

Low-energy local density of states of the 1D Hubbard model

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Low energy properties of fermionic systems in 1D

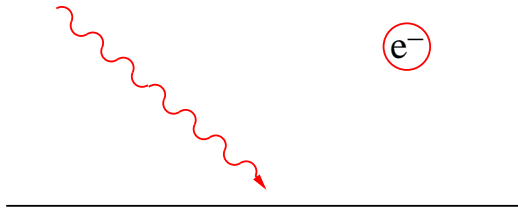
- Strong correlations, interactions dominant, universal behavior
- no single-particle picture possible, excitations collective bosonic modes

Luttinger liquid

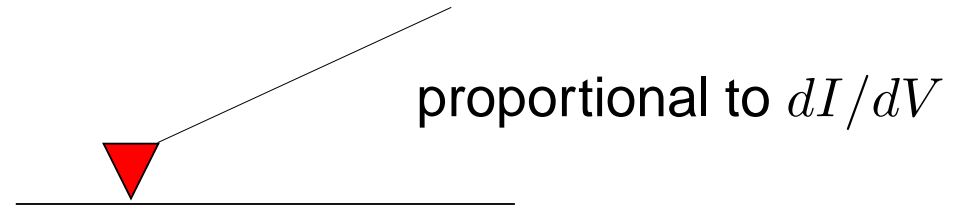
- Spin- and charge excitations decouple

Tunneling in quantum wires

Photo emission



Scanning tunneling spectroscopy



Density of states:

$$\begin{aligned}\rho(\omega) &= -\frac{1}{\pi} \text{Im} \int_0^\infty e^{i\omega t} G^r(t) dt \\ &= \sum_m |\langle \omega_m | \psi^\dagger | 0 \rangle|^2 \delta(\omega - \omega_m)\end{aligned}$$

Luttinger liquid:

$$\rho(\omega) \propto \omega^\alpha$$

Momentum resolved tunneling experiments in 1D

Experimental signatures of spin-charge separation

- Semiconductor hetero-structures

Auslaender, Steinberg, Yacoby, Tserkovnyak, Halperin, Baldwin, Pfeiffer, and West, *Science* 308, 88 (2005)

Jompol, Ford, Griffiths, Farrer, Jones, Anderson, Ritchie, Silk, and Schofield, *Science* 325 (2009)

- Quasi one-dimensional crystals

Kim, Koh, Rotenberg, Oh, Eisaki, Motoyama, Uchida, Tohyama, Maekawa, Shen, and Kim, *Nature Phys.* 2 (2006)

- Self-organized atomic chains

(Segovia, Purdie, Hengsberger, and Baer, *Nature* 402 (1999))

Scanning tunneling spectroscopy in 1D

● Carbon nanotubes

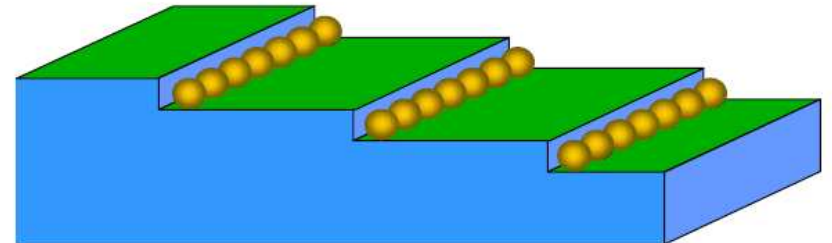
Lee, Eggert, Kim, Kahng, Shinorara, and Kuk, *Phys. Rev. Lett.* 93 (2004)

Venema, Wildöer, Janssen, Tans, Tuinstra, Kouwenhoven, and Dekker, *Science* 283 (1999)

Lemay, Janssen, van den Hout, Mooij, Bronikowski, Willis, Smalley, Kouwenhoven, and Dekker, *Nature* 412 (2001)

● Self-organized atomic gold chains

Blumenstein, Schäfer, Mietke, Meyer, Dollinger, Lochner, Cui, Patthey, Matzdorf, and Claessen, *Nature Phys.* 7 (2011)



Signatures of power law density of states \rightarrow Luttinger liquid behavior

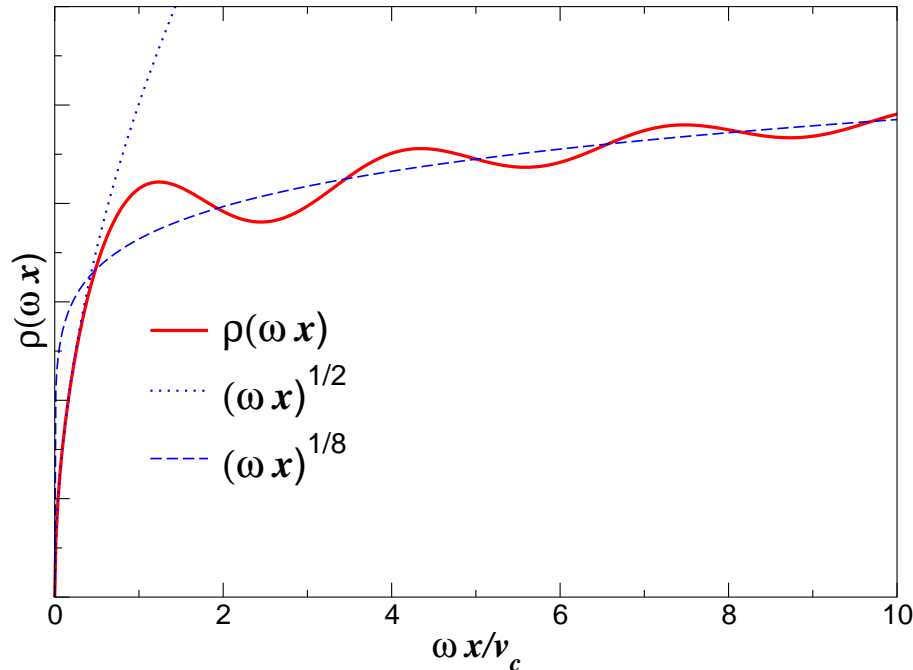
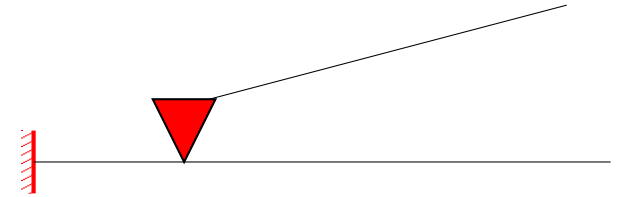
Outline

- Local density of states of interacting fermions in 1D
 - Luttinger liquid: power laws
here: Effects of boundaries and finite system sizes
 - DMRG: lattice model of spinless fermions
Spectral weight of individual excitations
 - Bosonization: Recursion formula
 - DMRG: Hubbard model

Luttinger liquid with impurity

Local density of states

$$\begin{aligned} \rho(\omega, x) &= \sum_m |\langle \omega_m | \psi^\dagger(x) | 0 \rangle|^2 \delta(\omega - \omega_m) \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{i\omega t} \langle \psi(x, t) \psi^\dagger(x, 0) \rangle dt \end{aligned}$$

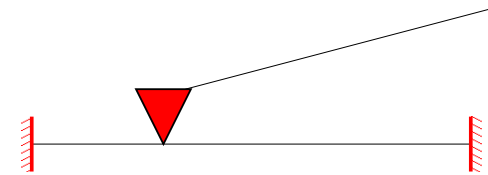


strong depletion for small energies and at the boundary

(here $K_s = 1, K_c = \frac{1}{2}$)

Finite Luttinger liquid with boundaries

$$\begin{aligned} \rho(\omega, x) &= \sum_m |\langle \omega_m | \psi^\dagger(x) | 0 \rangle|^2 \delta(\omega - \omega_m) \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{i\omega t} \langle \psi(x, t) \psi^\dagger(x, 0) \rangle dt \end{aligned}$$

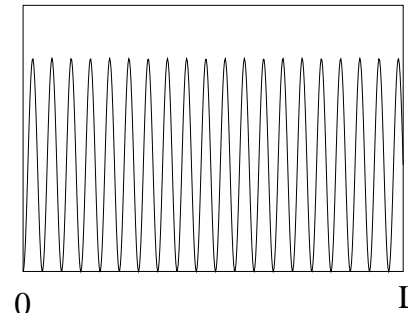


Free fermions:

single particle wave function:

$$\rho(\omega_m, x) = |\Psi_m(x)|^2$$

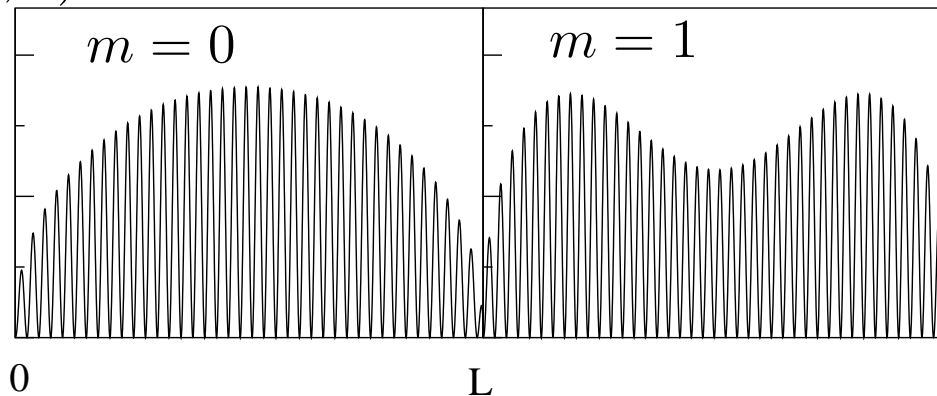
$\rho(\omega_0, x)$



Interacting:

Bosonization

$\rho(\omega_m, x)$



Anfuso, Eggert, Phys. Rev. B

68 (2003)

Exact lattice model: local density of states in DMRG

$$H = -t \sum_{x=1}^{L-1} \left(\psi_x^\dagger \psi_{x+1} + \psi_{x+1}^\dagger \psi_x \right) + U \sum_{x=1}^{L-1} n_x n_{x+1}, \quad n_x = \psi_x^\dagger \psi_x - 1/2$$

- Approach 1: Dynamical DMRG and tDMRG
 - + entire spectrum
 - energy levels not resolvable

Jeckelmann, arXiv:1111.6545

- Approach 2: transition matrix elements in DMRG
 - + energy levels resolvable
 - only for low energy excitations

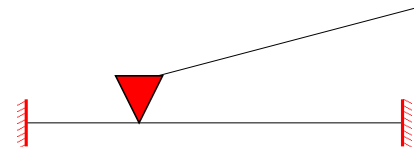
Schneider, Struck, Bortz, and Eggert, Phys. Rev. Lett. 101, 206401 (2008)

Söffing, Schneider, and Eggert, arXiv:1204.0003

$$\rho(\omega, x) = \sum_m |\langle \omega_m | \psi_x^\dagger | 0 \rangle|^2 \delta(\omega - \omega_m) = -\frac{1}{\pi} \text{Im} \int_0^\infty e^{i\omega t} G^r(x, t) dt$$

Understanding of individual excitations

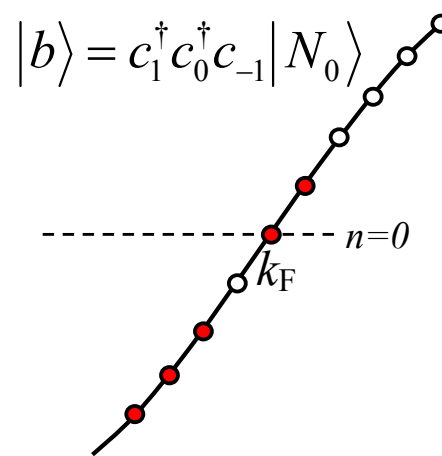
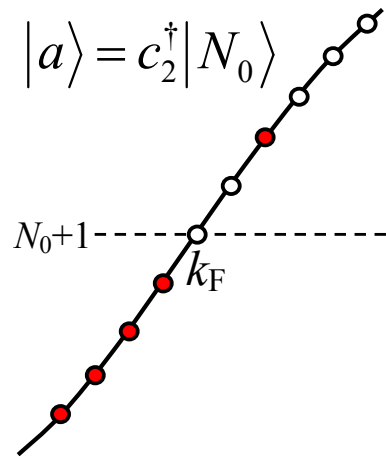
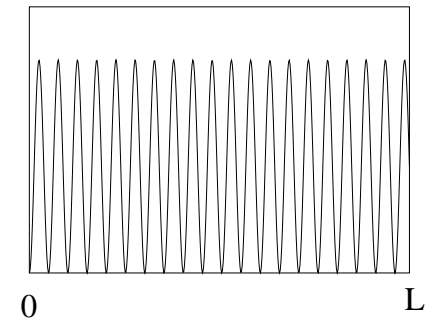
Free fermions



$$H = \sum_k \epsilon(k) c_k^\dagger c_k \quad k = \frac{\pi}{L}n \quad n = 0, 1, 2, \dots$$

$$|\langle N_0 | c_n \psi_x^\dagger | N_0 \rangle|^2 = \frac{2}{L} |\sin(k_F + k_n)x|^2$$

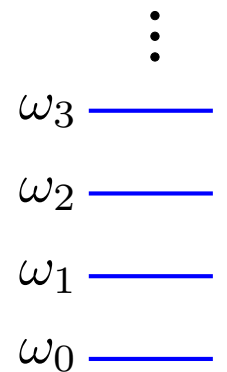
$$|\langle a | \psi_x^\dagger | N_0 \rangle|^2$$



lattice model

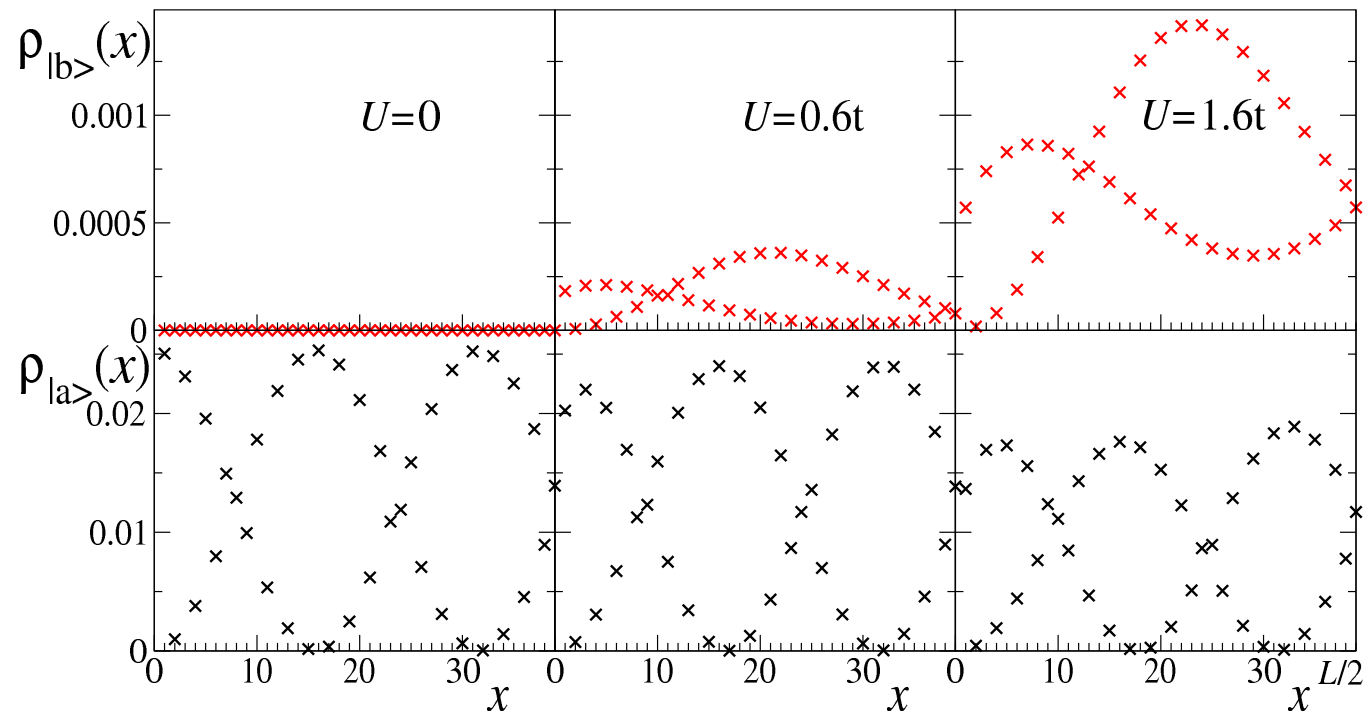
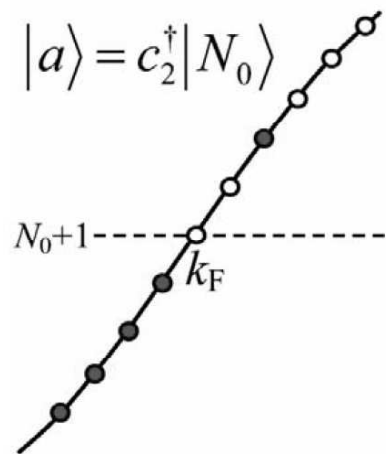
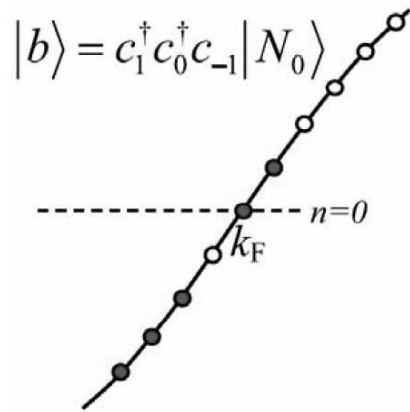
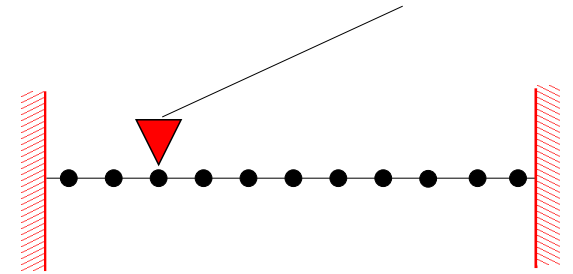


effective theory

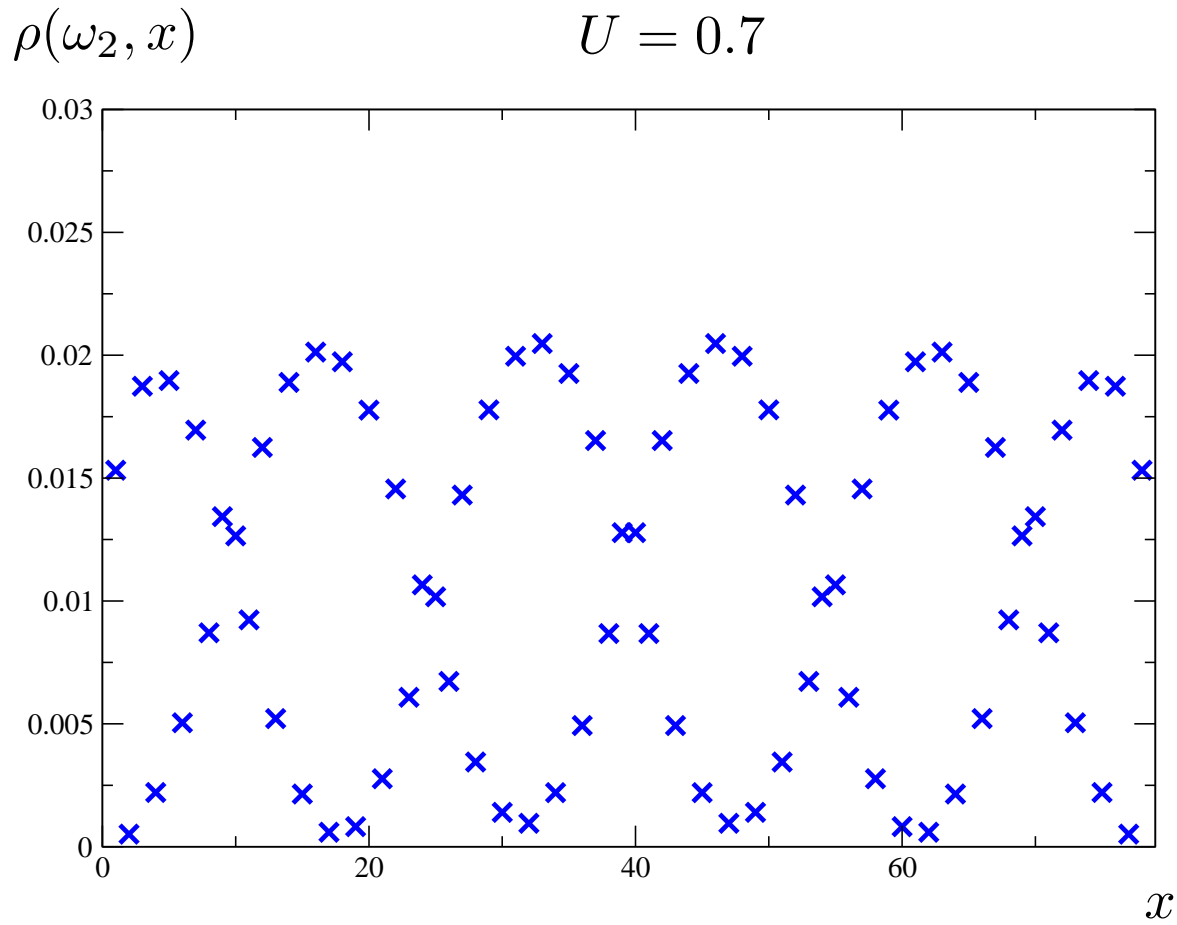


Local density of states: DMRG results

$$H = -t \sum_{x=1}^{L-1} \left(\psi_x^\dagger \psi_{x+1} + \psi_{x+1}^\dagger \psi_x \right) + U \sum_{x=1}^{L-1} n_x n_{x+1}$$



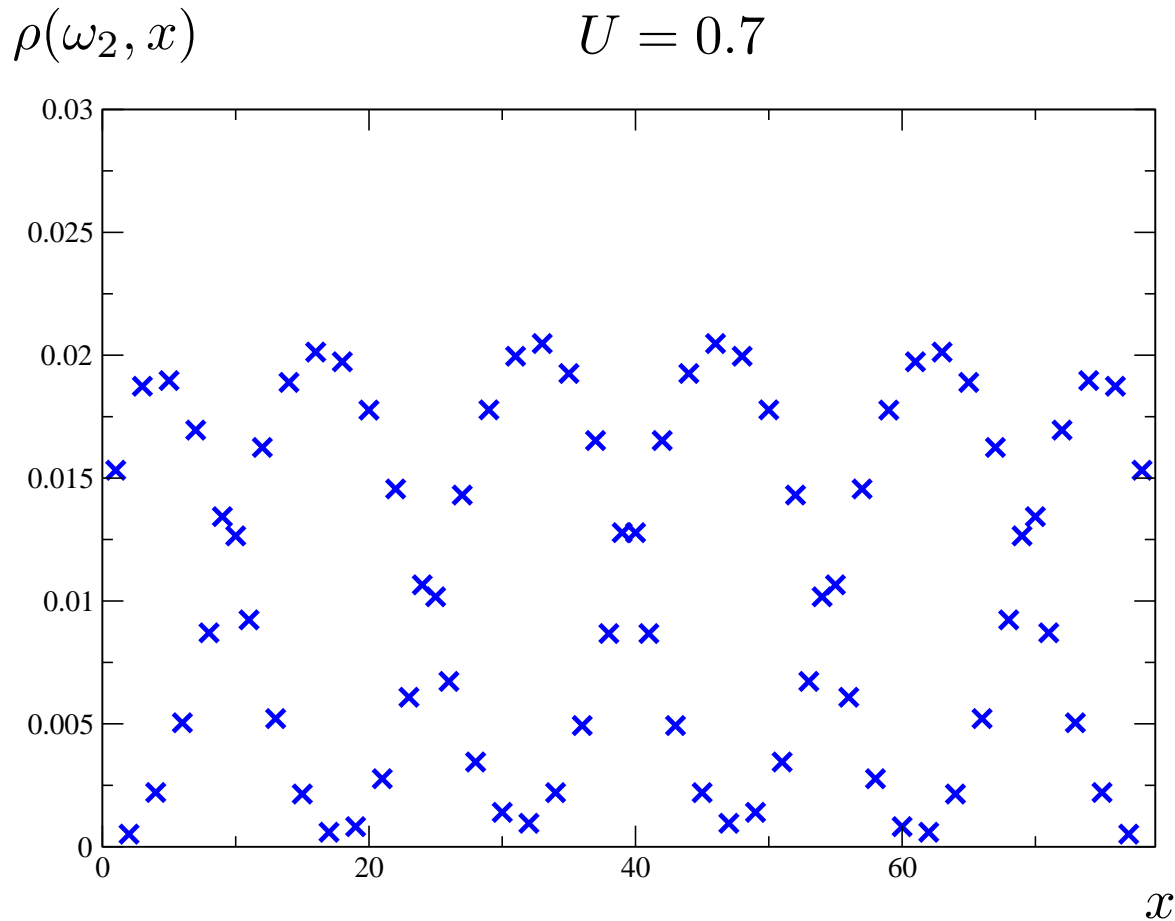
Local density of states: DMRG results



× DMRG

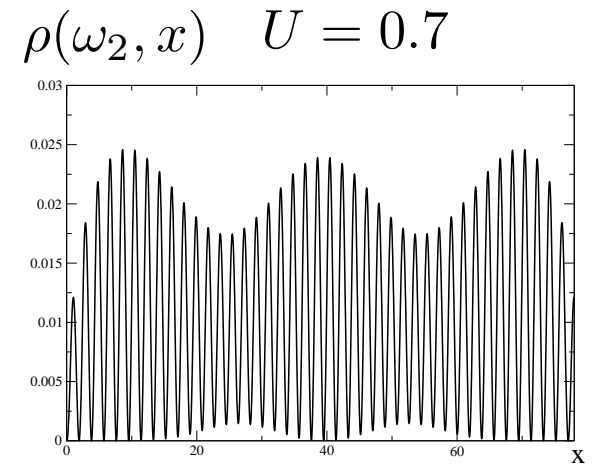
$$\rho(\omega_2, x) = |\langle a | \psi^\dagger(x) | 0 \rangle|^2 + |\langle b | \psi^\dagger(x) | 0 \rangle|^2$$

Local density of states: DMRG results

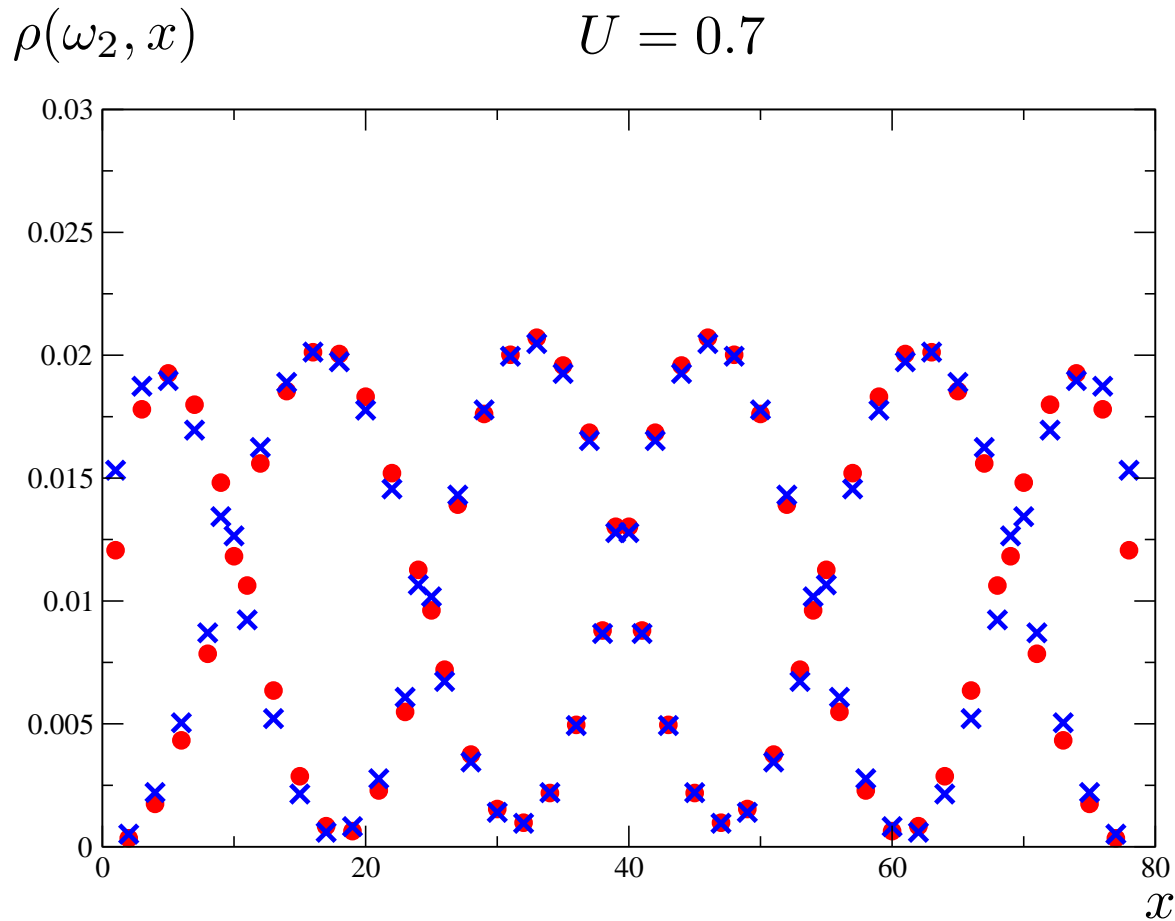


\times DMRG

$$\rho(\omega_2, x) = |\langle a | \psi^\dagger(x) | 0 \rangle|^2 + |\langle b | \psi^\dagger(x) | 0 \rangle|^2$$

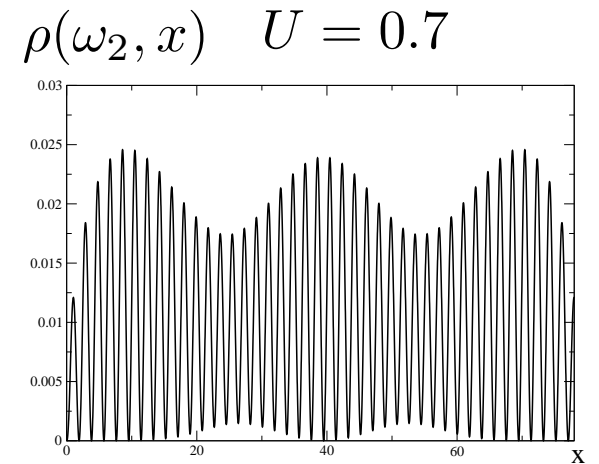


Local density of states: DMRG results

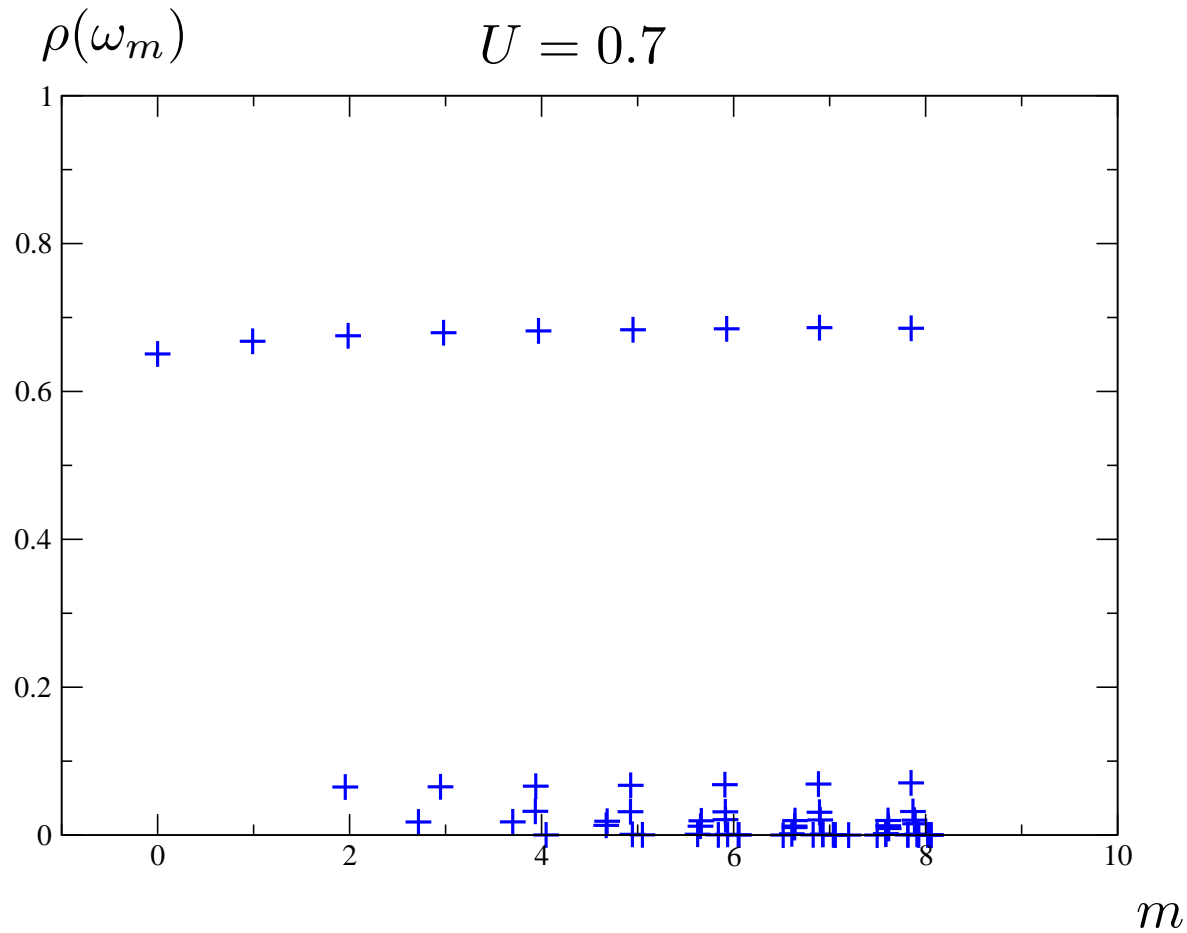


× DMRG ● Bosonization

$$\rho(\omega_2, x) = |\langle a | \psi^\dagger(x) | 0 \rangle|^2 + |\langle b | \psi^\dagger(x) | 0 \rangle|^2$$



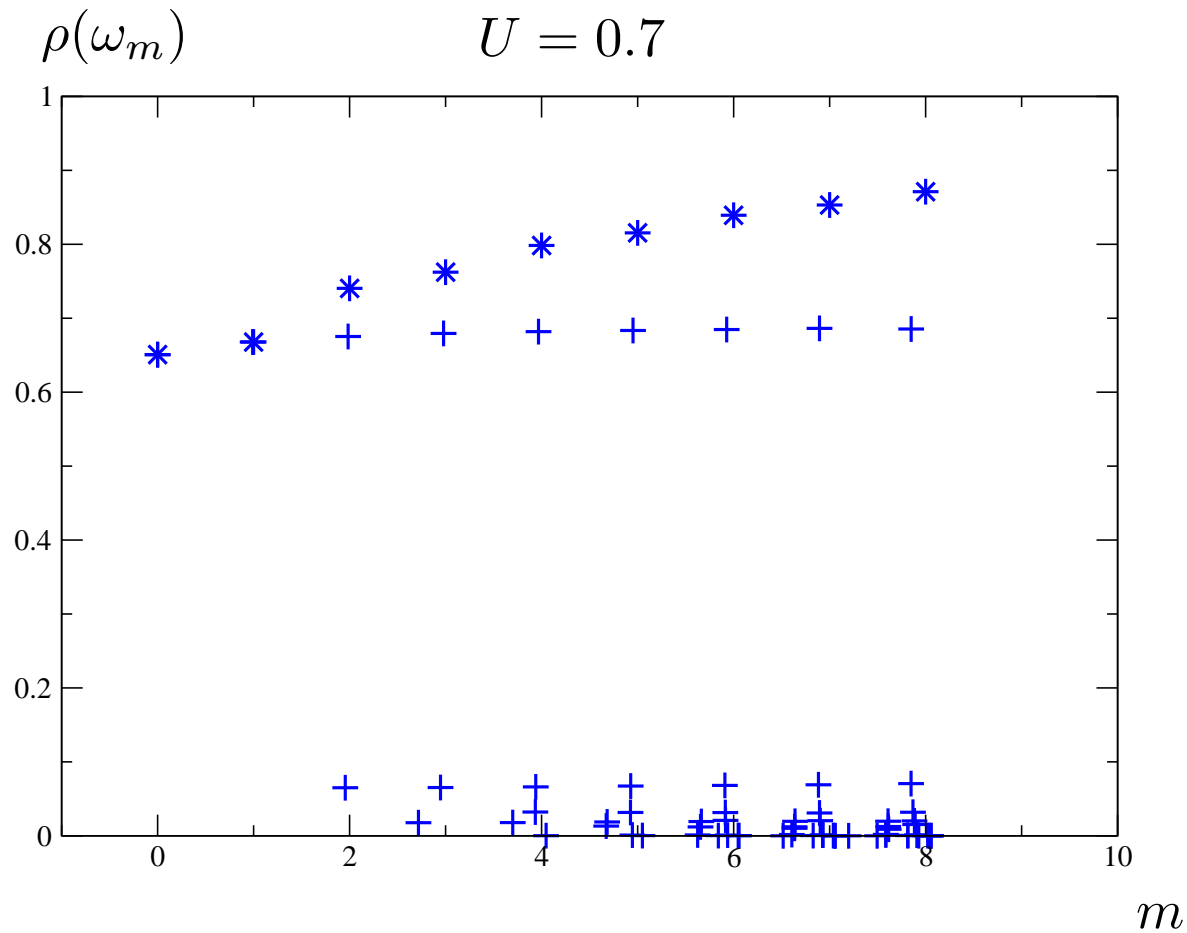
Density of states: position integrated



$$\rho(\omega_m) = \sum_x \rho(\omega_m, x)$$

+ DMRG

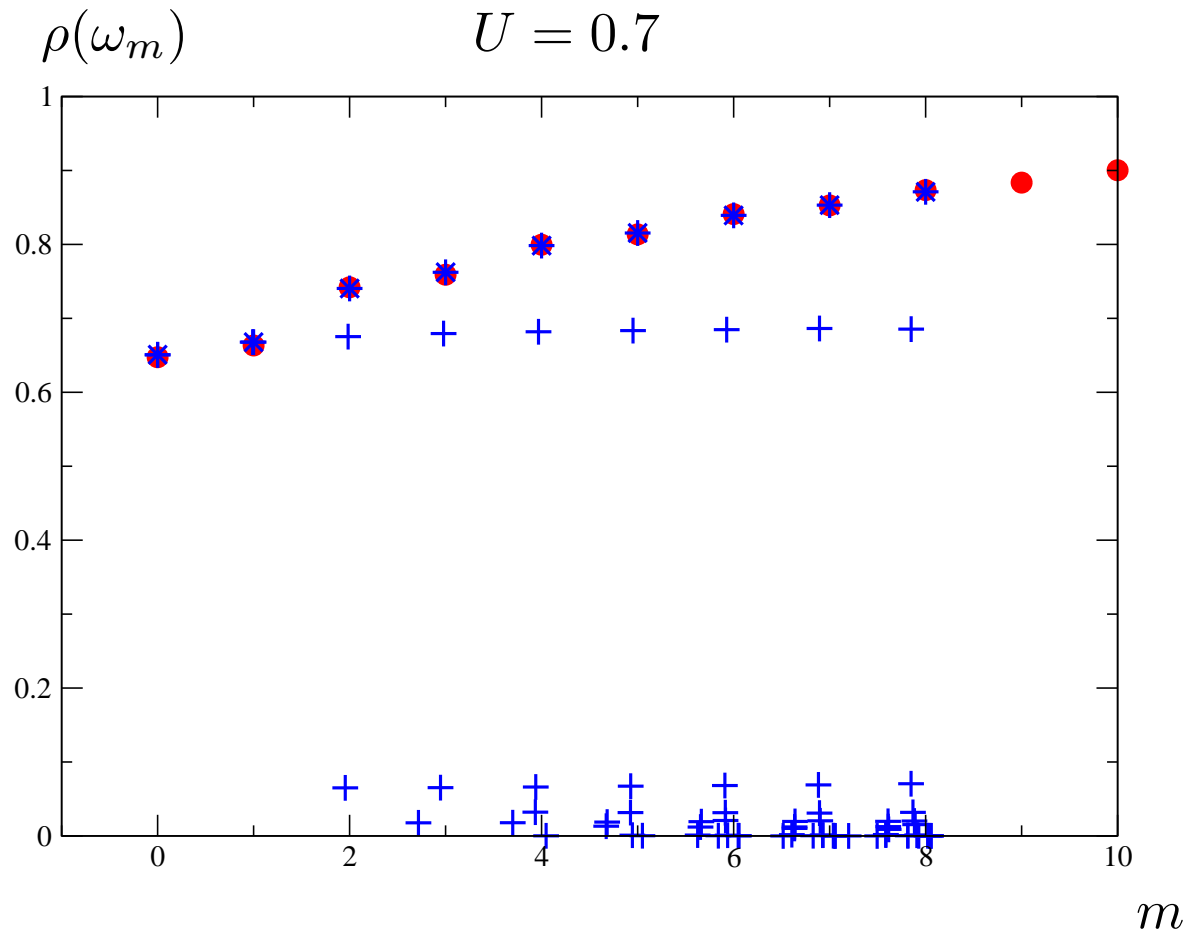
Density of states: position integrated



$$\rho(\omega_m) = \sum_x \rho(\omega_m, x)$$

- + DMRG
- * DMRG summiert

Density of states: position integrated



$$\rho(\omega_m) = \sum_x \rho(\omega_m, x)$$

- + DMRG
- * DMRG summed
- Bosonization

Recursive method for the density of states

- $\rho(\omega, x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{i\omega t} \langle \psi(x, t) \psi^\dagger(x, 0) \rangle dt$

- Correlation functions in standard bosonization

$$\langle \psi_R(x, t) \psi_R^\dagger(x, 0) \rangle = |c|^2 \exp \left(\sum_{\ell=1}^{\infty} \frac{1}{\ell} e^{-i\ell\Delta\omega t} \gamma_\ell(x) \right)$$

$$\psi_R^\dagger(x, t) := c(x) \exp \left[i \sum_{\ell=1}^{\infty} \frac{1}{\sqrt{\ell}} e^{i\ell\Delta\omega t} A_\ell^\dagger(x) \right] \exp \left[i \sum_{\ell=1}^{\infty} \frac{1}{\sqrt{\ell}} e^{-i\ell\Delta\omega t} A_\ell(x) \right]$$

$$\gamma_\ell(x) = [A_\ell(x), A_\ell^\dagger(x)], \quad A_\ell(x) = \alpha(K) e^{ik_\ell x} b_\ell^R - \beta(K) e^{-ik_\ell x} b_\ell^L$$

- Finite systems

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} dt e^{i\omega t} \langle \psi_R(x, t) \psi_R^\dagger(x, 0) \rangle = \sum_m \rho_m \delta(\omega - m\Delta\omega)$$

$$\rho_m = \frac{1}{m} (\rho_{m-1} \gamma_1 + \rho_{m-2} \gamma_2 + \dots + \rho_1 \gamma_{m-1} + \rho_0 \gamma_m) \quad \text{mit} \quad \rho_0 = |c|^2$$

Schneider and Eggert, Phys. Rev. Lett. 104 (2010)

earlier recursive approach: Schönhammer and Meden, Phys. Rev. B 47 (1993)

Spinless fermions with periodic b. c.

Density of states: $\rho_m = \frac{1}{m}(\rho_{m-1}\gamma_1 + \rho_{m-2}\gamma_2 + \cdots + \rho_1\gamma_{m-1} + \rho_0\gamma_m)$

- Commutator mode independent

$$\gamma = \frac{1}{2} \left(\frac{1}{K} + K \right) \quad \text{Luttinger-parameter } K$$

- Recursion formula exacty solvable

$$\rho_m = |c|^2 \frac{\Gamma(m + \gamma)}{\Gamma(\gamma)\Gamma(m + 1)} \approx |c|^2 \frac{1}{\Gamma(\gamma)} m^{\gamma-1} \quad \text{well known power law}$$

- in general $\gamma_\ell(x)$ mode and x dependent

Spinful fermions with open b. c.

Luttinger liquid picture:

States described by integer spin and charge quantum numbers $\{m_s, m_c\}$

Energies: $\omega_{m_s, m_c} = (m_s v_s + m_c v_c) \frac{\pi}{L+1}$ with $v_s \leq v_c$

Density of states:

$$\rho_{m_s, m_c}(x) = |c_x|^2 \left[\rho_{s, m_s}^{uni}(x) \rho_{c, m_c}^{uni}(x) - \cos(2k_F x) \rho_{s, m_s}^{osc}(x) \rho_{c, m_c}^{osc}(x) \right]$$

Calculate recursively, e.g. :

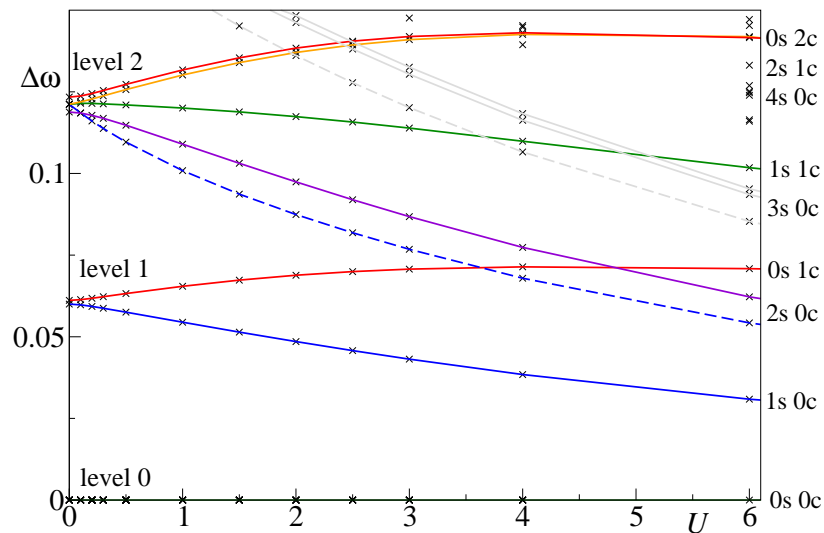
$$\rho_{c, m_c}^{uni}(x) = \frac{1}{m_c} \sum_{\ell=1}^{m_c} \rho_{c, m_c - \ell}^{uni}(x) \gamma_{c, \ell}^{uni}(x)$$

$$\gamma_{c, \ell}^{uni}(x) = (1/K_c + K_c)/4 + (1/K_c - K_c) \cos(2k_\ell x)$$

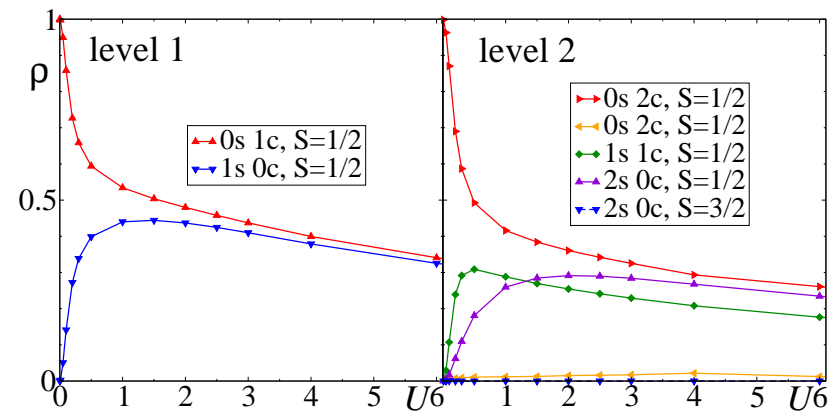
Comparison to DMRG results (1/3)

Hubbard model: $H = -t \sum_{\sigma, x=1}^{L-1} (\psi_{\sigma, x}^{\dagger} \psi_{\sigma, x+1} + \text{h.c.}) + U \sum_{x=1}^L n_{\uparrow, x} n_{\downarrow, x}$

Energies $\Delta\omega$



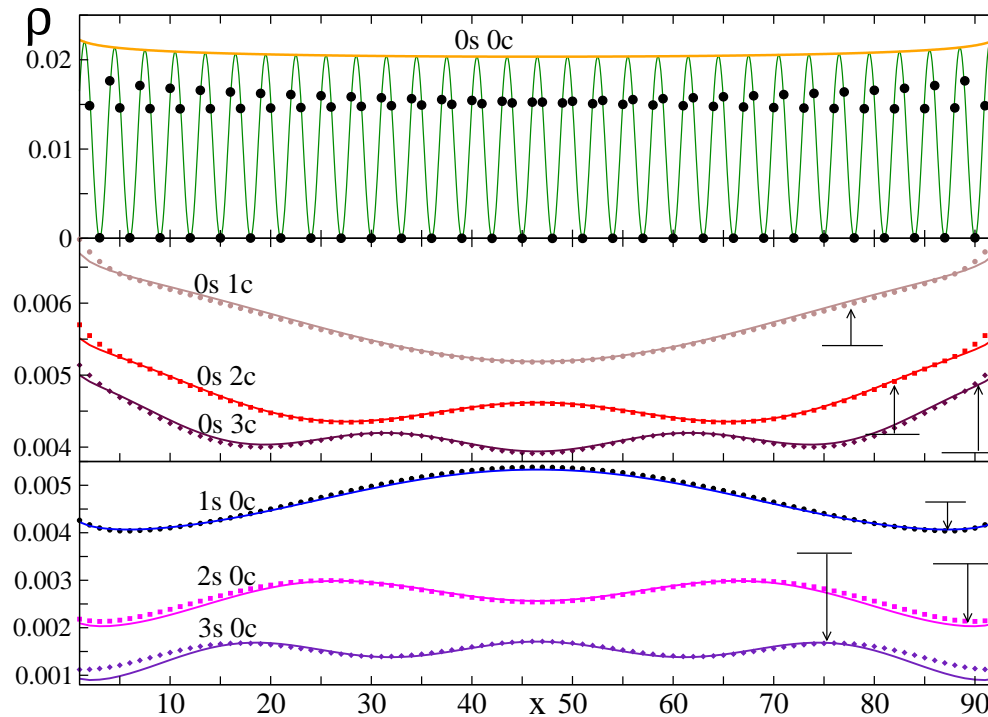
Total density of states



Parameter: $N_{\uparrow} = N_{\downarrow} + 1 = 31$ and $L = 90$

Comparison to DMRG results (2/3)

Local density of states:



$$N_{\uparrow} = N_{\downarrow} + 1 = 31$$

$$L = 92$$

$$U = 1$$

Lines: predictions for $K_c = 0.9081$ and $K_s = 1.16$ adjusted by shifts

Local density of states *increases* near boundary

Comparison to DMRG results (3/3)

Local density of states does *not* fit predictions by theory:

- Theory curves must be shifted down for charge and up for spin modes (competition of energy scales: band curvature vs interaction)
- Luttinger parameter K_s must be chosen considerably larger than unity → attractive behavior in the spin

Boundary exponent $\alpha_B = (1/K_s + 1/K_c)/2 - 1$ may become negative

Similar observations: Schuricht, Andergassen, and Meden preprint arXiv:1111.7174,

Andergassen, Enss, Meden, Metzner, Schollwöck, and Schönhammer, Phys. Rev. B 73 (2006),

Meden, Metzner, Schollwöck, Schneider, Stauber, and Schönhammer, Eur. Phys. J. B 16 (2000),

Schönhammer, Meden, Metzner, Schollwöck, and Gunnarsson, Phys. Rev. B 61, (2000)

- Multiplicative corrections to $G^r(x, t)$ due to marginal irrelevant operator?

Summary

- Local density of states for individual energy levels by DMRG
- Recursive formula: simple calculation of the density of states
- Numerical results in agreement with bosonization for spinless fermions
- Large deviations for the Hubbard model
→ effective negative boundary exponent