

Universality in the equilibration of isolated systems after a small quench

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Preliminaries

Equilibration? $\rho(t) \Rightarrow \rho_{eq}$

1. Isolated, **finite** system

2. Prepare initial state ρ_0

3. Evolve with H : $\rho(t) = e^{-iHt} \rho_0 e^{iHt}$

4. Monitor observable A :

$$a(t) = \langle A(t) \rangle = \text{Tr} A \rho(t)$$

• Observation window $[0, T]$

• Time average

$$\bar{f} = \int_0^T f(t) \frac{dt}{T}$$

$$P_A(a) da = \text{Prob}(\langle A(t) \rangle \in [a, a + da], t \in [0, T])$$

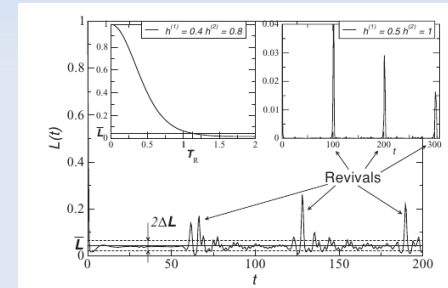
$$P_A(a) = \overline{\delta(a - a(t))}$$

Equilibration = concentration of $P_A(a)$

1. No strong convergence $\|\rho(t) - \rho_{eq}\| = \text{const}$

2. For finite systems no weak convergence

3. Stochastic convergence



Equilibration

$$\overline{a(t)} = \overline{\text{Tr} A \rho(t)} = \text{Tr} A \bar{\rho} \quad \Rightarrow \quad \rho_{eq} = \bar{\rho}$$

$$\Delta^2 a = \overline{a(t)^2} - \bar{a}^2 \leq \text{Ran}_A^2 \text{Tr} \bar{\rho}^2$$

+ "gap"
non-degeneracy

$$E_i - E_j = E_l - E_m \\ \Rightarrow i=j, l=m \vee i=l, j=m$$

Chebyshev's inequality

$$\text{Prob}(|a(t) - \bar{a}| \geq \epsilon) \leq \frac{\Delta^2 a}{\epsilon^2}$$

Reimann, PRL (2008)

$$\text{Tr} \bar{\rho}^2 = \overline{L(t)}$$

$$L(t) = |\langle \Psi | e^{-itH} | \Psi \rangle|^2 = \exp 2 \sum_{n=1} \langle H^n \rangle_c \frac{(-t^2)^n}{2n!} = e^{V_f(t)}$$

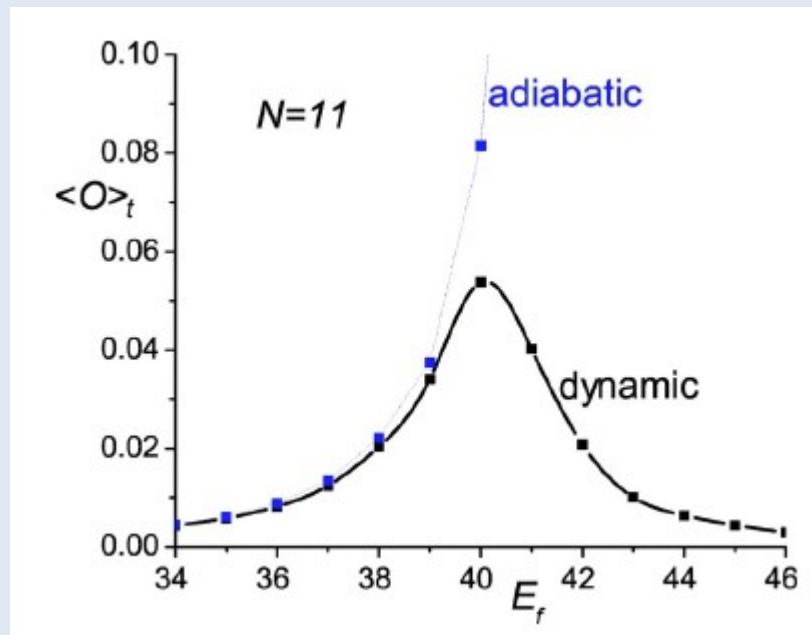
$$\Rightarrow \text{Tr} \bar{\rho}^2 \propto e^{-\alpha V}$$

Free systems?

Quench: dynamical detection of QPT's

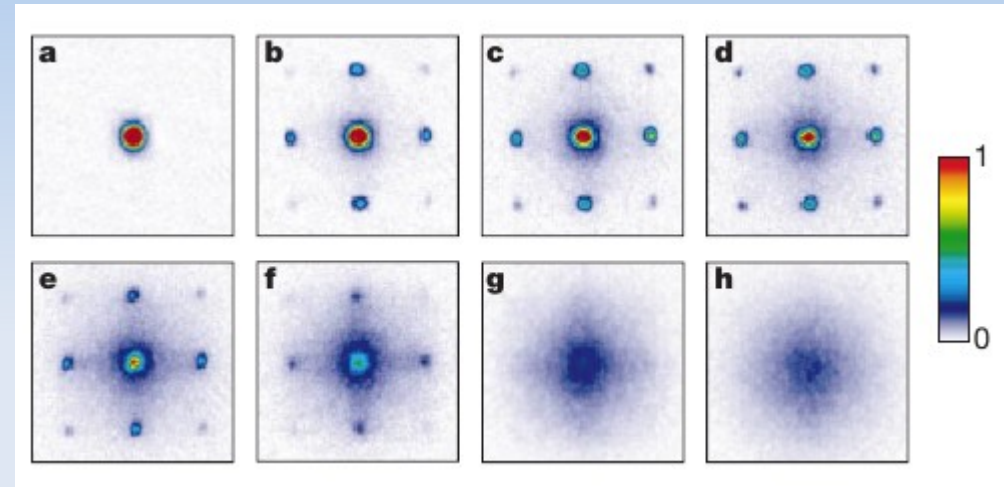
$$H \rightarrow |\Psi_0\rangle$$
$$H' = H + \delta\lambda V$$

Theoretical description



Sengupta, Powell, Sachdev, PRA (2004)

superfluid-Mott transition: experiment



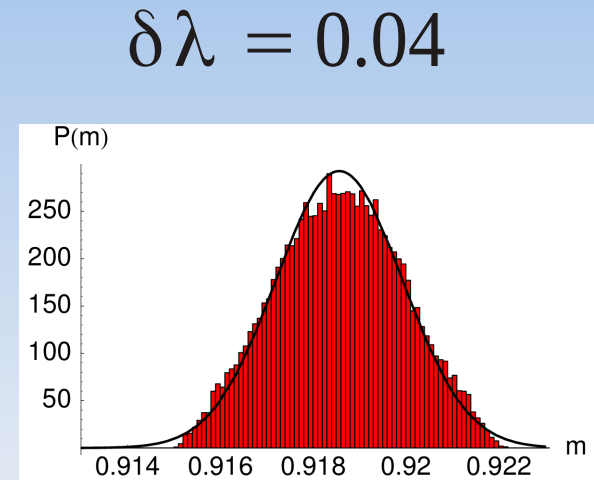
Greiner et al., Nature (2002)

Small* quench: full statistics

- Small quench, off-critical

$$L \gg \xi$$

$$L = 12$$

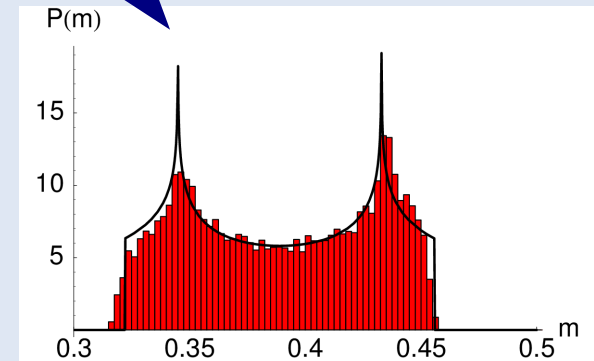


Bi-modal:
phase transition

- Small quench, quasi-critical

$$\xi \gg L$$

$$L = 16$$



(*) Small: $\delta\lambda^2 \chi_F \ll 1 \Rightarrow \delta\lambda \ll \frac{1}{L^{1/\nu}}$

Small quench: CLT

$$\langle A(t) \rangle = \sum_{n,m} A_{n,m} \overline{c(E_n)} c(E_m) e^{-it(E_m - E_n)} \quad c(E_n) = \langle E_n, \Psi_0 \rangle$$

$$\approx \bar{A} + \sum_{n>0} F_n \cos(t(E_n - E_0)) \quad F_n \approx 2A_{n,0} c(E_n)$$

rational
independence



$a(t)$ sum of
independent variables

$$\Delta^2 a = \frac{1}{2} \sum_{n>0} F_n^2 \leq 2[\langle 0|A^2|0 \rangle - \langle 0|A|0 \rangle^2] \propto L$$

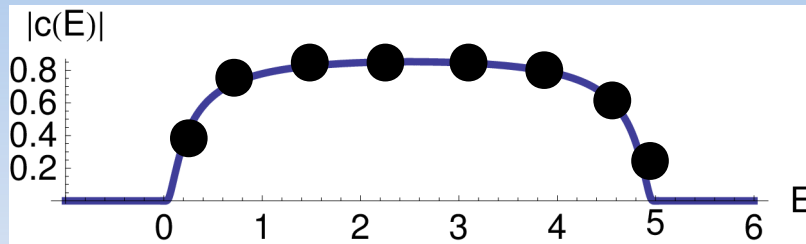


"Generally"
CLT

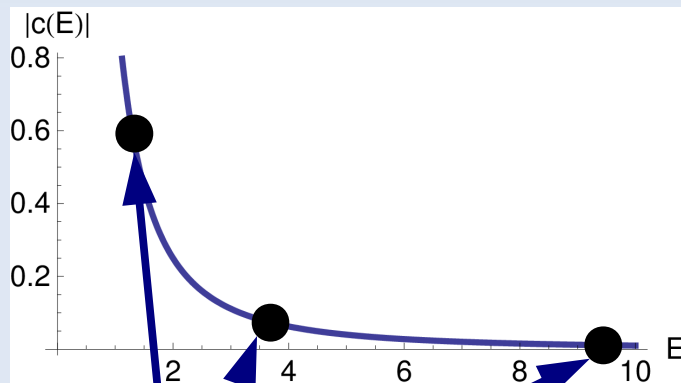
Look at $F_n, c(E_n)$

Small quench: explanation

Off-critical distribution



Critical case distribution



Allowed values

Scaling prediction
at criticality:

$$c(E) \sim E^{-1/\nu}$$

sum rule

$$\sum_E |c(E)|^2 = 1$$

Only few $|E\rangle$ states
contribute:
Bi-modal distribution

Small quench: critical case

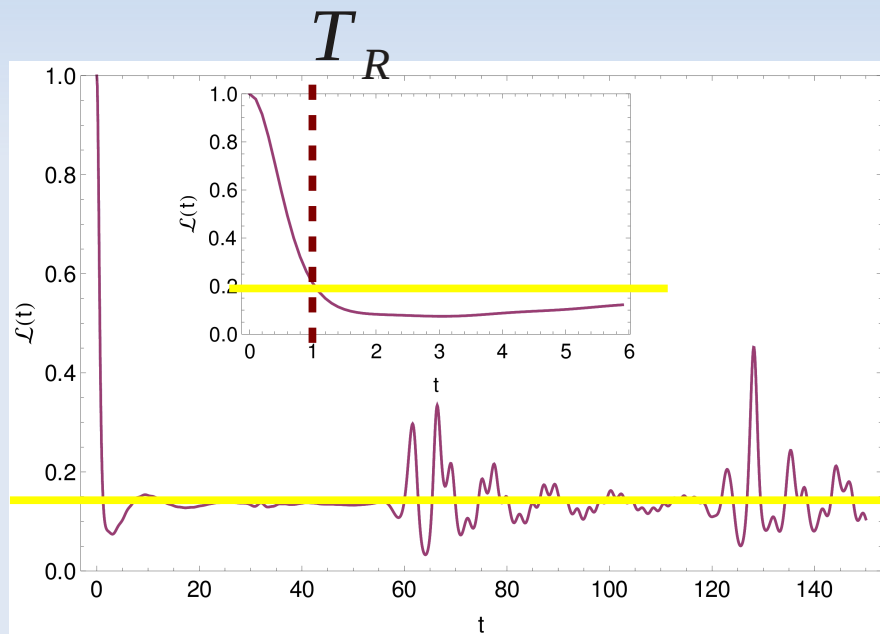
- Q: How to break CLT?
- A: most $F_n \rightarrow 0$

(Quantum) critical points are ***more*** stable against perturbations

Relaxation time



(Talk by Michael Pustilnik)



$$\langle A(T_R) \rangle := \bar{A}$$

Loschmidt echo

$$L(T_R) := \bar{L}$$

Relaxation time II

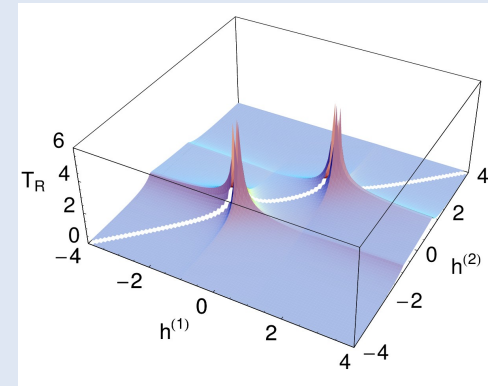
Loschmidt echo: short time - cumulant expansion

**Generally
(+small quench off-critical)**

$$L(t) = e^{-\sigma^2 t^2} \quad \bar{L} = e^{-\alpha L^d}, \quad \sigma^2 \sim L^d \Rightarrow T_R = O(L^0)$$

small quench criticality

$$\bar{L} \simeq e^{-2\delta\lambda^2\chi_F}, \quad \sigma^2 \sim L^{2(d-\Delta)}, \quad \chi_F \sim L^{2(d+\zeta-\Delta)} \Rightarrow T_R = O(L^\zeta)$$



$$T_R \sim |\lambda - \lambda_c|^{-\zeta \nu}$$

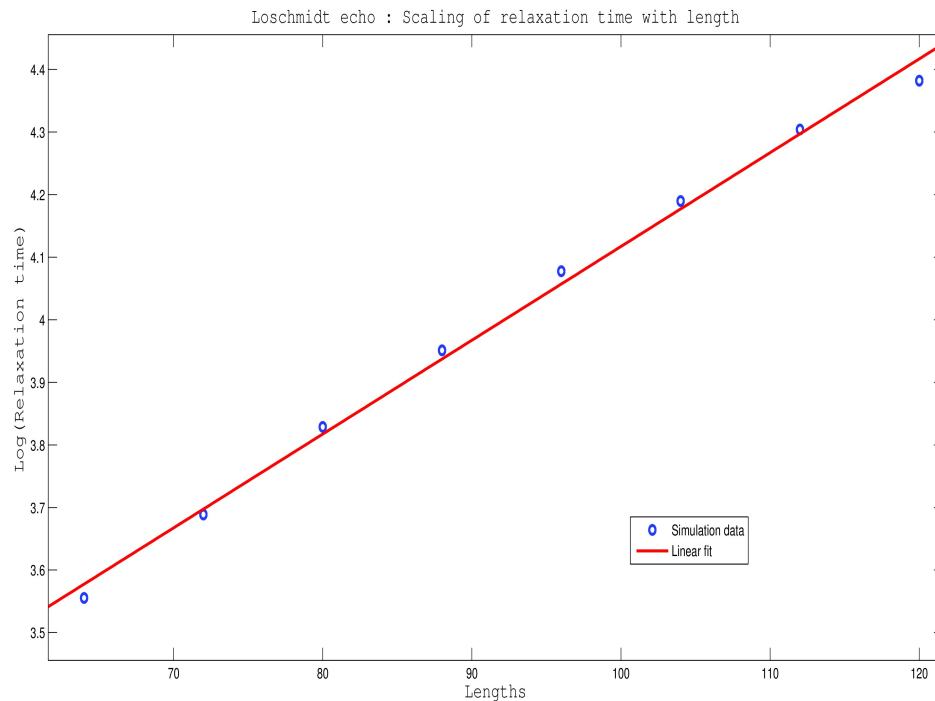
Relaxation time: Random Systems

~ Inguscio, Modugno, LENS

$$H = \sum_x (c_x^\dagger c_{x+1} + c_{x+1}^\dagger c_x) - \mu_x c_x^\dagger c_x$$

Loschmidt echo

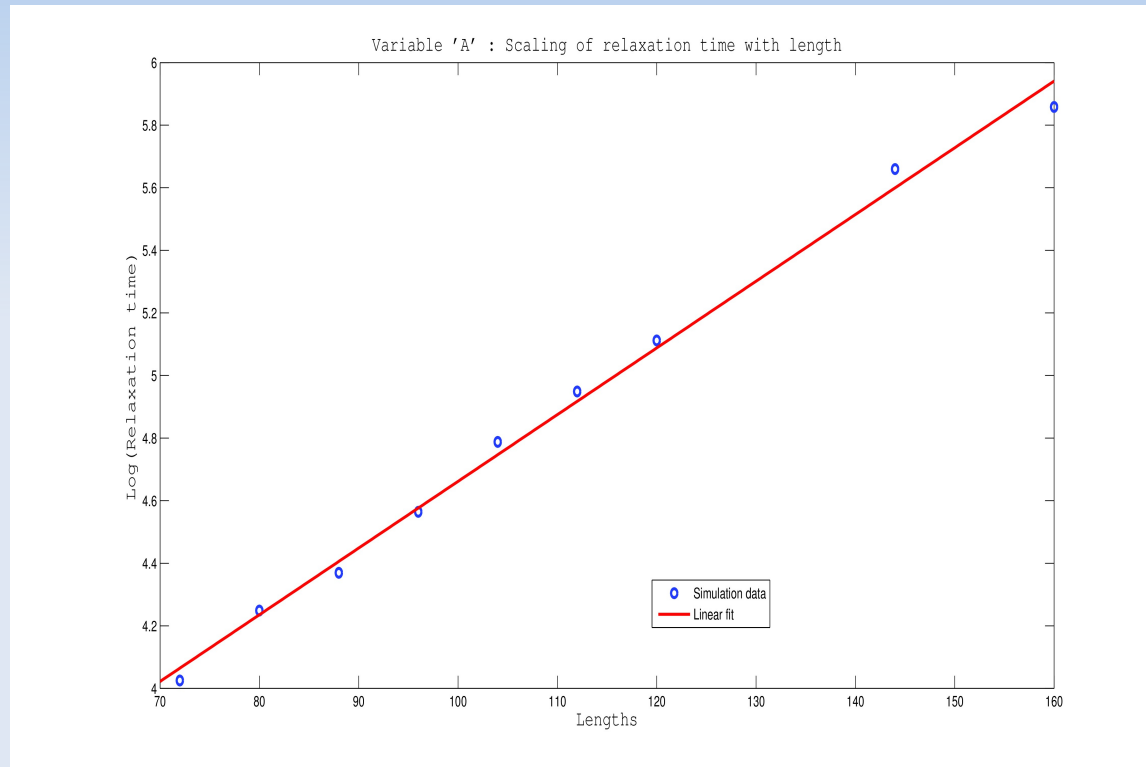
Random field



$$E[L(t)]$$

Relaxation time: Random Systems

Number operator $E[\langle N_l(t) \rangle]$



$$T_{\text{Relax}} \sim e^{\alpha L}$$

Equilibration & Integrability

$$a(t) = \langle A(t) \rangle \rightarrow P_A(a) = \overline{\delta(a - a(t))}$$

Generally, for A extensive:

$$\overline{a(t)} \propto V \quad \Delta^2 a \leq O(e^{-\alpha V})$$

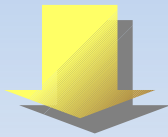
Integrable systems (free Fermions)

$$\begin{aligned} a(t) &= \text{Tr}(A e^{-itM} R e^{itM}) \\ &= \sum_{k,q} \underbrace{A_{q,k} R_{k,q}}_{F_{k,q}/2} e^{-it(\epsilon_k - \epsilon_q)} \end{aligned}$$

$$\begin{aligned} H &= \sum c_x^\dagger M_{x,y} c_y \\ A &= \sum c_x^\dagger A_{x,y} c_y \\ R_{y,x} &= \langle c_x^\dagger c_y \rangle \end{aligned}$$

Stat-mech parallel

Rational independence



$$\overline{e^{\lambda a(t)}} = \sum_{\theta's} e^{\lambda E(\theta's)} = e^{f(\lambda)V}$$

$$E(\theta's) = \sum_{k,q} F_{k,q} \cos(\theta_k - \theta_q)$$

$$F_{k,q} = F(|k - q|)$$



All cumulants
extensive:
CLT

$$Z = \frac{(a(t) - \bar{a})}{\sqrt{V}} \quad \text{Gaussian}$$

Classical XY model
on lattice $F_{i,j}$
(infinite temperature)

Equilibration & Integrability

Generally, for A extensive:

$$\overline{a(t)} \propto V \quad \Delta^2 a \leq O(e^{-\alpha V})$$

Integrable systems (free Fermions)

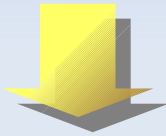
$$\overline{a(t)} \propto V \quad \Delta^2 a = O(V)$$

Gaussian(poor) equilibration

Loschmidt echo

$$[R, M] = 0$$

$$L(t) = \prod_k (1 - \alpha_k \sin^2(t \epsilon_k / 2))$$



$$Z = \frac{\log L(t) - \overline{\log L(t)}}{\sqrt{L}} \quad \text{Gaussian, } \Rightarrow L(t) \text{ Log-Normal}$$

For general models (RI spectrum), work in progress

$$\mu_n(L(t)) = f(\text{Tr}(\bar{\rho}^{2k}))$$

- Applies to XY model
- Generalizes to thermal quenches
- Generalizes to Ulman Fidelity

Curiosity: Riemman zeta

$$\zeta(\sigma + it) = \text{Tr}(e^{-itH} \rho_\sigma), \quad \rho_\sigma = e^{-\sigma H}$$

H primon gas:
free bosons

$$Z := \log |\zeta(\sigma_0 + it)|^2$$

Satisfies
CLT

Very similar
to Loschmidt Echo)

Conclusions

- Finite systems
 - Look at full time statistics
 - Small quench: a tool to detect criticality, engineer "*new quantum states of matter*"
 - Relaxation time can be defined
 - Integrability & equilibration:
integrable systems concentrate less
- } Ingredients

Thank you