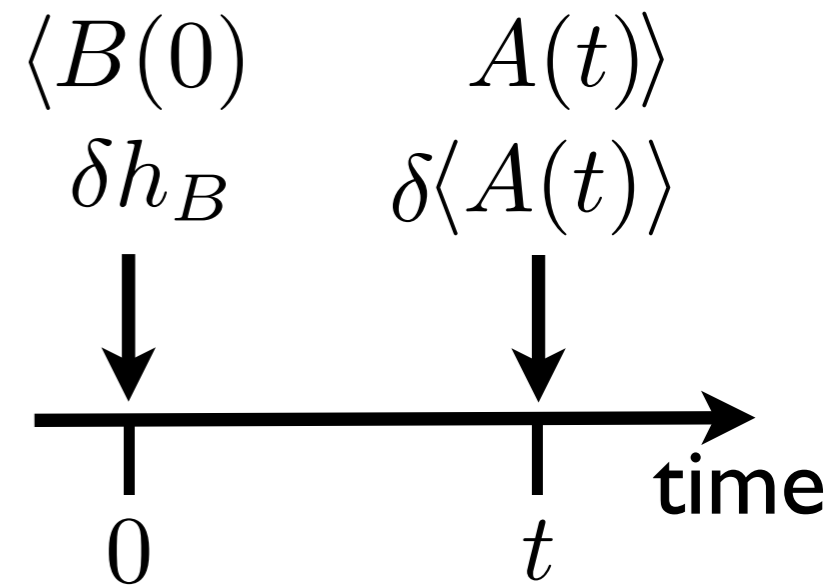
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$$-\frac{d}{dt} C_{AB}(t) \propto T R_{AB}(t) \quad \forall A, B$$



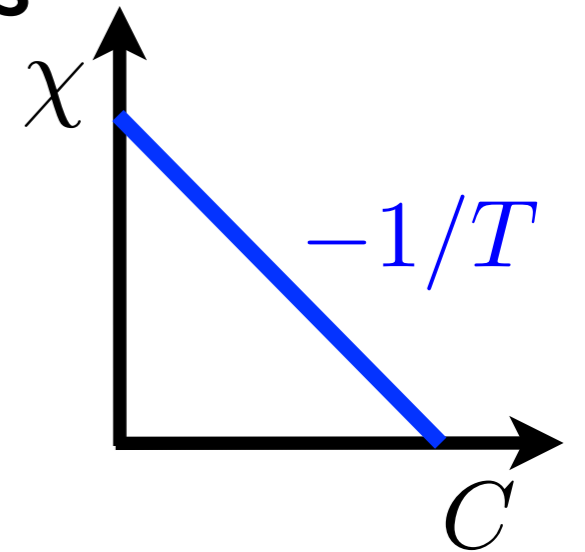
Non-equilibrium \Rightarrow ~~FDR~~

Temperature ?

Effective temperature

Dissipative (classical) dynamics

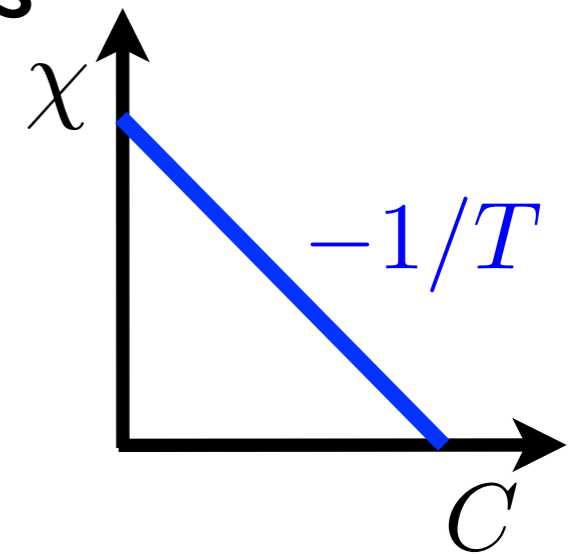
$$\chi(t - t_w) \equiv \int_{t_w}^t dt' R(t, t') = \frac{1 - C(t - t_w)}{k_B T}$$



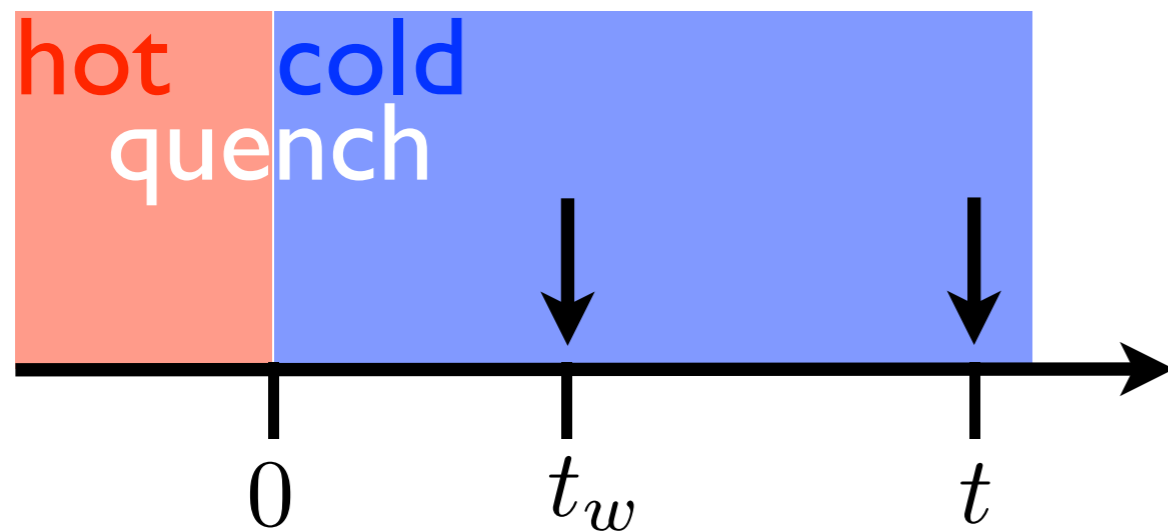
Effective temperature

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Glasses display **slow dyn. & aging**

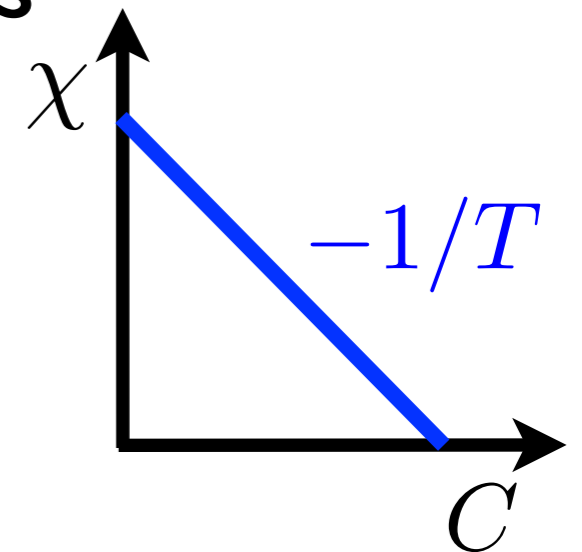


$$\chi(t, t_w) \quad C(t, t_w)$$

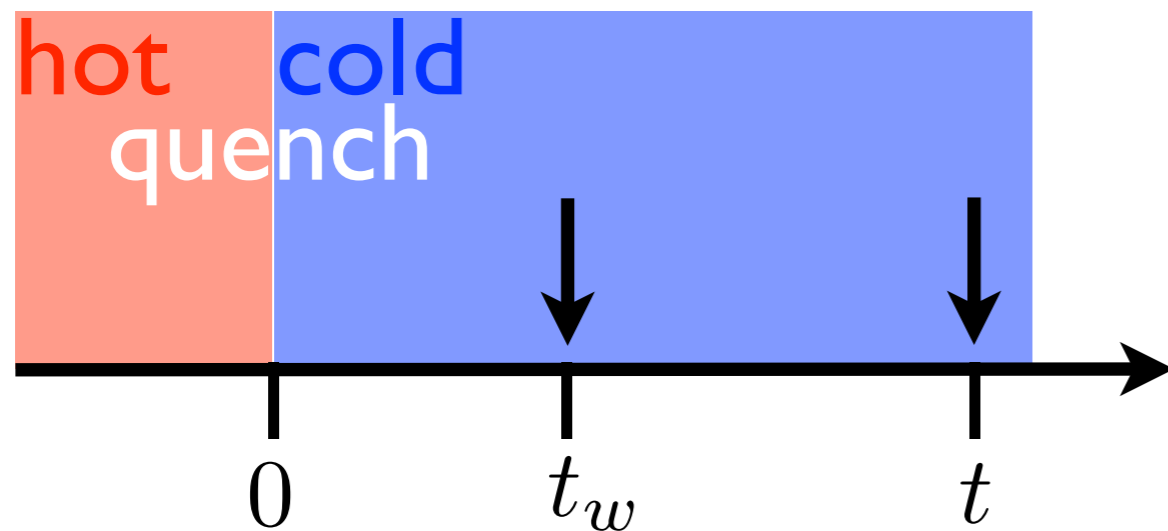
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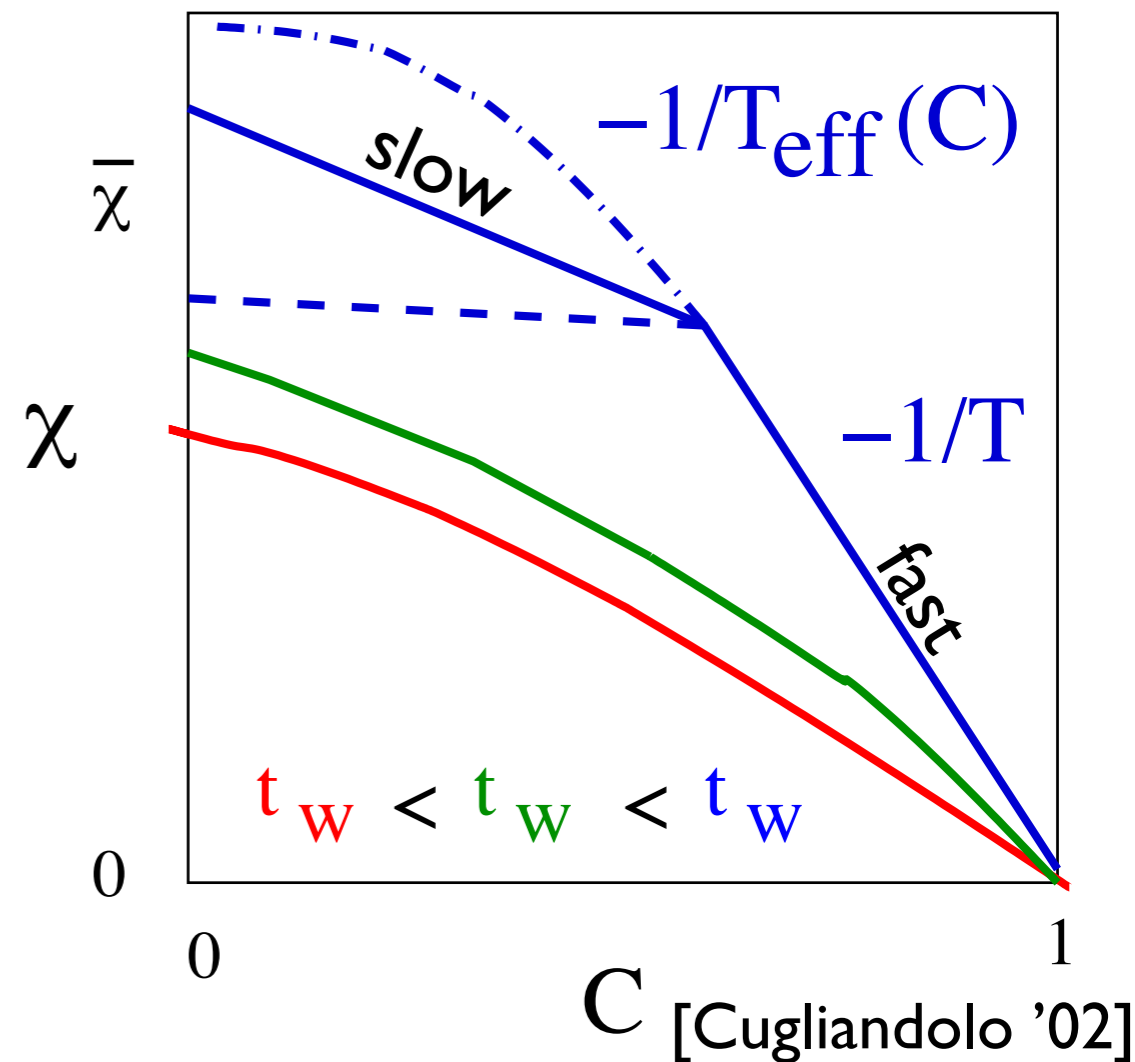
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$$\chi(t, t_w) \quad C(t, t_w)$$

Partial equilibration $\rightarrow T_{\text{eff}}$

2 time scale regimes

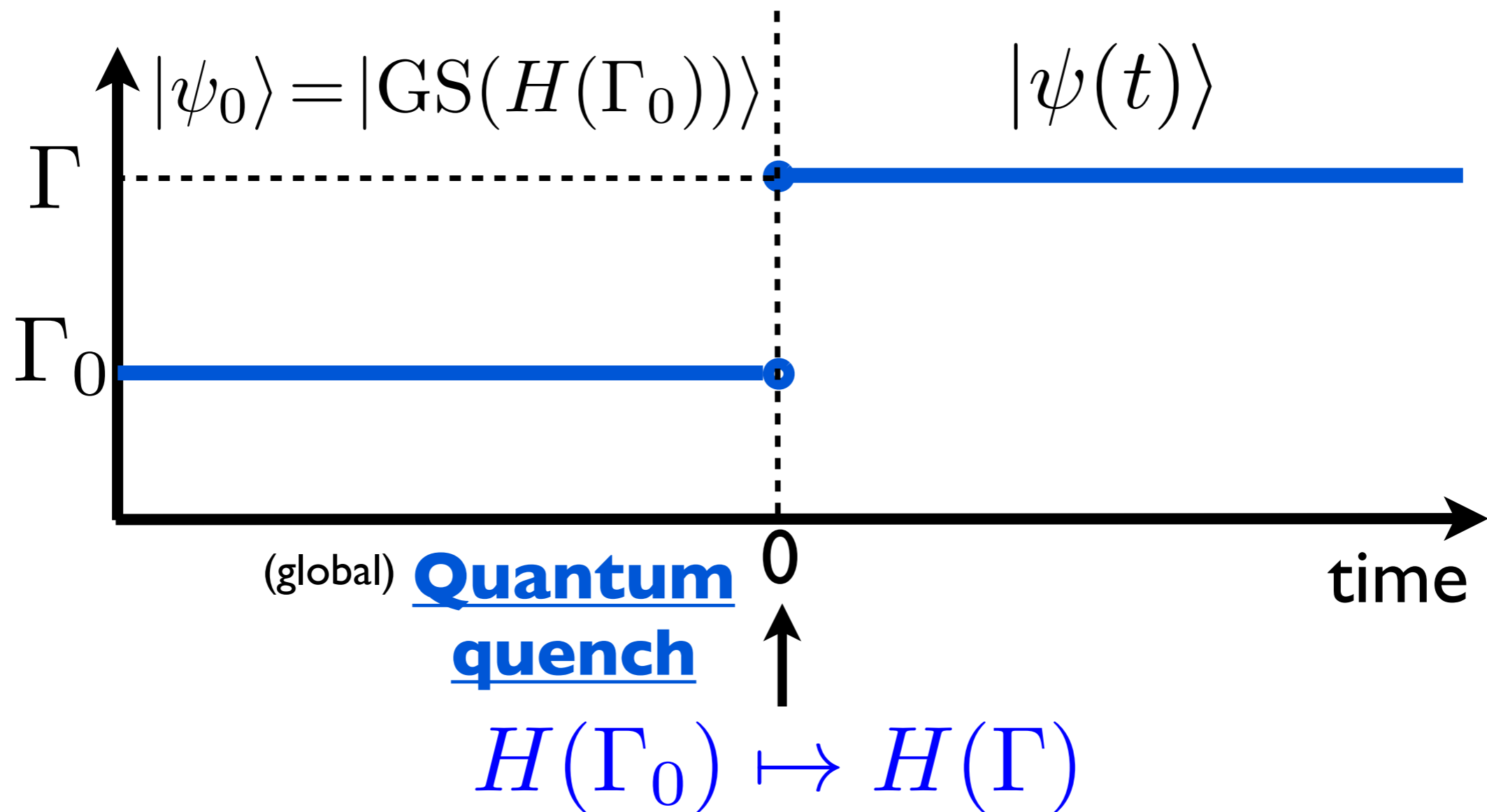


C [Cugliandolo '02]

Quantum dynamics of isolated systems

Unitary

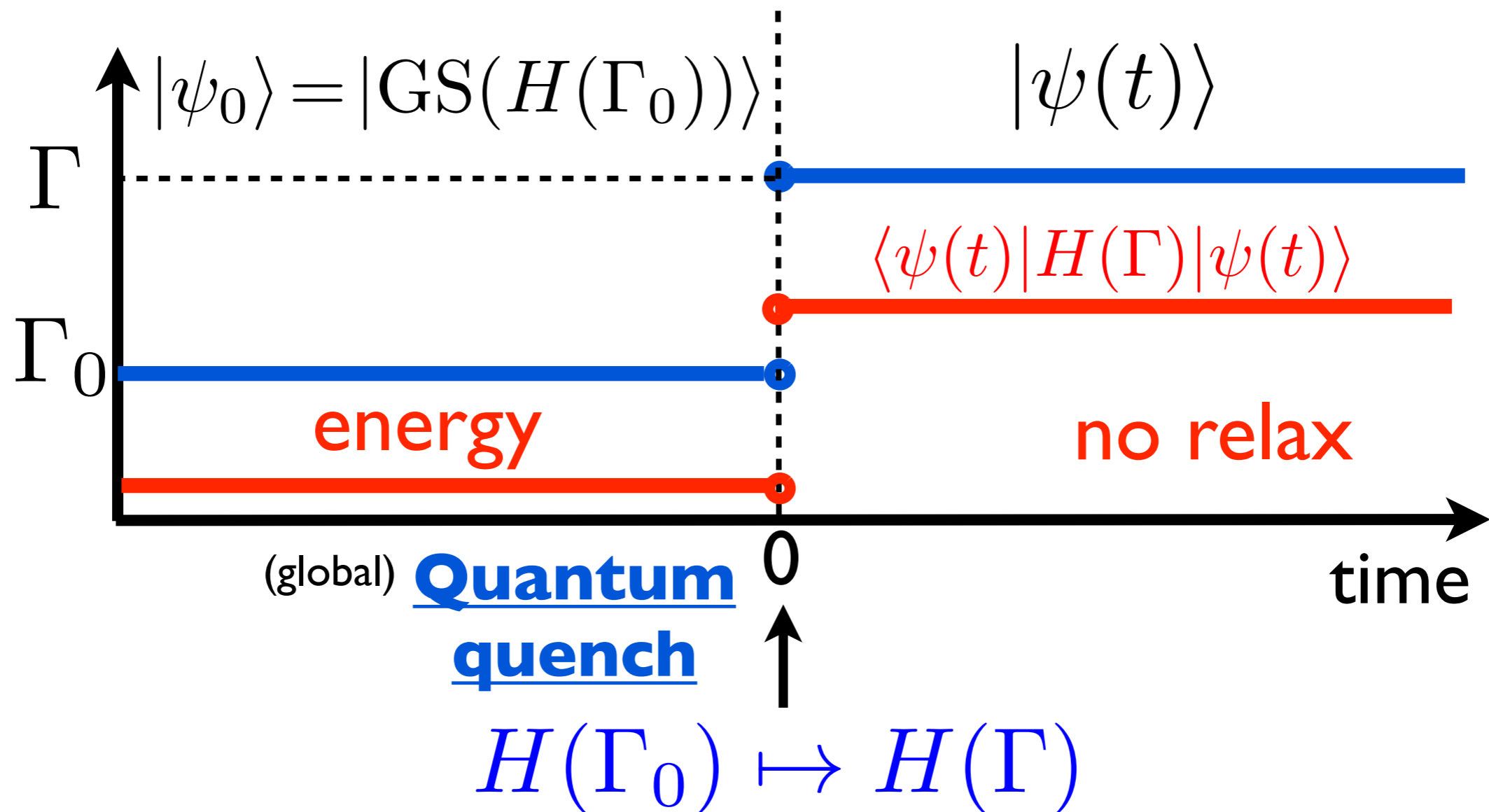
$$|\psi(t)\rangle = e^{-iH(\Gamma)t} |\psi_0\rangle$$



Quantum dynamics of isolated systems

Unitary

$$|\psi(t)\rangle = e^{-iH(\Gamma)t} |\psi_0\rangle$$



Issues...

- **Dynamical transitions**

Biroli, Sciolla - Schirò, Fabrizio - Gambassi, Calabrese - ...

- **Long-lived non-equilibrium states**

Kinoshita, Wenger, Weiss - Kollath, Lauchli, Altman - ...

- **Relaxation of inhomogeneity, aging**

Shutz, Trimper - Iglói, Rieger - Carleo, Becca, Schirò, Fabrizio - ...

- **Thermalization**

Srednicki - Biroli, Kollath, Lauchli - Rigol, Dunjko, Olshanii - Berges - Calabrese, Cardy - Rossini, Silva, Mussardo, Santoro - Calabrese, Essler, Fagotti - ...

stationary state?

*dimensionality?
cons. laws?*

integrability?

Issues...

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generic \Rightarrow *Gibbs (self-bath)*
T ?

effective temp.s ?

stationary state?

*dimensionality?
cons. laws?*

integrability?

Effective temp.s

Quench

$$L \rightarrow \infty$$

$$\langle \psi(t) | \mathcal{O} | \psi(t) \rangle \xrightarrow{t \rightarrow \infty} \langle \mathcal{O} \rangle_{\mathcal{Q}}$$

(? \exists)

Gibbs

$$\rho \propto e^{-H(\Gamma)/T}$$

$$\text{Tr}[\rho \mathcal{O}]$$

Effective temp.s

Quench

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$$\langle \psi(t) | \mathcal{O} | \psi(t) \rangle \xrightarrow{t \rightarrow \infty} \langle \mathcal{O} \rangle_{\mathcal{Q}} \quad (=)$$

(? \exists)

Gibbs

$$\rho \propto e^{-H(\Gamma)/T}$$

$$\text{Tr}[\rho \mathcal{O}] \quad \Rightarrow \quad T$$

Effective temp.s

Quench

$$L \rightarrow \infty$$

$$\langle \psi(t) | \mathcal{O} | \psi(t) \rangle \xrightarrow{t \rightarrow \infty} \langle \mathcal{O} \rangle_{\mathcal{Q}} \quad =$$

(? \exists)

$$H(\Gamma)$$

Gibbs

$$\rho \propto e^{-H(\Gamma)/T}$$

$$\text{Tr}[\rho \mathcal{O}] \Rightarrow T$$

$$\Rightarrow T^E$$

(Γ, Γ_0)

Effective temp.s

Quench

$$L \rightarrow \infty$$

$$\langle \psi(t) | \mathcal{O} | \psi(t) \rangle \xrightarrow{t \rightarrow \infty} \langle \mathcal{O} \rangle_{\mathcal{Q}} \quad (= \text{?} \exists)$$

$$H(\Gamma)$$

$$\mathcal{O} \longrightarrow f_{\mathcal{Q}}(\Gamma, \Gamma_0)$$

Gibbs

$$\rho \propto e^{-H(\Gamma)/T}$$

$$\text{Tr}[\rho \mathcal{O}] \Rightarrow T$$

$$f_G(T^E(\Gamma, \Gamma_0)) \Rightarrow T^E(\Gamma, \Gamma_0)$$

?
≐

Effective temp.s

Quench

$$L \rightarrow \infty$$

$$\langle \psi(t) | \mathcal{O} | \psi(t) \rangle \xrightarrow{t \rightarrow \infty} \langle \mathcal{O} \rangle_{\mathcal{Q}} \quad (? \exists)$$

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$$\text{Tr}[\rho \mathcal{O}] \Rightarrow T$$



$$T^E$$

$$(\Gamma, \Gamma_0)$$

$$f_G(T^E(\Gamma, \Gamma_0))$$



$$T^{\mathcal{O}}$$

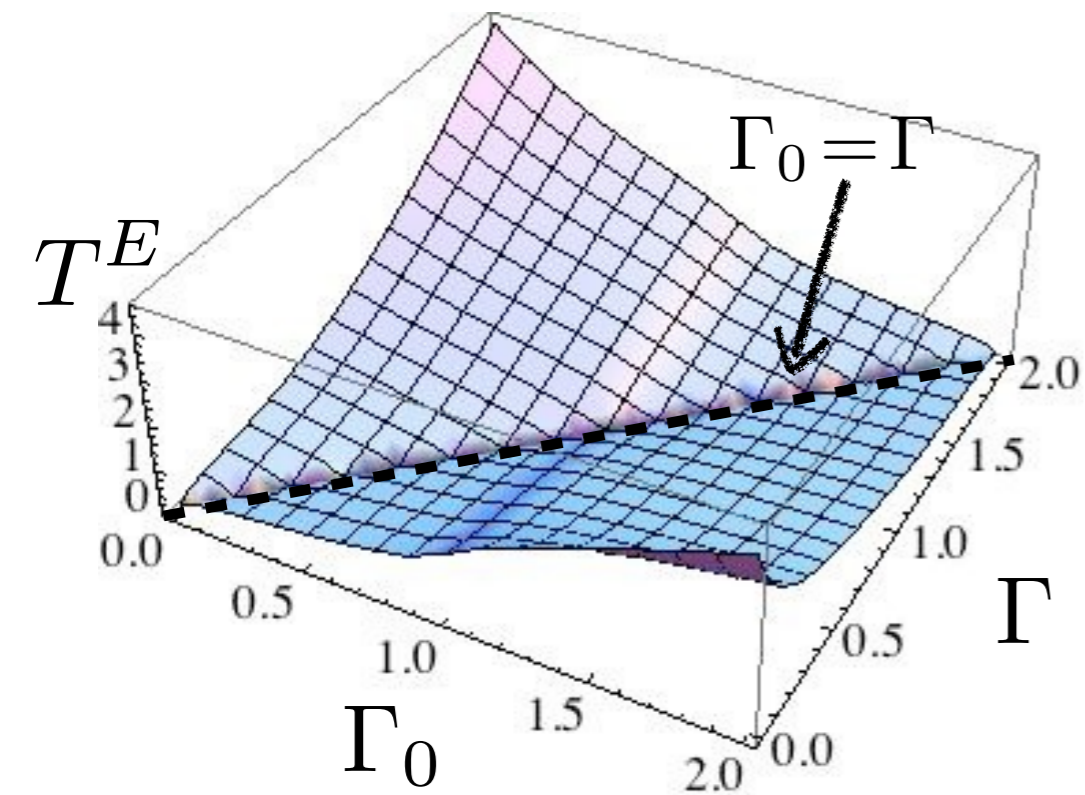
corr. length, coherence time

$$\langle \sigma_i \sigma_j \rangle_{\mathcal{Q}} \quad \langle \psi_0 | \sigma_i(0) \sigma_i(t) | \psi_0 \rangle$$

integrability - non-locality?

Rossini, Silva, Mussardo, Santoro

Effective temp.s



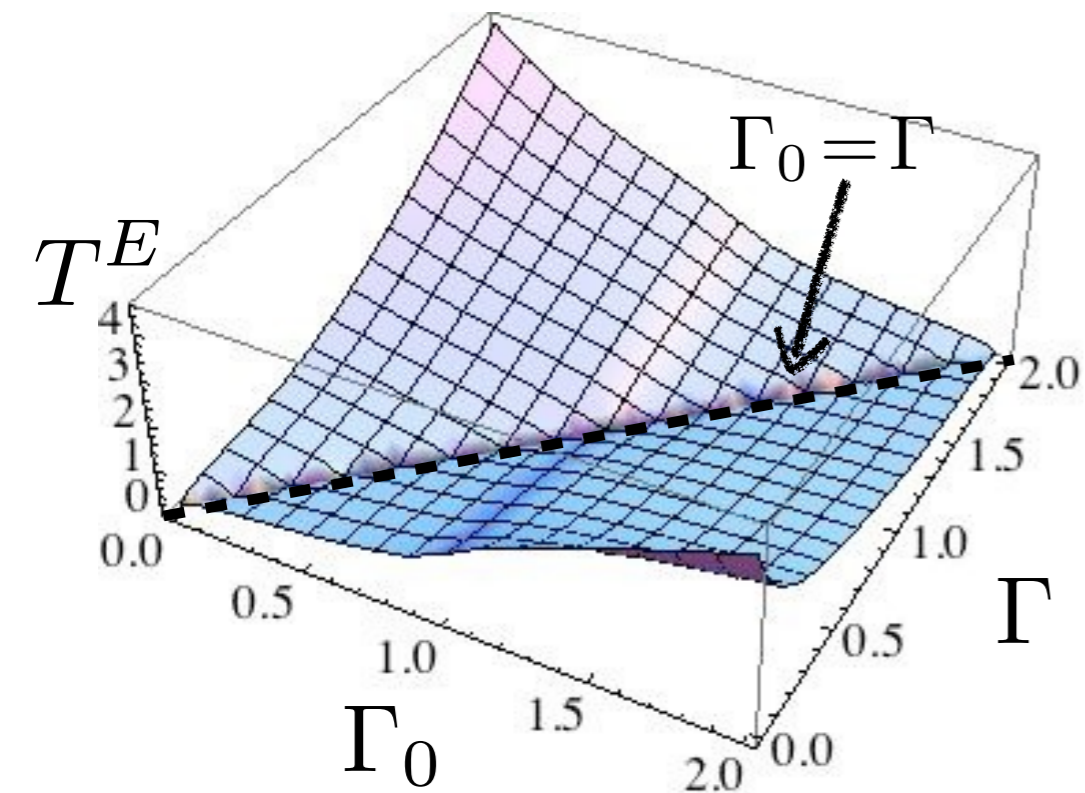
1d-lattice model

$$\Gamma, \Gamma_0 > 0$$

Effective temp.s

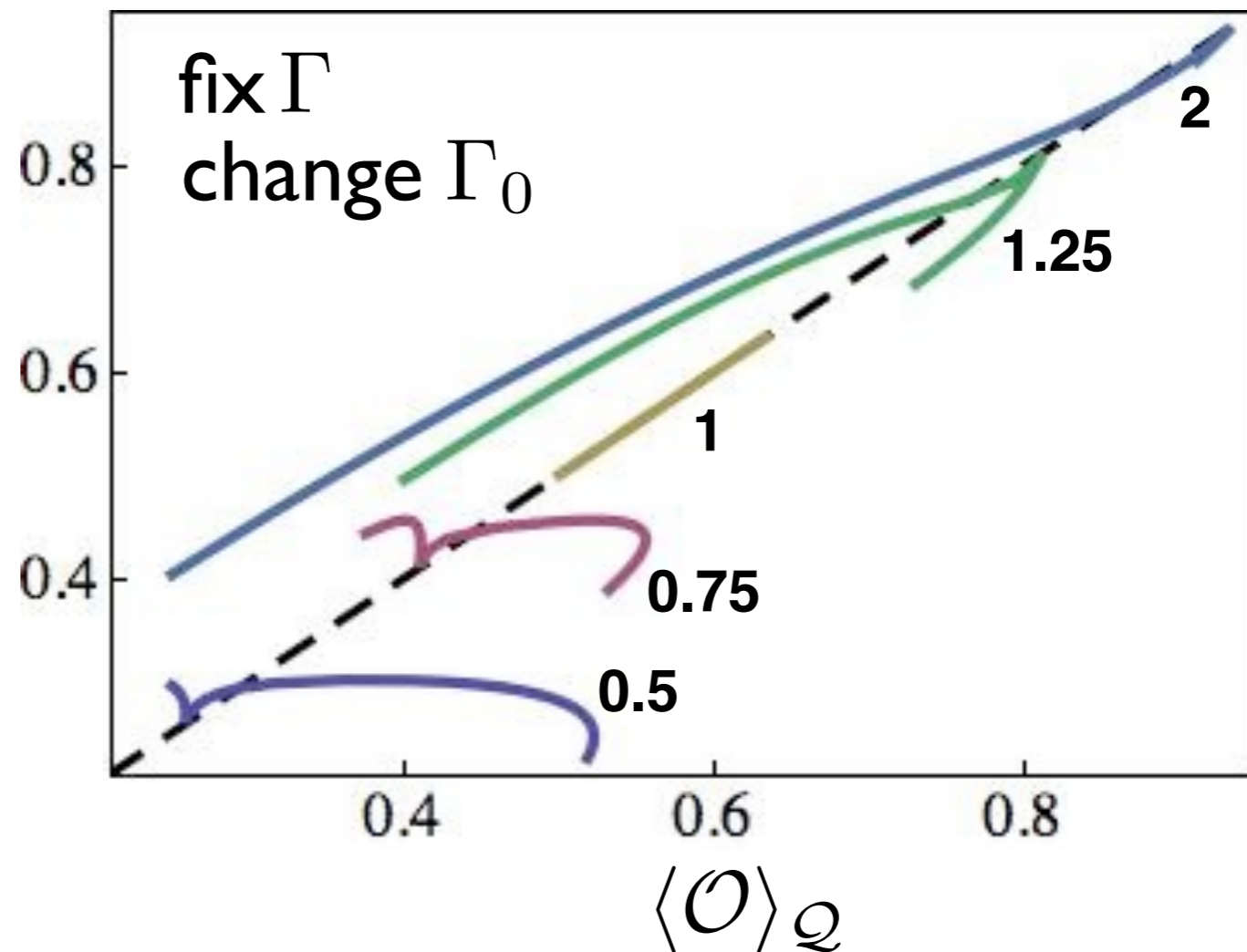
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observable \mathcal{O}

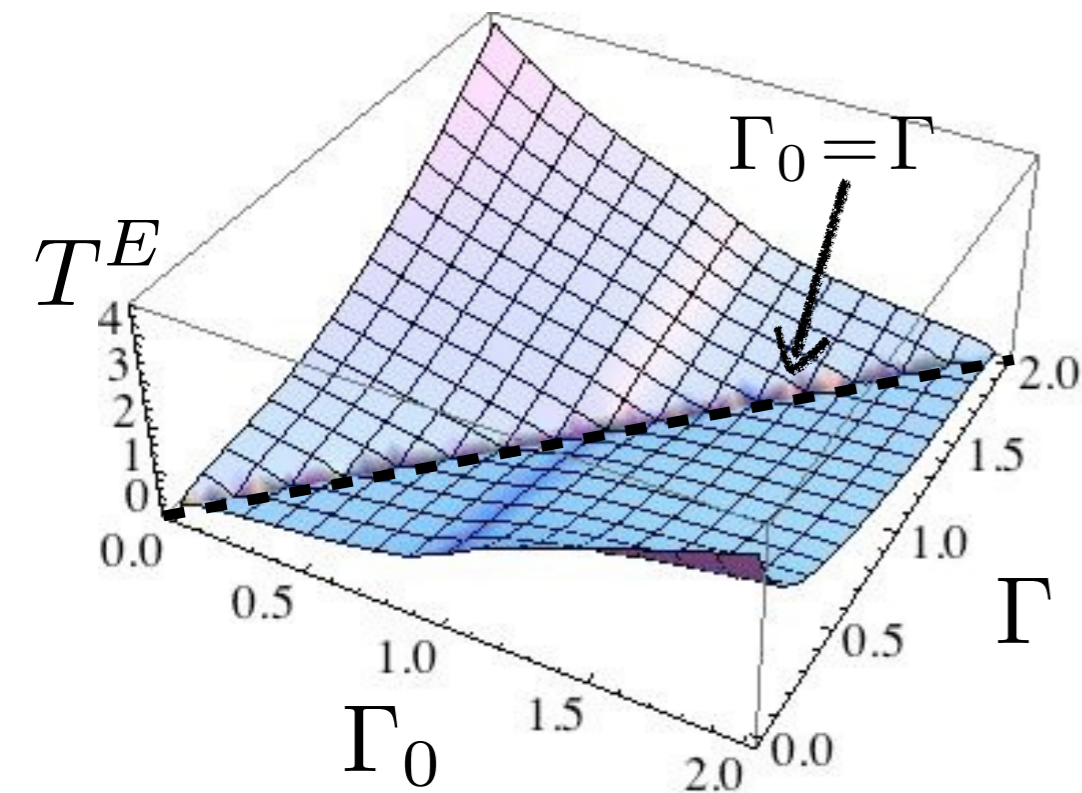
$$\langle \mathcal{O} \rangle_{\text{Gibbs}} \\ @ T = T^E(\Gamma_0, \Gamma)$$



Effective temp.s

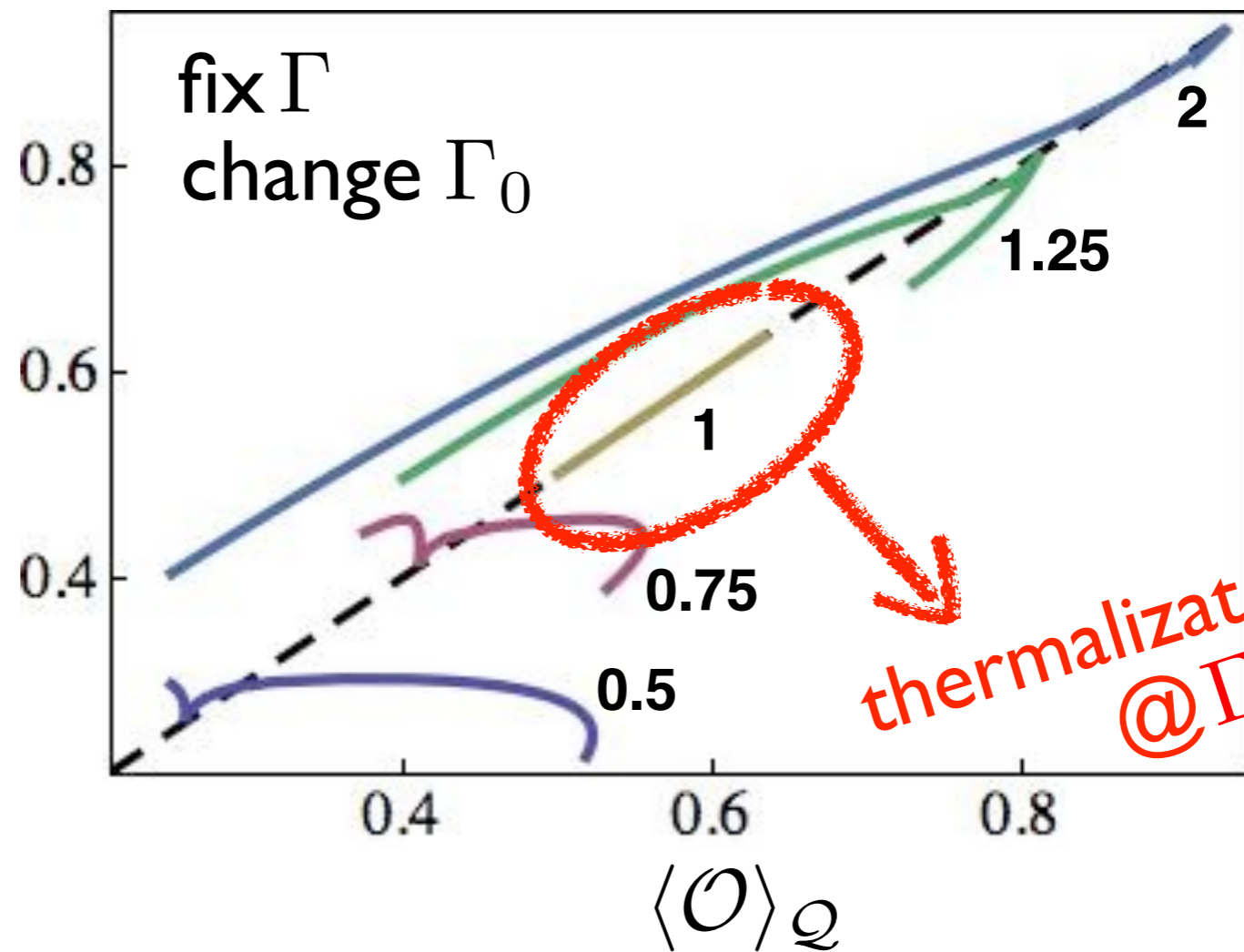
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Effective temp.s

Quench

Gibbs

$$\rho \propto e^{-H(\Gamma)/T}$$

$$U(t) = e^{-iHt}$$

quantum FDT

+

$$\frac{\hbar \operatorname{Im} R(\omega)}{\operatorname{Re} C(\omega)} = \operatorname{th} \left(\frac{\hbar\omega}{2T} \right)$$

Effective temp.s

Quench

$$\langle \psi_0 | \mathcal{O}(t+t_0) \mathcal{O}(t_0) | \psi_0 \rangle \xrightarrow{t_0 \rightarrow \infty}$$

$$\theta(t) [\quad , \quad] \longrightarrow R(t)$$

$$\{ \quad , \quad \} \longrightarrow C_+(t)$$

stationary regime

Gibbs

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$$\frac{\hbar \operatorname{Im} R(\omega)}{\operatorname{Re} C(\omega)} = \operatorname{th} \left(\frac{\hbar \omega}{2T} \right)$$

$$\frac{\hbar \operatorname{Im} R(\omega)}{C_+(\omega)} = \operatorname{th} \left(\frac{\hbar \omega}{2T_{\text{eff}}(\omega)} \right)$$

measure!

*O-dependent ?
dynamics !*

Effective temp.s

Quench

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stationary regime

Gibbs

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quantum FDT

$$\frac{\hbar \operatorname{Im} R(\omega)}{\operatorname{Re} C(\omega)} = \operatorname{th} \left(\frac{\hbar \omega}{2T} \right)$$

$$\frac{\hbar \operatorname{Im} R(\omega)}{C_+(\omega)} = \operatorname{th} \left(\frac{\hbar \omega}{2T_{\text{eff}}(\omega)} \right)$$

equilibrium



$$T_{\text{eff}}(\omega) = T^E$$

O-independent

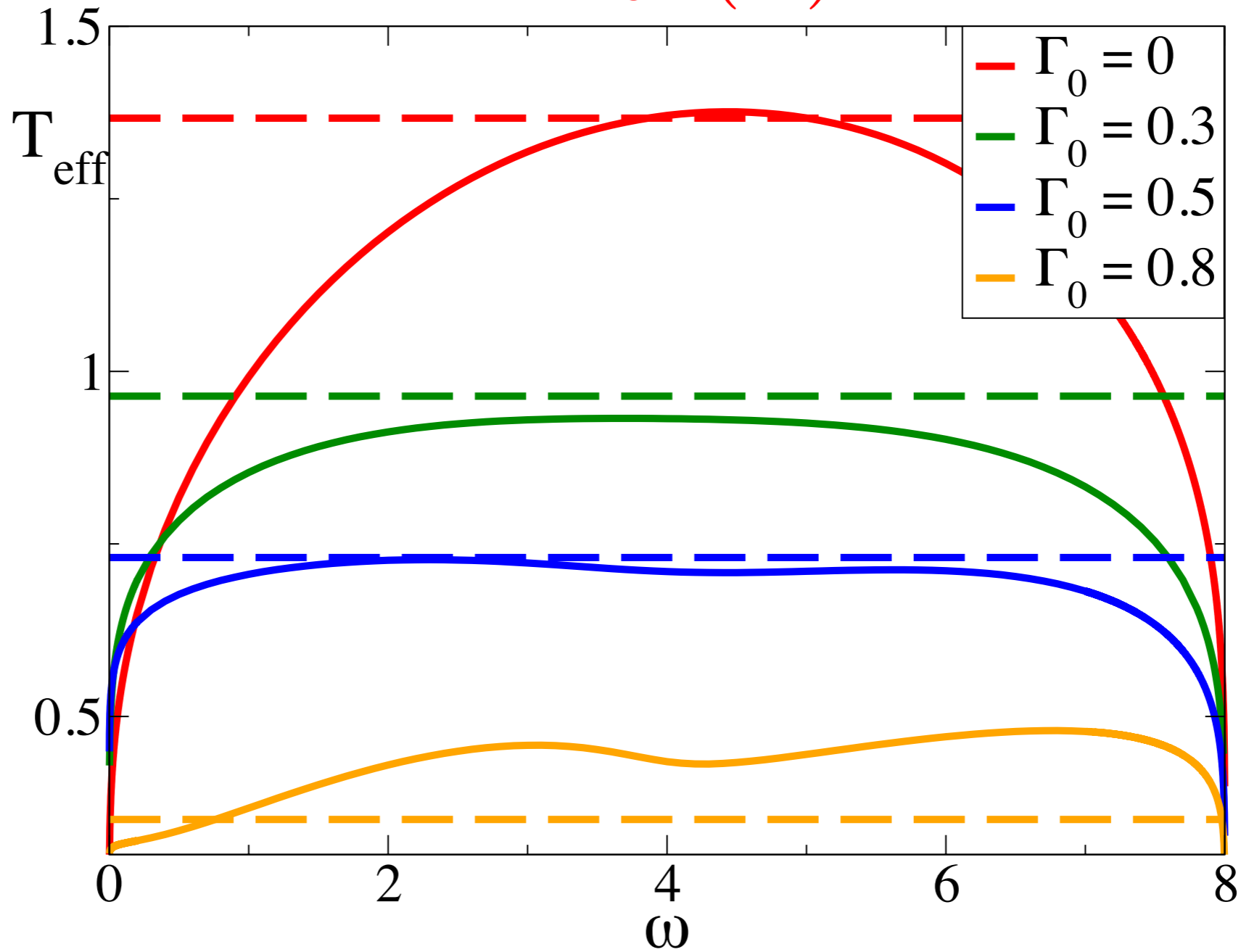
*O-dependent ?
dynamics !*

measure!

Effective temp.s

focus on $\Gamma = 1$
observable \mathcal{O}

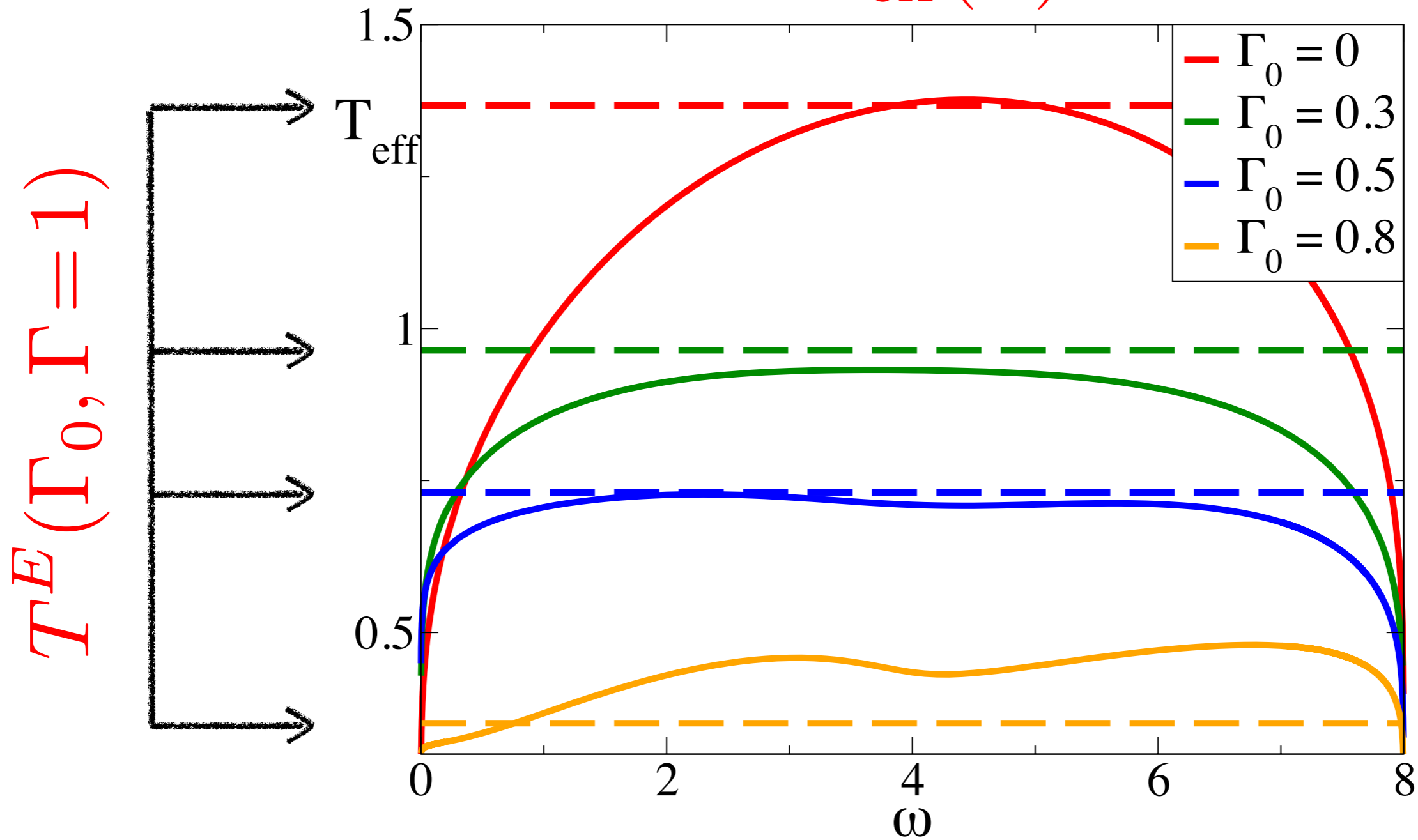
$$T_{\text{eff}}(\omega)$$



Effective temp.s

focus on $\Gamma=1$
observable \mathcal{O}

$T_{\text{eff}}(\omega)$

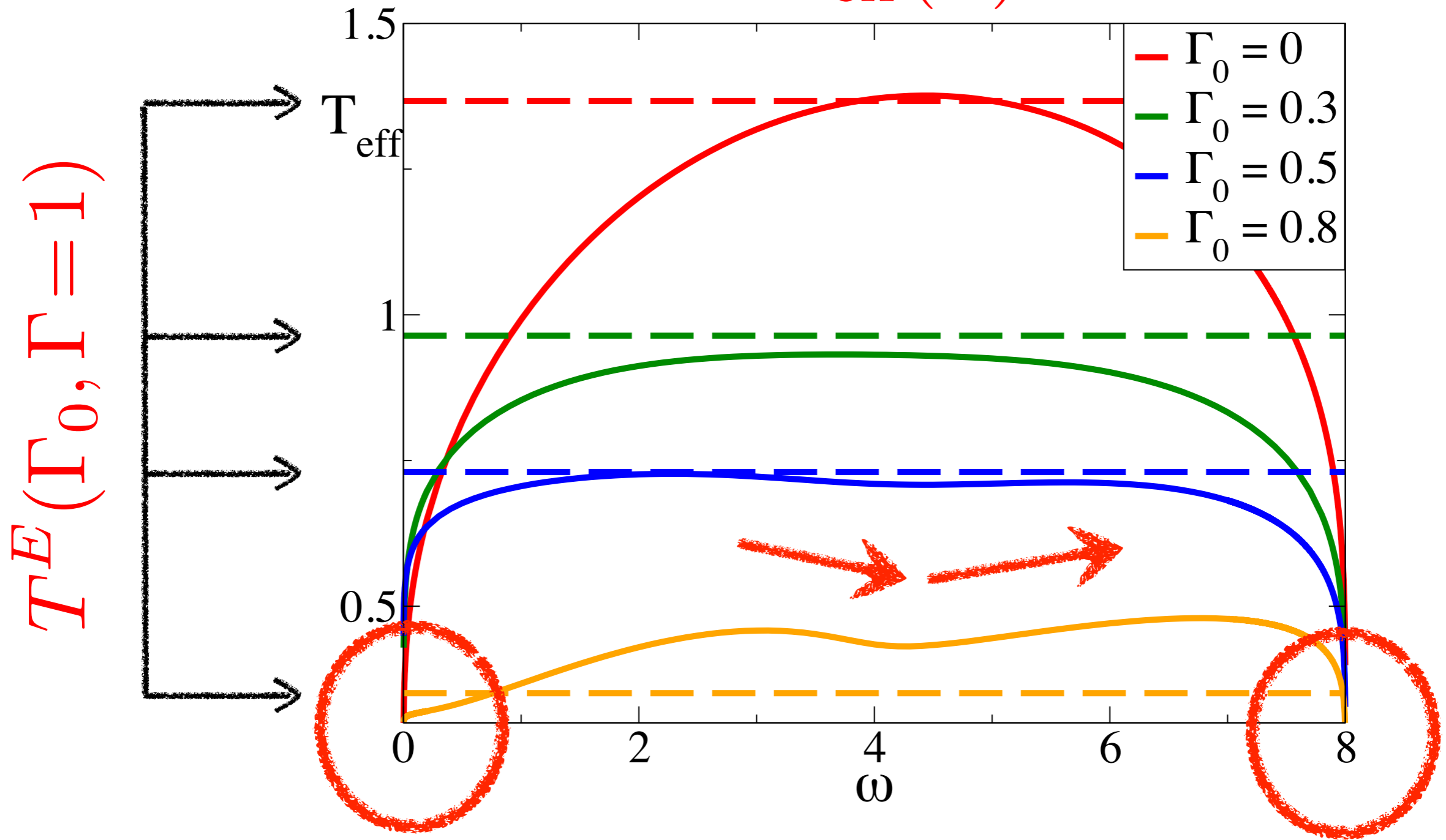


$T^E(\Gamma_0, \Gamma=1)$

Effective temp.s

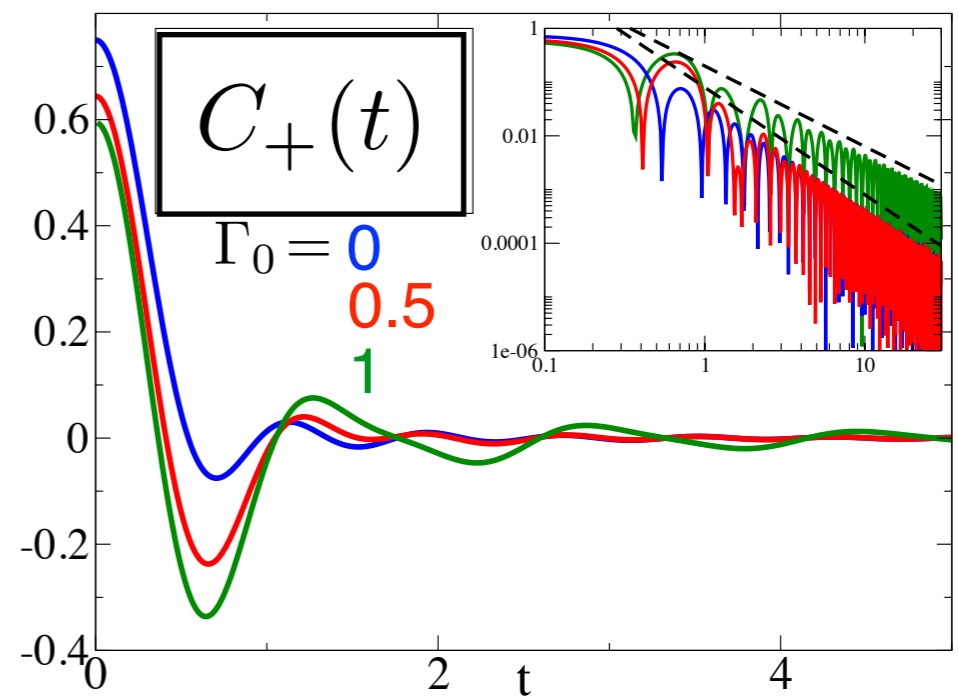
focus on $\Gamma = 1$
observable \mathcal{O}

$T_{\text{eff}}(\omega)$



$R(t), C_+(t)$ invar. $\Gamma_0 \leftrightarrow 1/\Gamma_0$

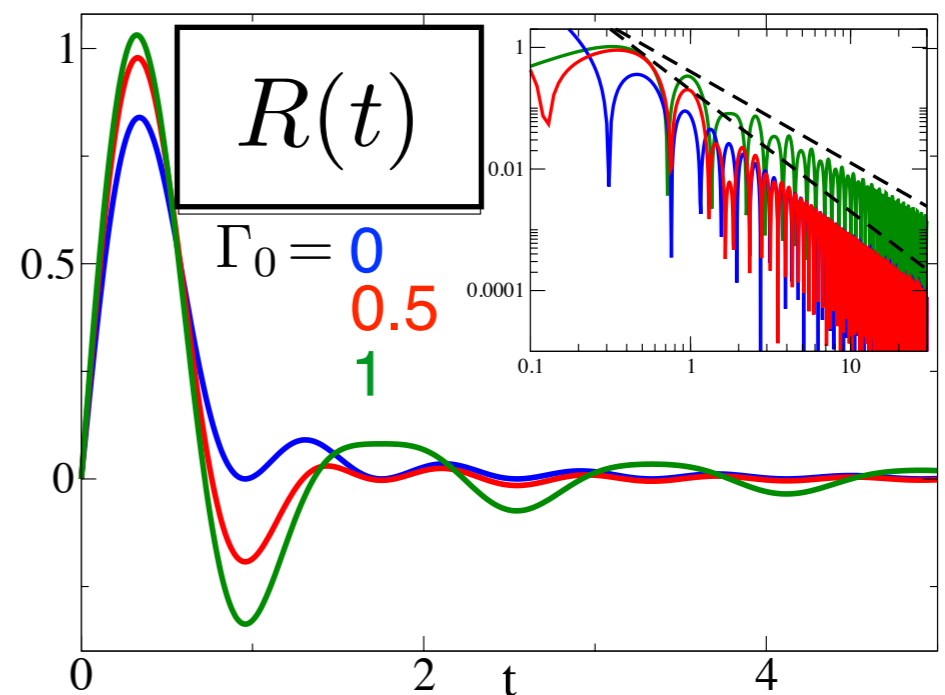
focus on $\Gamma = 1$
observable \mathcal{O}



$$C_+(t) = -\frac{1}{4\pi t^2} \cos(8t) + O(t^{-3}) \quad \Gamma_0 = 1 : \sim t^{-3/2}$$

$$R(t) = \frac{1}{4\pi t^2} [\Upsilon(\Gamma_0) - \sin(8t)] + O(t^{-3}) \quad \Upsilon(\Gamma_0) = \left(\frac{1 - \Gamma_0}{1 + \Gamma_0} \right)^2$$

eq. @ $\Gamma \neq 0$
 $C_+(t), R(t) \propto \frac{1}{t}$



focus on $\Gamma = 1$
observable \mathcal{O}

slow
relaxation

$$C_+(t) = \frac{1}{4\pi t^2} \cos(8t) + O(t^{-3})$$

$$\Gamma_0 = 1 : \sim t^{-3/2}$$

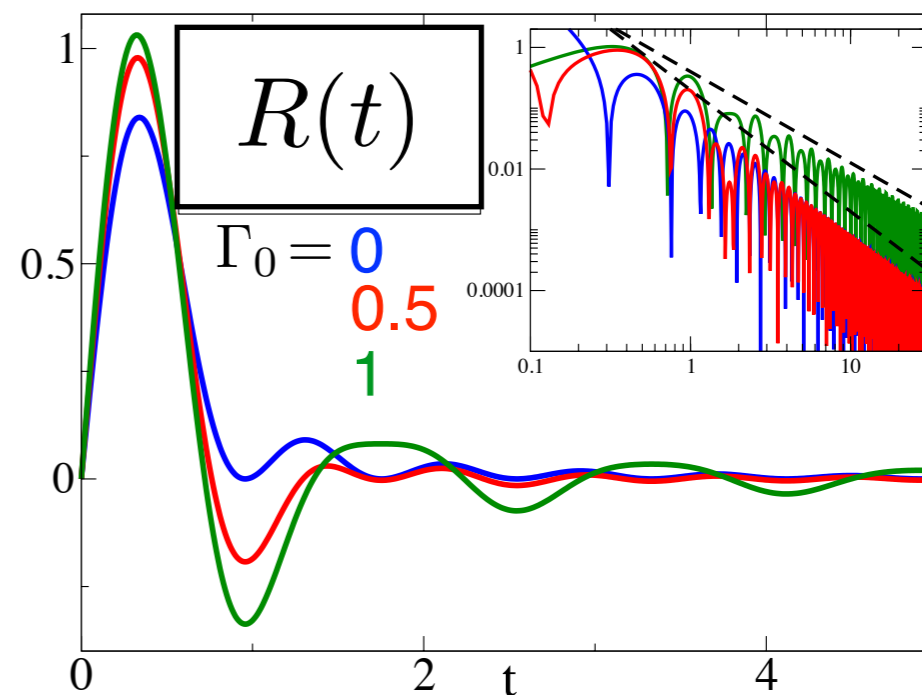
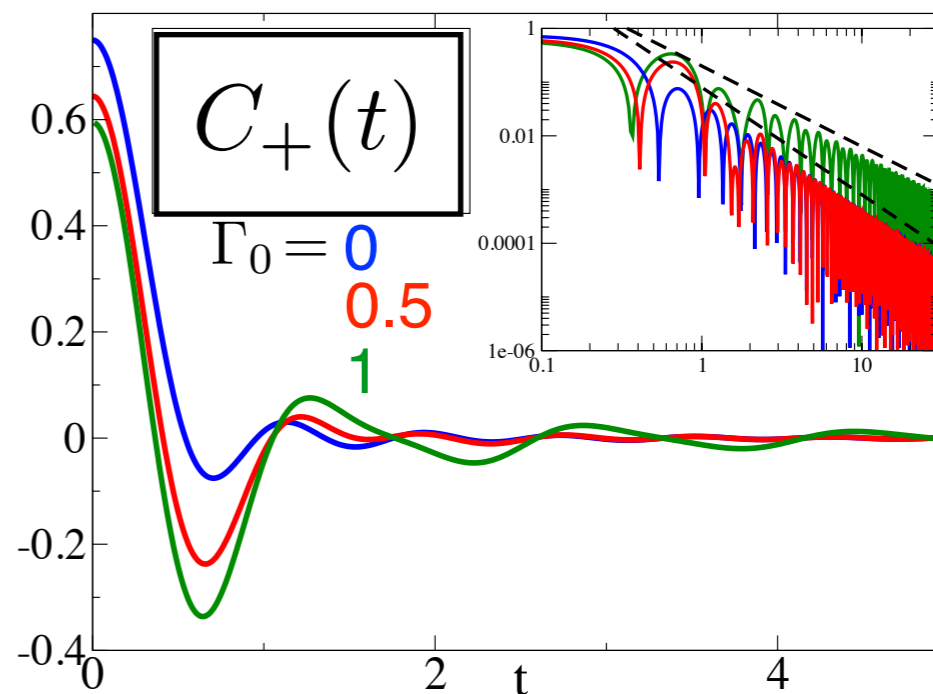
$$R(t) = \frac{1}{4\pi t^2} [\Upsilon(\Gamma_0) - \sin(8t)] + O(t^{-3})$$

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eq. @ $\Gamma \neq 0$

$$C_+(t), R(t) \propto \frac{1}{t}$$

no coher. time
no match



Effective temp.s

$$\langle H \rangle_Q = \langle H \rangle_{\text{Gibbs}} \implies T^E$$

static

observable \mathcal{O}

$$\langle \mathcal{O} \rangle_Q = \langle \mathcal{O} \rangle_{\text{Gibbs}@T=T^E} \quad \rightsquigarrow \text{thermalization ?}$$

Effective temp.s

$$\langle H \rangle_{\mathcal{Q}} = \langle H \rangle_{\text{Gibbs}} \implies T^E$$

static

observable \mathcal{O}

$$\langle \mathcal{O} \rangle_{\mathcal{Q}} = \langle \mathcal{O} \rangle_{\text{Gibbs}@T=T^E} \quad \rightsquigarrow \text{thermalization ?}$$

dynamics

$$R^{\mathcal{O}}(t), C_{+}^{\mathcal{O}}(t) \quad \text{Gibbs FDT} \implies T_{\text{eff}}^{\mathcal{O}}(\omega)$$

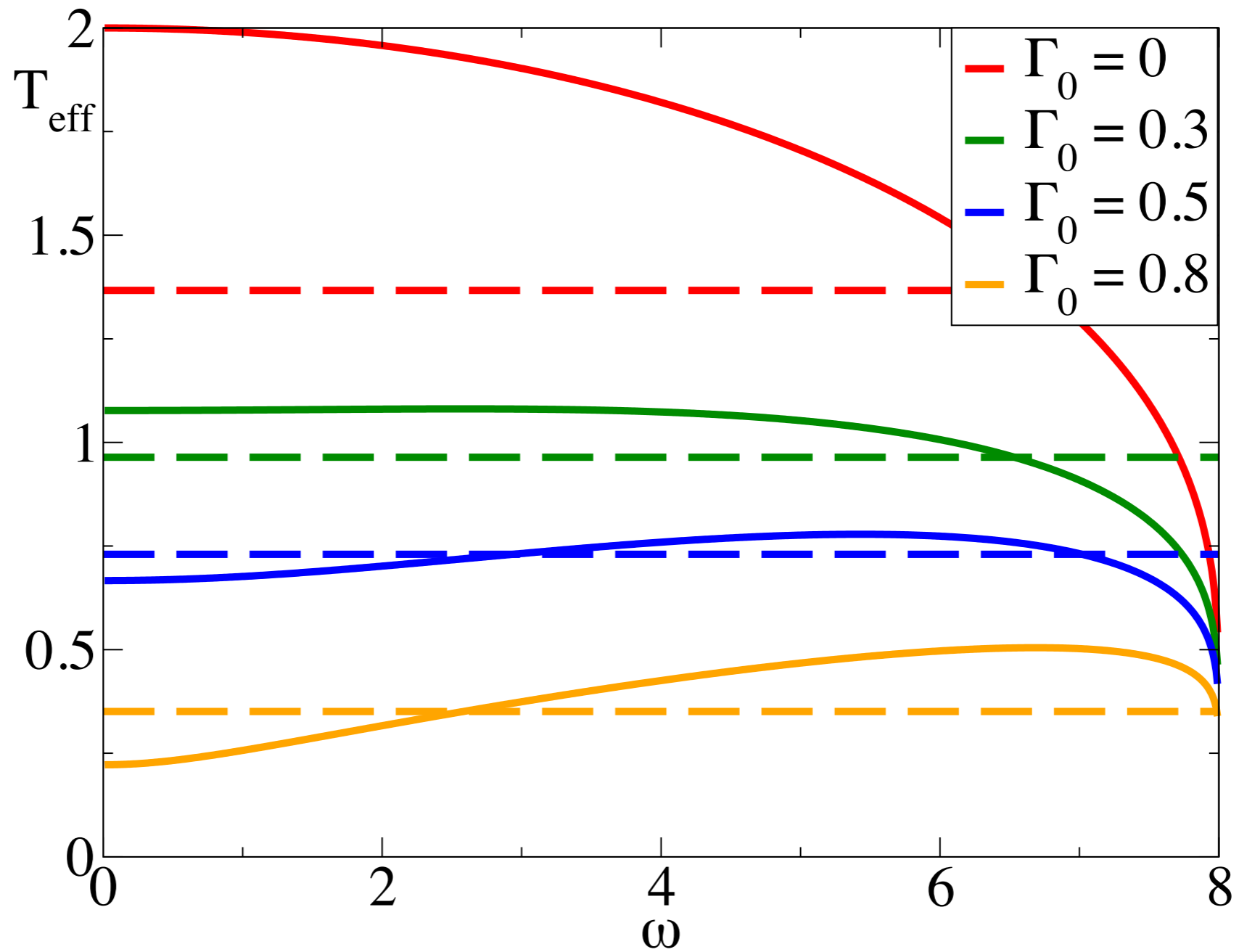
...not really...

...generic ?

Effective temp.s

focus on $\Gamma = 1$
observable \mathcal{O}'
 $\mathcal{O}' \neq \mathcal{O}$

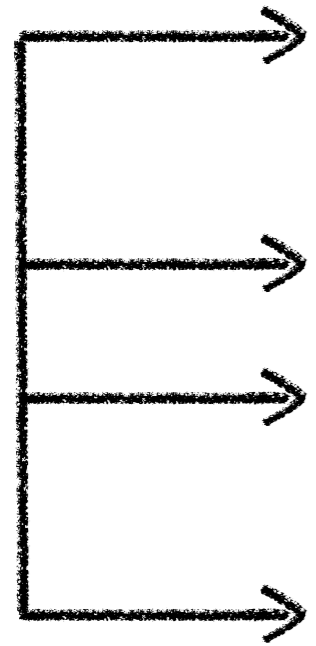
$$T_{\text{eff}}(\omega)$$



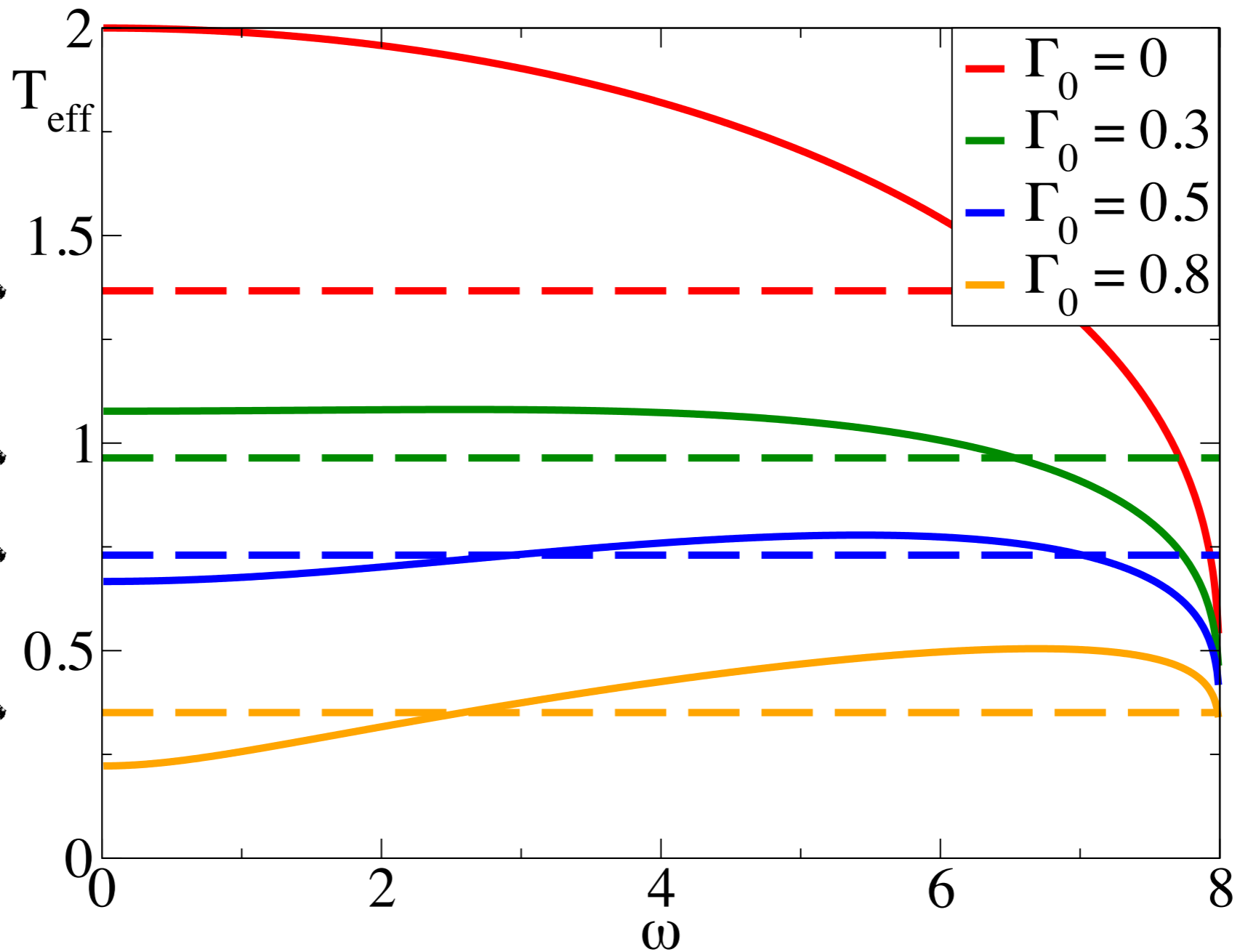
Effective temp.s

focus on $\Gamma = 1$
observable \mathcal{O}'
 $\mathcal{O}' \neq \mathcal{O}$

$T^E(\Gamma_0, \Gamma = 1)$



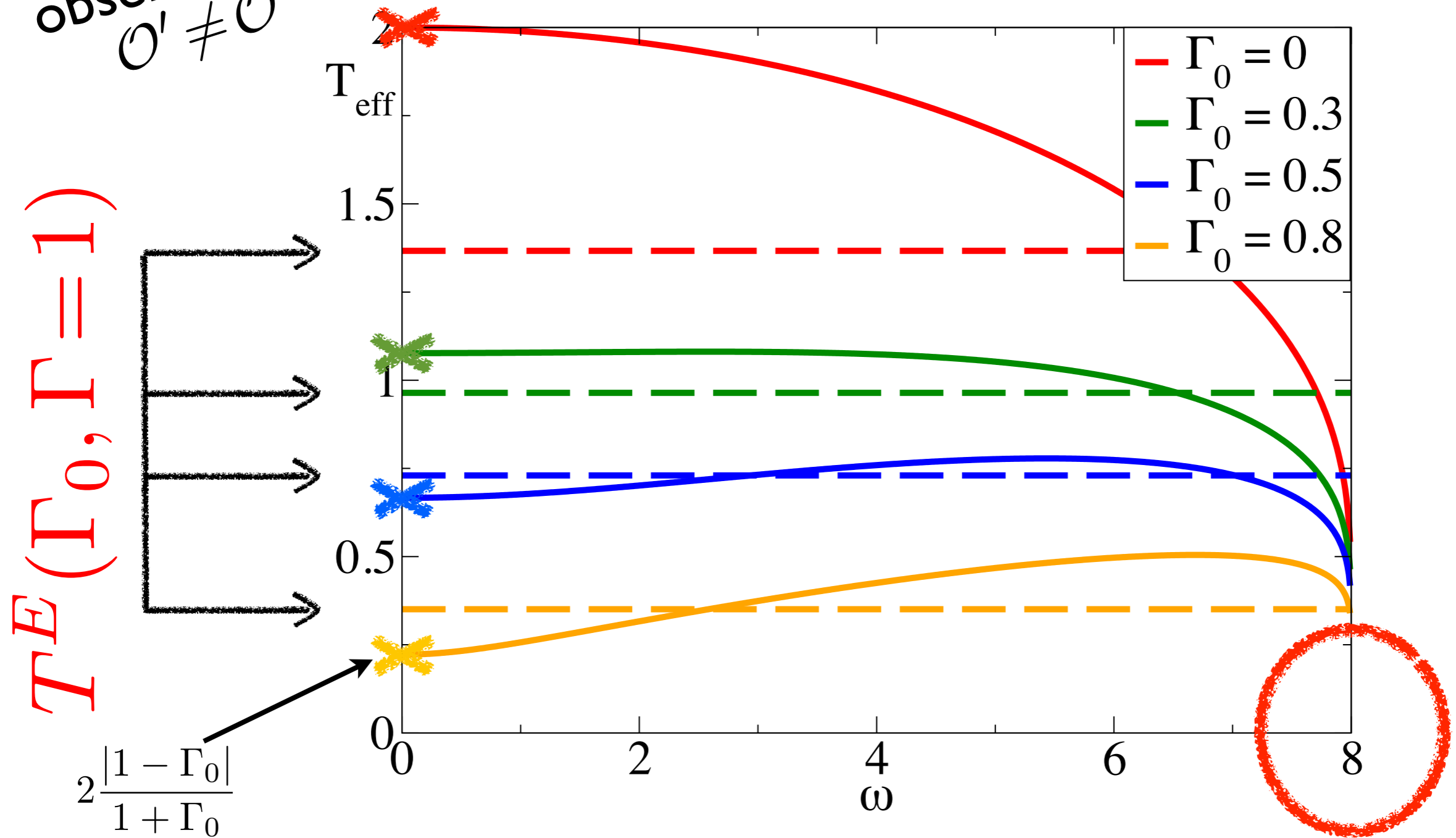
$T_{\text{eff}}(\omega)$



Effective temp.s

focus on $\Gamma = 1$
observable \mathcal{O}'
 $\mathcal{O}' \neq \mathcal{O}$

$T_{\text{eff}}(\omega)$



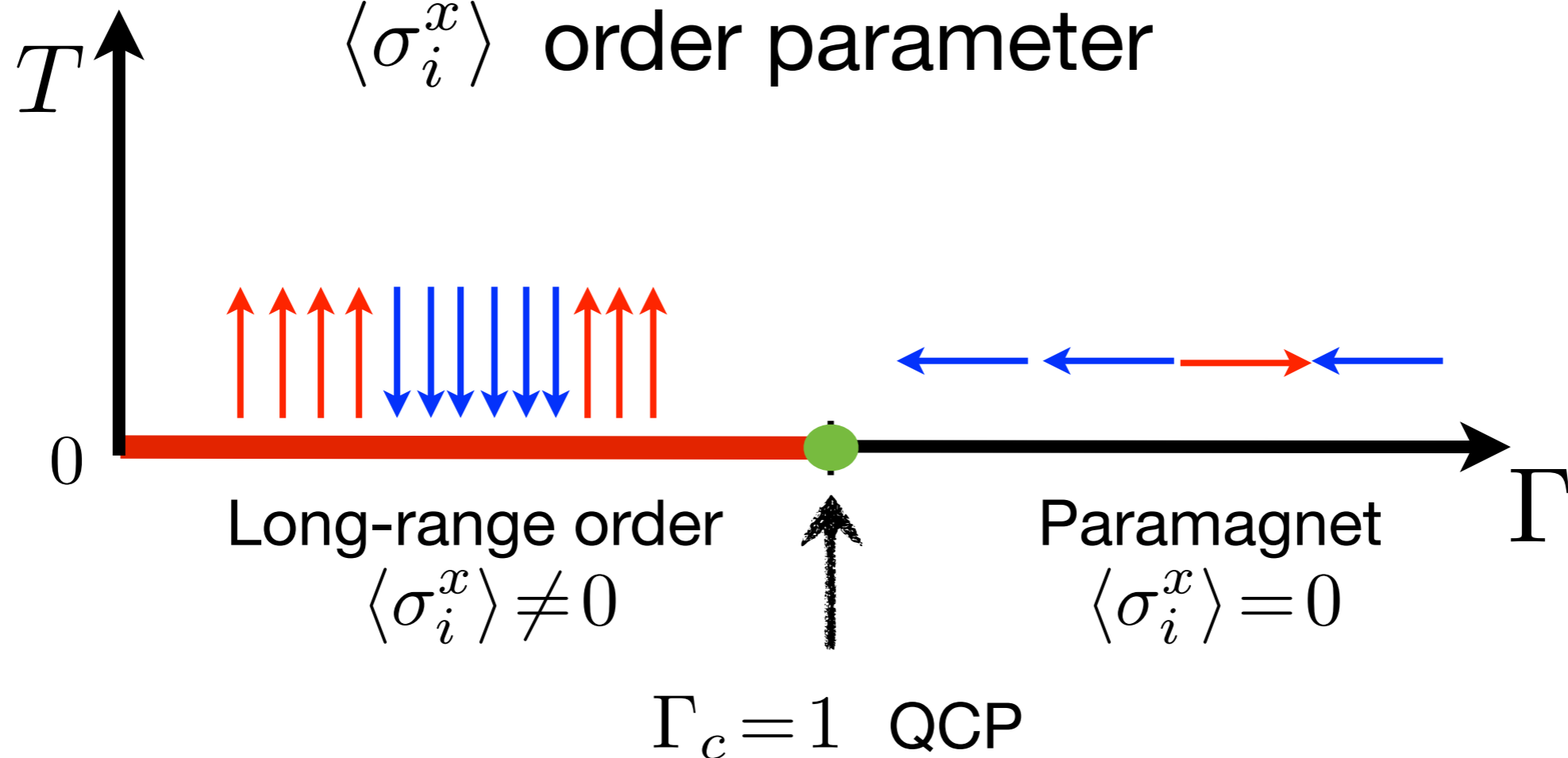
also for $\mathcal{O}' \neq \mathcal{O}$: $R(t), C_+(t)$ invar. $\Gamma_0 \leftrightarrow 1/\Gamma_0$

Transverse field Ising chain

$$H(\Gamma) = -J \sum_i [\sigma_i^x \sigma_{i+1}^x + \Gamma \sigma_i^z]$$

$T=0$ QPT

$\langle \sigma_i^x \rangle$ order parameter



Transverse field Ising chain

$$H(\Gamma) = -J \sum_i [\sigma_i^x \sigma_{i+1}^x + \Gamma \sigma_i^z]$$

Diagonalization: fermions

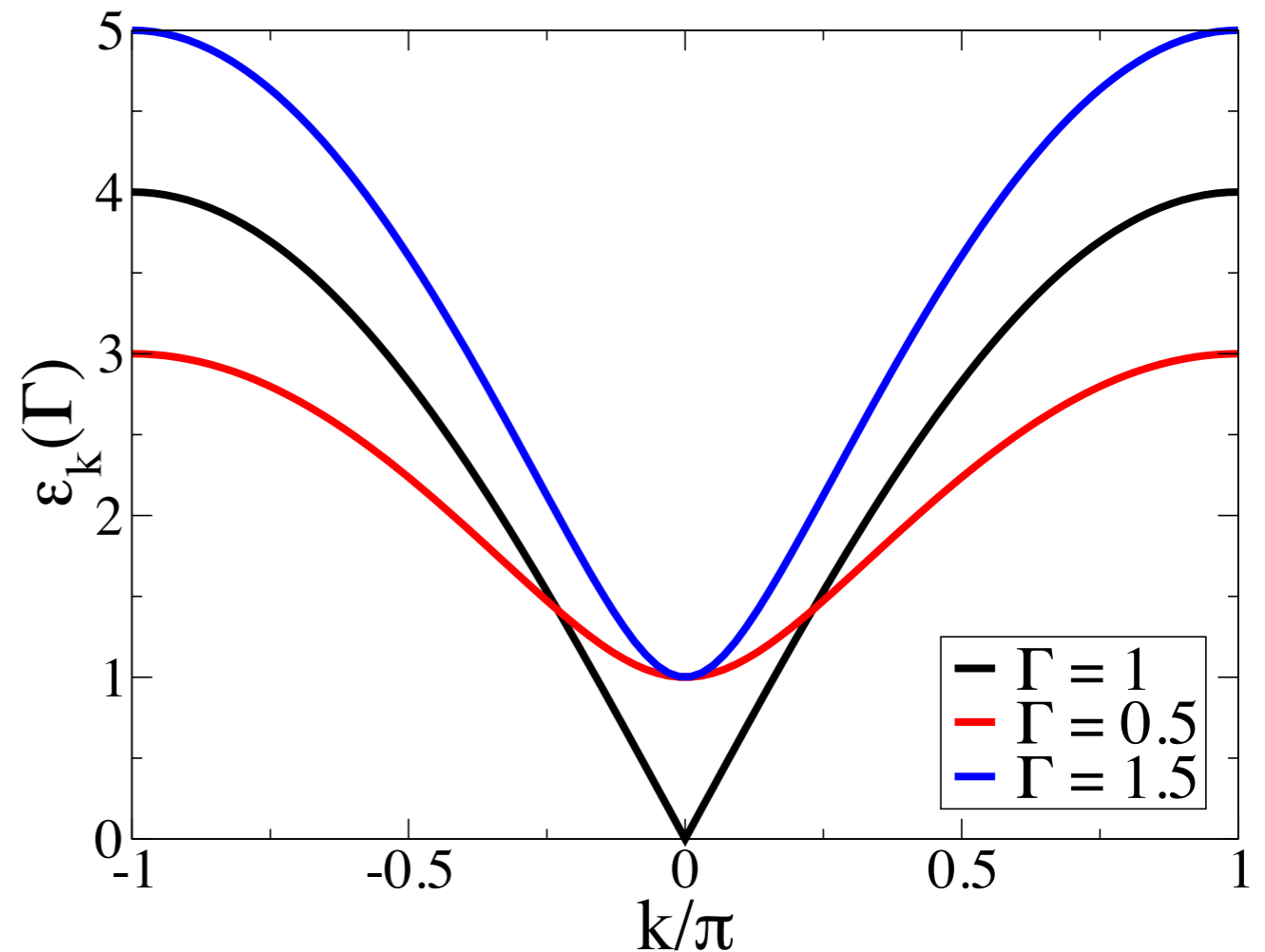
1. Jordan-Wigner:
$$c_i = \left(\prod_{j=1}^{i-1} \sigma_j^z \right) \sigma_i^- \quad \begin{aligned} \{c_i^\dagger, c_j\} &= \delta_{i,j} \\ \{c_i, c_j\} &= \{c_i^\dagger, c_j^\dagger\} = 0 \end{aligned}$$
2. Fourier transform (PBC):
$$c_k = \frac{1}{\sqrt{L}} \sum_j e^{-ikj} c_j$$
3. Bogoliubov rotation:
$$c_k = u_k \eta_k + v_k \eta_{-k}^\dagger$$

Spectrum

$$H(\Gamma) = \sum_k \epsilon_k(\Gamma) (\eta_k^\dagger \eta_k - 1/2)$$

$$\epsilon_k(\Gamma) = 2\sqrt{1 + \Gamma^2 - 2\Gamma \cos k}$$

$$\text{GS: } \eta_k |0\rangle = 0$$



Dynamics

Heisenberg picture:

$$\mathcal{O}(t) = e^{iHt} \mathcal{O} e^{-iHt}$$

Dyn. of excitations:

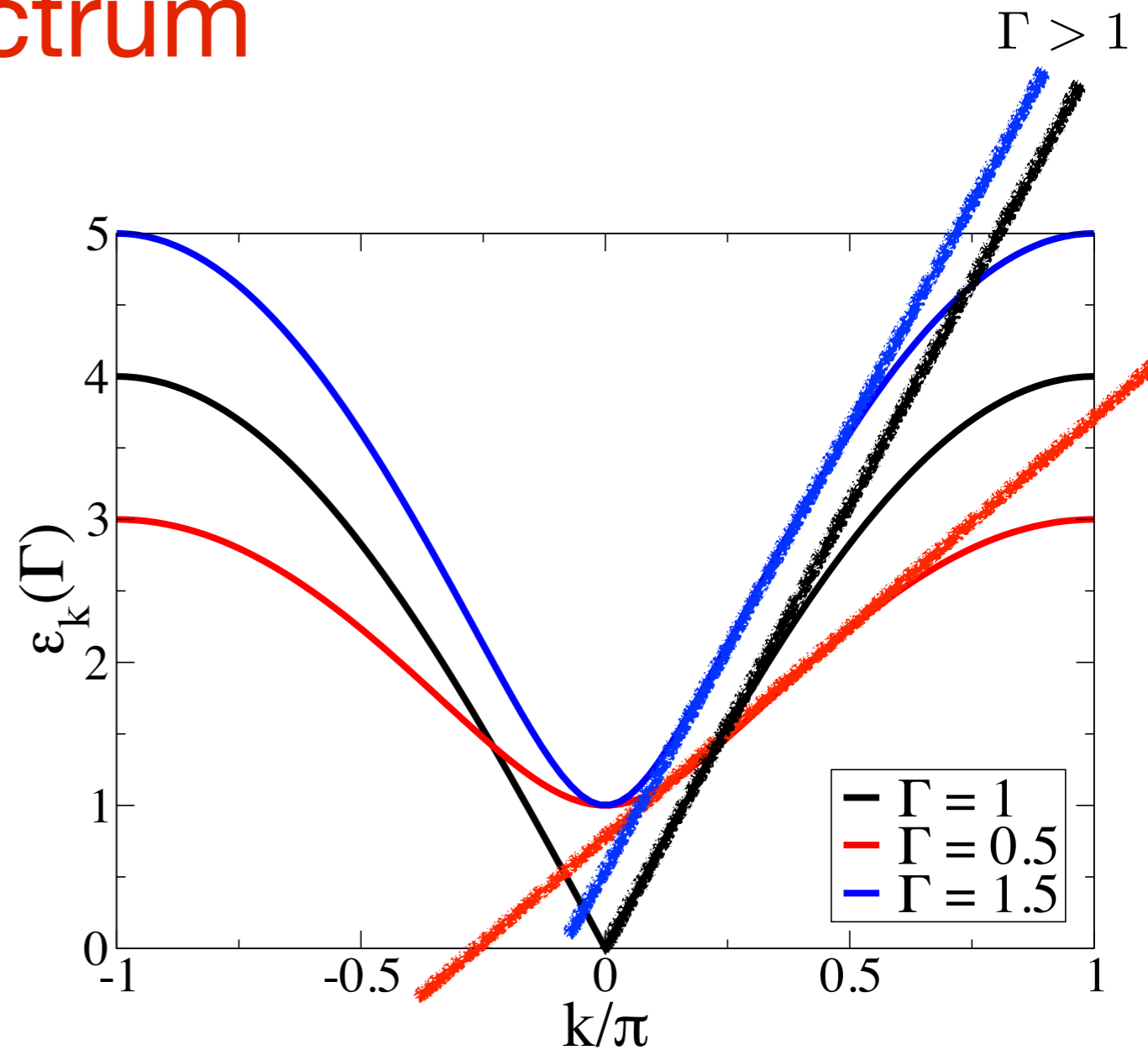
$$\eta_k(t) = e^{-i\epsilon_k(\Gamma)t} \eta_k \implies n_k \equiv \eta_k^\dagger \eta_k \text{ conserved}$$

Spectrum

$$H(\Gamma) = \sum_k \epsilon_k(\Gamma) (\eta_k^\dagger \eta_k - 1/2)$$

$$\epsilon_k(\Gamma) = 2\sqrt{1 + \Gamma^2 - 2\Gamma \cos k}$$

GS: $\eta_k |0\rangle = 0$



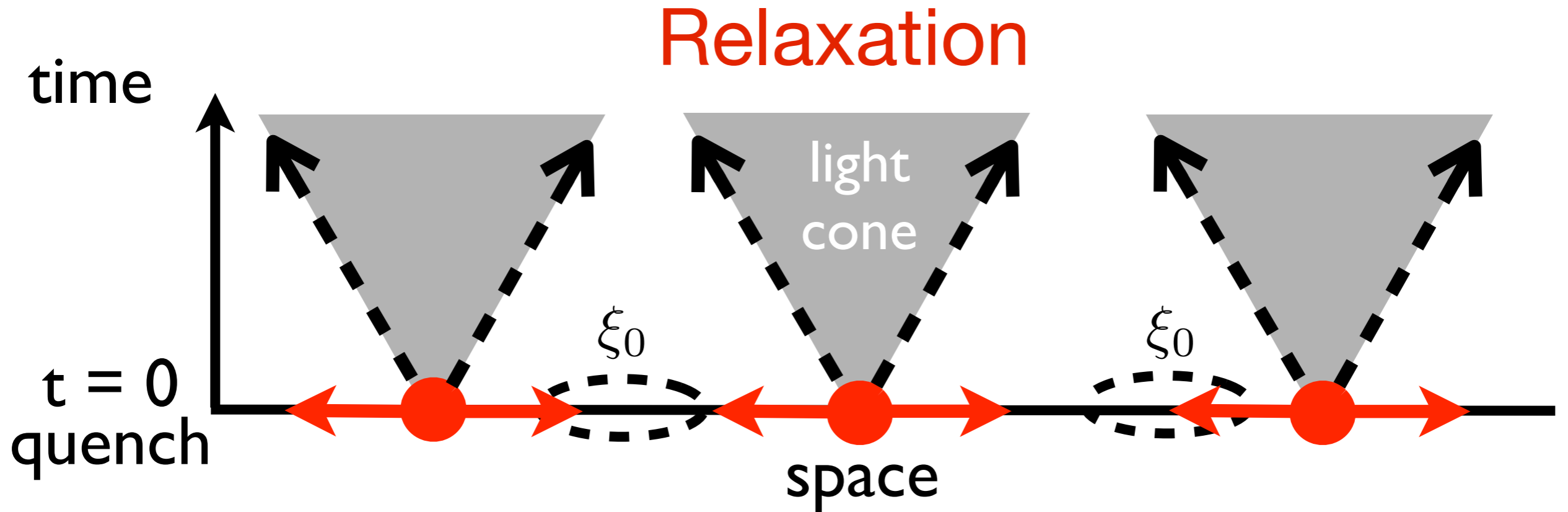
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$$\mathcal{O}(t) = e^{iHt} \mathcal{O} e^{-iHt}$$

Dyn. of excitations:

$$\eta_k(t) = e^{-i\epsilon_k(\Gamma)t} \eta_k \implies n_k \equiv \eta_k^\dagger \eta_k \text{ conserved}$$



- $|\psi_0\rangle = |\text{GS}(H(\Gamma_0))\rangle$ has extensive energy
- Ballistic excit.s, plane waves
- Finite systems \rightarrow recurrence

(eq.) Sachdev, Young '97

Igloi, Rieger, 2000, >2010

Calabrese, Cardy, >2006

Rossini, Mussardo, Santoro, Silva, Suzuki, >2009

Calabrese, Essler, Fagotti, >2011

observables: $\mathcal{O} \equiv \sigma_i^z$ ($\propto c_i^\dagger c_i$ density of fermions)

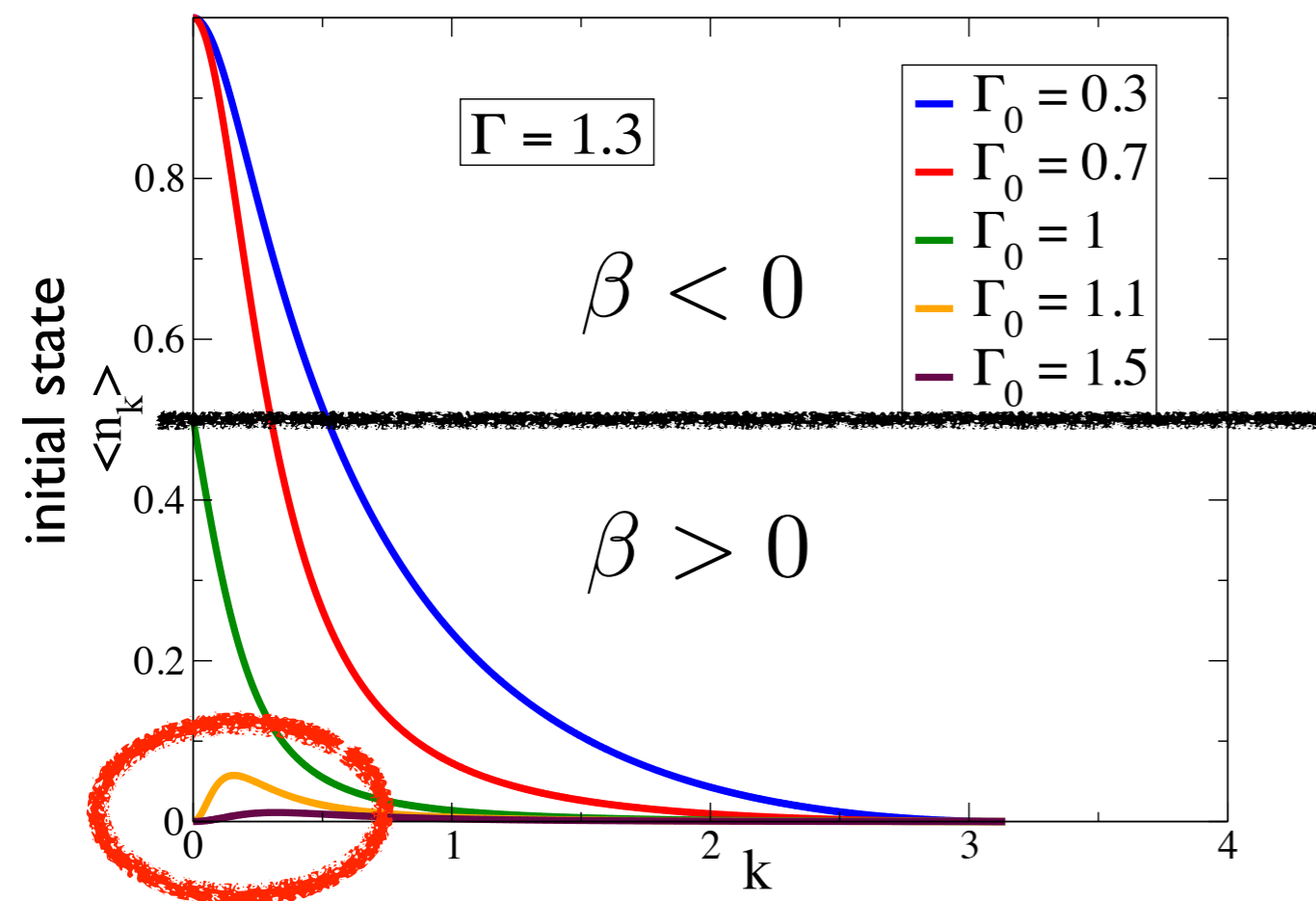
$$\mathcal{O}' \equiv M(t) = \frac{1}{L} \sum_i \sigma_i^z \quad C, R \propto L \times \langle [M, M]_{\pm} \rangle$$

Why critical quenches?

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Why critical quenches?



$$\langle n_k \rangle \leftrightarrow \frac{1}{1 + e^{\beta \epsilon_k(\Gamma)}}$$

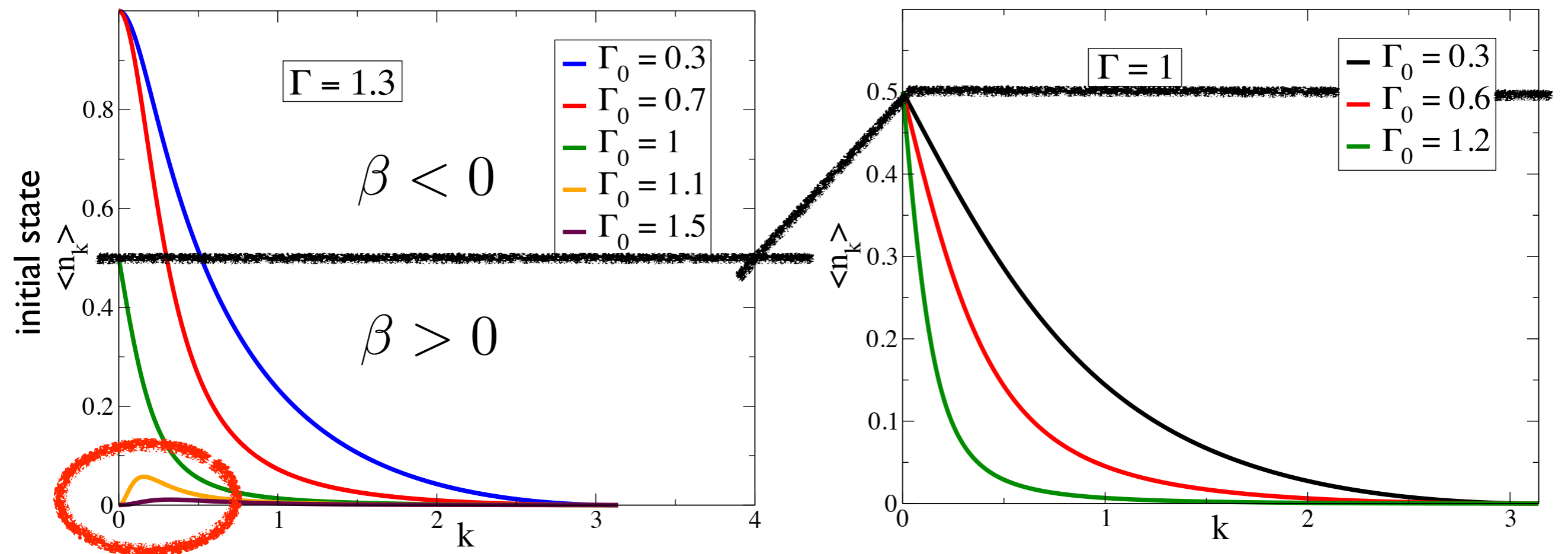
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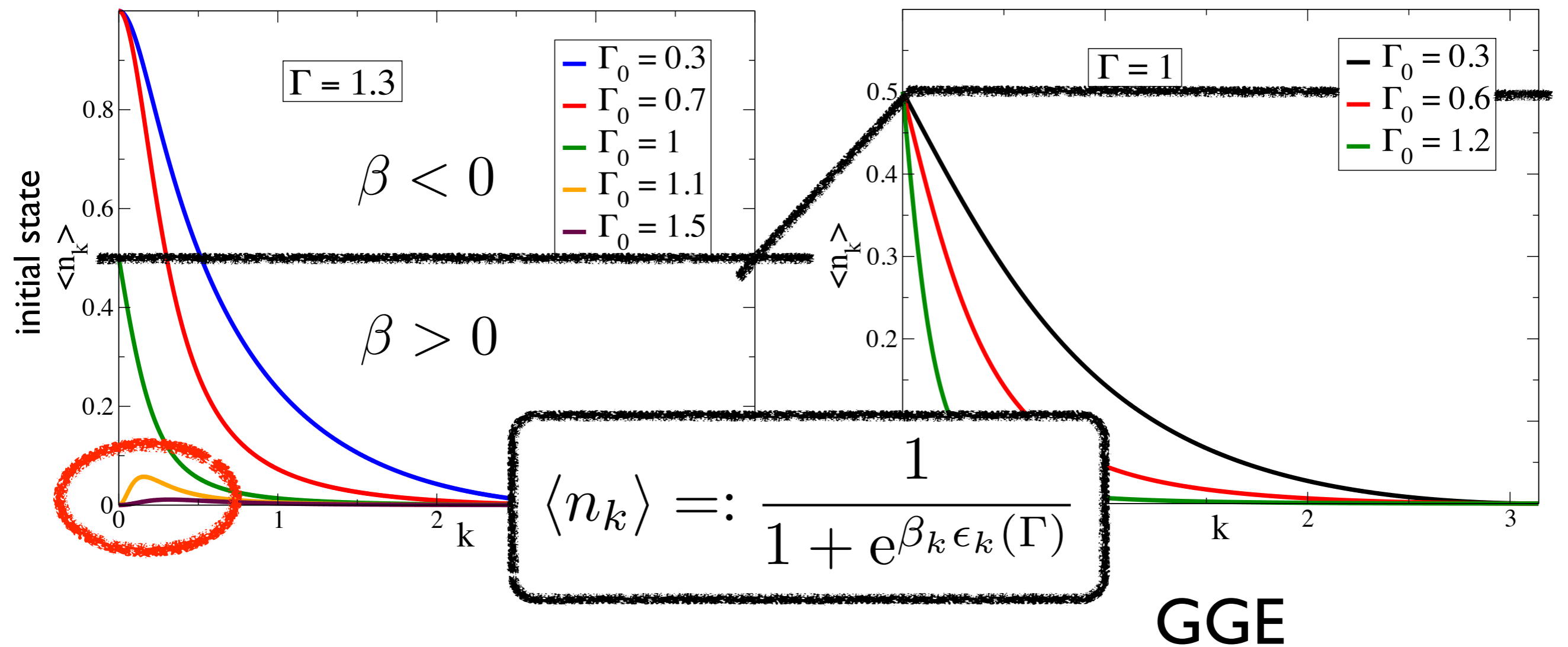
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Why critical quenches?



because....

- $\{\langle \psi_0 | n_k | \psi_0 \rangle\}_k$ “more” thermal
- σ_i^z “special”
- no lattice $\Rightarrow \epsilon_k \sim |k|$ CFT

[Calabrese, Cardy '06]

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[Calabrese, Cardy '06]

No Gibbs....

... GGE

[...-Calabrese, Essler, Fagotti-Blass, Rieger, Igloi]

what can we learn?
effective “Gibbs”?

$$\beta_k^{-1} = T_{\text{eff}}^M (\omega = 2\epsilon_k(\Gamma))$$

“measure” temp.s
of GGE

➤ local vs. non-local ?

σ_i^z, M

σ_i^x

[Rossini et al. 2010]

Role of “locality”

Transverse magnetization

$$\sigma_j^z = 2c_j^\dagger c_j - 1$$

Order parameter

$$\sigma_j^x = \exp \left[i\pi \sum_{l=1}^{j-1} c_l^\dagger c_l \right] (c_j^\dagger + c_j)$$

[Rossini, Suzuki, Mussardo, Santoro, Silva '09]

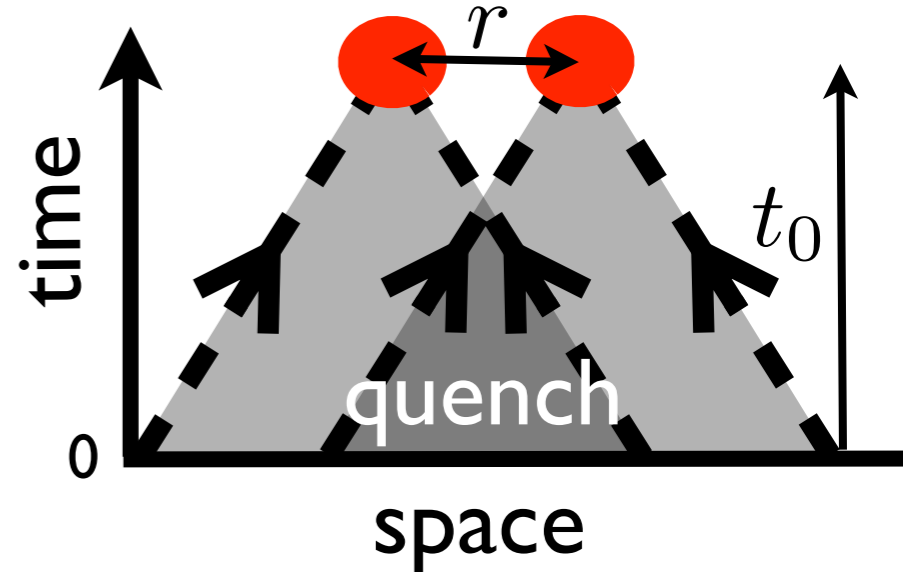
focus on $\Gamma=1$

Order parameter

$\langle \sigma^x(t) \rangle \xrightarrow[t \rightarrow \infty]{d > 1?} 0 \quad \forall \Gamma \neq \Gamma_0$
 $\sim e^{-t/\tau}$

thermal behav.
 @ $T \neq 0$

$\langle \sigma_i^x(t_0) \sigma_{i+r}^x(t_0) \rangle \underset{(t_0=\infty)}{\sim} e^{-r/\xi}$



$$\tau^{-1} = - \int_0^\pi \frac{dk}{\pi} \epsilon'_k(\Gamma) \ln |1 - 2\langle n_k \rangle|$$

$$\xi^{-1} = \dots \quad 1 \quad \dots$$

GGE

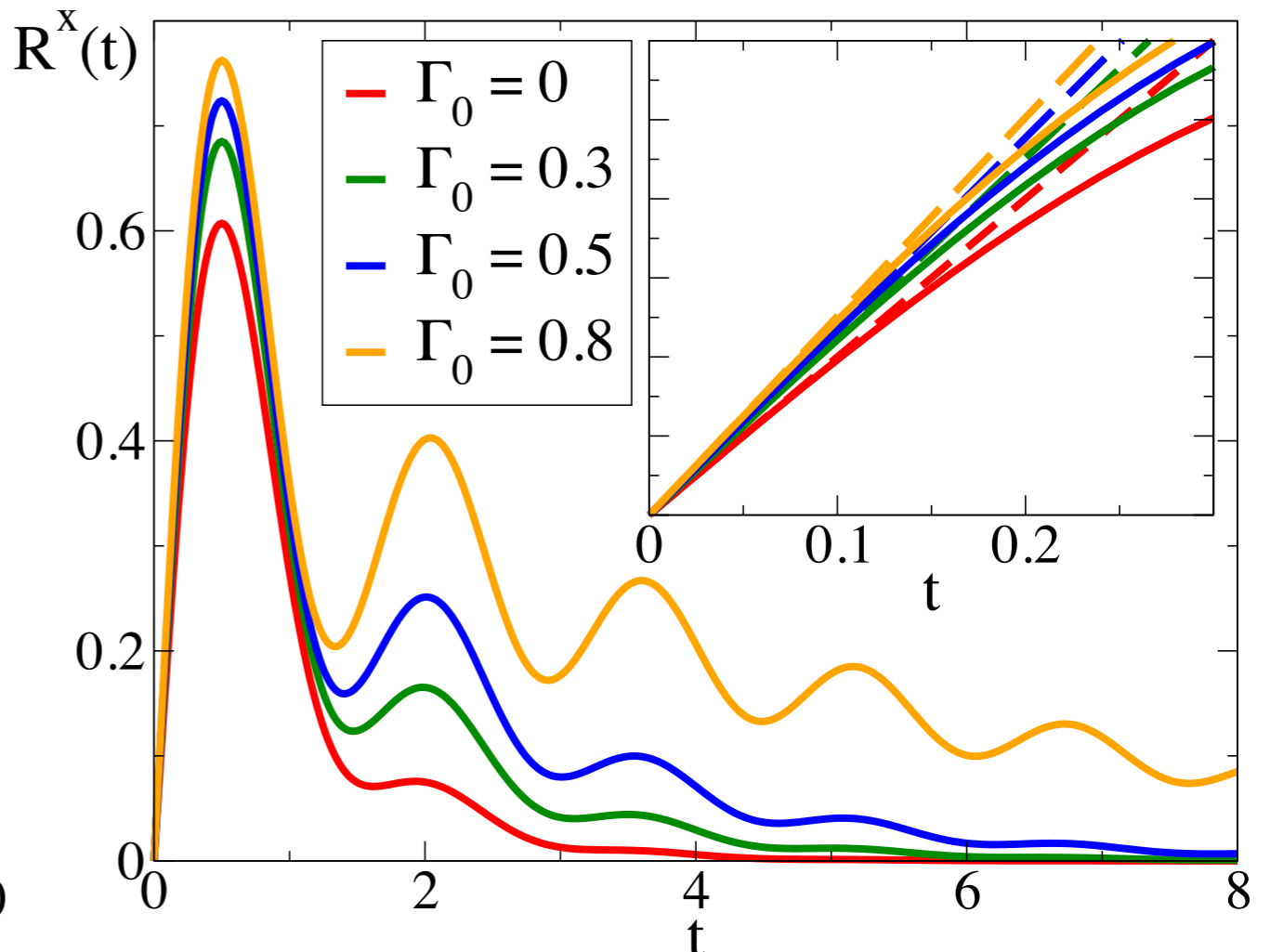
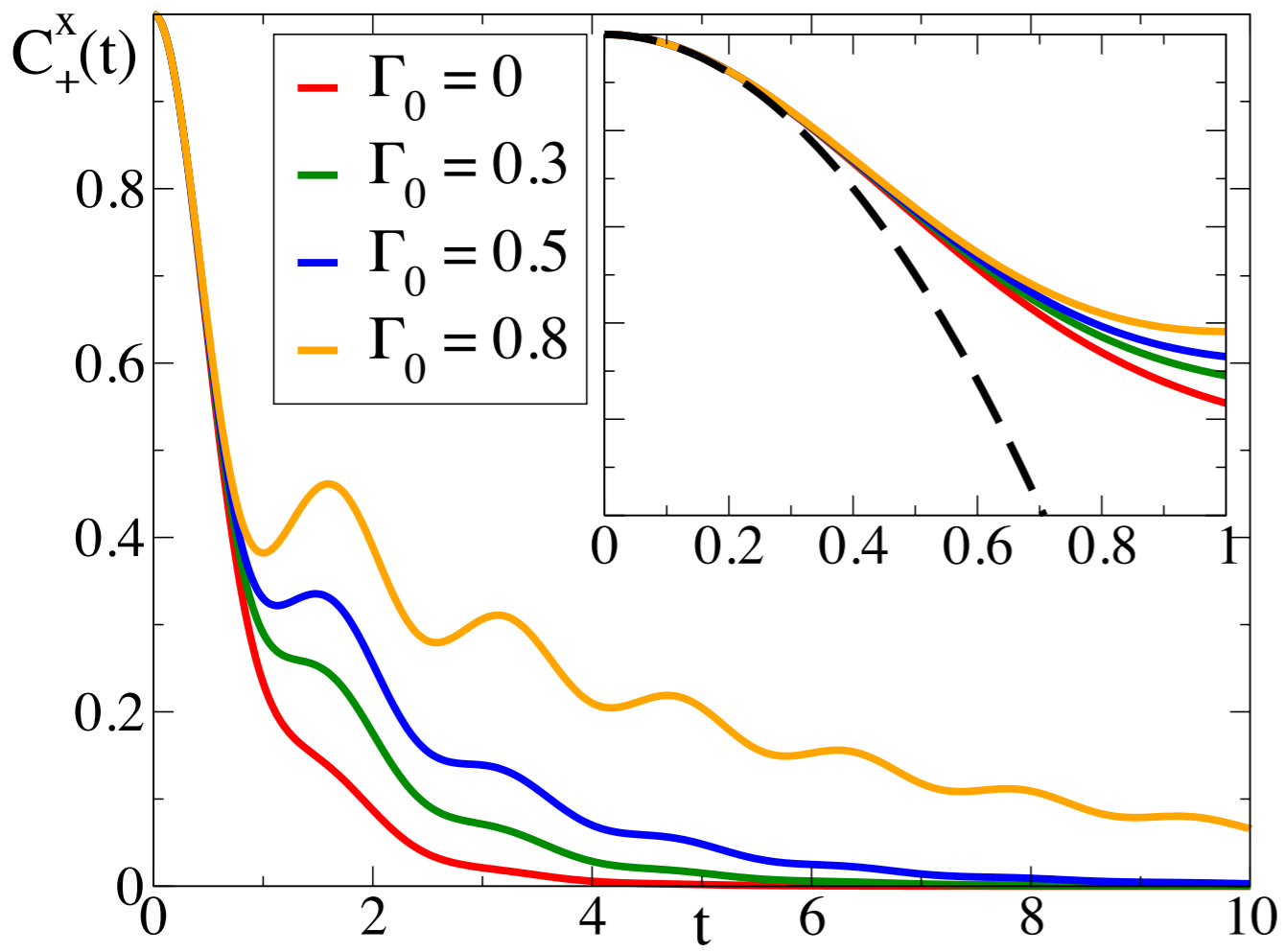
$\langle \sigma_i^x(t+t_0) \sigma_i^x(t_0) \rangle \sim e^{-t/\tau} \quad \tau \simeq \tau_{\text{eq}}(T = T^E)$

...we focus on the stationary regime

focus on $\Gamma = 1$

Order parameter $r = 0$

$$C^x(t \gtrsim 4) \simeq A_C(\Gamma_0) e^{-t/\tau(\Gamma_0)} \left[1 + \frac{a_C \approx 0.21}{\sqrt{t}} \sin(4t + \phi) \right]$$
$$A_C/A_R = 1.210(5)$$
$$R^x(t \gtrsim 4) \simeq A_R(\Gamma_0) e^{-t/\tau(\Gamma_0)} \left[1 - \frac{a_R \approx 0.52}{\sqrt{t}} \cos(4t + \phi) \right]$$



focus on $\Gamma = 1$

Order parameter $r = 0$

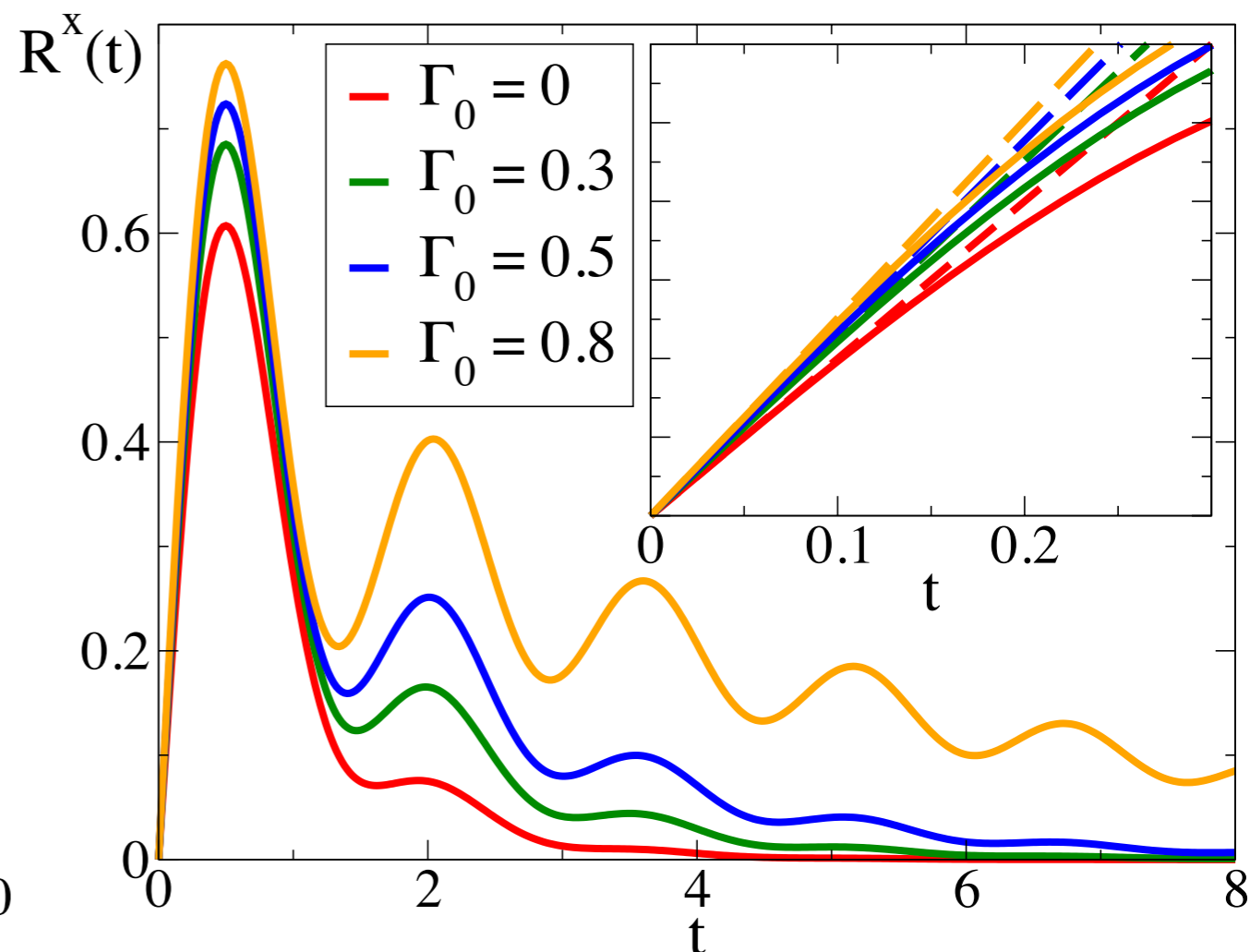
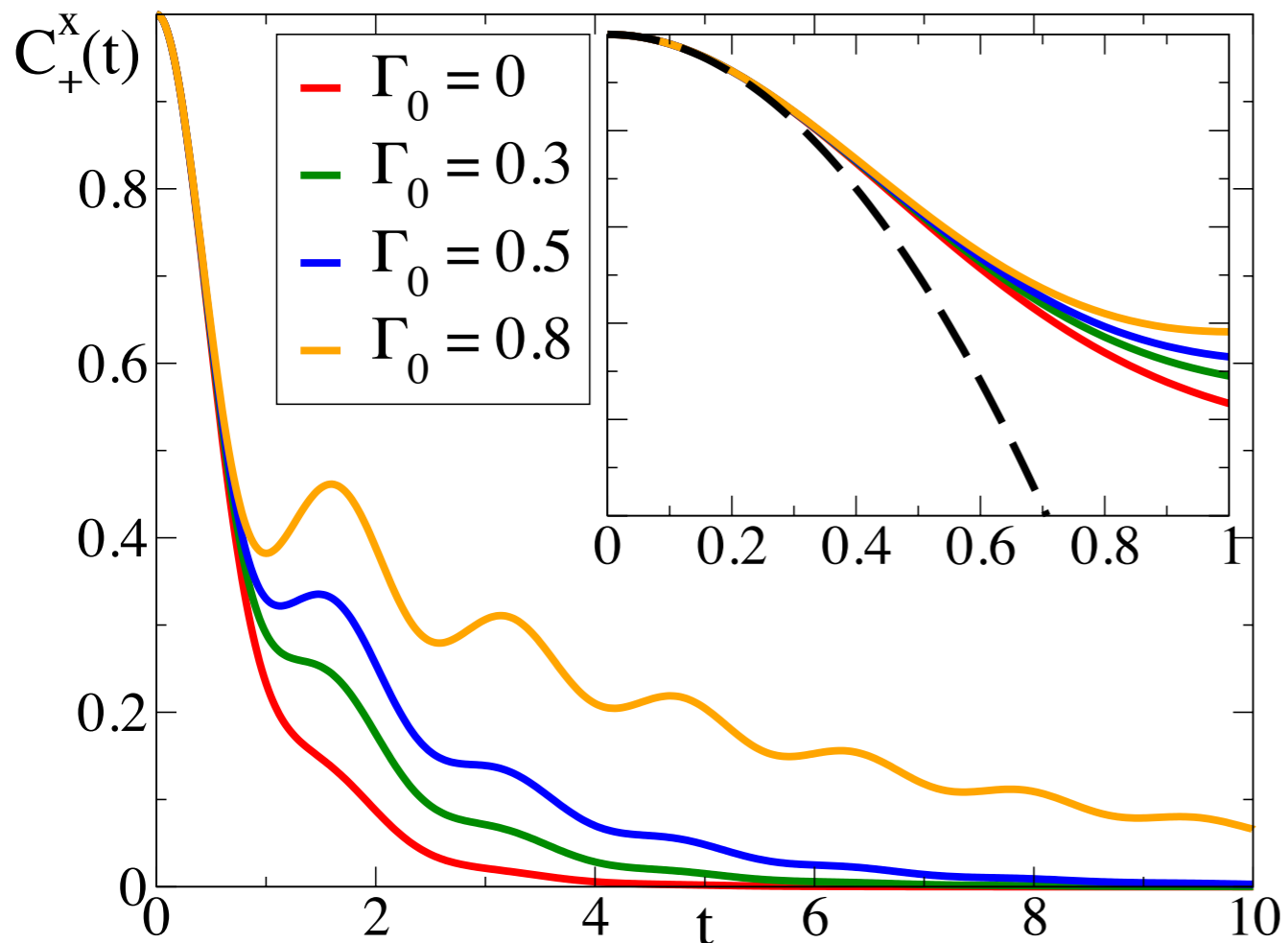
**fast
relaxation**

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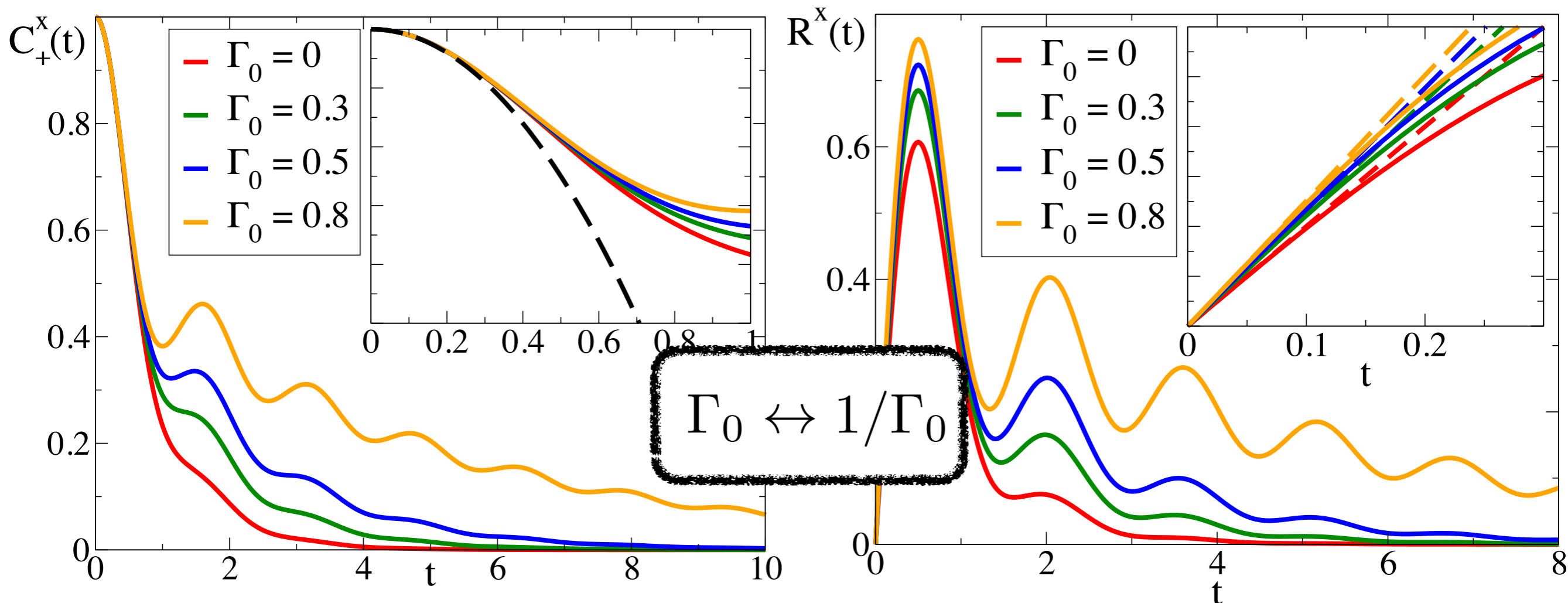
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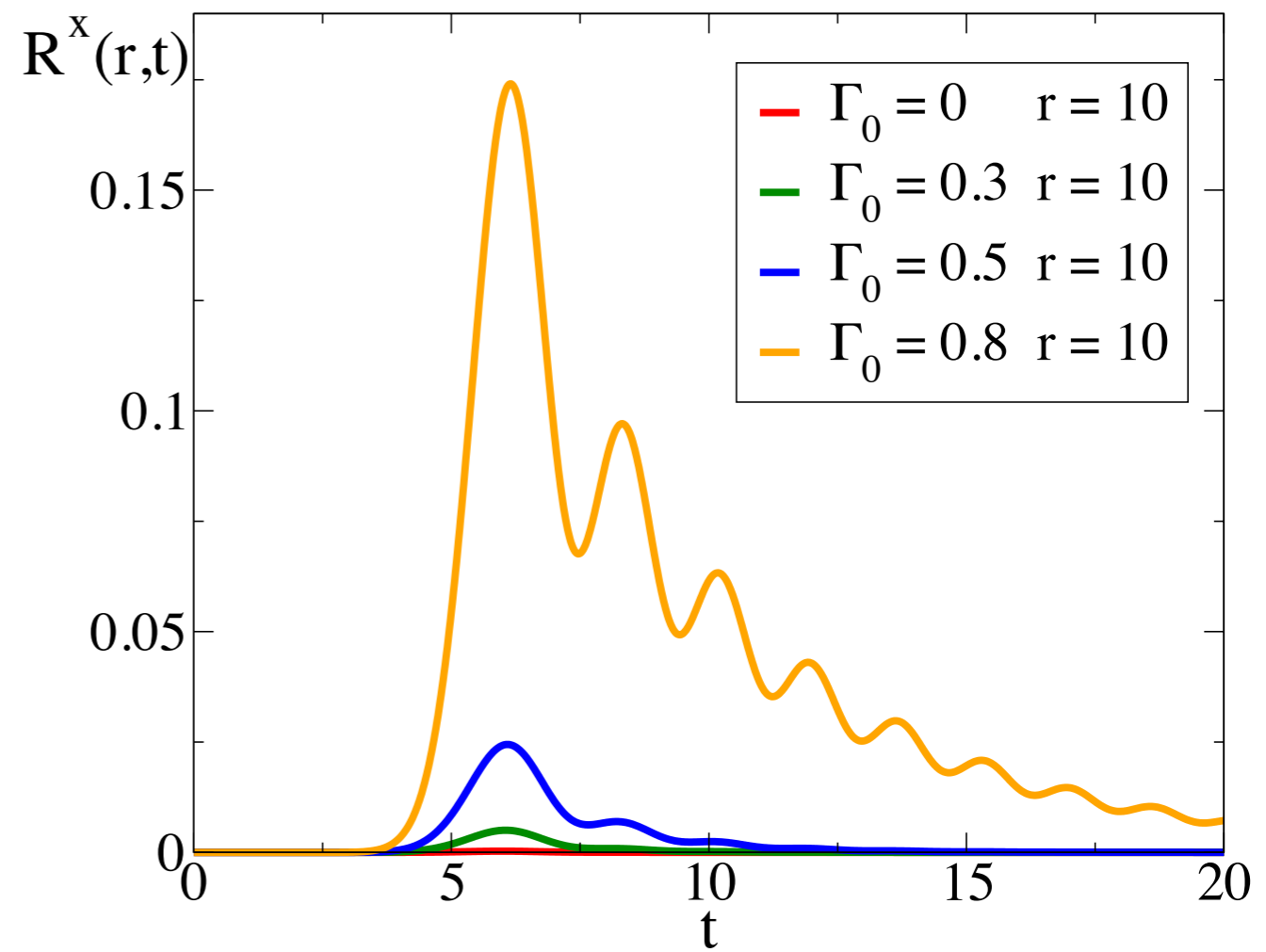
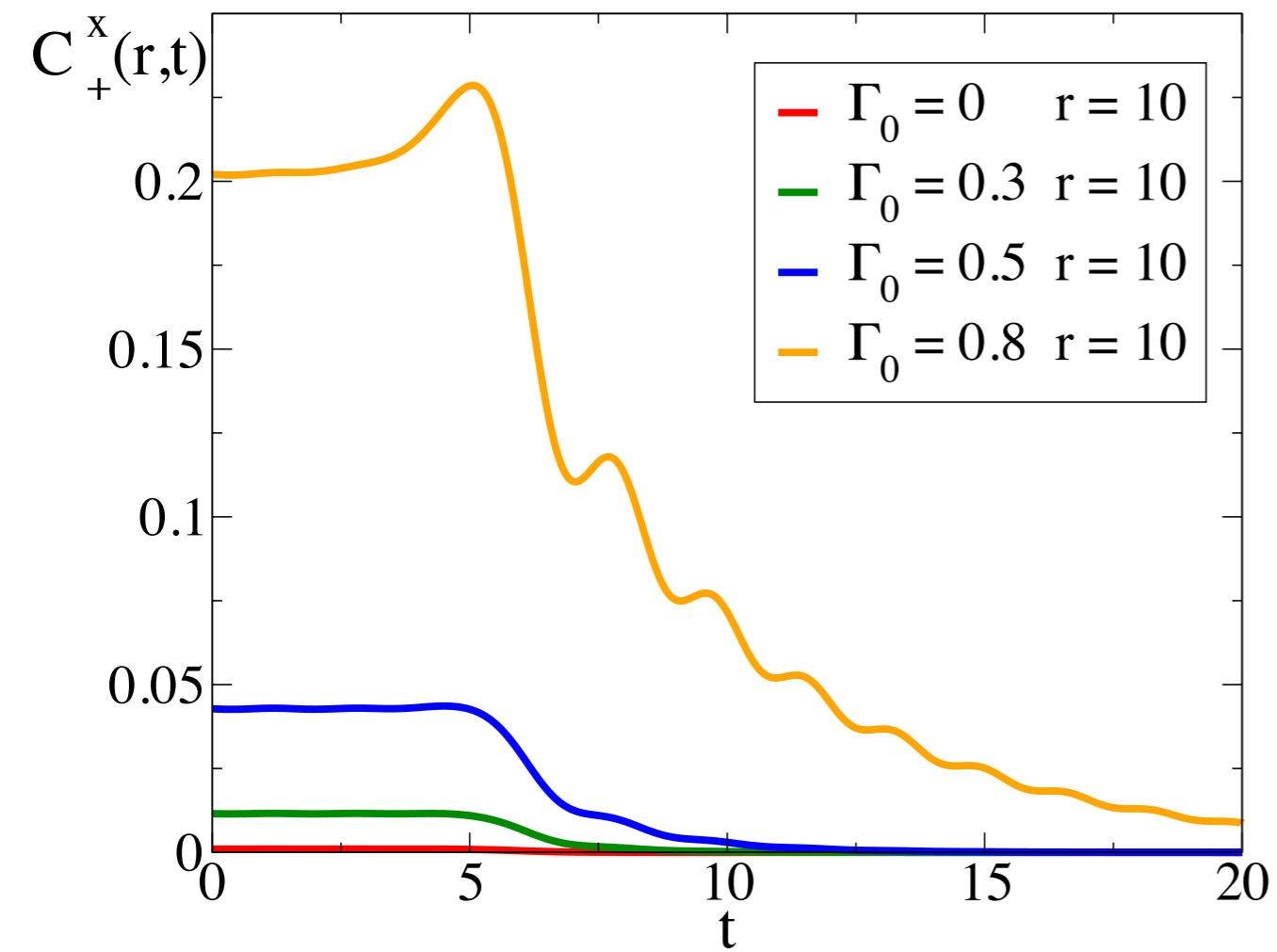
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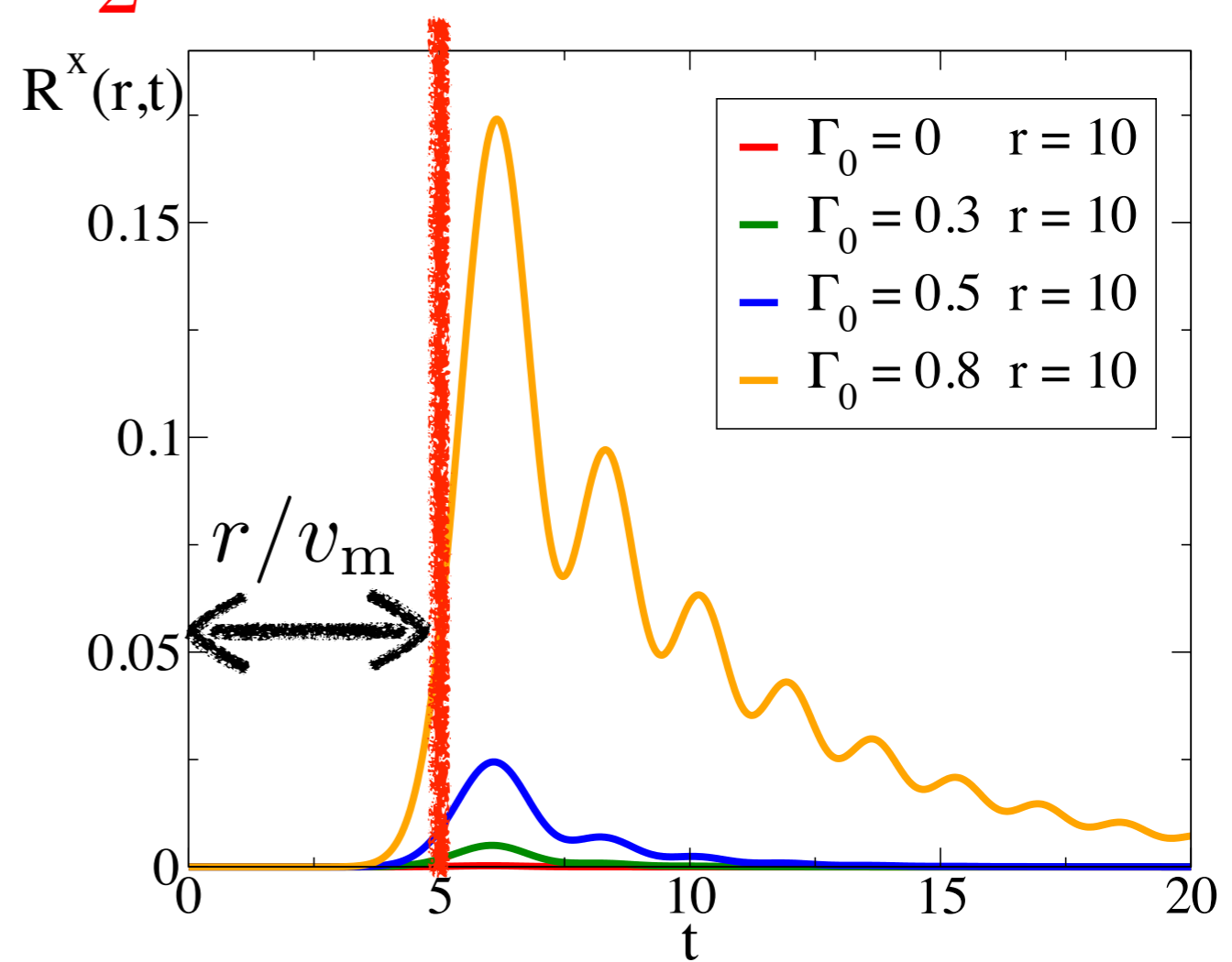
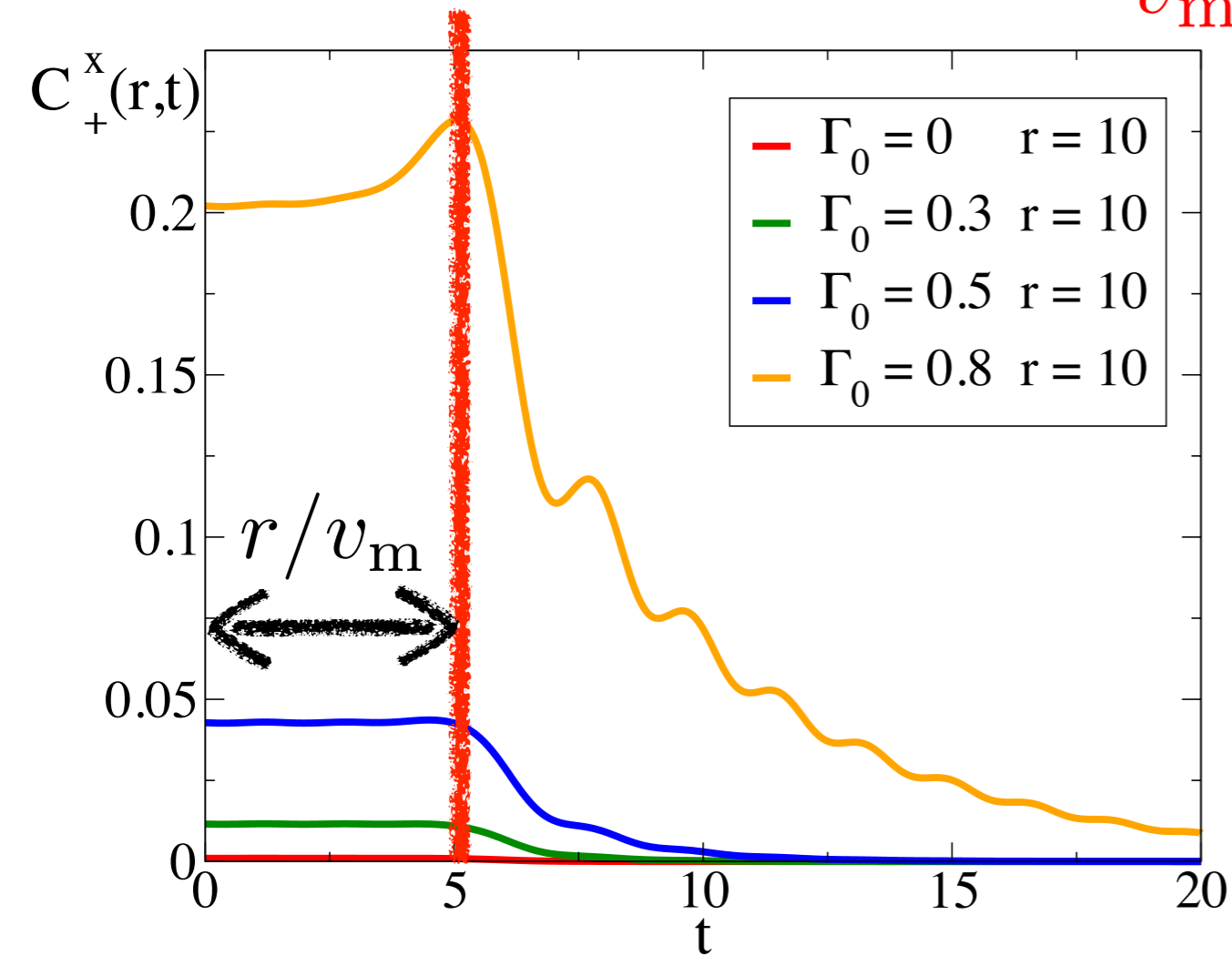
Order parameter $r \neq 0$



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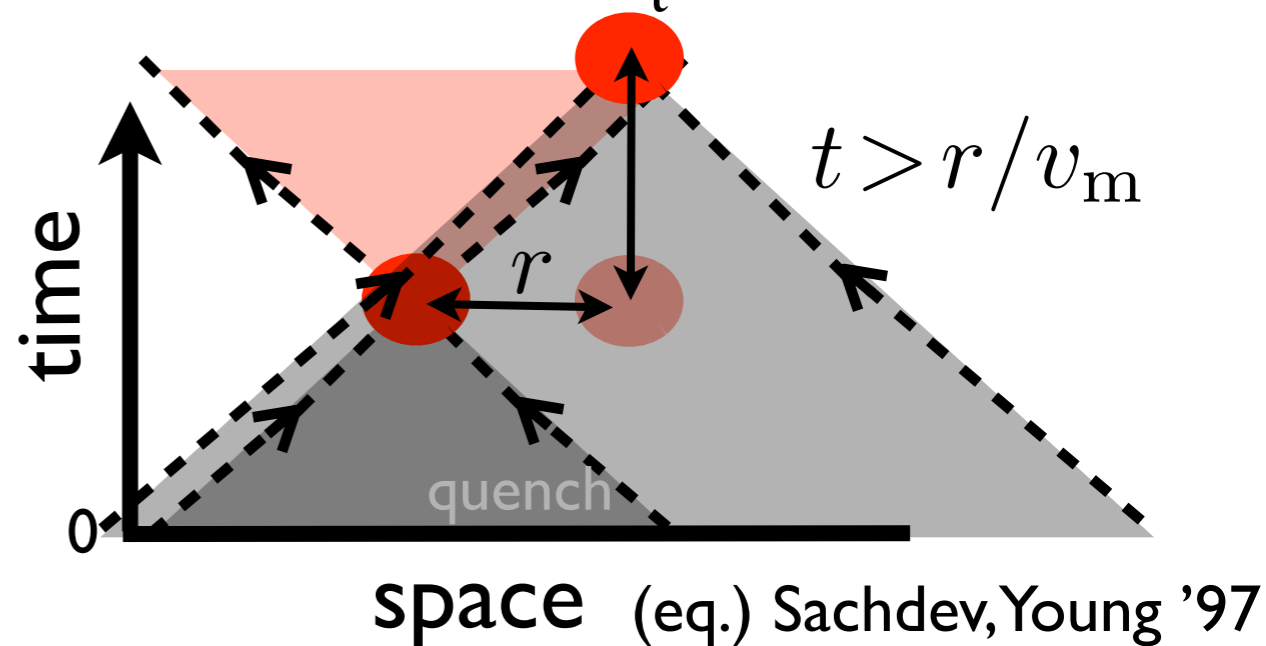
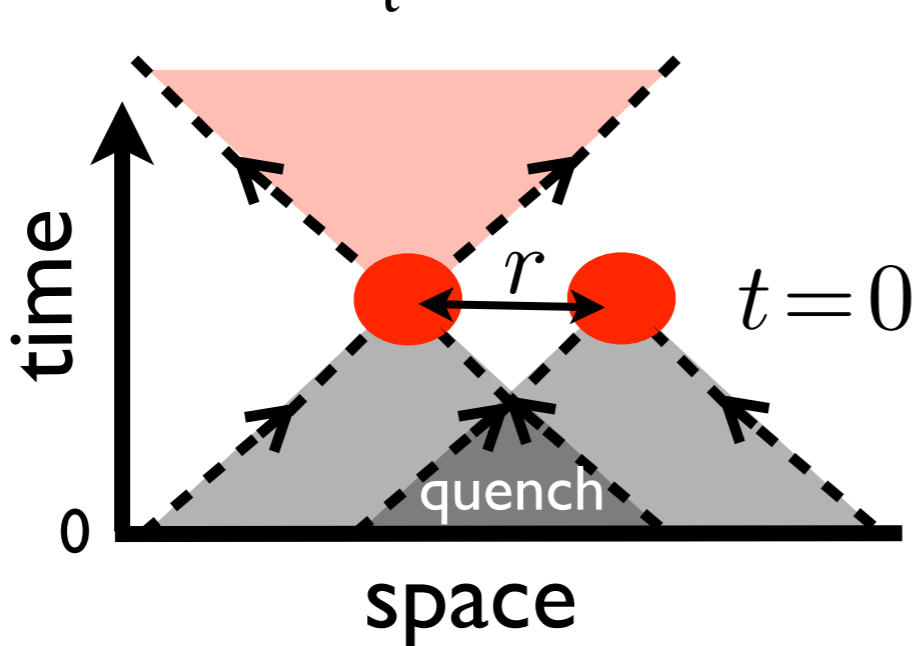
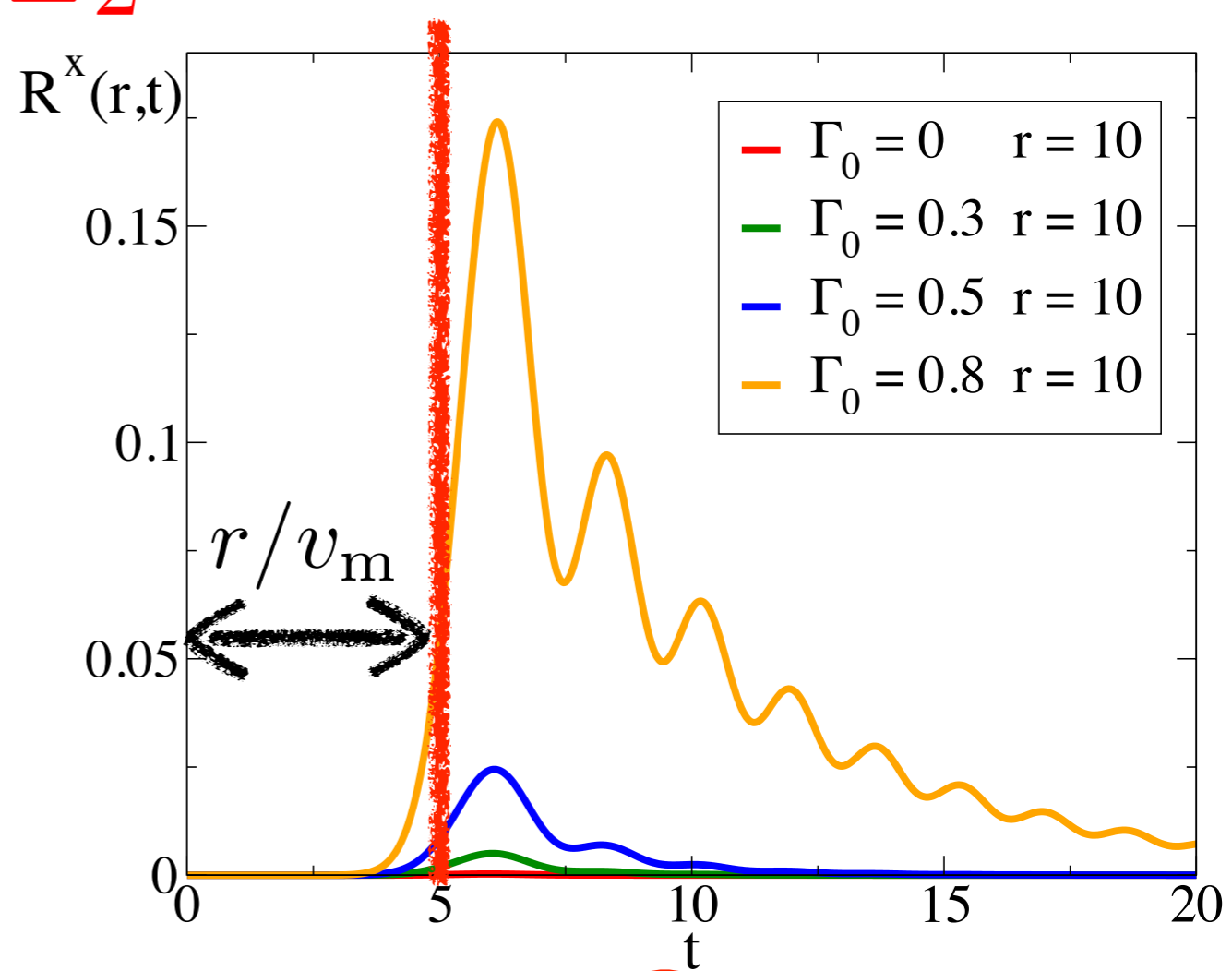
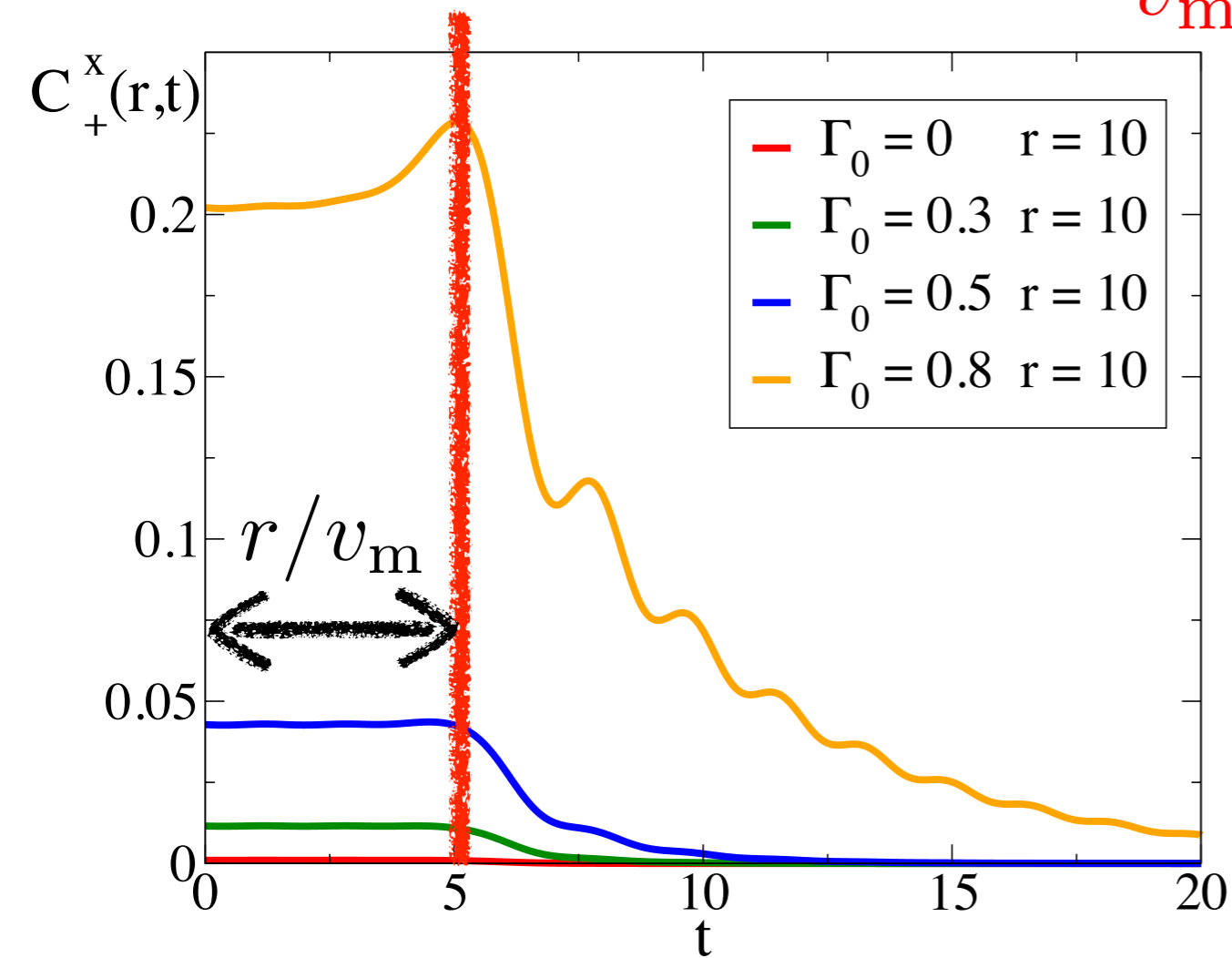
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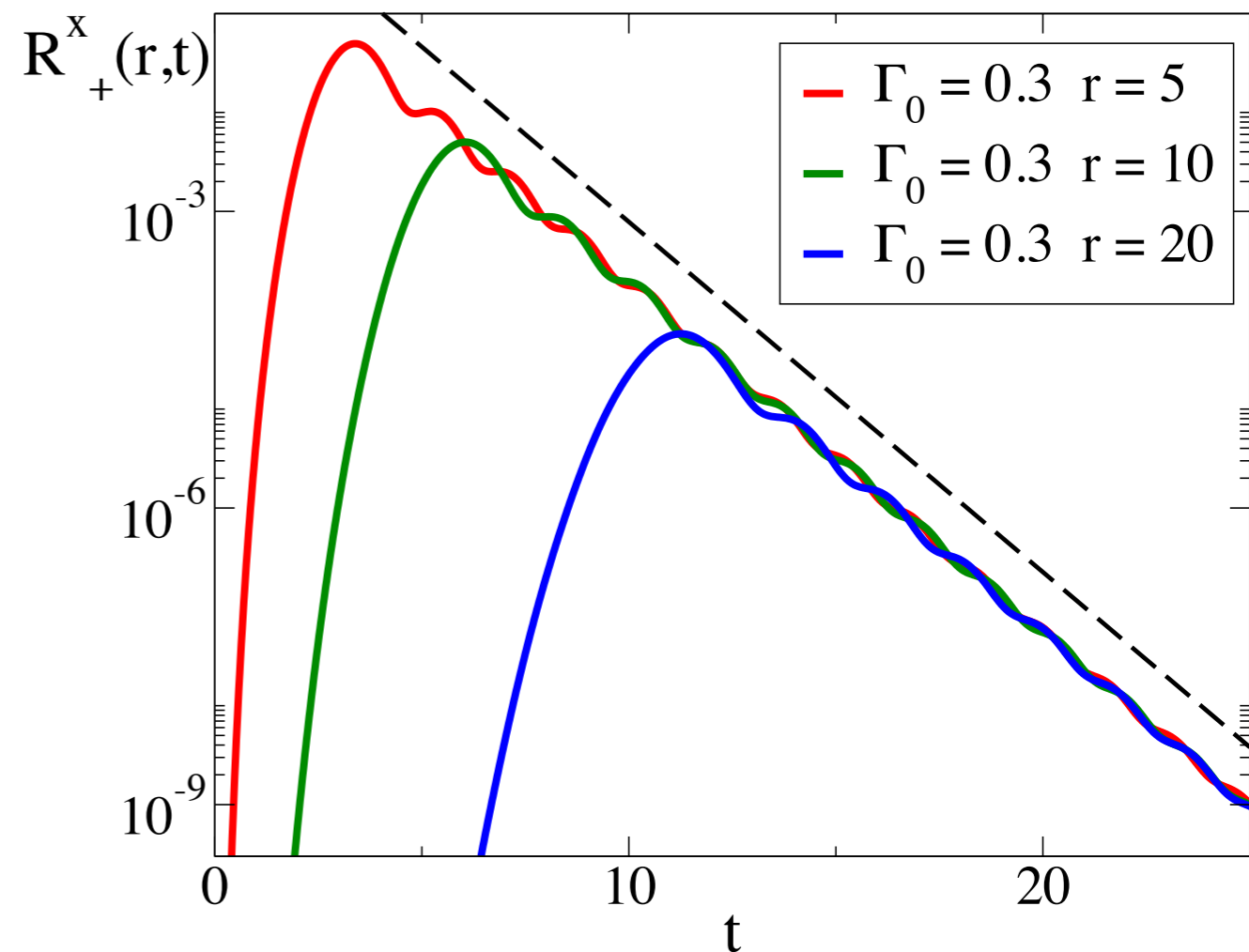
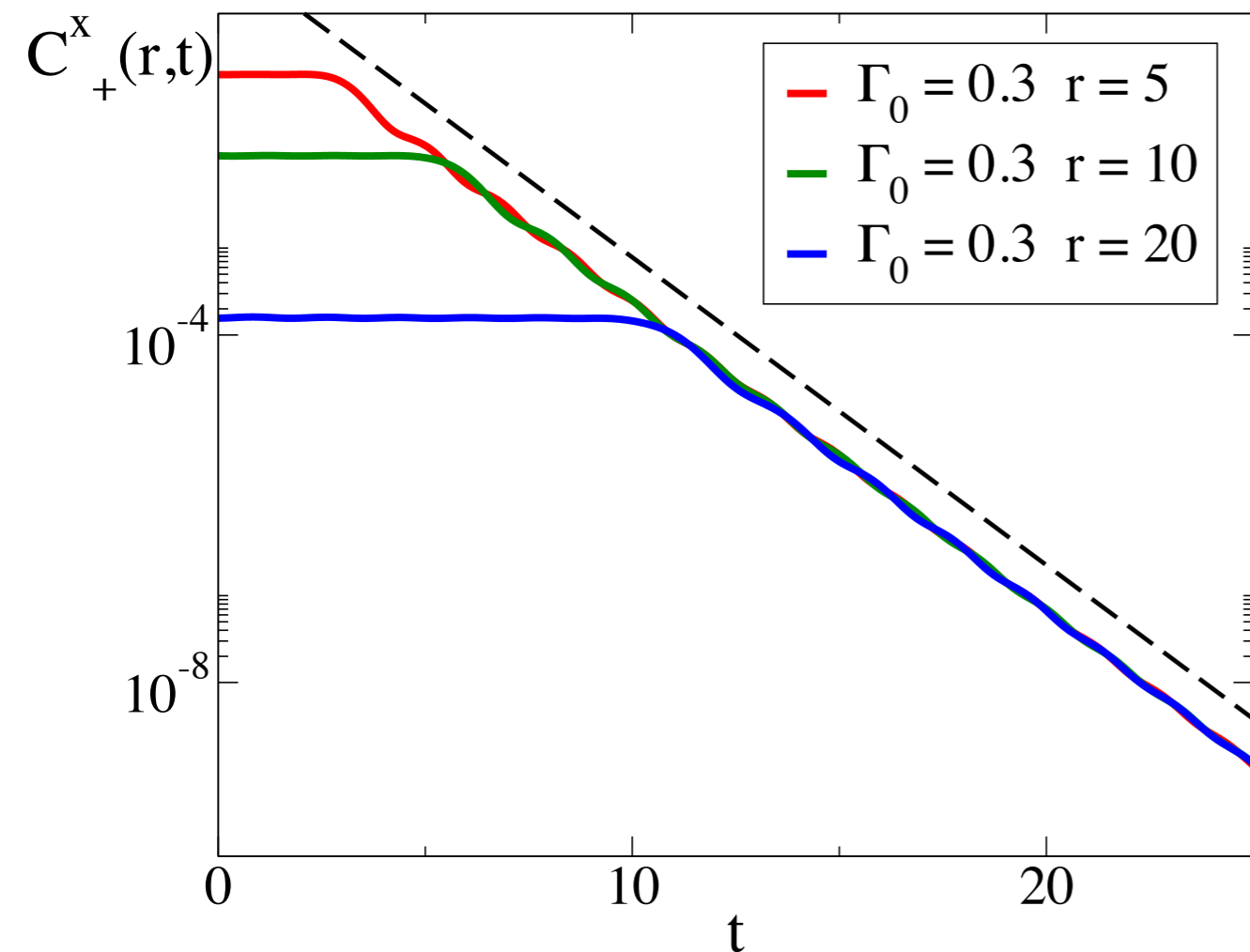
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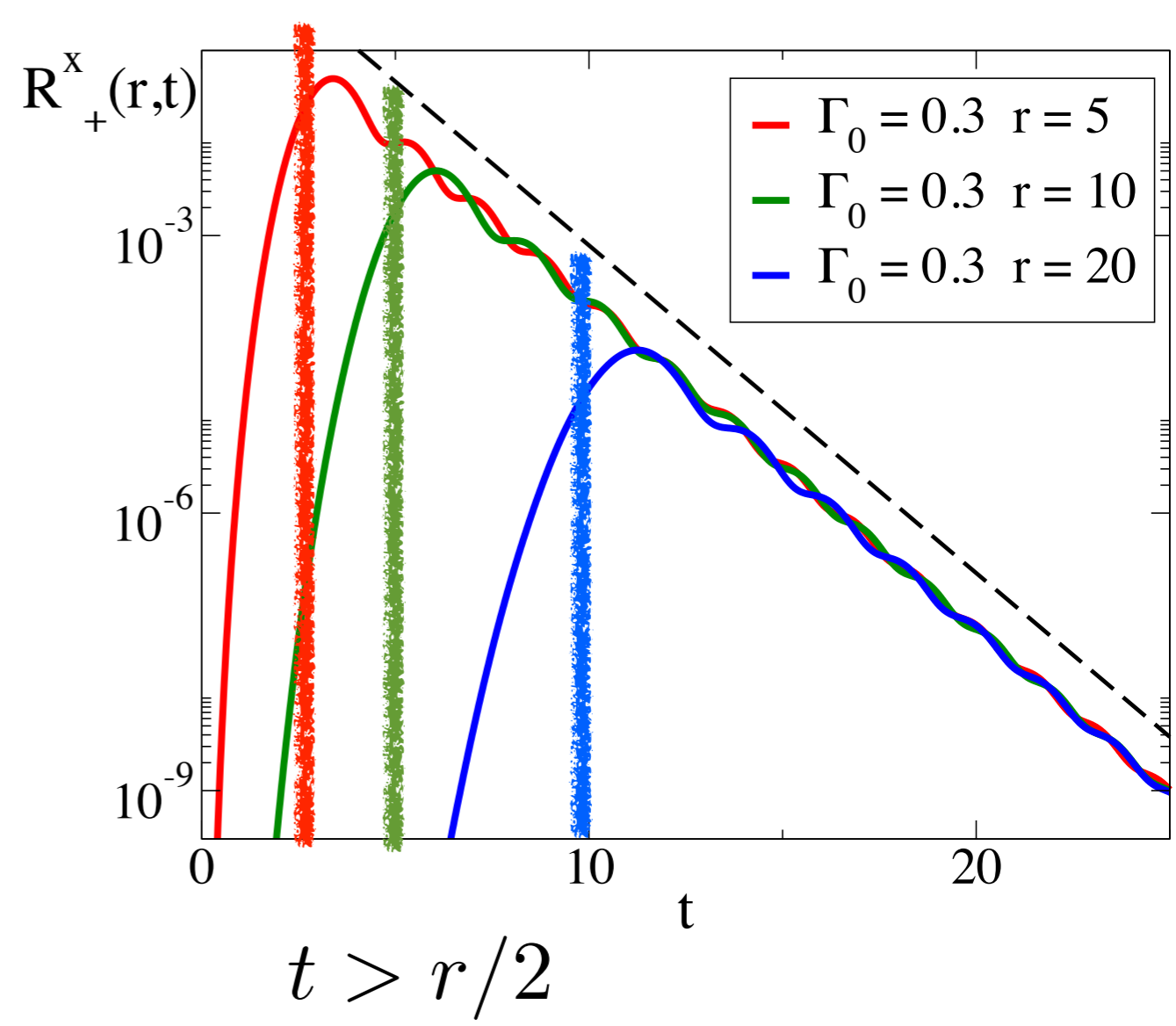
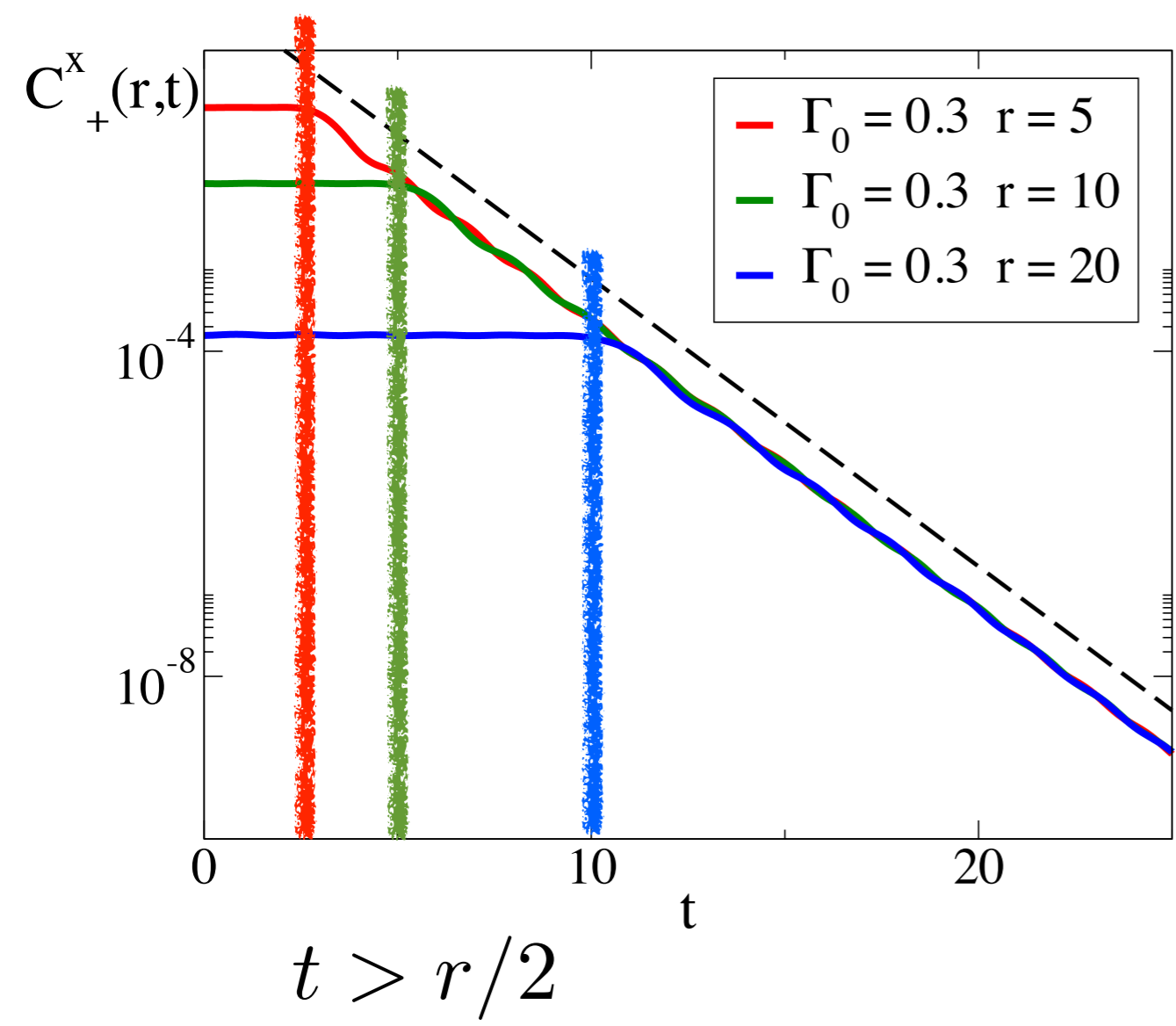


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Order parameter

FDT - effective temp.

$$\hbar \operatorname{Im} R(\omega) = \operatorname{th} \left(\frac{\beta \hbar \omega}{2} \right) \operatorname{Re} C(\omega)$$

$$R(t) = \frac{2i}{\hbar} \sum_{\text{odd } n=1}^{\infty} c_n \left(\frac{i\beta \hbar}{2} \right)^n \frac{d^n C(t)}{dt^n}$$

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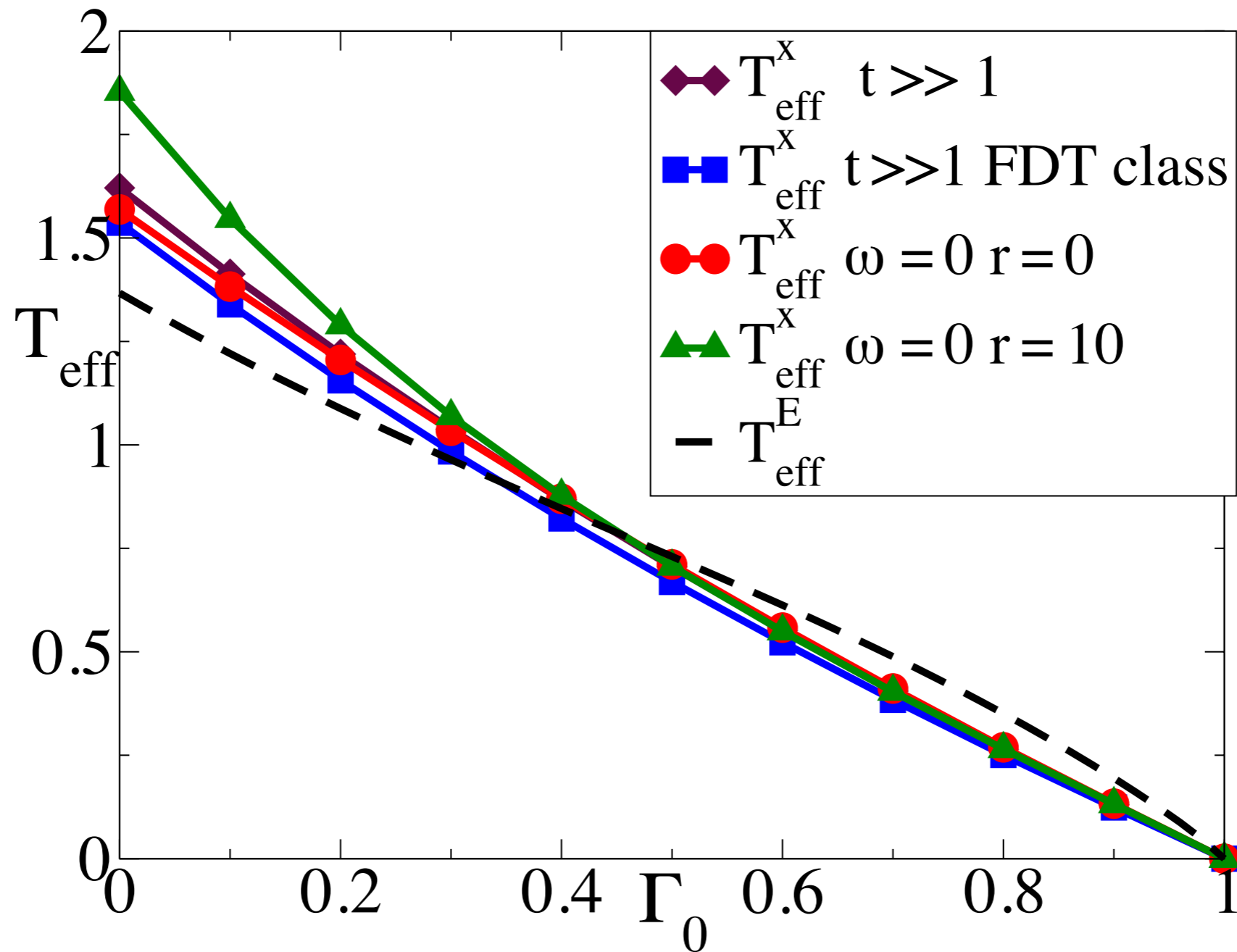
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$$t \rightarrow \infty \quad \Longrightarrow \quad \frac{\hbar A_R}{2A_C} = \tan \left(\frac{\hbar}{2\tau T_{\text{eff}}^x} \right) \quad \Longrightarrow \quad T_{\text{eff}}^x$$

Order parameter

FDT - effective temp.



Take-home
message:

Conclusions

- Effective temperatures via FDRs in *generic* **closed quantum** systems
 - “Dynamic” probe of thermal behavior beyond expectation values or time/space length scales
- Ising:
- Dynamics strongly depends on the observable
 - Correlations spread ballistically - light cones
 - Symmetries of stationary crit. dynamics
 - No thermal behavior - role of effective temp.s?

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- Effective temperatures via FDRs in *generic* **closed quantum** systems
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Thank you!