

# Strong (light) Higgs dynamics with MFV

Juan Yepes

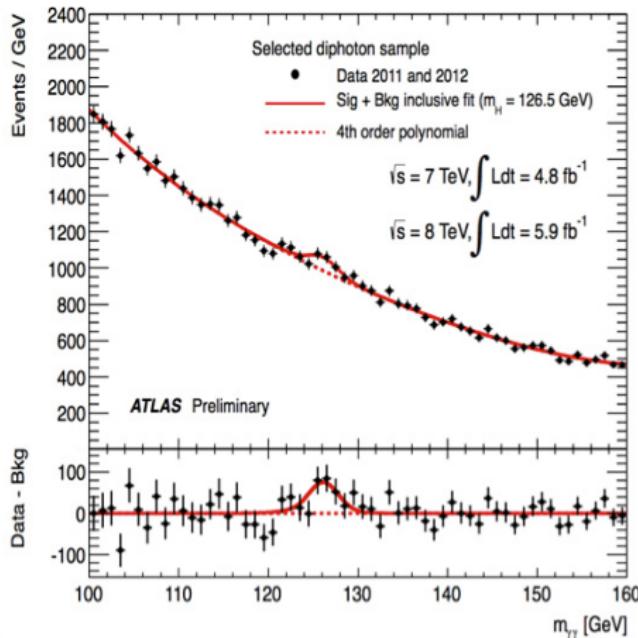


WHAT IS  $\nu$ ?-GGI Workshop-2012, Florence

Alonso, Gavela, Merlo, Rigolin & JY, JHEP 1206 (2012) 076, [hep-ph/1201.1511](#)

July 10, 2012

# ATLAS

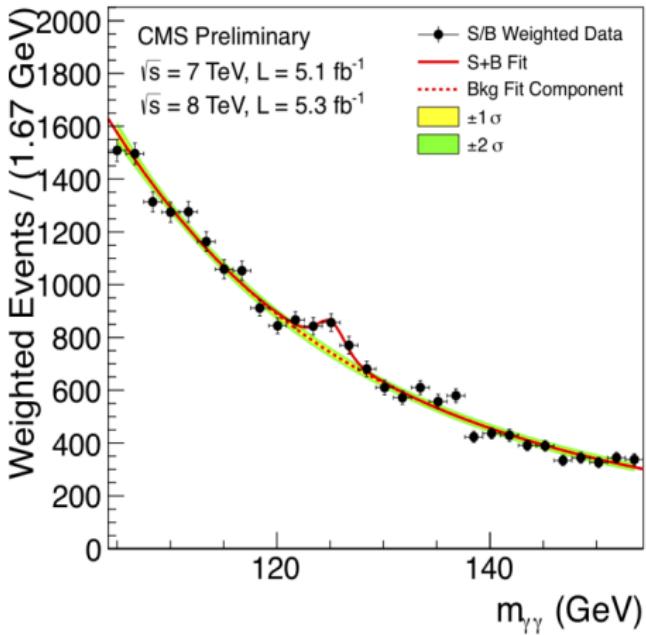


$123 < m_H(\text{GeV}) < 130$

Not excluded

$123 < m_H(\text{GeV}) < 130$  @  
95% C.L.

Excess  $\sim 126.5 \text{ GeV}$  @  $5\sigma$



$122.5 < m_H(\text{GeV}) < 127$

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Global

Local

$$SU(2)_L \times SU(2)_R \rightarrow \mathbf{U} \sim (2, 2) \quad SU(2)_L \times U(1)_Y \rightarrow L(x) \mathbf{U}(x) R^\dagger(x),$$

$$L(x) = e^{i\vec{\epsilon}_L(x) \cdot \vec{\tau}/2}, \quad R(x) = e^{i\epsilon_Y(x)\tau_3/2}$$

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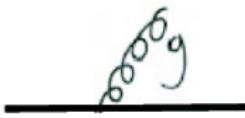
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Covariant derivative

$$\mathcal{D}_\mu \mathbf{U} \equiv \partial_\mu \mathbf{U} + \frac{i g}{2} \tau_i W_\mu^i \mathbf{U} - \frac{i g'}{2} \mathbf{U} \tau_3 B_\mu.$$

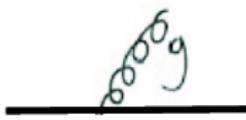
# What changes if the Higgs has strong interacting dynamics?

As in strongly interacting QCD  $\alpha_s \sim 1 \Rightarrow$  unsuppressed multiple gluon emission



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Strongly interacting Higgs  $\lambda \sim 1 \Rightarrow$  unsuppressed longitudinal  $W - Z$  components emission



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Building blocks:  $\mathbf{U}$ ,  $\mathcal{D}_\mu \mathbf{U}$ ,  $\mathbf{T} = \mathbf{U} \tau_3 \mathbf{U}^\dagger$ ,  $\mathbf{V}_\mu = (\mathcal{D}_\mu \mathbf{U}) \mathbf{U}^\dagger$

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Tower of operators contain, e.g.,

$$\frac{v^2}{4} \text{Tr}[\mathbf{V}^\mu \mathbf{V}_\mu] \dots$$

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⇒ **STILL STRONG DYNAMICS**

Giudice, Grojean, Pomarol & Rattazzi '07

$$\Rightarrow \quad a_i \mathcal{O}_i \left[ \rho + a \frac{h}{f} + b \frac{h^2}{f^2} + \dots \right], \quad \rho = \frac{v}{f}$$

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We focus on the 1st term for  $v \approx f$ .

Now go to **FLAVOR** ⇒

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$$\mathcal{G}_F = SU(3)_{Q_L} \times SU(3)_{U_R} \times SU(3)_{D_R}$$

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recovering  $\mathcal{G}_f$

$$\implies Y_U \sim (3, \bar{3}, 1) \quad \& \quad Y_D \sim (3, 1, \bar{3})$$

$$\mathcal{L} = \mathcal{L}_{SM} + a_i \frac{\mathcal{O}_i^{d=6}}{\Lambda_f^2} + \dots$$

$$\mathcal{O}^{d=6} \sim c_{\alpha\beta} \bar{\psi}_\alpha \gamma^\mu \psi_\beta (\Phi^\dagger \mathcal{D}_\mu \Phi), \quad c_{\alpha\beta} c_{\gamma\delta} \bar{\psi}_\alpha \psi_\beta \bar{\psi}_\gamma \psi_\delta$$

$c \sim YY^\dagger$ , D'Ambrosio, Giudice, Isidori & Strumia '02

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Minimally flavour violating dimension six operator	main observables	$\Lambda_f$ [TeV]
		-    +
$\mathcal{O}_0 = \frac{1}{2}(\bar{Q}_L \lambda_{FC} \gamma_\mu Q_L)^2$	$\epsilon_K, \Delta m_{B_d}$	6.4    5.0
$\mathcal{O}_{F1} = H^\dagger \left( \bar{D}_R \lambda_d \lambda_{FC} \sigma_{\mu\nu} Q_L \right) F_{\mu\nu}$	$B \rightarrow X_s \gamma$	9.3    12.4
$\mathcal{O}_{G1} = H^\dagger \left( \bar{D}_R \lambda_d \lambda_{FC} \sigma_{\mu\nu} T^a Q_L \right) G_{\mu\nu}^a$	$B \rightarrow X_s \gamma$	2.6    3.5
$\mathcal{O}_{\ell 1} = (\bar{Q}_L \lambda_{FC} \gamma_\mu Q_L)(\bar{L}_L \gamma_\mu L_L)$	$B \rightarrow (X)\ell\bar{\ell}, \quad K \rightarrow \pi\nu\bar{\nu}, (\pi)\ell\bar{\ell}$	3.1    2.7    *
$\mathcal{O}_{\ell 2} = (\bar{Q}_L \lambda_{FC} \gamma_\mu \tau^a Q_L)(\bar{L}_L \gamma_\mu \tau^a L_L)$	$B \rightarrow (X)\ell\bar{\ell}, \quad K \rightarrow \pi\nu\bar{\nu}, (\pi)\ell\bar{\ell}$	3.4    3.0    *
$\mathcal{O}_{H1} = (\bar{Q}_L \lambda_{FC} \gamma_\mu Q_L)(H^\dagger i D_\mu H)$	$B \rightarrow (X)\ell\bar{\ell}, \quad K \rightarrow \pi\nu\bar{\nu}, (\pi)\ell\bar{\ell}$	1.6    1.6    *
$\mathcal{O}_{q5} = (\bar{Q}_L \lambda_{FC} \gamma_\mu Q_L)(\bar{D}_R \gamma_\mu D_R)$	$B \rightarrow K\pi, \quad \epsilon'/\epsilon, \dots$	$\sim 1$

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$$\Rightarrow \Lambda_f \sim \text{TeV}$$

## STRONG HIGGS DYNAMICS + MINIMAL FLAVOR VIOLATION

Non-linear expansion @  $d_\chi = 4$

$$\mathcal{O}_1 \sim \bar{\psi}_\alpha \gamma^\mu \{ \mathbf{U} \tau_3 \mathbf{U}^\dagger, (\mathcal{D}_\mu \mathbf{U}) \mathbf{U}^\dagger \} \psi_\beta,$$

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Linear expansion

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## Correspondences

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$\mathcal{O}_4$  is a CP-ODD op.! → Natural CP @ LO!!

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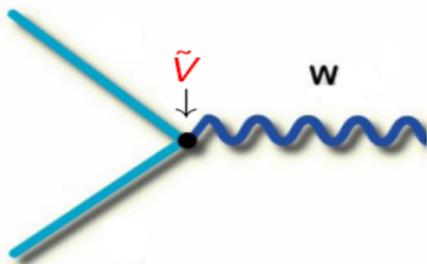
$$\delta \mathcal{L}_{d_\chi=4} = -\frac{g}{\sqrt{2}} [W^{\mu+} \bar{U}_L \gamma_\mu (a_W + i a_{CP}) (\mathbf{y}_U^2 V + V \mathbf{y}_D^2) D_L + h.c.] + \\ -\frac{g}{2c_W} Z^\mu [a_Z^u \bar{U}_L \gamma_\mu (\mathbf{y}_U^2 + V \mathbf{y}_D^2 V^\dagger) U_L + a_Z^d \bar{D}_L \gamma_\mu (\mathbf{y}_D^2 + V^\dagger \mathbf{y}_U^2 V) D_L]$$

$$a_Z^u \equiv a_1 + a_2 + a_3, \quad a_Z^d \equiv a_1 - a_2 - a_3, \\ a_W \equiv a_2 - a_3, \quad a_{CP} \equiv -a_4.$$

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$$\begin{aligned} \delta\mathcal{L}_{d_\chi=4} = & -\frac{g}{\sqrt{2}} [W^{\mu+}\bar{U}_L\gamma_\mu(\textcolor{red}{a_W} + i\textcolor{red}{a_{CP}})(\mathbf{y}_U^2 V + V\mathbf{y}_D^2) D_L + h.c.] + \\ & -\frac{g}{2c_W} Z^\mu [\textcolor{red}{a_Z^u} \bar{U}_L\gamma_\mu (\mathbf{y}_U^2 + V\mathbf{y}_D^2 V^\dagger) U_L + \textcolor{red}{a_Z^d} \bar{D}_L\gamma_\mu (\mathbf{y}_D^2 + V^\dagger\mathbf{y}_U^2 V) D_L] \end{aligned}$$

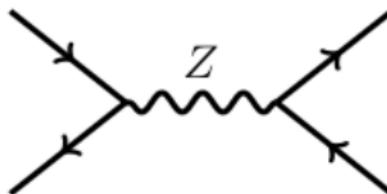
$$\begin{aligned} \textcolor{red}{a}_Z^u &\equiv a_1 + a_2 + a_3, & \textcolor{red}{a}_Z^d &\equiv a_1 - a_2 - a_3, \\ \textcolor{red}{a}_W &\equiv a_2 - a_3, & \textcolor{red}{a}_{CP} &\equiv -a_4. \end{aligned}$$



$$\tilde{V}_{ij} = V_{ij} \left[ 1 + (\textcolor{red}{a_W} + i \textcolor{red}{a_{CP}})(y_{u_i}^2 + y_{d_j}^2) \right]$$

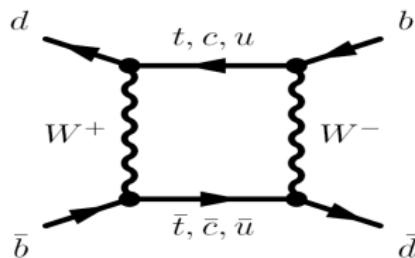
⇒ Impacts on  $\Delta F = 1$  &  $\Delta F = 2$  observables...

$$\Delta F = 1$$



Wilson coefficient modification  $\Rightarrow Q_{\bar{\nu}\nu}, Q_{9V}, Q_7 \dots$

$$\Delta F = 2$$



Modifications on  $\Rightarrow M_{12}^K, \varepsilon_K, M_{12}^{d,s}, A_{sl}^b$

# $\Delta F = 1$ observables

$$\frac{G_F \alpha}{2\sqrt{2}\pi s_W^2} V_{ti}^* V_{tj} \sum_n C_n Q_n + \text{h.c.},$$

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$$\frac{G_F \alpha}{2\sqrt{2}\pi s_W^2} V_{ti}^* V_{tj} \sum_n C_n \mathcal{Q}_n + \text{h.c.}, \quad C_n = C_n^{SM}$$

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$$\frac{G_F \alpha}{2\sqrt{2}\pi s_W^2} V_{ti}^* V_{tj} \sum_n C_n \mathcal{Q}_n + \text{h.c.}, \quad C_n = C_n^{SM} + \mathcal{C}_n^{NP}$$

# $\Delta F = 1$ observables

$$\frac{G_F \alpha}{2\sqrt{2}\pi s_W^2} V_{ti}^* V_{tj} \sum_n C_n \mathcal{Q}_n + \text{h.c.}, \quad C_n = C_n^{SM} + \textcolor{blue}{C_n^{NP}}$$

FCNC operators basis

$$\begin{aligned} \mathcal{Q}_{\bar{\nu}\nu} &= \bar{d}_i \gamma_\mu (1 - \gamma_5) d_j \bar{\nu} \gamma^\mu (1 - \gamma_5) \nu, & \mathcal{Q}_7 &= e_q \bar{d}_i \gamma_\mu (1 - \gamma_5) d_j \bar{q} \gamma^\mu (1 + \gamma_5) q, \\ \mathcal{Q}_{9V} &= \bar{d}_i \gamma_\mu (1 - \gamma_5) d_j \bar{\ell} \gamma^\mu \ell, & \mathcal{Q}_9 &= e_q \bar{d}_i \gamma_\mu (1 - \gamma_5) d_j \bar{q} \gamma^\mu (1 - \gamma_5) q, \\ \mathcal{Q}_{10A} &= \bar{d}_i \gamma_\mu (1 - \gamma_5) d_j \bar{\ell} \gamma^\mu \gamma_5 \ell, & \mathcal{Q}_{7\gamma} &= \frac{m_i}{g^2} \bar{d}_i (1 - \gamma_5) \sigma_{\mu\nu} d_j (e F^{\mu\nu}), \\ && \mathcal{Q}_{8G} &= \frac{m_i}{g^2} \bar{d}_i (1 - \gamma_5) \sigma_{\mu\nu} T^a d_j (g_s G_a^{\mu\nu}). \end{aligned}$$

Wilson coefficient modifications

$$C_n^{NP} \left\{ \begin{array}{ll} \sim y_t^2 a_Z^d, & n = \bar{\nu}\nu, 9V, \dots, 9 \\ 0, & n = 7\gamma, 8G \end{array} \right.$$

## $\Delta F = 1$ observables

Operator	Observable	Bound (@ 95% C.L.)
$\mathcal{O}_{9V}$	$B \rightarrow X_s I^+ I^-$	$-0.811 < a_Z^d < 0.232$
$\mathcal{O}_{10A}$	$B \rightarrow X_s I^+ I^- , B \rightarrow \mu^+ \mu^-$	$-0.050 < a_Z^d < 0.009$
$\mathcal{O}_{\bar{\nu}\nu}$	$K^+ \rightarrow \pi^+ \bar{\nu}\nu$	$-0.044 < a_Z^d < 0.133$

## $\Delta F = 2$ observables

$\Delta F = 2$ -effective hamiltonian

$$\mathcal{H}_{\text{eff}}^{\Delta F=2} = \frac{G_F^2 M_W^2}{4\pi^2} C(\mu) Q,$$

$$Q = (\bar{d}_i^\alpha \gamma_\mu P_L d_j^\alpha)(\bar{d}_i^\beta \gamma^\mu P_L d_j^\beta)$$

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Mixing amplitudes:

$$M_{12}^K = \frac{\langle \bar{K}^0 | \mathcal{H}_{\text{eff}}^{\Delta S=2} | K^0 \rangle^*}{2 m_K},$$

$$M_{12}^q = \frac{\langle \bar{B}_q^0 | \mathcal{H}_{\text{eff}}^{\Delta B=2} | B_q^0 \rangle^*}{2 m_{B_q}} \quad q = d, s,$$

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Either  $K$  or  $B$ -system,

$$M_{12}$$

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Either  $K$  or  $B$ -system,

$$M_{12} = (M_{12})_{SM}$$

## $\Delta F = 2$ observables

$\Delta F = 2$ -effective hamiltonian

$$\mathcal{H}_{\text{eff}}^{\Delta F=2} = \frac{G_F^2 M_W^2}{4\pi^2} C(\mu) Q,$$

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Either  $K$  or  $B$ -system,

$$M_{12} = (M_{12})_{SM} + (M_{12})_{NP}$$

# $\Delta F = 2$ observables

## Neutral kaon oscillation

$$\Delta M_K = 2 \left[ \text{Re}(M_{12}^K)_{SM} + \text{Re}(\textcolor{blue}{M}_{12}^K)_{NP} \right],$$

$$\varepsilon_K = \frac{\kappa_\epsilon e^{i\varphi_\epsilon}}{\sqrt{2}(\Delta M_K)_{\text{exp}}} \left[ \text{Im} \left( M_{12}^K \right)_{SM} + \text{Im} \left( \textcolor{blue}{M}_{12}^K \right)_{NP} \right]$$

## $\Delta F = 2$ observables

### Neutral kaon oscillation

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Neglecting all contributions proportional to  $y_{u,d,s}$  and  $y_c^n$  with  $n > 2$ :

$$\begin{aligned} (M_{12}^K)_{NP} \sim & \eta_2 \lambda_t^2 \mathcal{O} \left( y_t^2 a_W, y_t^4 a_{CP}^2, y_t^4 (a_Z^d)^2 \right) \\ & + \eta_1 \lambda_c^2 \mathcal{O} (y_c^2 a_W) \\ & + 2 \eta_3 \lambda_t \lambda_c \mathcal{O} (y_t^2 a_W, y_t^4 a_{CP}^2) \end{aligned}$$

$\Delta F = 2$  observables

Neutral meson oscillation

Mixing amplitude

$$M_{12}^q = (M_{12}^q)_{\text{SM}} \mathcal{C}_{B_q} e^{2i\varphi_{B_q}}$$

$B_{d,s}$ -mass differences

$$\Delta M_q = 2 |M_{12}^q| \equiv (\Delta M_q)_{\text{SM}} \mathcal{C}_{B_q}$$

$$\mathcal{C}_{B_d} = \mathcal{C}_{B_s} = \left| 1 + \mathcal{O} \left( y_t^2 a_W, y_t^4 (a_Z^d)^2 \right) + i \mathcal{O} \left( y_t^2 y_b^2 a_W a_{CP} \right) \right|$$

$$\varphi_{B_d} = \varphi_{B_s} \sim \mathcal{O} \left( y_t^2 y_b^2 a_W, a_{CP} \right)$$

# $\Delta F = 2$ observables

## Neutral meson oscillation

Mixing-induced CP asymmetries  $S_{\psi K_S}$  &  $S_{\psi \phi}$  for  
 $B_d^0 \rightarrow \psi K_S$  &  $B_s^0 \rightarrow \psi \phi$

$$S_{\psi K_S} = \sin(2\beta + 2\varphi_{B_d}), \quad S_{\psi \phi} = \sin(2\beta_s - 2\varphi_{B_s}),$$

UT-angles  $\beta$  &  $\beta_s$

$$\beta \equiv \arg \left( -\frac{V_{cb}^* V_{cd}}{V_{tb}^* V_{td}} \right), \quad \beta_s \equiv \arg \left( -\frac{V_{tb}^* V_{ts}}{V_{cb}^* V_{cs}} \right),$$

$R_{BR/\Delta M}$

$$R_{BR/\Delta M} \sim \frac{|1 + (\textcolor{red}{a_W} + i \textcolor{red}{a_{CP}}) y_b^2|^2}{\mathcal{C}_{B_d}}$$

## $\Delta F = 2$ observables

### $B$ -semileptonic CP-Asymmetry

$$A_{sl}^b = (0.594 \pm 0.022) a_{sl}^d + (0.406 \pm 0.022) a_{sl}^s ,$$

NP contributions

$$\Gamma_{12}^q = (\Gamma_{12}^q)_{SM} \tilde{\mathcal{C}}_{B_q} \quad \text{with} \quad \tilde{\mathcal{C}}_{B_q} = 1 + 2 \textcolor{red}{a_W} y_b^2 ,$$

$$a_{sl}^q = \left| \frac{(\Gamma_{12}^q)_{SM}}{(M_{12}^q)_{SM}} \right| \frac{\tilde{\mathcal{C}}_{B_q}}{\mathcal{C}_{B_q}} \sin(\phi_q + 2\varphi_{B_q}) ,$$

$\implies$

## $\varepsilon_K$ vs. $R_{BR/\Delta M}$

$a'$ 's from  $a_i \mathcal{O}_i$

$$R_{BR/\Delta M} = \frac{BR(B^+ \rightarrow \tau^+ \nu)}{\Delta M_{B_d}}$$

$$\begin{aligned} a_{CP} &= \pm 0.1 & \rightarrow & \delta \varepsilon_K \approx 1.1\%, & \delta R \approx -1.4\%, \\ a_W &= 0.1(-0.1) & \rightarrow & \delta \varepsilon_K \approx +26\%(-19\%), & \delta R \approx -25\%(+30\%), \\ a_Z^d &= \pm 0.1 & \rightarrow & \delta \varepsilon_K \approx 124\%, & \delta R \approx -62\%. \end{aligned}$$

## $\varepsilon_K$ vs. $R_{BR/\Delta M}$

$a'$ 's from  $a_i \mathcal{O}_i$

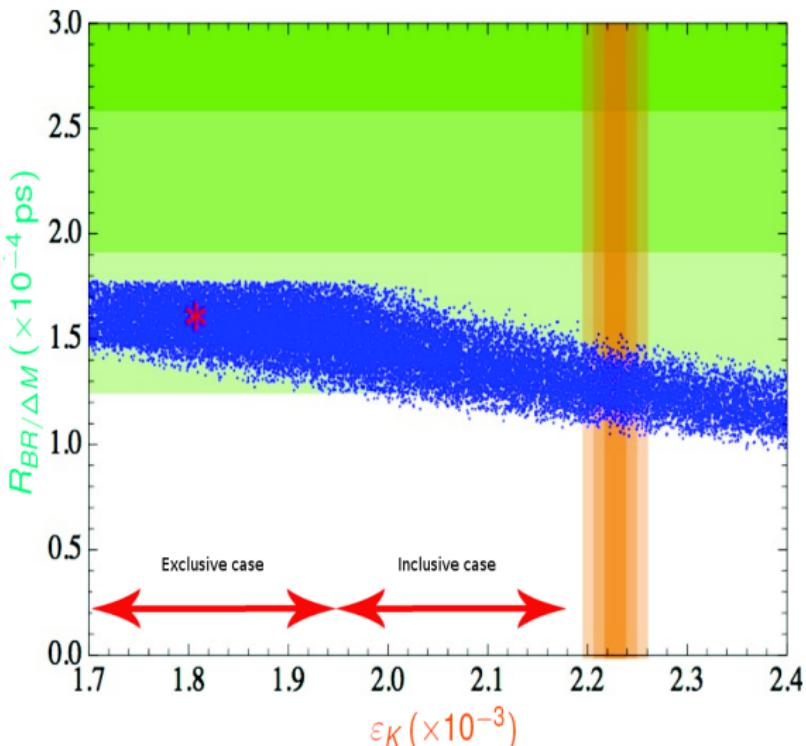
$$R_{BR/\Delta M} = \frac{BR(B^+ \rightarrow \tau^+ \nu)}{\Delta M_{B_d}}$$

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$$\varepsilon_K \uparrow (\approx \varepsilon_K^{\text{exp}} \text{ } \& \text{ } S_{\psi K_S} \approx S_{\psi K_S}^{\text{exp}}) \text{ } \& \text{ } R_{BR/\Delta M} \downarrow$$

$\Rightarrow$  SHD + MFV able to soften  $\varepsilon_K - S_{\psi K_S}$  anomaly,  
but not the SM tension for  $BR(B^+ \rightarrow \tau^+ \nu)$

$\varepsilon_K$  vs.  $R_{BR/\Delta M}$  from  $\mathcal{O}_4$

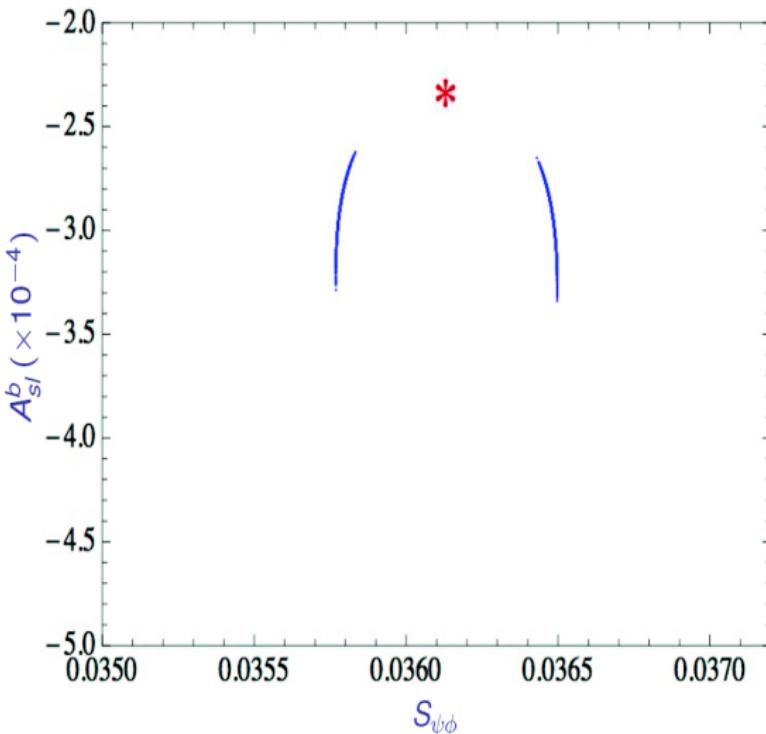


- : correlation
  - $\varepsilon_K - R_{BR/\Delta M}$
  - \*: SM values

Green & Orange:  
1, 2, 3 $\sigma$  exp. values

$a_{CP} \in [-1, 1]$

# $S_{\psi\phi}$ vs. $A_{sl}^b$ from $\mathcal{O}_4$



●: correlation  $S_{\psi\phi} - A_{sl}^b$

\*: SM values

Large values for  $a_{CP}$

# STRONG HIGGS DYNAMICS

+

# MINIMAL FLAVOR VIOLATION



- ▶ Different MFV phenomenology for the perturbative Higgs and the strong interacting regime, e.g.,  $\mathcal{O}_4$
- ▶ Natural  $\cancel{CP}(\mathcal{O}_4) @ \text{LO}!!$
- ▶  $\varepsilon_K - S_{\psi K_S}$  anomaly softened, still SM tension for  $BR(B^+ \rightarrow \tau^+ \nu)$
- ▶ Small  $\delta S_{\psi\phi}$  &  $\delta A_{sl}^b$  experimentally allowed

$d_\chi = 4$  ops. have effects also on  $b \rightarrow s\gamma\dots$

Going to  $d_\chi = 5\dots$

$$\bar{Q}_L \mathcal{O}_i \mathbf{U} Q_R$$

$$\begin{aligned}\mathcal{O}_1 &= \mathbf{V}_\mu \mathbf{V}^\mu, & \mathcal{O}_4 &= \mathbf{V}_\mu \mathbf{T} \mathbf{V}^\mu \mathbf{T} \\ \mathcal{O}_2 &= \mathbf{V}_\mu \mathbf{V}^\mu \mathbf{T}, & \mathcal{O}_5 &= \mathbf{T} \mathbf{V}_\mu \mathbf{T} \mathbf{V}^\mu \\ \mathcal{O}_3 &= \mathbf{V}_\mu \mathbf{T} \mathbf{V}^\mu, & \mathcal{O}_6 &= \mathbf{T} \mathbf{V}_\mu \mathbf{T} \mathbf{V}^\mu \mathbf{T}\end{aligned}$$

$$\bar{Q}_L \sigma^{\mu\nu} \mathcal{O}_i U Q_R$$

$$\begin{aligned}\mathcal{O}_7 &= [\mathbf{V}_\mu, \mathbf{V}_\nu], & \mathcal{O}_9 &= [\mathbf{V}_\mu \mathbf{T}, \mathbf{V}_\nu \mathbf{T}] \\ \mathcal{O}_8 &= [\mathbf{V}_\mu, \mathbf{V}_\nu] \mathbf{T}, & \mathcal{O}_{10} &= [\mathbf{V}_\mu \mathbf{T}, \mathbf{V}_\nu \mathbf{T}] \mathbf{T}\end{aligned}$$

Dipole-type

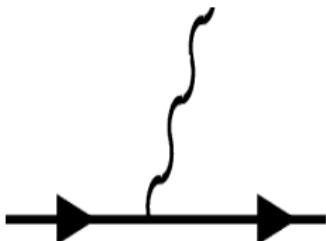
$$\begin{aligned}\mathcal{O}_{11} &= B_{\mu\nu}, & \mathcal{O}_{13} &= W_{\mu\nu}, & \mathcal{O}_{15} &= \mathbf{T} W_{\mu\nu} \\ \mathcal{O}_{12} &= B_{\mu\nu} \mathbf{T}, & \mathcal{O}_{14} &= W_{\mu\nu} \mathbf{T}, & \mathcal{O}_{16} &= \mathbf{T} W_{\mu\nu} \mathbf{T}\end{aligned}$$

⇒ Impact on  $\Delta F = 1$  observable...dipole-type operator

$$\frac{a'}{\Lambda_{NP}^2} H^\dagger \bar{D}_R Y_d \lambda_{FC} \sigma_{\mu\nu} Q_L F^{\mu\nu},$$

$$b \rightarrow s\gamma$$

From  $d_\chi = 5$



$$\frac{a_{7\gamma}}{f} = \frac{v}{2\sqrt{2}} \frac{a'}{\Lambda_{NP}^2} \leq \frac{10^{-4} - 10^{-3}}{\text{TeV}}$$

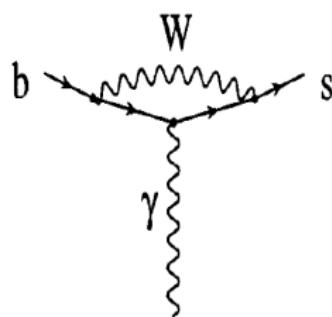
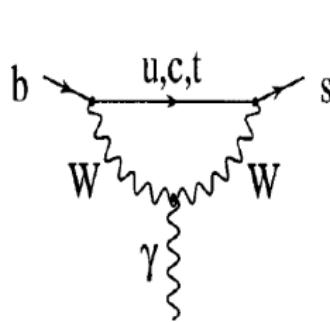
$$\Lambda_{NP} \sim 10 \text{TeV}$$

$$\alpha_{7\gamma} = c_W a_{B_{\mu\nu}} + s_W a_{W_{\mu\nu}^3}$$

$\Rightarrow$  Impact on  $\Delta F = 1$  observable...dipole-type operator

$$b \rightarrow s\gamma$$

From  $d_\chi = 4$



$$a_W, a_{CP}, a_Z^d \sim \mathcal{O}(1)$$

# Thanks!

## Tools

$$\mathcal{D}_\mu \mathbf{U} \equiv \partial_\mu \mathbf{U} + \frac{i g}{2} \tau_i W_\mu^i \mathbf{U} - \frac{i g'}{2} \mathbf{U} \tau_3 B_\mu$$

$$\mathbf{T} = \mathbf{U} \tau_3 \mathbf{U}^\dagger,$$

$$\mathbf{T} \rightarrow L \mathbf{T} L^\dagger,$$

$$\mathbf{V}_\mu = (\mathcal{D}_\mu \mathbf{U}) \mathbf{U}^\dagger,$$

$$\mathbf{V}_\mu \rightarrow L \mathbf{V}_\mu L^\dagger.$$

Lagrangian for interaction between the gauge fields and the scalar sector:

$$\mathcal{L} = -\frac{1}{4} W_{\mu\nu}^i W_i^{\mu\nu} - \frac{1}{4} B_{\mu\nu} B^{\mu\nu} + \frac{v^2}{4} \text{Tr}[\mathbf{V}_\mu \mathbf{V}^\mu] + \delta \mathcal{L},$$

$$\delta \mathcal{L}_{d_\chi=2} = a_{WZ} \frac{v^2}{4} \text{Tr}[\mathbf{T} \mathbf{V}_\mu] \text{Tr}[\mathbf{T} \mathbf{V}^\mu],$$

Non-linear Yukawa interactions:

$$\mathcal{L}_Y = \frac{v}{\sqrt{2}} \bar{Q}_L \gamma \mathbf{U} Q_R + \text{h.c.}, \quad Q_R = (u_R, d_R)$$

## Tools

$$\mathcal{Y} \equiv \begin{pmatrix} Y_U & 0 \\ 0 & Y_D \end{pmatrix} = \begin{pmatrix} V^\dagger \mathbf{y}_U & 0 \\ 0 & \mathbf{y}_D \end{pmatrix}$$

Spurion field  $\sim (8, 1, 1)$  for new FCNC effects

$$\lambda_F \equiv Y_U Y_U^\dagger + Y_D Y_D^\dagger = V^\dagger \mathbf{y}_U^2 V + \mathbf{y}_D^2$$

$d_\chi = 4$  ops.

$$\mathcal{O}_1 = \frac{i}{2} \bar{Q}_L \lambda_F \gamma^\mu \{\mathbf{T}, \mathbf{V}_\mu\} Q_L ,$$

$$\mathcal{O}_2 = i \bar{Q}_L \lambda_F \gamma^\mu \mathbf{V}_\mu Q_L ,$$

$$\mathcal{O}_3 = i \bar{Q}_L \lambda_F \gamma^\mu \mathbf{T} \mathbf{V}_\mu \mathbf{T} Q_L ,$$

$$\mathcal{O}_4 = \frac{1}{2} \bar{Q}_L \lambda_F \gamma^\mu [\mathbf{T}, \mathbf{V}_\mu] Q_L .$$

## Effective Low-Energy Lagrangian

$$\begin{aligned} \delta \mathcal{L}_{d_\chi=4} &= -\frac{g}{\sqrt{2}} \left[ W^{\mu+} \bar{U}_L \gamma_\mu (a_W + i a_{CP}) \left( \mathbf{y}_U^2 V + V \mathbf{y}_D^2 \right) D_L + h.c. \right] + \\ &\quad -\frac{g}{2c_W} Z^\mu \left[ a_Z^u \bar{U}_L \gamma_\mu \left( \mathbf{y}_U^2 + V \mathbf{y}_D^2 V^\dagger \right) U_L + a_Z^d \bar{D}_L \gamma_\mu \left( \mathbf{y}_D^2 + V^\dagger \mathbf{y}_U^2 V \right) D_L \right] \end{aligned}$$

$$\begin{aligned} a_Z^u &\equiv a_1 + a_2 + a_3 , & a_Z^d &\equiv a_1 - a_2 - a_3 , \\ a_W &\equiv a_2 - a_3 , & a_{CP} &\equiv -a_4 . \end{aligned}$$

# Non Unitarity and CP Violation

$$\tilde{V}_{ij} = V_{ij} \left[ 1 + (a_W + i a_{CP})(y_{u_i}^2 + y_{d_j}^2) \right]$$

$$\begin{aligned}\sum_k \tilde{V}_{ik}^* \tilde{V}_{jk} &\simeq \delta_{ij} + \left[ 2 a_W y_t^2 + (a_W^2 + a_{CP}^2) y_t^4 \right] \delta_{it} \delta_{jt} \\ \sum_k \tilde{V}_{ki}^* \tilde{V}_{kj} &\simeq \delta_{ij} + \left[ 2 a_W y_t^2 + (a_W^2 + a_{CP}^2) y_t^4 \right] V_{ti}^* V_{tj}\end{aligned}$$

$$\begin{aligned}\arg \left( -\frac{\tilde{V}_{ik}^* \tilde{V}_{il}}{\tilde{V}_{jk}^* \tilde{V}_{jl}} \right) &= \arg \left( -\frac{V_{ik}^* V_{il}}{V_{jk}^* V_{jl}} \right) + a_{CP} [2 a_W (y_{uj}^2 - y_{ui}^2) (y_{dl}^2 - y_{dk}^2) + \\ &- (3 a_W^2 - a_{CP}^2) (y_{uj}^2 - y_{ui}^2) (y_{dl}^2 - y_{dk}^2) (y_{ui}^2 + y_{uj}^2 + y_{dk}^2 + y_{dl}^2)] + O(a^4),\end{aligned}$$

$$\arg \left( -\frac{\tilde{V}_{tb}^* \tilde{V}_{td}}{\tilde{V}_{ub}^* \tilde{V}_{ud}} \right) \simeq \alpha + 2 y_b^2 y_t^2 a_W a_{CP},$$

$$\arg \left( -\frac{\tilde{V}_{cb}^* \tilde{V}_{cd}}{\tilde{V}_{tb}^* \tilde{V}_{td}} \right) \simeq \beta - 2 y_b^2 y_t^2 a_W a_{CP},$$

$$\arg \left( -\frac{\tilde{V}_{ub}^* \tilde{V}_{ud}}{\tilde{V}_{cb}^* \tilde{V}_{cd}} \right) \simeq \gamma - 2 y_c^2 y_b^2 a_W a_{CP} \simeq \gamma.$$

# $\Delta F = 1$ observables

FCNC-effective lagrangian

$$\frac{G_F \alpha}{2\sqrt{2}\pi s_W^2} V_{ti}^* V_{tj} \sum_n C_n \mathcal{Q}_n + \text{h.c.},$$

Wilson coefficient  $C_n$ :

$$C_n = C_n^{SM} + C_n^{NP}$$

FCNC operators basis

$$\begin{aligned} \mathcal{Q}_{\bar{\nu}\nu} &= \bar{d}_i \gamma_\mu (1 - \gamma_5) d_j \bar{\nu} \gamma^\mu (1 - \gamma_5) \nu, & \mathcal{Q}_7 &= e_q \bar{d}_i \gamma_\mu (1 - \gamma_5) d_j \bar{q} \gamma^\mu (1 + \gamma_5) q, \\ \mathcal{Q}_{9V} &= \bar{d}_i \gamma_\mu (1 - \gamma_5) d_j \bar{\ell} \gamma^\mu \ell, & \mathcal{Q}_9 &= e_q \bar{d}_i \gamma_\mu (1 - \gamma_5) d_j \bar{q} \gamma^\mu (1 - \gamma_5) q, \\ \mathcal{Q}_{10A} &= \bar{d}_i \gamma_\mu (1 - \gamma_5) d_j \bar{\ell} \gamma^\mu \gamma_5 \ell, & \mathcal{Q}_{7\gamma} &= \frac{m_i}{g_Z^2} \bar{d}_i (1 - \gamma_5) \sigma_{\mu\nu} d_j (e F^{\mu\nu}), \\ && \mathcal{Q}_{8G} &= \frac{m_i}{g_s^2} \bar{d}_i (1 - \gamma_5) \sigma_{\mu\nu} T^a d_j (g_s G_a^{\mu\nu}). \end{aligned}$$

Leading NP contributions non-linear MFV operators:

$$\begin{aligned} C_{\nu\bar{\nu}}^{NP} &= -\kappa y_t^2 a_Z^d, & C_7^{NP} &= +2\kappa s_W^2 y_t^2 a_Z^d, \\ C_{9V}^{NP} &= \kappa (1 - 4s_W^2) y_t^2 a_Z^d, & C_9^{NP} &= -2\kappa c_W^2 y_t^2 a_Z^d, \\ C_{10A}^{NP} &= -\kappa y_t^2 a_Z^d, & C_{7\gamma}^{NP} &= C_{8G}^{NP} = 0. \end{aligned}$$

$$B^+ \rightarrow \tau^+ \nu$$

$B^+ \rightarrow \tau^+ \nu$ -tree-level charged current process.

$$BR(B^+ \rightarrow \tau^+ \nu) = \frac{G_F^2 m_{B^+} m_\tau^2}{8\pi} \left(1 - \frac{m_\tau^2}{m_{B^+}^2}\right)^2 F_{B^+}^2 |V_{ub}|^2 \left|1 + (a_W + i a_{CP}) y_b^2\right|^2 \tau_{B^+},$$

$F_{B^+}$  is  $B$  decay constant.

# $\Delta F = 2$ observables

$\Delta F = 2$ -effective hamiltonian

$$\mathcal{H}_{\text{eff}}^{\Delta F=2} = \frac{G_F^2 M_W^2}{4\pi^2} C(\mu) Q$$

$Q$  neutral meson mixing operator:

$$Q = (\bar{d}_i^\alpha \gamma_\mu P_L d_j^\alpha)(\bar{d}_i^\beta \gamma^\mu P_L d_j^\beta)$$

Mixing amplitudes  $M_{12}^i$  ( $i = K, d, s$ ):

$$M_{12}^K = \frac{\langle \bar{K}^0 | \mathcal{H}_{\text{eff}}^{\Delta S=2} | K^0 \rangle^*}{2 m_K}, \quad M_{12}^q = \frac{\langle \bar{B}_q^0 | \mathcal{H}_{\text{eff}}^{\Delta B=2} | B_q^0 \rangle^*}{2 m_{B_q}},$$

with  $q = d, s$ . For the  $K$  system,  $M_{12}^K = (M_{12}^K)_{SM} + (M_{12}^K)_{NP}$ . Neglecting all contributions proportional to  $y_{u,d,s}$  and  $y_c^n$  with  $n > 2$ :

$$\begin{aligned} (M_{12}^K)_{SM} &= R_K \left[ \eta_2 \lambda_t^2 S_0(x_t) + \eta_1 \lambda_c^2 S_0(x_c) + 2 \eta_3 \lambda_t \lambda_c S_0(x_c, x_t) \right]^*, \\ (M_{12}^K)_{NP} &= R_K \left[ \eta_2 \lambda_t^2 \left( y_t^2 (2 a_W + y_t^2 a_{CP}^2) G(x_t) + \frac{(4\pi y_t^2 a_Z^d)^2}{g^2} \right) + 2 \eta_1 \lambda_c^2 a_W y_c^2 G(x_c) + \right. \\ &\quad \left. + 2 \eta_3 \lambda_t \lambda_c \left( y_t^2 (2 a_W + a_{CP}^2 y_t^2) H(x_t, x_c) + 2 a_W y_c^2 H(x_c, x_t) \right) \right]^* \end{aligned}$$

$$R_K \equiv \frac{G_F^2 M_W^2}{12\pi^2} F_K^2 m_K \hat{B}_K$$

# $\Delta F = 2$ observables

## Neutral kaon oscillation

$$\Delta M_K = 2 \left[ \text{Re}(M_{12}^K)_{SM} + \text{Re}(M_{12}^K)_{NP} \right],$$

$$\varepsilon_K = \frac{\kappa_\epsilon e^{i\varphi_\epsilon}}{\sqrt{2}(\Delta M_K)_{\text{exp}}} \left[ \text{Im} \left( M_{12}^K \right)_{SM} + \text{Im} \left( M_{12}^K \right)_{NP} \right]$$

## Neutral meson oscillation

$$M_{12}^q = (M_{12}^q)_{SM} \mathcal{C}_{Bq} e^{2i\varphi_{Bq}},$$

NP effects from  $\mathcal{C}_{B_{d,s}}$  and  $\varphi_{B_{d,s}}$

$$M_{12}^q = R_{Bq} \left[ \lambda_t^2 S_0(x_t) \right]^*, \quad \text{with} \quad R_{Bq} \equiv \frac{G_F^2 M_W^2}{12\pi^2} F_{Bq}^2 m_{Bq} \hat{B}_{Bq} \eta_B,$$

$B_{d,s}$ -mass differences

$$\Delta M_q = 2 |M_{12}^q| \equiv (\Delta M_q)_{SM} \mathcal{C}_{Bq},$$

$$\mathcal{C}_{B_d} = \mathcal{C}_{B_s} = \left| 1 + 2 a_W \left( y_t^2 \frac{G(x_t)}{S_0(x_t)} + y_b^2 \right) + \frac{(4\pi y_t^2 a_Z^d)^2}{g^2 S_0(x_t)} + 2 i a_W a_{CP} y_t^2 y_b^2 \frac{G(x_t)}{S_0(x_t)} \right|.$$

# $\Delta F = 2$ observables

## Neutral meson oscillation

Mixing-induced CP asymmetries  $S_{\psi K_S}$  &  $S_{\psi \phi}$  for  $B_d^0 \rightarrow \psi K_S$  &  $B_s^0 \rightarrow \psi \phi$

$$S_{\psi K_S} = \sin(2\beta + 2\varphi_{B_d}), \quad S_{\psi \phi} = \sin(2\beta_s - 2\varphi_{B_s}),$$

UT-angles  $\beta$  &  $\beta_s$

$$\beta \equiv \arg \left( -\frac{V_{cb}^* V_{cd}}{V_{tb}^* V_{td}} \right), \quad \beta_s \equiv \arg \left( -\frac{V_{tb}^* V_{ts}}{V_{cb}^* V_{cs}} \right),$$

New phases

$$\varphi_{B_d} = \varphi_{B_s} = 2 a_W a_{CP} y_t^2 y_b^2 \frac{G(x_t)}{S_0(x_t)}.$$

## $R_{BR/\Delta M}$

$$R_{BR/\Delta M} = \frac{3\pi \tau_{B+}}{4\eta_B \hat{B}_{B_d} S_0(x_t) M_W^2} \frac{m_\tau^2}{|V_{tb}^* V_{td}|^2} \left( 1 - \frac{m_\tau^2}{m_{B_d}^2} \right)^2 \frac{|1 + (a_W + i a_{CP}) y_b^2|^2}{C_{B_d}}$$

# $\Delta F = 2$ observables

## $B$ -semileptonic CP-Asymmetry

$$A_{sl}^b \equiv \frac{N_b^{++} - N_b^{--}}{N_b^{++} + N_b^{--}} ,$$

$N_b^{++}$  &  $N_b^{--}$  number of events containing two  $\mu^+$  or  $\mu^-$ . In  $p\bar{p}$  colliders, such events can only arise through  $B_d^0 - \bar{B}_d^0$  or  $B_s^0 - \bar{B}_s^0$  mixings.

$$A_{sl}^b = (0.594 \pm 0.022) a_{sl}^d + (0.406 \pm 0.022) a_{sl}^s ,$$

$$a_{sl}^d \equiv \left| \frac{(\Gamma_{12}^d)_{SM}}{(M_{12}^d)_{SM}} \right| \sin \phi_d = (5.4 \pm 1.0) \times 10^{-3} \sin \phi_d ,$$

$$a_{sl}^s \equiv \left| \frac{(\Gamma_{12}^s)_{SM}}{(M_{12}^s)_{SM}} \right| \sin \phi_s = (5.0 \pm 1.1) \times 10^{-3} \sin \phi_s ,$$

$$\phi_d \equiv \arg \left( - (M_{12}^d)_{SM} / (\Gamma_{12}^d)_{SM} \right) = -4.3^\circ \pm 1.4^\circ ,$$

$$\phi_s \equiv \arg \left( - (M_{12}^s)_{SM} / (\Gamma_{12}^s)_{SM} \right) = 0.22^\circ \pm 0.06^\circ .$$

NP contributions

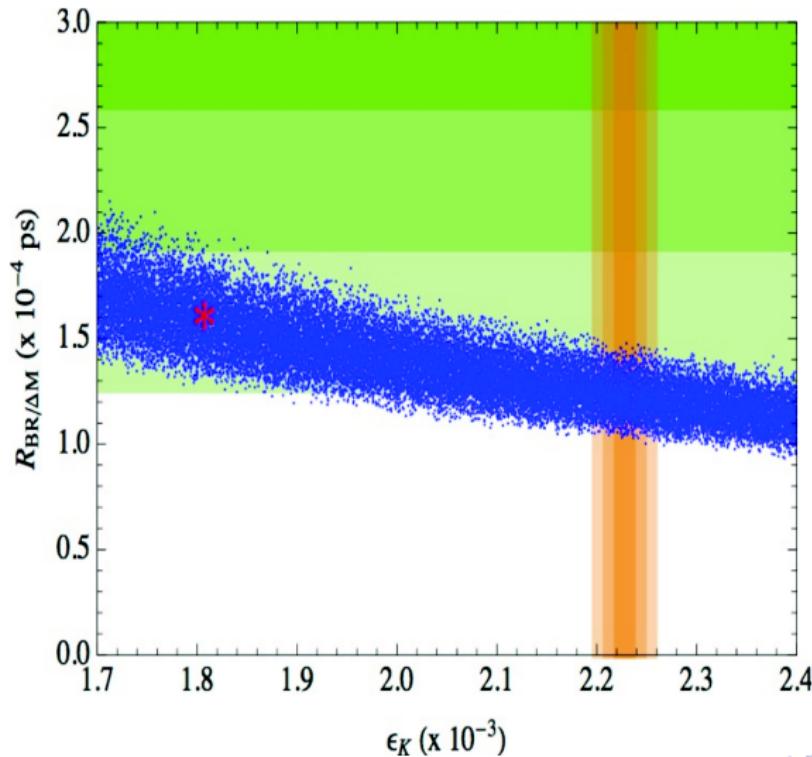
$$\Gamma_{12}^q = (\Gamma_{12}^q)_{SM} \tilde{C}_{Bq} \quad \text{with} \quad \tilde{C}_{Bq} = 1 + 2 a_W y_b^2 ,$$

$$a_{sl}^q = \left| \frac{(\Gamma_{12}^q)_{SM}}{(M_{12}^q)_{SM}} \right| \frac{\tilde{C}_{Bq}}{C_{Bq}} \sin \left( \phi_q + 2\varphi_{Bq} \right) ,$$

## Impact on the observables of specific parameter values

Parameter	$\delta\varepsilon_K$	$\delta R_{BR/\Delta M}$	$\delta A_{sl}^b$
$a_{CP} = 0.1(-0.1)$	$\approx 1.1\%$	$\approx -1.4\%$	$\approx 1.1\%(-1.6\%)$
$a_W = 0.1(-0.1)$	$\approx +26\%(-19\%)$	$\approx -25\%(+30\%)$	$\approx +33\%(-23\%)$
$a_Z^d = \pm 0.1$	$\approx 124\%$	$\approx -62\%$	$\approx 160\%$

$\varepsilon_K$  vs.  $R_{BR/\Delta M}$  from  $\mathcal{O}_{1,2,3}$



- correlation
  - $\varepsilon_K - R_{BR/\Delta M}$
  - \*: SM values

Green & Orange:  
 $1, 2, 3\sigma$  exp. values

$a_W \in [-1, 1]$ ,  
 $a_Z^d \in [-0.1, 0.1]$

and  $a_{CP} = 0$

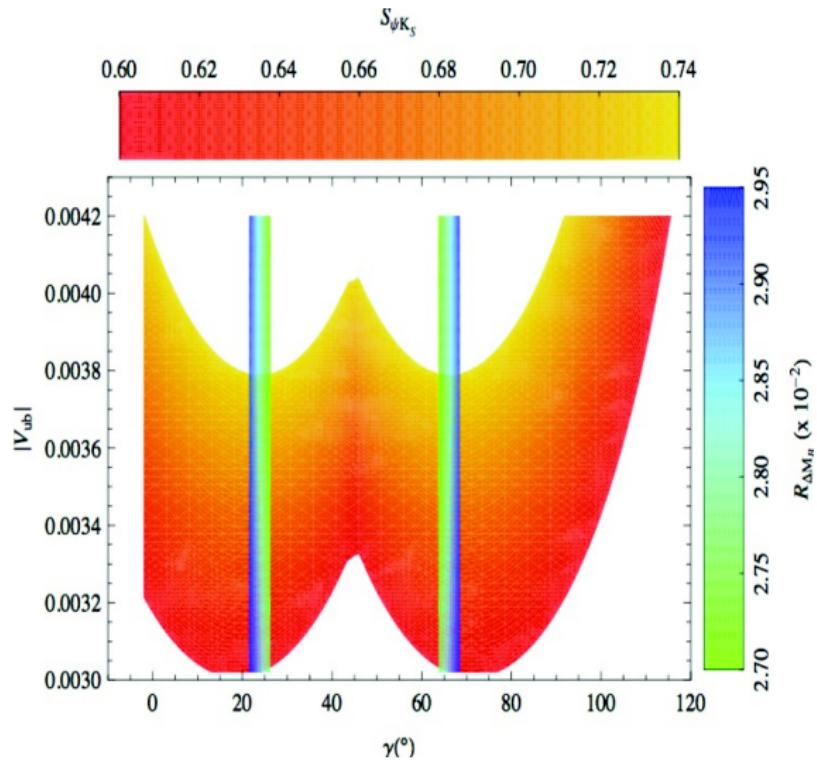
# Phenomenological Analysis

## Input parameters and the SM analysis

$G_F = 1.16637(1) \times 10^{-5}$ GeV $^{-2}$	[30]	$m_{B_d} = 5279.5(3)$ MeV	[30]
$M_W = 80.399(23)$ GeV	[30]	$m_{B_s} = 5366.3(6)$ MeV	[30]
$s_W^2 \equiv \sin^2 \theta_W = 0.23116(13)$	[30]	$F_{B_d} = 205(12)$ MeV	[32]
$\alpha(M_Z) = 1/127.9$	[30]	$F_{B_s} = 250(12)$ MeV	[32]
$\alpha_s(M_Z) = 0.1184(7)$	[30]	$\hat{B}_{B_d} = 1.26(11)$	[32]
$m_u(2\text{ GeV}) = 1.7 \div 3.1$ MeV	[30]	$\hat{B}_{B_s} = 1.33(6)$	[32]
$m_d(2\text{ GeV}) = 4.1 \div 5.7$ MeV	[30]	$F_{B_d} \sqrt{\hat{B}_{B_d}} = 233(14)$ MeV	[32]
$m_s(2\text{ GeV}) = 100^{+30}_{-20}$ MeV	[30]	$F_{B_s} \sqrt{\hat{B}_{B_s}} = 288(15)$ MeV	[32]
$m_c(m_c) = (1.279 \pm 0.013)$ GeV	[36]	$\xi = 1.237(32)$	[32]
$m_b(m_b) = 4.19^{+0.18}_{-0.06}$ GeV	[30]	$\eta_B = 0.55(1)$	[37, 38]
$M_t = 172.9 \pm 0.6 \pm 0.9$ GeV	[30]	$\Delta M_d = 0.507(4)$ ps $^{-1}$	[30]
$m_K = 497.614(24)$ MeV	[30]	$\Delta M_s = 17.77(12)$ ps $^{-1}$	[30]
$F_K = 156.0(11)$ MeV	[32]	$\sin(2\beta)_{b \rightarrow c\bar{s}} = 0.673(23)$	[30]
$\hat{B}_K = 0.737(20)$	[32]	$\phi_s^{\psi\phi} = 0.55^{+0.38}_{-0.36}$	[39, 40]
$\kappa_\epsilon = 0.923(6)$	[41]	$\phi_s^{\psi\phi} = 0.03 \pm 0.16 \pm 0.07$	[42]
$\varphi_\epsilon = (43.51 \pm 0.05)^\circ$	[43]	$R_{\Delta M_B} = (2.85 \pm 0.03) \times 10^{-2}$	[30]
$\eta_1 = 1.87(76)$	[44]	$A_{sl}^b = (-0.787 \pm 0.172 \pm 0.093) \times 10^{-2}$	[34]
$\eta_2 = 0.5765(65)$	[37]	$ V_{us}  = 0.2252(9)$	[30]
$\eta_3 = 0.496(47)$	[45]	$ V_{cb}  = (40.6 \pm 1.3) \times 10^{-3}$	[30]
$\Delta M_K = 0.5292(9) \times 10^{-2}$ ps $^{-1}$	[30]	$ V_{ub}^{\text{incl.}}  = (4.27 \pm 0.38) \times 10^{-3}$	[30]
$ \epsilon_K  = 2.228(11) \times 10^{-3}$	[30]	$ V_{ub}^{\text{excl.}}  = (3.38 \pm 0.36) \times 10^{-3}$	[30]
$\tau_{B^\pm} = (1641 \pm 8) \times 10^{-3}$ ps	[30]	$ V_{ub}^{\text{comb.}}  = (3.89 \pm 0.44) \times 10^{-3}$	[30]
$BR(B^+ \rightarrow \tau^+ \nu) = (1.65 \pm 0.34) \times 10^{-4}$	[30]	$\gamma = (73^{+22}_{-25})^\circ$	[30]

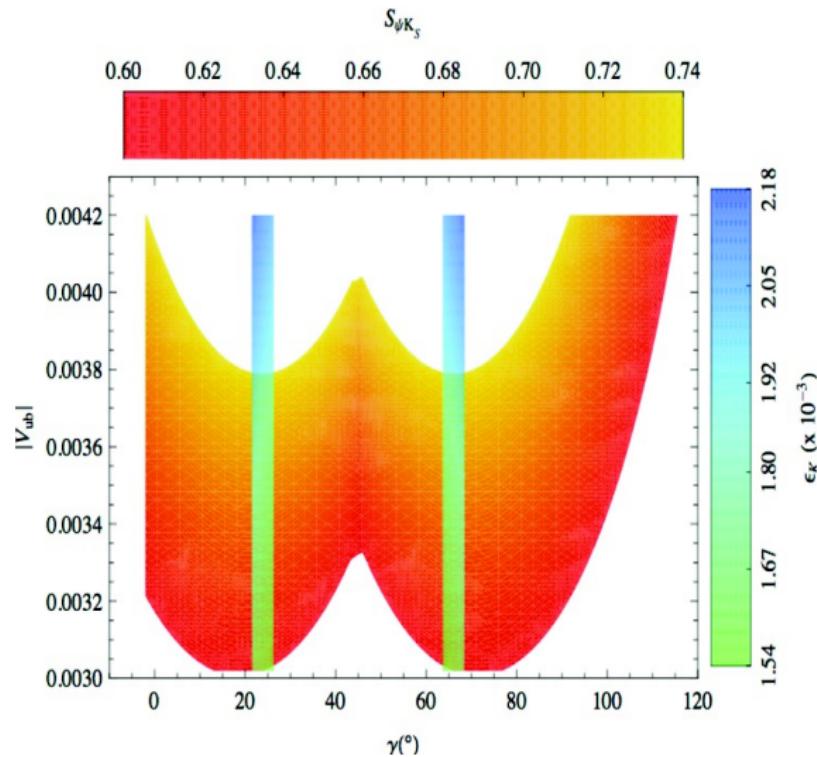
# Phenomenological Analysis

$|V_{ub}| - \gamma$  parameter space



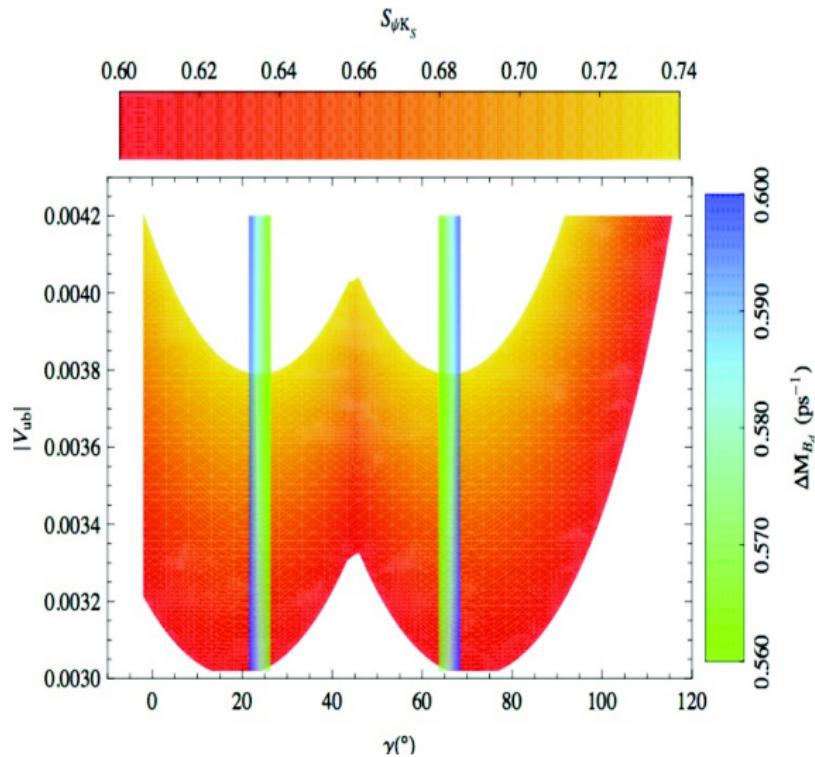
# Phenomenological Analysis

εK



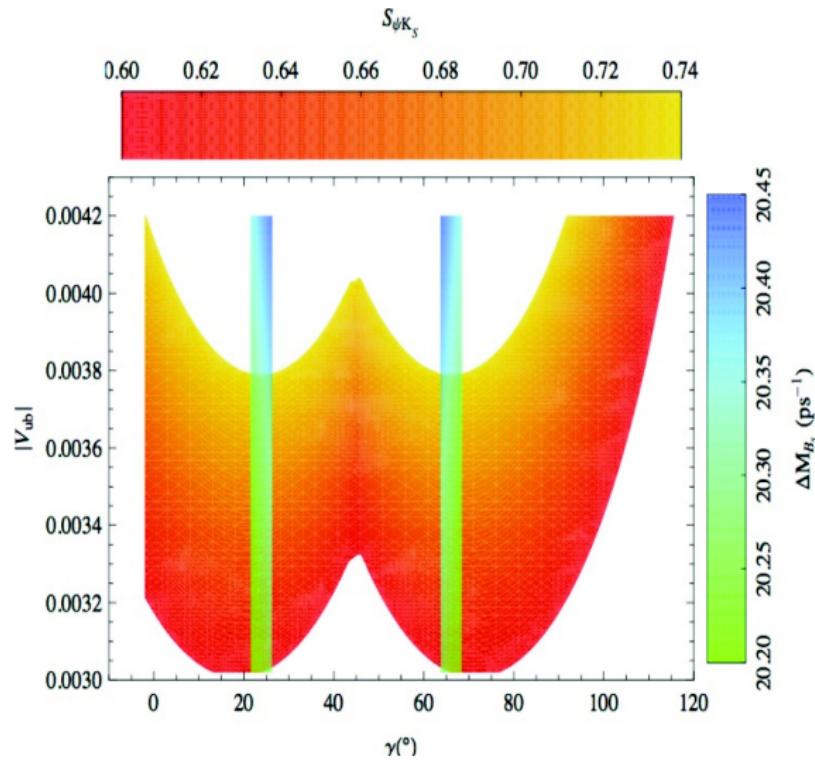
# Phenomenological Analysis

$\Delta M_{B_d}$



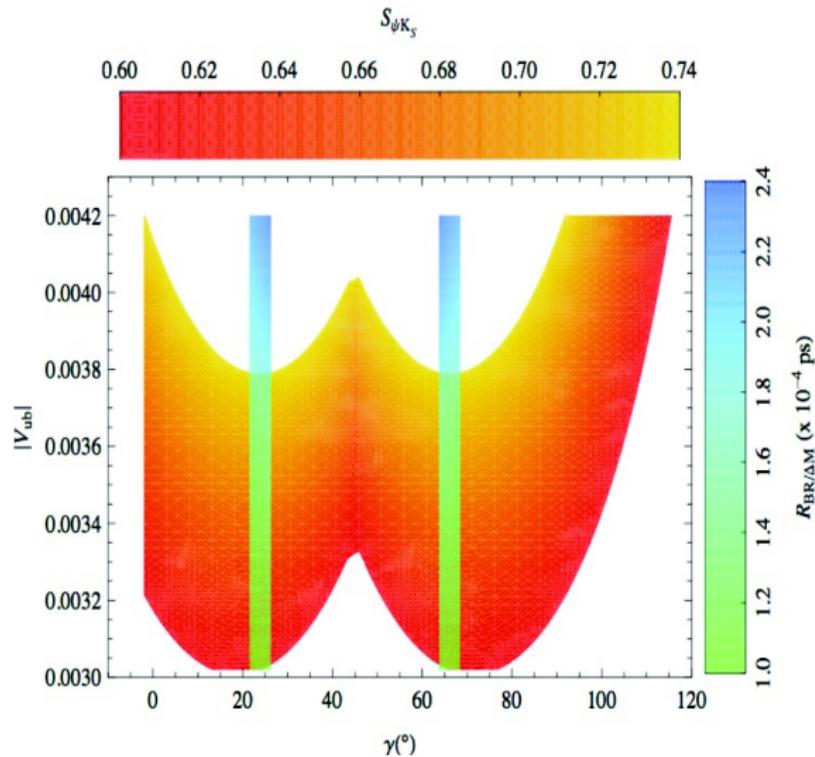
# Phenomenological Analysis

$$\Delta M_{B_s}$$



# Phenomenological Analysis

$$R_{BR/\Delta M}$$



# Phenomenological Analysis

In order to illustrate the features of the MFV scenario with a strong interacting Higgs sector, the numerical analysis of the following sections will be presented choosing as reference point,  
 $(|V_{ub}|, \gamma) = (3.5 \times 10^{-3}, 66^\circ)$ , corresponding to  $S_{\psi K_S} \simeq 0.692$  and  $R_{\Delta M_B} \simeq 2.83 \times 10^{-2}$ . For this point

$$\varepsilon_K = 1.8 \times 10^{-3}, \quad R_{BR/\Delta M} = 1.6 \times 10^{-4} \text{ ps}.$$

$$S_{\psi\phi} = 0.036.$$

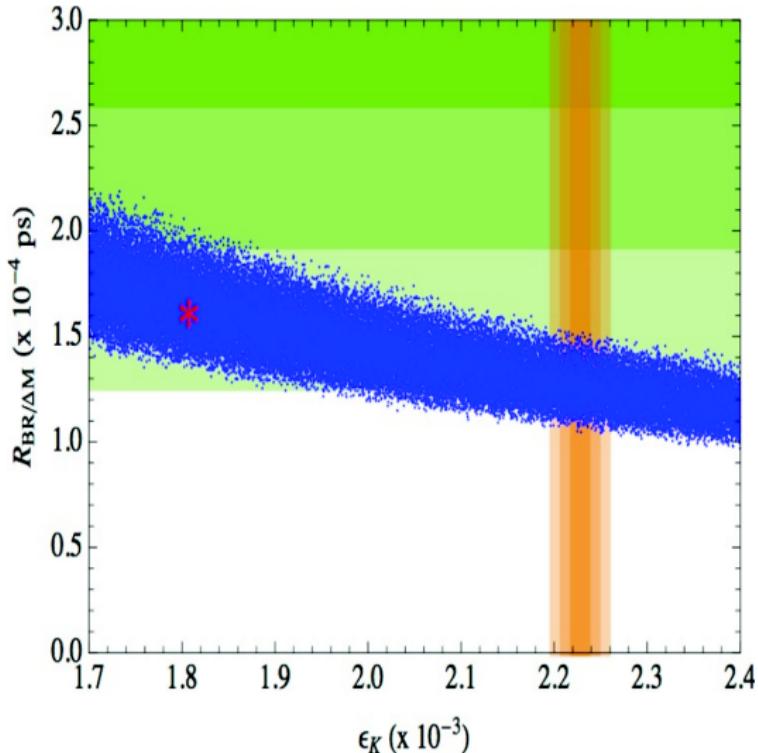
$$A_{sl}^b = -2.3 \times 10^{-4} \quad (a_{sl}^d = -4.0 \times 10^{-4}, \quad a_{sl}^s = 1.9 \times 10^{-5}),$$

$$a_{CP} = 0.1(-0.1) \longrightarrow \delta A_{sl}^b \approx 1.1\%(1.6\%)$$

$$a_W = 0.1(-0.1) \longrightarrow \delta A_{sl}^b \approx 33\%(-23\%)$$

$$a_Z^d = \pm 0.1 \longrightarrow \delta A_{sl}^b \approx 160\%.$$

$\epsilon_K$  vs.  $R_{BR/\Delta M}$  from all  $O_i$



●: correlation

$\epsilon_K - R_{BR/\Delta M}$

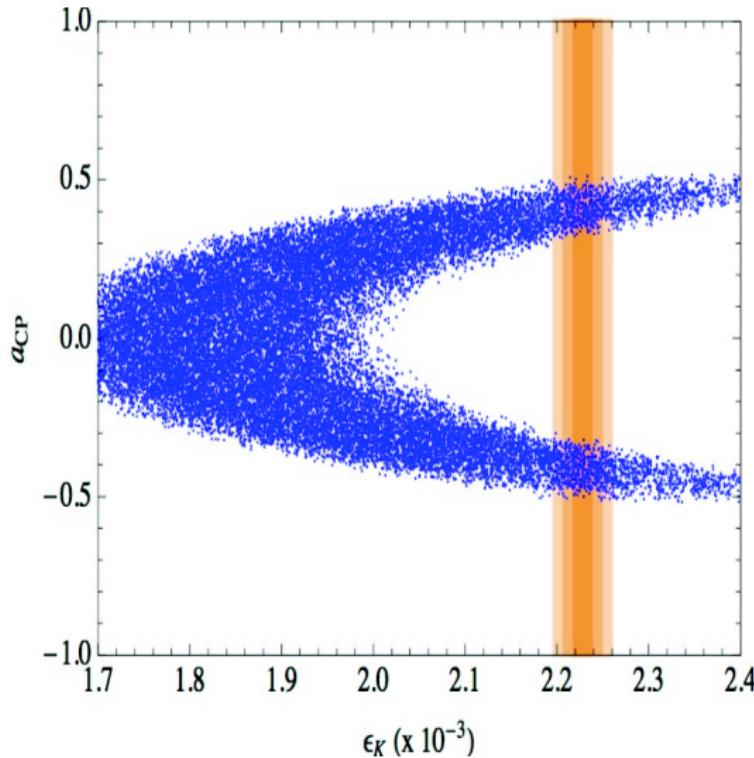
\*: SM values

Green & Orange: 1, 2, 3 $\sigma$   
exp. values

$a_W, a_{CP} \in [-1, 1]$ ,  
 $a_Z^d \in [-0.1, 0.1]$

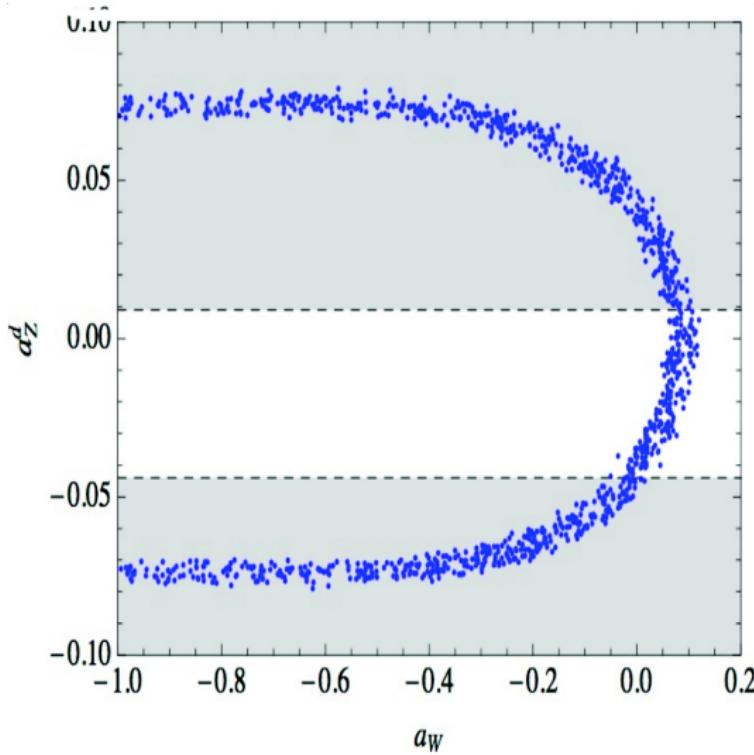
From  $\mathcal{O}_4$

$a_{CP}$  from  $\mathcal{O}_4$



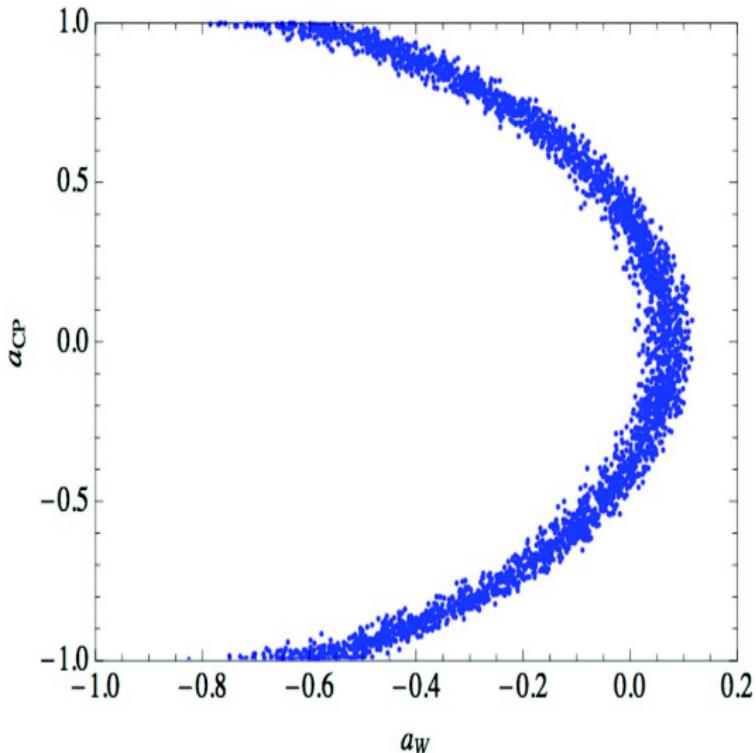
From  $\mathcal{O}_{1,2,3}$

$a_W - a_Z^d$  from  $\mathcal{O}_{1,2,3}$



*Obs* inside its  $3\sigma$  exp. value

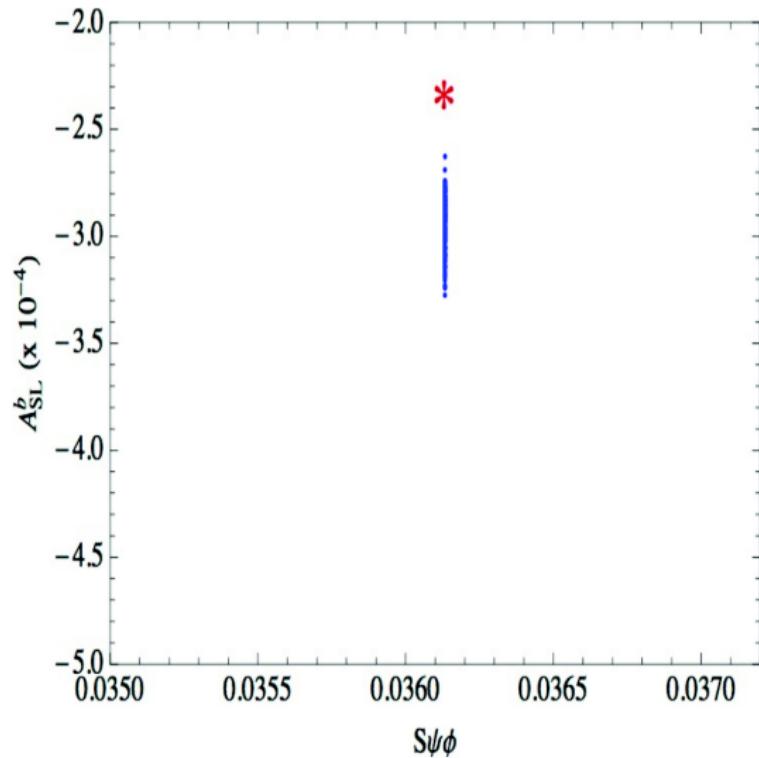
From all  $\mathcal{O}_i$



$a_W - a_{CP}$  and  
 $a_Z^d \in [-0.044, 0.009]$   
from  $\mathcal{O}_i$

*Obs* inside its  $3\sigma$  exp.  
value

## $S_{\psi\phi}$ vs. $A_{sI}^b$ from $\mathcal{O}_{1,2,3}$



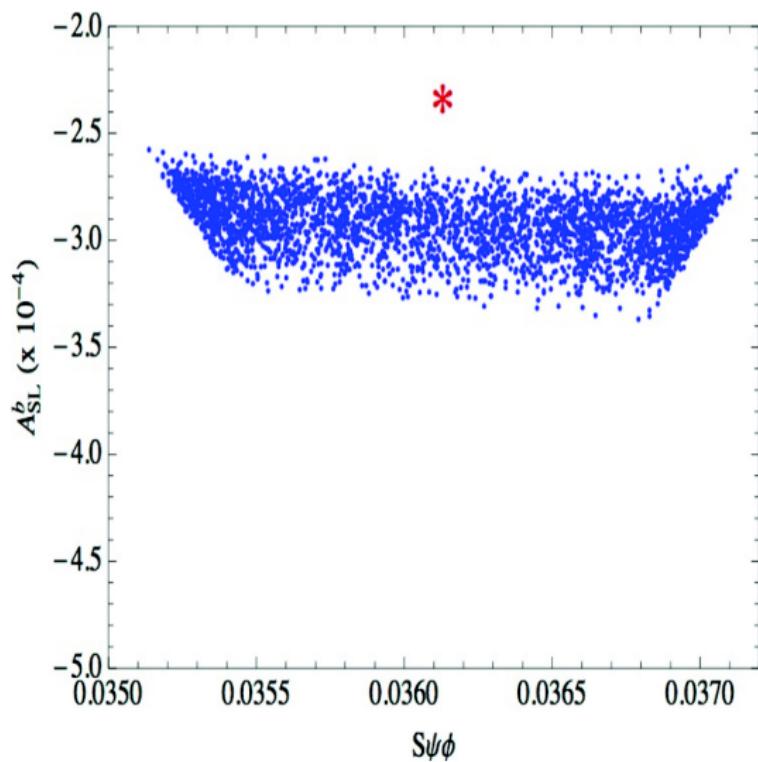
•: correlation  $S_{\psi\phi} - A_{sI}^b$

\*: SM values

$a_W \in [-1, 1]$ ,  $a_{CP} = 0$ ,

$a_Z^d \in [-0.044, 0.009]$

## $S_{\psi\phi}$ vs. $A_{sl}^b$ from all $\mathcal{O}_i$



●: correlation  $S_{\psi\phi} - A_{sl}^b$

\*: SM values

$a_W, a_{CP} \in [-1, 1]$ ,  
 $a_Z^d \in [-0.044, 0.009]$