





FONDS NATIONAL SUISSE Schweizerischer Nationalfonds Fondo nazionale svizzero Swiss National Science Foundation

Fundamental questions in interacting systems

> spreading of correlations

coherent dynamics in a Mottinsulator



- preparation of complex states, e.g. BEC formation
- Iimit of 'speed of information'

destruction of coherence by dissipation: interplay between dissipation and interaction



- heating in optical lattices
- Interaction can fight dissipation?
- can dissipation be used to prepare complex quantum states?

Relativistic quantum mechanics



De Lahaye, Science (1998)

relativistic: correlation propagation maximally with speed of light non-relativistic: no such maximal velocity build in

Lieb and Robinson Bound

The Finite Group Velocity of Quantum Spin Systems

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finite group velocity for spreading of correlations outside of a 'light-cone' only exponentially small changes

not yet proven in generality for bosonic unbounded systems counter example has been constructed

Bravy , Hastings, Nachtergaele, Sims, Eisert ...

Quasi-particle interpretation



what is the velocity of quasi-particles?



Calabrese and Cardy (2006)

& other specific models

(Igloi, Rieger, Laeuchli, Cazalilla, Muramatsu, Manmana, Mathey, Schollwöck, Eisert,...)

How fast can correlations spread in an optical lattice?



How do correlations spread?

Theory: **Experiments**: decrease or increase of ratio U/J sudden decrease or increase of the optical lattice height Superfluid Phase Mott-Insulating phase quench U/J quasi-particle picture DMRG exact solution

M. Chenau (MPQ Munich), M. Endres, T. Fukuhara, C. Weitenberg, P. Schauß, C. Gross, I. Bloch (MPQ Munich, LMU Munich) S. Kuhr (Strathclyde, Glasgow)

P. Barmettler, D. Poletti, C. Kollath

M. Cheneau et al., Nature (2012) P. Barmettler et al, PRA (2012)

Strongly interacting bosons



only important for low final interation strength

P. Barmettler et al. in preparation
Batista and Ortiz PRL 2001 for spin1 -chains
bosonic approach see E. Altman and Auerbach and S. Huber et al.

Equilibrium local observables



short range quantities well described both by Fermionic quasi-particles and perturbation theory in J/U

Dynamic of Fermionic quasi-particles

large interaction limit:

symmetric coherent superposition (no change of density!)



first order perturbation in J/U (more complicated expression for quasiparticles):

$$\begin{aligned} |\psi(t)\rangle = |\bar{n}\rangle + i\frac{2\sqrt{2}J}{U}\sum_{k}\sin(k)c^{\dagger}_{k,+}c^{\dagger}_{-k,-}|\bar{n}\rangle \\ &-i\frac{2\sqrt{2}J}{U}\sum_{k}\sin(k)e^{i6J\cos(k)t/\hbar}c^{\dagger}_{k,+}c^{\dagger}_{-k,-}|\bar{n}\rangle \\ &=\sum_{j,d}\frac{(-i)^{d}\hbar d}{3Jt}\mathcal{J}_{d}(6Jt/\hbar)c^{\dagger}_{j,+}c^{\dagger}_{j+d,-}|\bar{n}\rangle \end{aligned}$$

Dispersion of quasi-particles



Real space imaging



parity distribution in real space

parity correlations

$$C_d(t) \propto \left\langle n_j^{even} n_{j+d}^{even} \right\rangle - \left\langle n_j^{even} \right\rangle \left\langle n_{j+d}^{even} \right\rangle$$

Parity distribution: Superfluid to Mott-insulator



parity distribution in real space flourescence imaging:





Bakr et al. (2009)

Real space imaging

equal time parity correlation between sites of distance d

$$C_{d}(t) \propto \left\langle n_{j}^{even} n_{j+d}^{even} \right\rangle - \left\langle n_{j}^{even} \right\rangle \left\langle n_{j+d}^{even} \right\rangle$$



Propagation of correlations in Mott-insulator



M. Cheneau et al, Nature (2012)

 $\mathcal{C}_{d}(t)/\mathcal{C}_{d}^{\mathsf{max}}$

0

Cone-like spreading?



linear spreading only exponentially small corrections before (Besselfunction)

Deviations from light cone propagation

$$C_{d\gg1} \approx -\left(\frac{2d^{2/3}2^{1/3}J\hbar}{3Ut}\right)^2 \operatorname{Ai}^2\left(-(2/d)^{1/3}(6Jt/\hbar - d)\right)$$



$$\begin{aligned} \mathbf{v}_{d} = & \frac{\hbar}{(t_{\text{peak}}(d+1) - t_{\text{peak}}(d))} \\ = & \mathbf{v}_{\infty} \left(1 - \frac{z_{0}}{2^{1/3}3} d^{-2/3} \right) + \mathcal{O}(d^{-5/3}) \end{aligned}$$

algebraic corrections to light cone (perturbation theory)

width $\sim d^1/3$

Velocity at long times



algebraic corrections to linear spreading corrections for small distances t~d+d^1/3

Velocity of correlation spreading



velocity:

generic maximal velocity in bounded model continuous across equilibrium phase transition (at phase transition difficult to determine)

sf->MI larger velocities found (origin higher bands?)





Fundamental questions in interacting systems

coherent spreading of correlations



- light-cone like dynamics
- well explained by quasi-particle picture
- generic max. velocity in system with bounded spectra

open questions:

- influence of higher bands
- faster quantum channels?
- approach for higher dimensional systems

interplay between dissipation and interaction



- heating in optical lattices
- Interaction can fight dissipation?
- can dissipation be used to prepare complex quantum states?

Theoretical model for dissipative coupling

dissipative effects:

scattering with thermal atoms flourescence scattering with light fields



described by Markovian master equation:

H system Hamiltonian

closed system dynamics

$$i\hbar\partial_{t}\hat{
ho} = [\hat{H},\hat{
ho}] + i\mathcal{L}\left(\hat{
ho}
ight)_{a}$$

dissipative dynamics

F. Gerbier and Y. Castin (2010) S. Pichler et al (2010)

D. Poletti, JS. Bernier, A. Georges, C. Kollath, arXiv (2012)

Double well: lowest band approximation

two-site Bose-Hubbard model (lowest band only) N atoms



start with ground state and apply dissipation at time t=0

how long does the initial coherence survive?

 $C = \langle \hat{b}_1^{\dagger} \hat{b}_2 + \hat{b}_2^{\dagger} \hat{b}_1 \rangle$

Localization effects



Interaction saves coherence

Γ/**J=1**, **N=60**



Evolution of coherence for large interaction



Mapping to classical diffusion equation (large U, N)

neglecting fast modes $(t > 1/\Gamma)$ (adiabatic elimination) and performing large N limit

$$\frac{d}{dx}\left(D(x)\frac{d}{dx}\phi(x,\tilde{t})\right) = \frac{\hbar U^2}{2J^2\Gamma}\frac{\partial}{\partial\tilde{t}}\phi(x,\tilde{t})$$
$$\downarrow$$
$$D(x) = \frac{x(1-x)}{(2x-1)^2}, \quad \tilde{t} = t/N^2$$

corresponds to diagonal density matrix elements

slowly diffusing states at the boundary correspond to large imbalance



Diffusion of initial state (non-Brownian)





Slow evolution towards totally mixed state



ground state close to balanced

build up of strongly imbalanced states energetically very costly

Diffusion of initial state (non-Brownian)



Scaling of coherence

Continuum limes: $\tilde{t} = t / N^2$

$$C(t) = f(t / N^2)$$

hand-waving: C(t)/N should not depend on N

 $C(t) / N \sim \frac{1}{\sqrt{t}}$

for large times $t >> \hbar / \lambda_0$ can feel discreteness

$$C(t)/N \propto e^{-\lambda_0 t/\hbar}$$



Interaction saves coherence

Γ/**J=1**, **N=60**



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open questions:

- faster quantum channels?
- implications for preparation of complex states in cold gases?
- approach for higher dimensional systems?

interplay: dissipation and interaction



- 3 regimes for decay of coherence
- slow algebraic decay due to interaction blocking
- slow diffusive states in configuration space

open questions:

- can interaction be used to prevent dissipation?
- implications for heating of complex states in cold gases?



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