



# $S_3$ and the angle $\theta_{13}$

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# Outline

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1. *Different strategies to go Beyond SM*
2.  *$S_3$  as a family symmetry*
3. *A particular realization of  $S_3$*
4. *Quark mixing*
5. *Lepton mixing: the angle  $\theta_{13}$*
6. *Outlook*

# Beyond SM

- *Fermion masses and CP violation*
- *Neutrino physics*
- *Hierarchy problem (EW and Planck scale)*
- $a_\mu$ : *magnetic anomalous moment of muon*
- *Dark matter*

# Beyond SM

- *Fermion masses and CP violation*
- *Neutrino physics*
- *Hierarchy problem (EW and Planck scale)*
- $a_\mu$ : *magnetic anomalous moment of muon*
- *Dark matter*

*Early LHC: It may be that flavour is decoupled from the hierarchy problem . . .*

## *Bottom up approach*

1. *Propose an effective model for fermion mixing at the EW scale*
2. *Later, address the rest of the problems*

# $S_3$ as a family symmetry

- *Rescued from old ideas:*

*Pakvasa 77, Wyler, Derman, Frere 78*

*A. and M. Mondragon, et. al 1998, Kubo, Branco, Paes, etc*

- *Two main classes*

- $S_3$  extension (that is, must be broken)

$$S_3 \supset S_{3L} \times S_{3R} \supset S_3^{\text{diag.}} \supset S_2$$

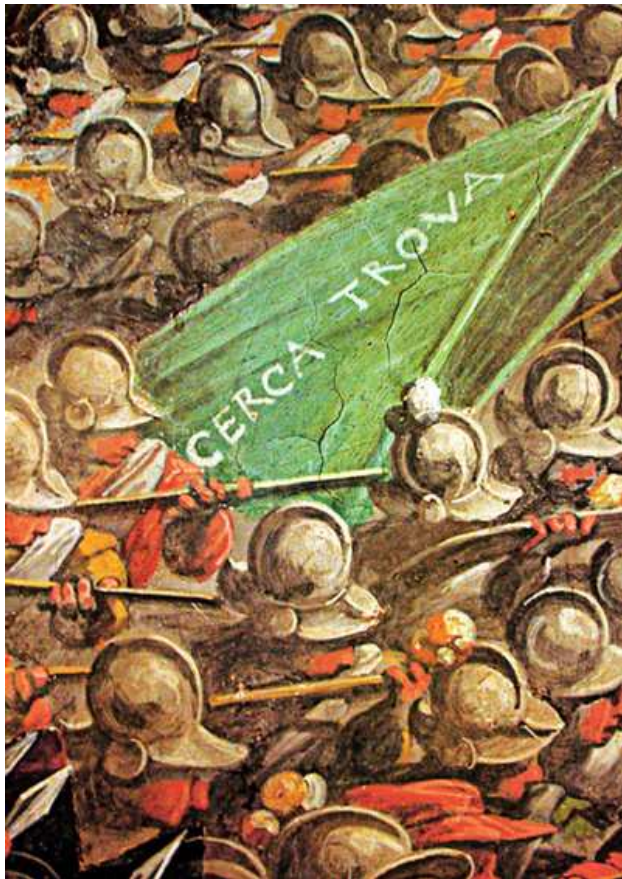
- $S_3$ -SM:  $S_3$  realized as an exact symmetry, possible due to the addition of *Higgs doublets*.

- *Fertile patch: keep exploring* (Review in 1205.4755)

*Time for advertisement: it is worth to keep exploring: "Look-for" then "find" (cerca-trova)*

# Cerca trova

*Maurizio Seracini: lost painting beneath one of the walls of a room at Palazzo Vecchio*



*Detail of Giorgio Vasari's wall painting (Battle of Marciano) at Palazzo Vecchio*

# Cerca trovi: Davinci's Battaglia di Anghiari



*After 30 years of research, Seracini may have found the, for long time considered lost, painting of Davinci's La battaglia di Anghiari (left: a copy of Rubens, right: an sketch of Davinci for the painting)*

# A particular realization of $S_3$ : $S_3$ -SM

To appear soon (A. Mondragón, F. González-Canales, U. J. Saldaña and L. V-S)

1.  $S_3$  is considered to be a global symmetry of the theory, even after the Higgs mechanism → the invariance under permutations of the 3 families can be seen as a hidden symmetry that makes its appearance only through the observed mass spectrum
2. The scalar sector: four Higgs  $SU(2)_L$  doublets belonging to the  $S_3$  reducible representation,  $4 = \mathbf{2} \oplus \mathbf{1}_S \oplus \mathbf{1}_A$ 
  - Four effective  $SU(2)_L$  doublets
  - Three effective  $SU(2)_L$  doublets
3. What is so far this construction achieved?
  - Through the solution of 9 equations with 9 variables (in each sector), one can get exact formulas of mixing angles in terms of the physical mass ratios of fermions
  - Excelent fits of quark and lepton mixing
4. Next steps: reduce the number of parameters to be predictive and explore EWSB



- *The assignment of fermion families in  $S_3$  is suggested by the observed mass hierarchy*

$$\begin{pmatrix} f_{I(L,R)} \\ f_{II(L,R)} \end{pmatrix} \sim \mathbf{2}; \quad f_{III(L,R)} \equiv g_{L,R} \sim \mathbf{1}_S,$$

$$f_{IIIL} = (b_L, t_L) \quad f_{IIIR} = t_R,$$

$$\begin{pmatrix} f_{IL} \\ f_{IIL} \end{pmatrix} = \begin{pmatrix} (u_L, d_L) \\ (c_L, s_L) \end{pmatrix}, \quad \begin{pmatrix} f_{IR} \\ f_{IIR} \end{pmatrix}_{f=u} = \begin{pmatrix} u_R \\ c_R \end{pmatrix}, \quad \begin{pmatrix} f_{IR} \\ f_{IIR} \end{pmatrix}_{f=d} = \begin{pmatrix} d_R \\ s_R \end{pmatrix},$$

■ *Higgs sector*

$$\begin{pmatrix} H_{1W} \\ H_{2W} \end{pmatrix} \sim \mathbf{2}; \quad H_{SW} \sim \mathbf{1}_S; \quad H_{AW} \sim \mathbf{1}_A.$$

■ *Yukawa interactions*

$$\begin{aligned} \mathcal{L}_{Y_f} = & Y_1^f (\bar{g}_{SW} g_{SR} H_{SW}) + \frac{1}{\sqrt{2}} Y_2^f (\bar{f}_{1W} f_{1R} + \bar{f}_{2W} f_{2R}) H_{SW} + \\ & \frac{1}{2} Y_3^f \left[ (\bar{f}_{1W} H_{2W} + \bar{f}_{2W} H_{1W}) f_{1R} + (\bar{f}_{1W} H_{1W} - \bar{f}_{2W} H_{2W}) f_{2R} \right] + \\ & \frac{1}{\sqrt{2}} Y_4^f (\bar{f}_{1W} f_{2R} - \bar{f}_{2W} f_{1R}) H_{AW} + \frac{1}{\sqrt{2}} Y_5^f (\bar{f}_{1W} H_{1W} + \bar{f}_{2W} H_{2W}) g_{SR} + \\ & \frac{1}{\sqrt{2}} Y_6^f [\bar{g}_{SW} (H_{1W} f_{1R} + H_{2W} f_{2R})] + h.c., \quad l, f = d, \\ & f = u, \nu, \quad H_{\bullet W} \rightarrow i\sigma_2 H_{\bullet W}^* \end{aligned}$$

- Just as in the SM, Higgs acquire VEVs and CP violation comes from complex Yukawa couplings

$$\langle 0|H_{1W}|0\rangle \equiv w_1, \quad \langle 0|H_{2W}|0\rangle \equiv w_2, \quad \langle 0|H_{SW}|0\rangle \equiv v_S, \quad \langle 0|H_{AW}|0\rangle \equiv v_A.$$

- Effective mass Lagrangian after EWSB (same for leptons)

$$\mathcal{L}_q = -(\bar{u}_L, \bar{c}_L, \bar{t}_L) \mathcal{M}_{S_3}^u \begin{pmatrix} u_R \\ c_R \\ t_R \end{pmatrix} - (\bar{d}_L, \bar{s}_L, \bar{b}_L) \mathcal{M}_{S_3}^d \begin{pmatrix} d_R \\ s_R \\ b_R \end{pmatrix} + h.c.,$$

$$\mathcal{M}_{S_3}^f = \begin{pmatrix} \sqrt{2}Y_2^f v_S + Y_3^f w_2 & Y_3^f w_1 + \sqrt{2}Y_4^f v_A & \sqrt{2}Y_5^f w_1 \\ Y_3^f w_1 - \sqrt{2}Y_4^f v_A & \sqrt{2}Y_2^f v_S - Y_3^f w_2 & \sqrt{2}Y_5^f w_2 \\ \sqrt{2}Y_6^f w_1 & \sqrt{2}Y_6^f w_2 & 2Y_1^f v_S \end{pmatrix}.$$

- We know that if the matrix looks like

$$\mathcal{M}^f = \begin{pmatrix} 0 & A^f & 0 \\ A^{f*} & |B^f| & C^f \\ 0 & C^{f*} & |D^f| \end{pmatrix} \quad (\text{Think Fritzsch})$$

*it produces a good phenomenology (easy to obtain in these scenarios)*

- So, let us just make a rotation

$$\mathcal{M}^f \equiv \mathcal{R}(\theta)_{12} \mathcal{M}_{S_3}^f \mathcal{R}(\theta)_{12}^T$$

*and identify when we can then achieve a matrix with two texture zeroes*

$$Y_5^f = Y_6^{f*}, \quad \arg(Y_4^f) = \pm \frac{\pi}{2}, \quad \arg(Y_1^f) = \arg(Y_2^f) = \arg(Y_3^f) = 0$$

$$\sqrt{2} Y_2^f v_s = -Y_3^f w_2 \cos^2 \theta (1 - 3 \tan^2 \theta)$$

- We do not alter the mass matrices as given by the  $S_3$  symmetry if we perform the same rotation in both, up and down sectors and another different rotation in up and neutrino sector
- Once we identify this, we can easily and exactly diagonalize the mass matrices, writing

$$\mathcal{M}^f = \mathcal{P}_f^\dagger \bar{\mathcal{M}}^f \mathcal{P}_f,$$

$$\bar{\mathcal{M}}^f = m_3^f \begin{pmatrix} 0 & \sqrt{\frac{\tilde{m}_1^f \tilde{m}_2^f}{1 - \delta_f}} & 0 \\ \sqrt{\frac{\tilde{m}_1^f \tilde{m}_2^f}{1 - \delta_f}} & \tilde{m}_1^f - \tilde{m}_2^f + \delta_f & \sqrt{\frac{\delta_f}{1 - \delta_f}} \xi_1^f \xi_2^f \\ 0 & \sqrt{\frac{\delta_f}{1 - \delta_f}} \xi_1^f \xi_2^f & 1 - \delta_f \end{pmatrix};$$

$$\tilde{m}_i^f \equiv \frac{m_i^f}{m_3^f}, \quad \xi_1^f \equiv 1 - \tilde{m}_1^f - \delta_f, \quad \xi_2^f \equiv 1 + \tilde{m}_2^f - \delta_f,$$

## Quark mixing

- In this parameterization, we can obtain exact formulas for CKM elements

$$V_{ud}^{th} = \sqrt{\frac{\tilde{m}_c \tilde{m}_s \xi_1^u \xi_1^d}{\mathcal{D}_{1u} \mathcal{D}_{1d}}} + \sqrt{\frac{\tilde{m}_u \tilde{m}_d}{\mathcal{D}_{1u} \mathcal{D}_{1d}}} \left( \sqrt{(1 - \delta_u)(1 - \delta_d) \xi_1^u \xi_1^d} + \sqrt{\delta_u \delta_d \xi_2^u \xi_2^d} e^{i\phi_2} \right) e^{i\phi_1},$$

$$V_{us}^{th} = -\sqrt{\frac{\tilde{m}_c \tilde{m}_d \xi_1^u \xi_2^d}{\mathcal{D}_{1u} \mathcal{D}_{2d}}} + \sqrt{\frac{\tilde{m}_u \tilde{m}_s}{\mathcal{D}_{1u} \mathcal{D}_{2d}}} \left( \sqrt{(1 - \delta_u)(1 - \delta_d) \xi_1^u \xi_2^d} + \sqrt{\delta_u \delta_d \xi_2^u \xi_1^d} e^{i\phi_2} \right) e^{i\phi_1},$$

$$V_{ub}^{th} = \sqrt{\frac{\tilde{m}_c \tilde{m}_d \tilde{m}_s \delta_d \xi_1^u}{\mathcal{D}_{1u} \mathcal{D}_{3d}}} + \sqrt{\frac{\tilde{m}_u}{\mathcal{D}_{1u} \mathcal{D}_{3d}}} \left( \sqrt{(1 - \delta_u)(1 - \delta_d) \delta_d \xi_1^u} - \sqrt{\delta_u \xi_2^u \xi_1^d \xi_2^d} e^{i\phi_2} \right) e^{i\phi_1},$$

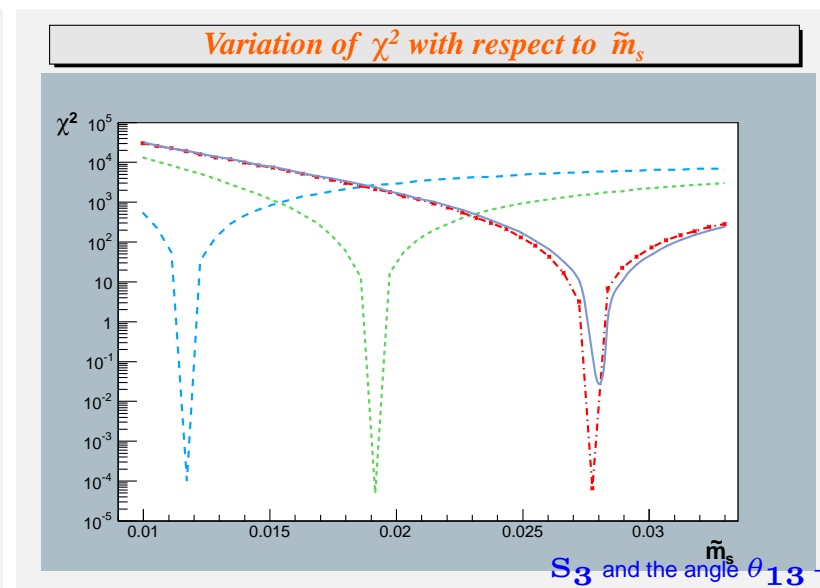
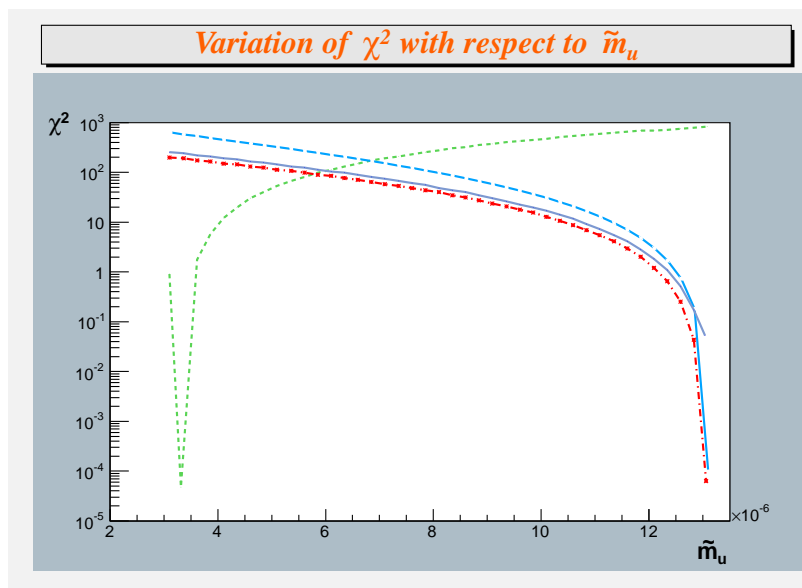
- Which allow us to make a precision  $\chi^2$  test

$$\chi^2 = \frac{(V_{ud}^{th} - V_{ud})^2}{\sigma_{V_{ud}}^2} + \frac{(V_{us}^{th} - V_{us})^2}{\sigma_{V_{us}}^2} + \frac{(V_{ub}^{th} - V_{ub})^2}{\sigma_{V_{ub}}^2} + \frac{(\mathcal{J}_q^{th} - \mathcal{J}_q)^2}{\sigma_{\mathcal{J}_q}^2},$$

$\tilde{m}_u$  and  $\tilde{m}_s$  have the biggest uncertainties, so have a big impact on the  $\chi^2$  value

	Mass ratios at $M_Z$
$\tilde{m}_u (M_Z)$	$0.0000081 \pm 0.0000025$
$\tilde{m}_c (M_Z)$	$0.0036 \pm 0.0004$
$\tilde{m}_d (M_Z)$	$0.000998 \pm 0.000188$
$\tilde{m}_s (M_Z)$	$0.021 \pm 0.006$

Fits with 3 or 4 Higgs doublets



## Lepton mixing

- See-saw  $M_{\nu L} = M_{\nu D} M_{\nu R}^{-1} M_{\nu D}^T$
- The mixing matrix has exactly the same form of the quark sector with the obvious replacements

$$\begin{aligned}
 V_{e1}^{th} &= \sqrt{\frac{\tilde{m}_\mu \tilde{m}_{\nu 2} f_{l1} f_{\nu 1}}{\mathcal{D}_{l1} \mathcal{D}_{\nu 1}}} + \sqrt{\frac{\tilde{m}_e \tilde{m}_{\nu 1}}{\mathcal{D}_{l1} \mathcal{D}_{\nu 1}}} \left( \sqrt{(1 - \delta_l)(1 - \delta_\nu)} f_{l1} f_{\nu 1} e^{i\Phi_1} + \sqrt{\delta_l \delta_\nu} f_{l2} f_{\nu 2} e^{i\Phi_2} \right), \\
 V_{e2}^{th} &= -\sqrt{\frac{\tilde{m}_\mu \tilde{m}_{\nu 1} f_{l1} f_{\nu 2}}{\mathcal{D}_{l1} \mathcal{D}_{\nu 2}}} + \sqrt{\frac{\tilde{m}_e \tilde{m}_{\nu 2}}{\mathcal{D}_{l1} \mathcal{D}_{\nu 2}}} \left( \sqrt{(1 - \delta_l)(1 - \delta_\nu)} f_{l1} f_{\nu 2} e^{i\Phi_1} + \sqrt{\delta_l \delta_\nu} f_{l2} f_{\nu 1} e^{i\Phi_2} \right), \\
 V_{e3}^{th} &= \sqrt{\frac{\tilde{m}_\mu \tilde{m}_{\nu 1} \tilde{m}_{\nu 2} \delta_\nu f_{l1}}{\mathcal{D}_{l1} \mathcal{D}_{\nu 3}}} + \sqrt{\frac{\tilde{m}_e}{\mathcal{D}_{l1} \mathcal{D}_{\nu 3}}} \left( \sqrt{\delta_\nu (1 - \delta_l)(1 - \delta_\nu)} f_{l1} e^{i\Phi_1} - \sqrt{\delta_e} f_{l2} f_{\nu 1} f_{\nu 2} e^{i\Phi_2} \right),
 \end{aligned}$$

- Neutrino masses (example of values with 3 Higgs doublets )

$$\tilde{m}_{\nu 1} = \sqrt{1 - \frac{(\Delta m_{32}^2 + \Delta m_{21}^2)}{m_{\nu 3}^2}}, \quad \tilde{m}_{\nu 2} = \sqrt{1 - \frac{\Delta m_{32}^2}{m_{\nu 3}^2}}.$$

$$m_{\nu 1} = \left( 3.22_{-0.39}^{+0.67} \right) \times 10^{-3} \text{ eV}, \quad m_{\nu 2} = \left( 9.10_{-0.13}^{+0.25} \right) \times 10^{-3} \text{ eV}, \quad m_{\nu 3} = \left( 4.92_{-0.22}^{+0.21} \right) \times 10^{-2} \text{ eV}.$$



## Mixing angles

- 3 Higgs doublets:  $M_{1R}$  and  $M_{2R}$  degenerate (our case has the same features as the one in 1205.4755)

$$\sin^2 \theta_{12}^{lh} \approx \frac{f_{\nu 2} \left\{ \frac{\tilde{m}_{\nu 1}}{\tilde{m}_{\nu 2}} + \frac{\tilde{m}_e}{\tilde{m}_\mu} (1 - \delta_\nu) + 2 \sqrt{\frac{\tilde{m}_{\nu 1}}{\tilde{m}_{\nu 2}} \frac{\tilde{m}_e}{\tilde{m}_\mu} (1 - \delta_\nu) \cos \Phi_1} \right\}}{(1 + \tilde{m}_{\nu 2})(1 - \delta_\nu) \left(1 + \frac{\tilde{m}_{\nu 1}}{\tilde{m}_{\nu 2}}\right) \left(1 + \frac{\tilde{m}_e}{\tilde{m}_\mu}\right)},$$

$$\sin^2 \theta_{23}^{lh} \approx \frac{\delta_\nu + \delta_e f_{\nu 2} - \sqrt{\delta_\nu \delta_e f_{\nu 2} \cos(\Phi_1 - \Phi_2)}}{\left(1 + \frac{\tilde{m}_e}{\tilde{m}_\mu}\right) (1 + \tilde{m}_{\nu 2})},$$

$$\sin^2 \theta_{13}^{lh} \approx \frac{\delta_\nu \left\{ \frac{\tilde{m}_e}{\tilde{m}_\mu} + \frac{\tilde{m}_{\nu 1} \tilde{m}_{\nu 2}}{(1 - \delta_\nu)} - 2 \sqrt{\frac{\tilde{m}_e}{\tilde{m}_\mu} \frac{\tilde{m}_{\nu 1} \tilde{m}_{\nu 2}}{(1 - \delta_\nu)} \cos \Phi_1} \right\}}{\left(1 + \frac{\tilde{m}_e}{\tilde{m}_\mu}\right) (1 + \tilde{m}_{\nu 2})}.$$

$$\theta_{12}^{lh} = \left(34.43_{-0.98}^{+0.85}\right)^\circ, \quad \theta_{23}^{lh} = \left(43.60_{-2.22}^{+1.97}\right)^\circ, \quad \theta_{13}^{lh} = \left(6.80_{-0.66}^{+0.95}\right)^\circ,$$

## The angle $\theta_{13}$

- 4 Higgs doublets:  $M_{1R}$  and  $M_{2R}$  not degenerate

$$m_{\nu_2} = 0.056 \text{ eV}, \quad m_{\nu_1} = 0.053 \text{ eV}, \quad m_{\nu_3} = 0.048 \text{ eV},$$
$$\delta_l = \pi/2, \quad \mu_0 = 0.049 \text{ eV}, \quad d = 8 \times 10^{-5} \text{ eV},$$

$$\sin^2 \theta_{13}^l \approx \frac{(\mu_0 + 2d - m_{\nu_3})(\mu_0 - m_{\nu_3})}{(m_{\nu_1} - m_{\nu_3})(m_{\nu_2} - m_{\nu_3})}.$$
$$\theta_{13}^l \approx 9.8^\circ$$

$$\text{Daya Bay} \rightarrow \theta_{13}^l \approx 9.1^\circ \quad d = \frac{2|\lambda||\mu_2^\nu|^2}{|\overline{M}|}, \quad \mu_2^\nu \text{ free parameter}$$

$$\lambda = \frac{1}{2} \left( \frac{M_2 - M_1}{M_1 + M_2} \right), \quad \overline{M} = 2 \frac{M_1 M_2}{M_2 + M_1}.$$

## Counting of free parameters

- *Quark sector: there are 10 observables*
- *Lepton sector: there are 10 + 2 observables*
- *Quark sector: (both models o.k. but different qualities of the fit)*
  - *3 Higgs-doublets: 9 free parameters*
  - *4 Higgs-doublets: 10 parameters*
- *Lepton sector:*
  - *3 Higgs-doublets:  $M_{1R}$  and  $M_{2R}$  almost degenerate,  $\theta_{13}^l$  **small**: 9 free parameters*
  - *4 Higgs-doublets:  $M_{1R}$  not  $M_{2R}$  degenerate : 10 free parameters*

$$\theta_{13}^l \approx 9.8^\circ$$

- *Status: We can accommodate observables to a great degree*

# Outlook

- *Try to reduce parameters*
- *FCNC under control, but precise values still to be computed*
- *Interesting contributions to the  $a_\mu$*
- *Understand the values of the VEVs of the Higgs  $\leftrightarrow$  EWSB*