

Adi and iso according to CMB and LSS

Vesa Muhonen
Helsinki Institute of Physics

In collaboration with J. Väliiviita, H. Kurki-Suonio and R. Keskitalo

GGI, Firenze, 23.10.2006

What?

We are considering events well after any process that generated the primordial perturbations.

- e.g., well after the end of inflation

In general there are the curvature perturbations \mathcal{R} and the entropy perturbations \mathcal{S} (can be several kinds of).

A general perturbation can be then divided into an adiabatic and an isocurvature mode

- adi: initially $\mathcal{R} \neq 0$ and $\mathcal{S} = 0$
- iso: initially $\mathcal{R} = 0$ and $\mathcal{S} \neq 0$

Adi and iso can be correlated since entropy perturbations can source curvature perturbations even on superhorizon scales.

Why?

We know that a simple adiabatic model is a very good fit to the data.

- Is isocurvature better constrained by the WMAP 3-year data?
 - revisit our earlier results
- How much isocurvature does the data allow?
 - it's not difficult to produce isocurvature
 - e.g., multi-field inflation

How?

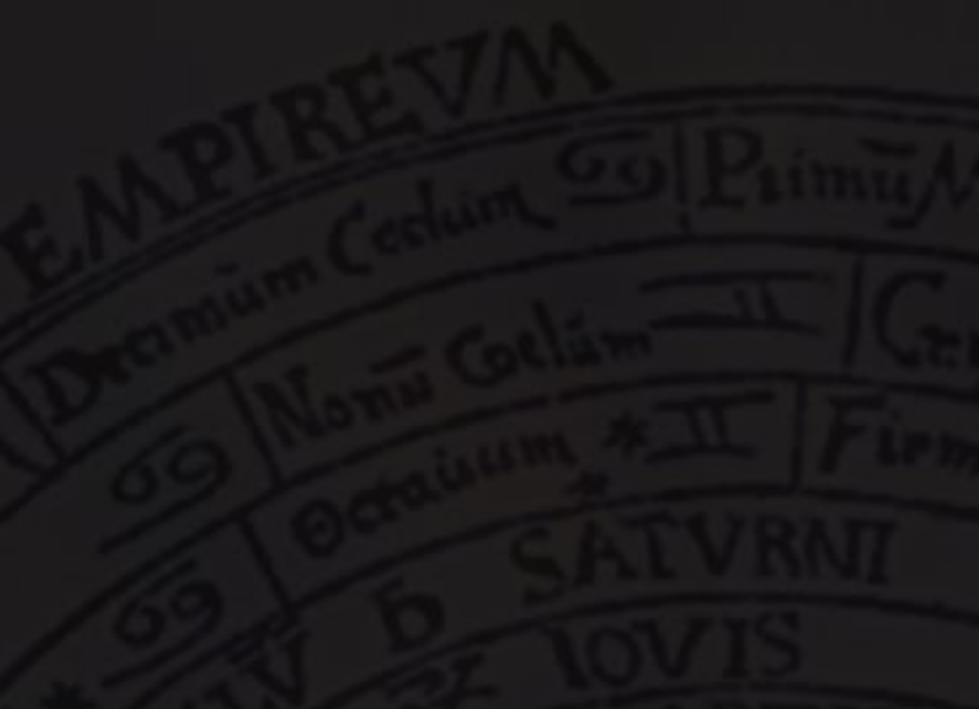
- We consider a spatially flat universe
 - dark energy is the cosmological constant
 - CDM isocurvature
- We use the CMB data from WMAP-3 with additional small scale data and LSS data from SDSS
- The total C_l is a sum of four components

$$\begin{aligned} C_l &= A^2 [(1 - \alpha)(1 - |\gamma|)\hat{C}_l^{\text{ad1}} + (1 - \alpha)|\gamma|\hat{C}_l^{\text{ad2}} \\ &\quad + \alpha\hat{C}_l^{\text{iso}} + \text{sign}(\gamma)\sqrt{\alpha(1 - \alpha)|\gamma|}\hat{C}_l^{\text{cor}}] \\ &\equiv C_l^{\text{ad1}} + C_l^{\text{ad2}} + C_l^{\text{iso}} + C_l^{\text{cor}}, \end{aligned}$$

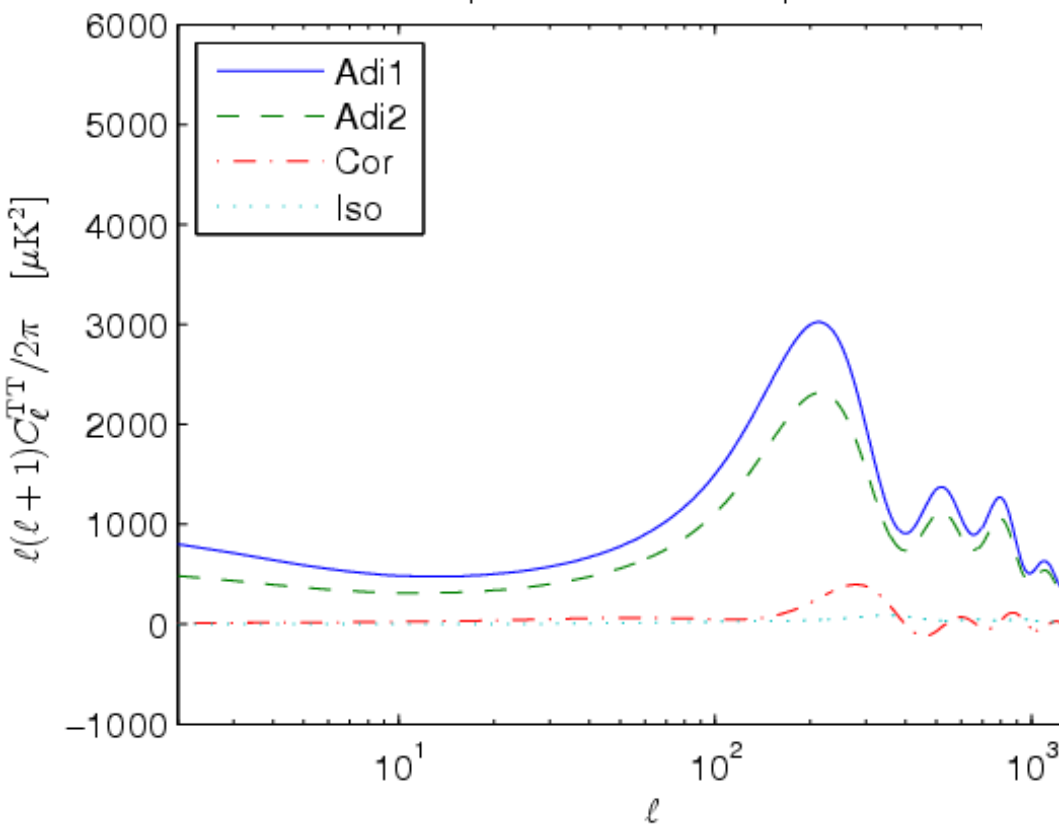
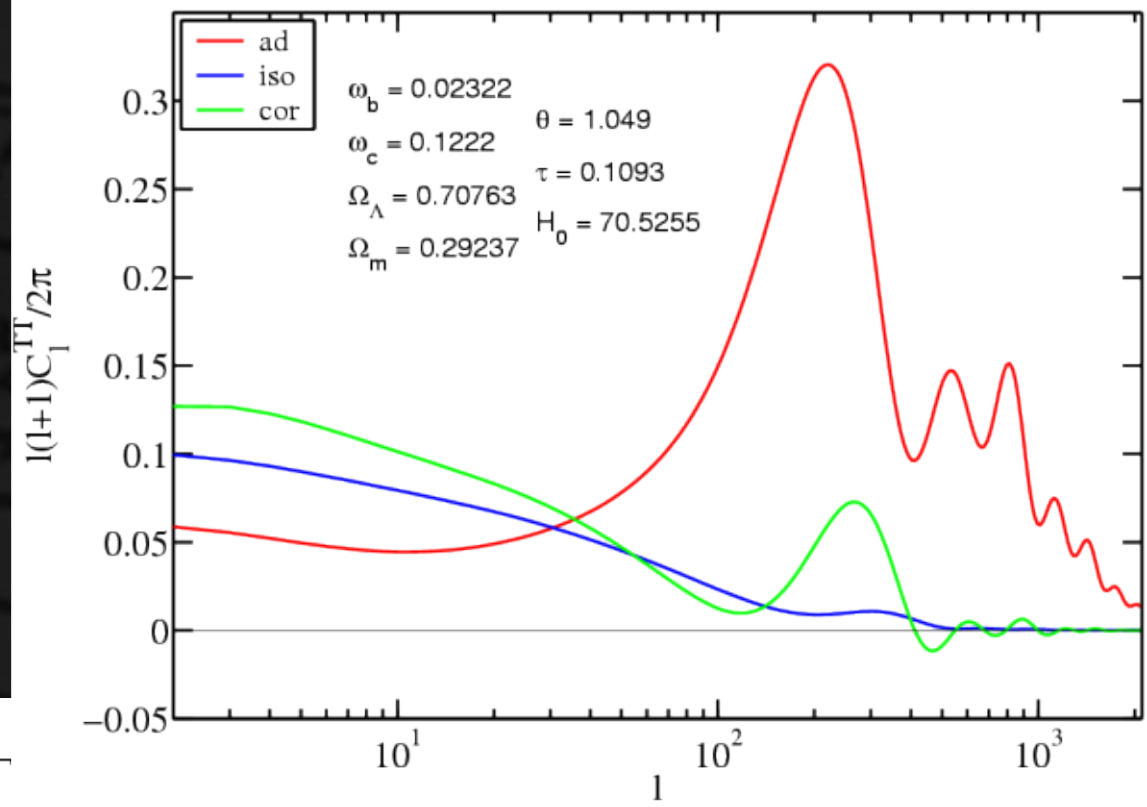
- In total there are 11 parameters

$$\omega_b, \omega_c, \theta, \tau, b, n_{\text{ad1}}, n_{\text{ad2}}, n_{\text{iso}}, \ln(10^{10} A^2), \alpha, \gamma$$

- Then we do a normal MCMC analysis



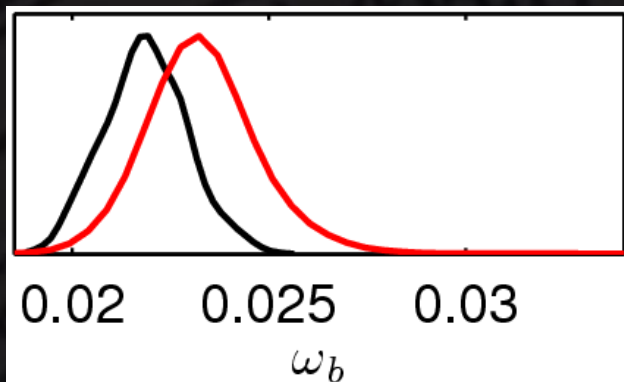
Components of the TT spectrum



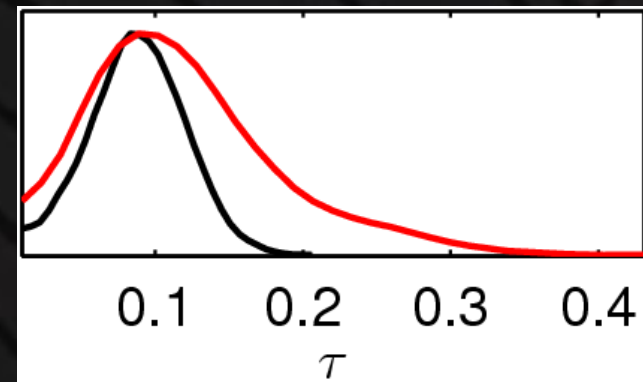
What do we find?

The isocurvature model is a slightly better fit to the data

- in terms of χ^2 the improvement is $\Delta\chi^2 \sim 10$

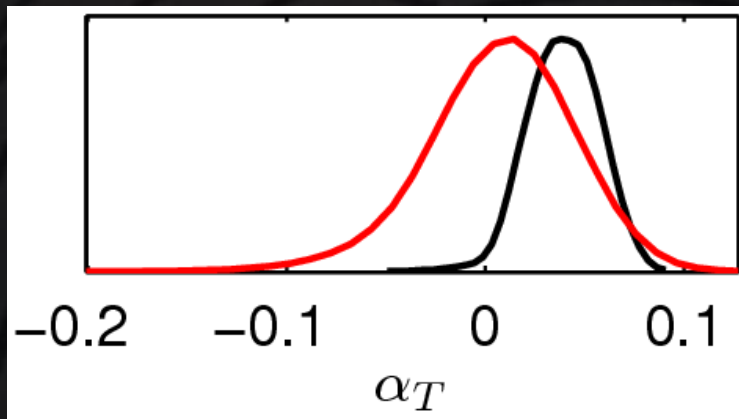
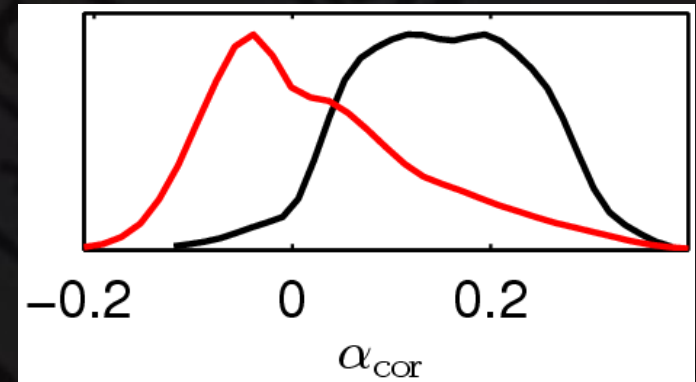
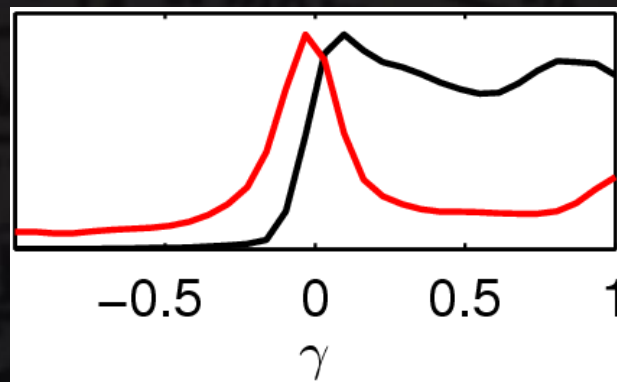
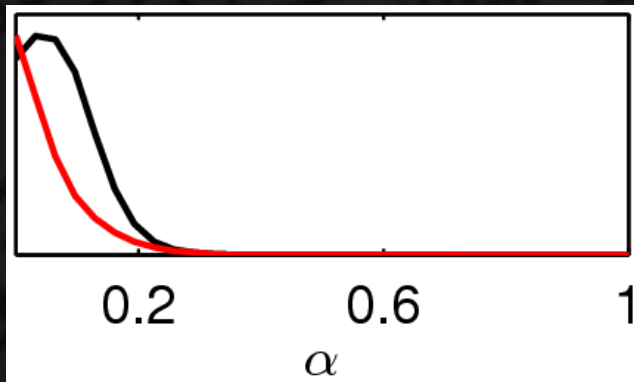


Slightly better fit
to BBN values.



Better constrained due to
WMAP polarization data.

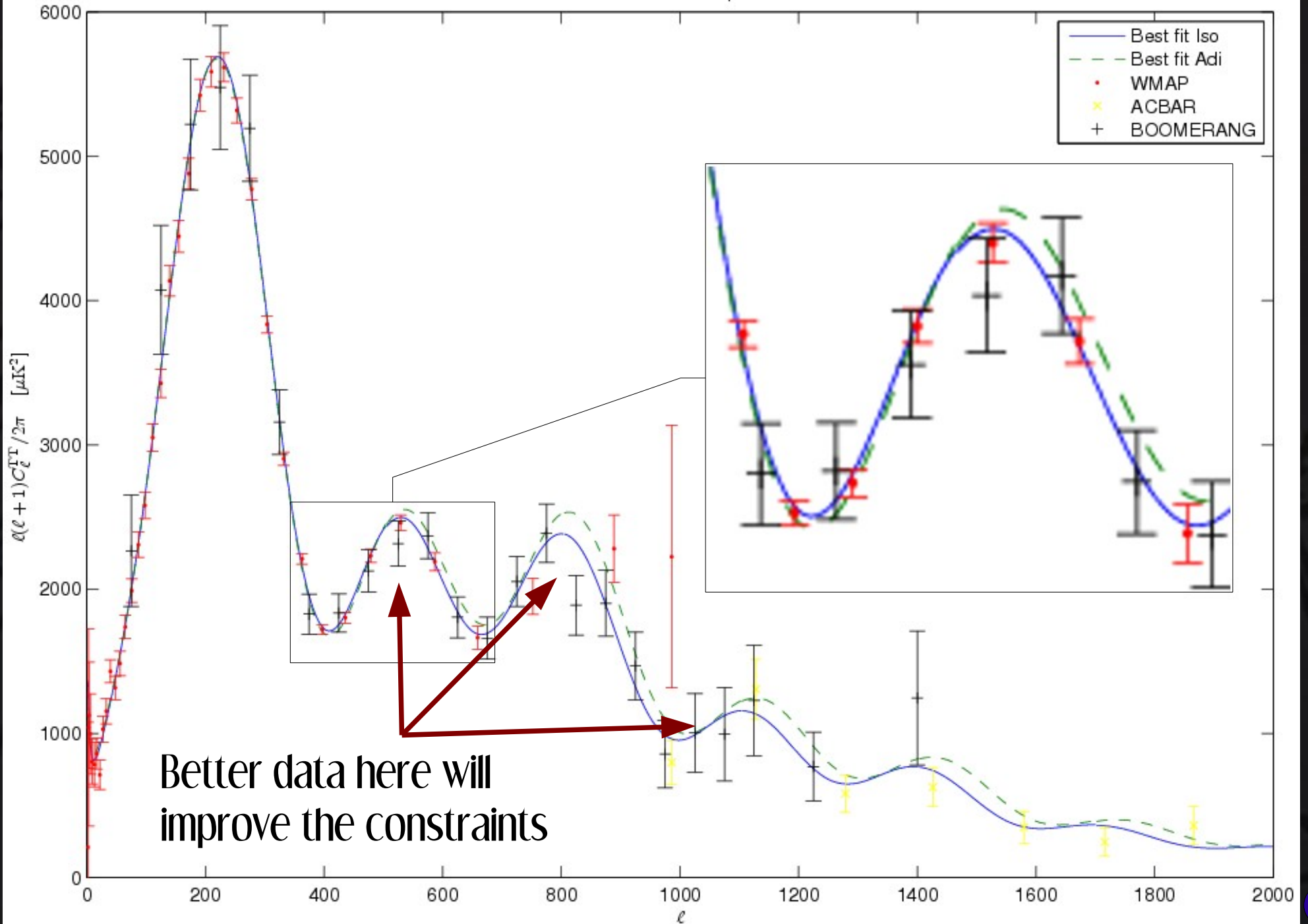
Surprisingly, the WMAP 3-year data does not lead to tighter constraints on the isocurvature parameters.



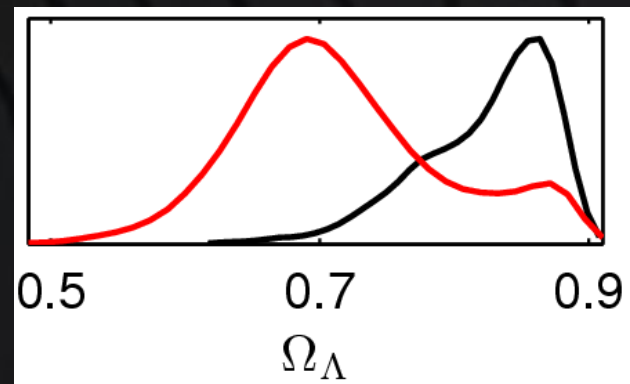
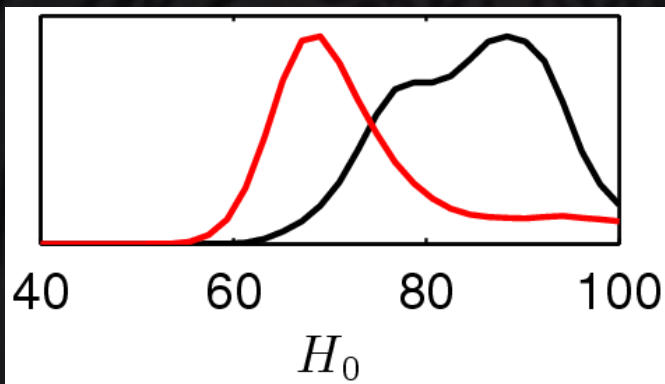
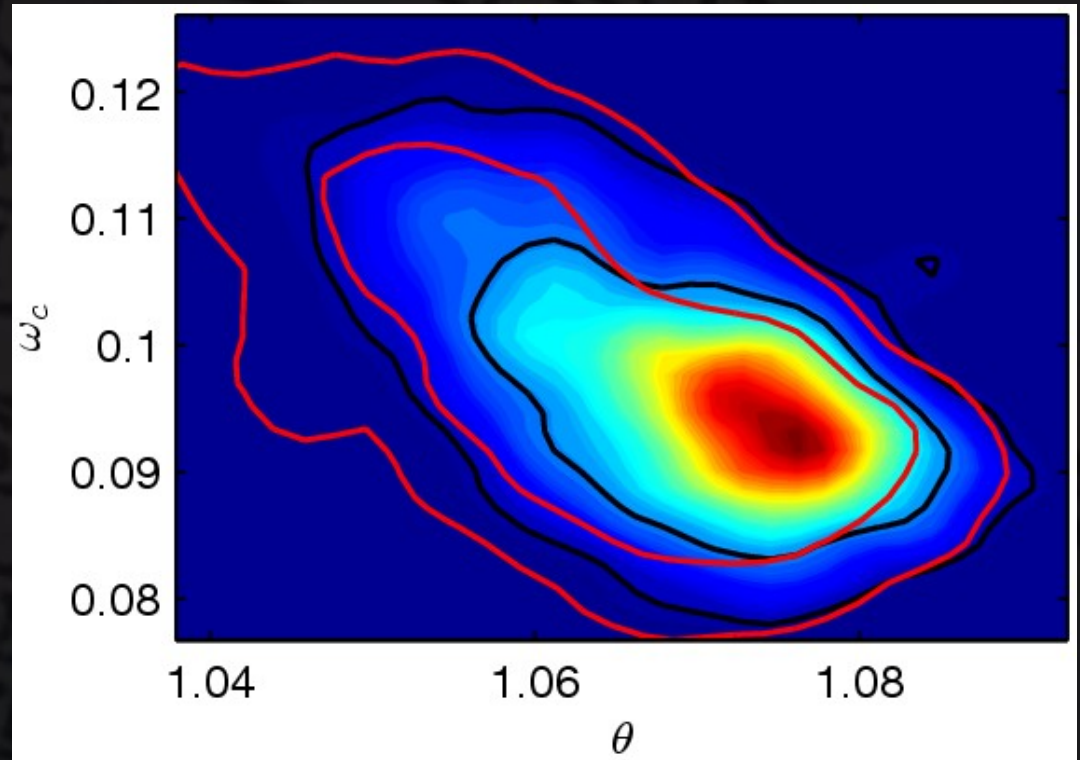
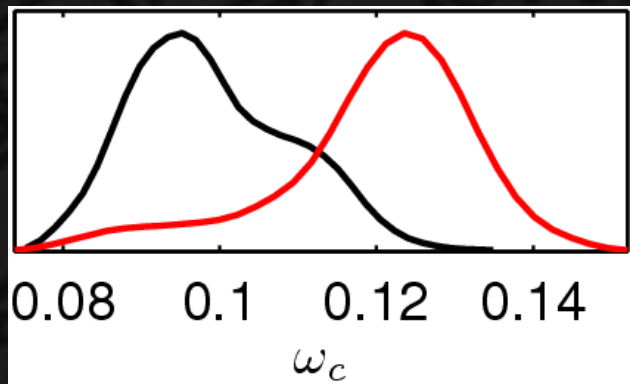
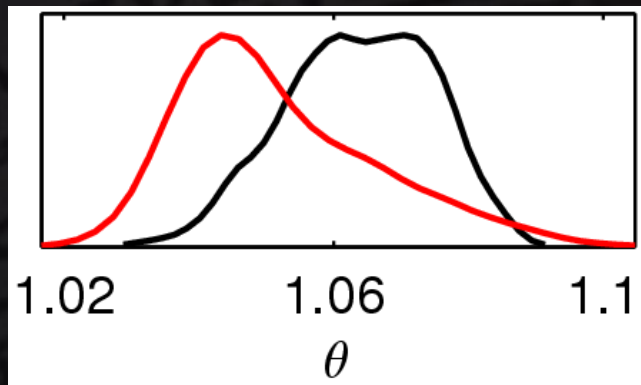
$$\begin{aligned} \alpha_T &= \frac{\langle (\delta T^{\text{non-ad}})^2 \rangle}{\langle (\delta T^{\text{total}})^2 \rangle} \\ &= \frac{\sum_{\ell} (2\ell + 1) (C_{\ell}^{\text{TTiso}} + C_{\ell}^{\text{TTcor}})}{\sum_{\ell} (2\ell + 1) C_{\ell}^{\text{TT}}} \end{aligned}$$

A non-adiabatic contribution $\sim 5\%$ is allowed by the data.

The combined TT spectrum



There are some effects, however,
on the other parameters...



Conclusions

- The CMB is dominantly adiabatic, but a small isocurvature component is clearly allowed
 - this is true even with the latest more accurate data
 - there might be a small feature that can be explained with iso
 - more accurate data on the 2nd and 3rd peak will give further constraints
- In the observed CMB spectra there can be $\sim 5\%$ non-adiabatic contribution

To calculate the CMB power spectra, one needs the curvature and entropy perturbations given deep in the radiation dominated era.

$$\begin{bmatrix} \mathcal{R}(\mathbf{k}) \\ \mathcal{S}(\mathbf{k}) \end{bmatrix}_{\text{rad}} = \begin{bmatrix} 1 & T_{\mathcal{R}\mathcal{S}}(k) \\ 0 & T_{\mathcal{S}\mathcal{S}}(k) \end{bmatrix} \begin{bmatrix} \mathcal{R}(\mathbf{k}) \\ \mathcal{S}(\mathbf{k}) \end{bmatrix}_*$$

A correlation between two random variables is given by:

$$\langle x(\mathbf{k})y^*(\mathbf{k}') \rangle = \frac{2\pi^2}{k^3} \mathcal{C}_{xy}(k) \delta^{(3)}(\mathbf{k} - \mathbf{k}')$$

$$\mathcal{C}_{\mathcal{R}\mathcal{R}}(t_{\text{rad}}, k) = \mathcal{P}_{\mathcal{R}}(t_*, k) + T_{\mathcal{R}\mathcal{S}}(k)^2 \mathcal{P}_{\mathcal{S}}(t_*, k)$$

$$\mathcal{C}_{\mathcal{R}\mathcal{S}}(t_{\text{rad}}, k) = T_{\mathcal{R}\mathcal{S}}(k) T_{\mathcal{S}\mathcal{S}}(k) \mathcal{P}_{\mathcal{S}}(t_*, k)$$

$$\mathcal{C}_{\mathcal{S}\mathcal{S}}(t_{\text{rad}}, k) = T_{\mathcal{S}\mathcal{S}}(k)^2 \mathcal{P}_{\mathcal{S}}(t_*, k),$$

Approximating the power spectra and the transfer functions by power laws leads to:

$$\mathcal{P}_{\mathcal{R}}(k) \equiv \mathcal{C}_{\mathcal{R}\mathcal{R}}(k) = A_r^2 \hat{k}^{n_{\text{ad}1}-1} + A_s^2 \hat{k}^{n_{\text{ad}2}-1}$$

$$\mathcal{P}_{\mathcal{S}}(k) \equiv \mathcal{C}_{\mathcal{S}\mathcal{S}}(k) = B^2 \hat{k}^{n_{\text{iso}}-1}$$

$$\mathcal{C}_{\mathcal{R}\mathcal{S}}(k) = \mathcal{C}_{\mathcal{S}\mathcal{R}}(k) = A_s B \hat{k}^{n_{\text{cor}}-1}$$

$$n_{\text{cor}} = \frac{1}{2}(n_{\text{iso}} + n_{\text{ad}2})$$

The total CMB angular power spectrum is now:

$$C_{\ell}^{ab} = 4\pi \sum_{xy} \int \frac{dk}{k} \mathcal{C}_{xy}(k) g_{x\ell}^a(k) g_{y\ell}^b(k)$$