The Galileo Galilei Institute for Theoretical Physics Arcetri, Florence, May 2012

Workshop: "New states of matter in and out of equilibrium"

Tensor network states that go beyond the boundary law for entanglement entropy

Guifre Vidal, Perimeter Institute



collaboration with Glen Evenbly, Caltech

Evenbly, Vidal, arxiv1205.0639 Evenbly, Vidal, arxiv120x.yyyy



MOTIVATION:

At low energies, "a many-body system may decouple into two (or several) sets of independent degrees of freedom"

Examples:

• 1D system: spin-charge separation



- 2D systems with 1D Fermi surface (or 1D Bose surface)
 - Fermi liquids
 - spin Bose metal



Haldane, Shankar, Swingle, Fisher,

•••

MOTIVATION:

Boundary law for entanglement entropy:

$$S_L \approx L^{D-1}$$

 $\overbrace{L}{}$

Systems with a Fermi/Bose surface are among the most entangled phases of quantum matter:

- Logarithmic violation: $S_L \approx L^{D-1} \log(L)$
- Beyond reach of tensor network states in D>1 dimensions:





 $S_L \approx L^{D-1}$

Outline



Glen Evenbly

• Introduction

Quantum circuits, simulatability and entanglement

MPS and TTN

• MERA

• branching MERA



branching MERA



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Quantum Circuit



Quantum Circuit

Can be used to *efficiently* encode many-body states:



time

Quantum Circuit as a many-body variational ansatz

Questions:

- 1) Cost of computing a local reduced density matrix
- 2) Entropy of a block of contiguous sites



time

























Example I:

 $w \approx 2aN$



inefficient

Example II:

$w \approx 2a \log(N)$



efficient

• Entanglement entropy of a block of contiguous sites

 $\left| 0 \right\rangle \left| 0 \right\rangle \left|$



• Entanglement entropy of a block of contiguous sites

 $\left| 0 \right\rangle \left| 0 \right\rangle \left|$



• Entanglement entropy of a block of contiguous sites



• Entanglement entropy of a block of contiguous sites





Example II:



Summary: Quantum Circuit as a many-body variational ansatz

Questions:

time

- Cost of computing a local reduced density matrix
- Entropy of a block of contiguous sites



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matrix product state MPS

tree tensor network

MPS: computational cost





scaling of entropy:

 $S(A) \approx const$

TTN: computational cost













 $c \approx N$ $S(A) \approx const$



 $c \approx \log(N)$ $S(A) \approx const$



cost

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MERA (multi-scale entanglement renormalization ansatz)



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MERA: entanglement entropy



 $n(A) \approx \log(L)$

scaling of entropy:

$$S(A) \approx \log(L)$$



cost

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MERA



branching MERA







branching MERA: computational cost past causal cone width: w' = 2w



MERA: entanglement entropy



 $n(A) \approx \log(L)$

scaling of entropy:

$$S(A) \approx \log(L)$$

ranching MERA: entanglement entropy



 $n(A) \approx 2\log(L)$

scaling of entropy:

$$S(A) \approx 2\log(L)$$

branching MERA



branching MERA



branching MERA: computational cost past causal cone width: w' = qw



cost of computing ho(A) :

 $c \approx q \exp(w)$

 $c \approx O(N)$

branching MERA: entanglement entropy





$$n(A) \approx O(L)$$

scaling of entropy:

$$S(A) \approx L$$

Conclusions



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 quantum circuits can be used to encode many-body states

let us add translation (+scale) invariance





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