

Experimental review on neutrinoless double beta decay

GGI Neutrino and Invisibles meeting
June 25, 2012

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**University of
Zurich^{UZH}**

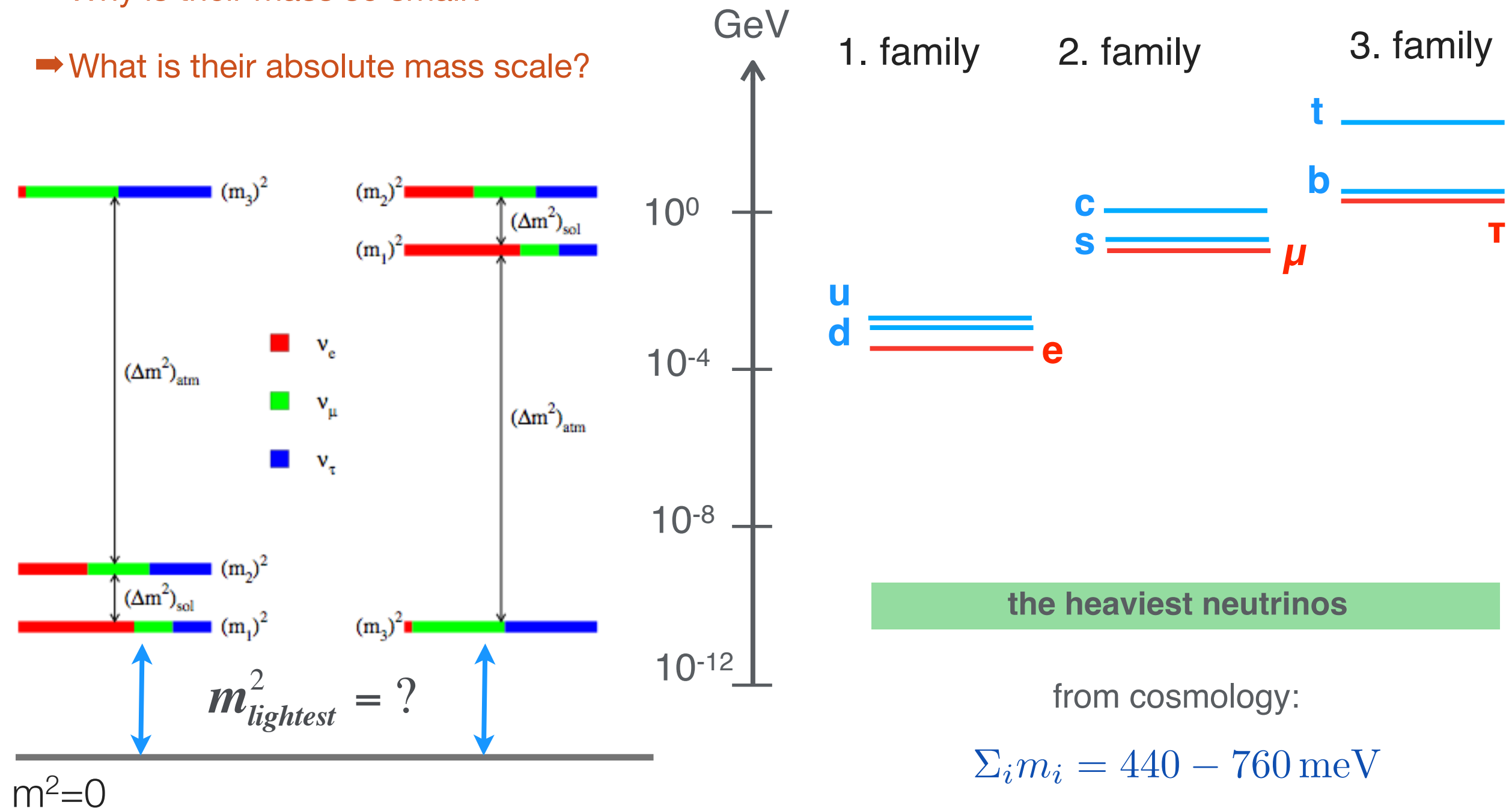


Neutrinos and masses of elementary particles

- Neutrinos: much lighter than other known particles

➔ Why is their mass so small?

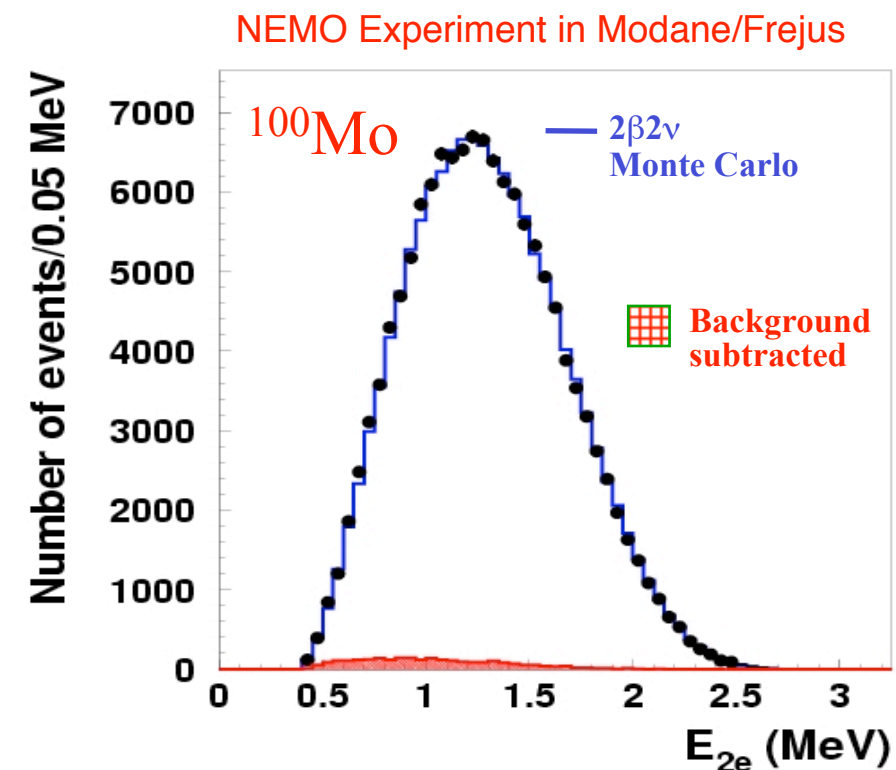
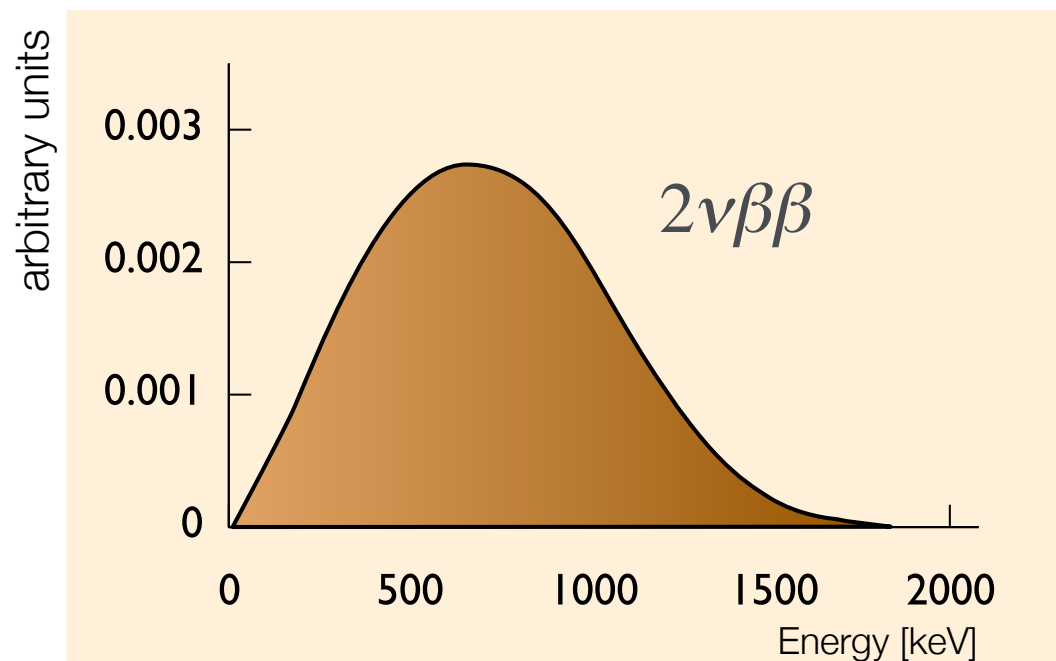
➔ What is their absolute mass scale?



Double beta decay

- The decay with emission of 2 neutrinos was observed in more than 10 different nuclei: ^{48}Ca , ^{76}Ge , ^{82}Se , ^{96}Zr , ^{100}Mo , ^{116}Cd , ^{128}Te , ^{130}Te , ^{136}Xe , ^{150}Nd , ^{238}U
- The observed energy spectrum of the two electrons is continuous, up to the Q-value

$$\Gamma^{2\nu} = \frac{1}{T_{1/2}^{2\nu}} = G^{2\nu}(Q, Z) |M^{2\nu}|^2 \quad G^{2\nu} \propto (G_F \cos\theta_C)^4 Q^7 \left(1 + \frac{Q}{2} + \frac{Q^2}{9} + \frac{Q^3}{90} + \frac{Q^4}{1980} \right)$$



$$Q = E_{e1} + E_{e2} + E_{\nu1} + E_{\nu2} - 2m_e$$

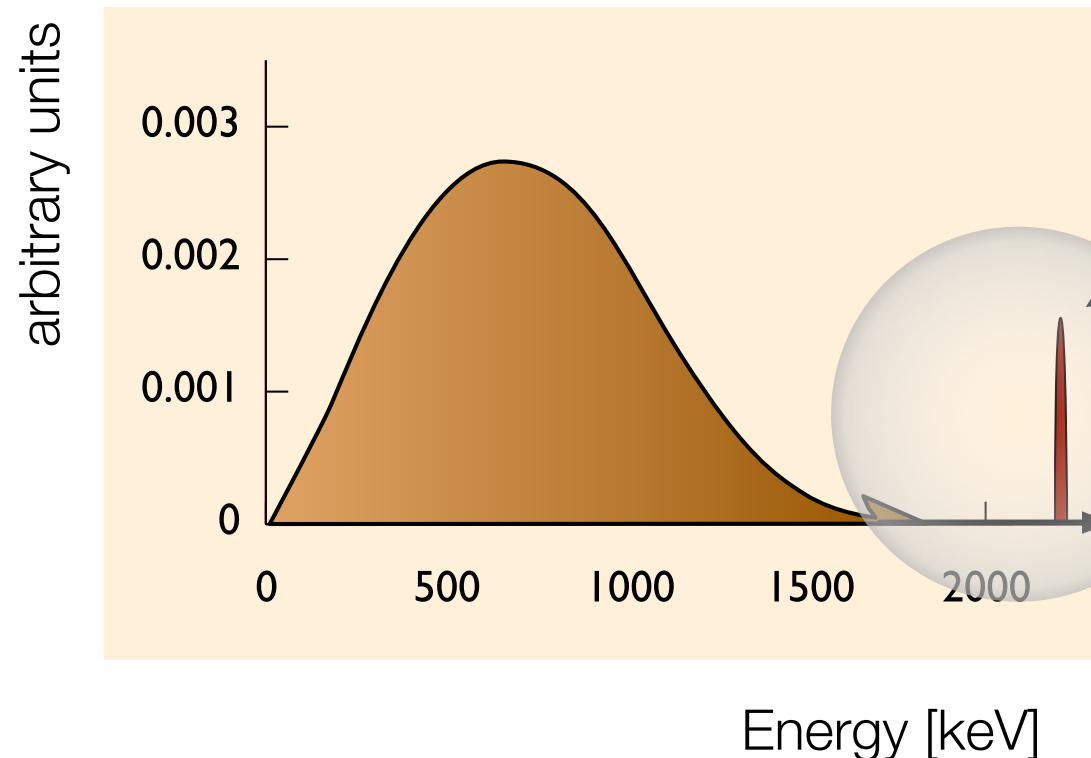
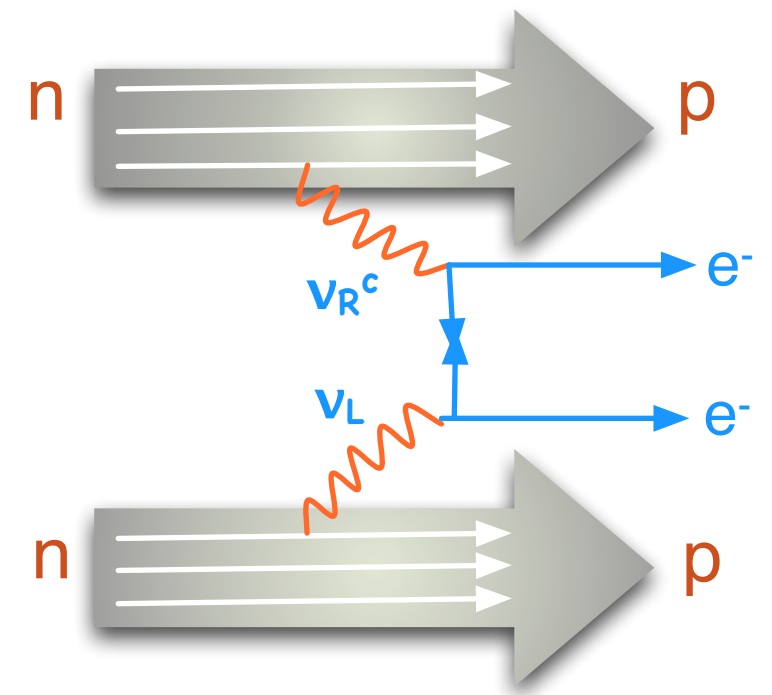
$^{100}\text{Mo}: T_{1/2} = 7.15 \times 10^{18} \text{ a}$

Neutrinoless double beta decay

- More interesting: the decay mode without emission of neutrinos (“forbidden” in the SM, since $\Delta L = 2$)

$$\Gamma^{0\nu} = \frac{1}{T_{1/2}^{0\nu}} = G^{0\nu}(Q, Z) |M^{0\nu}|^2 \frac{|m_{\beta\beta}|^2}{m_e^2}$$

$$G^{0\nu} \propto (G_F \cos\theta_C)^4 \cdot \left[\frac{Q^5}{30} - \frac{2Q^2}{3} + Q - \frac{2}{5} \right] \propto (G_F \cos\theta_C)^4 \cdot Q^5$$

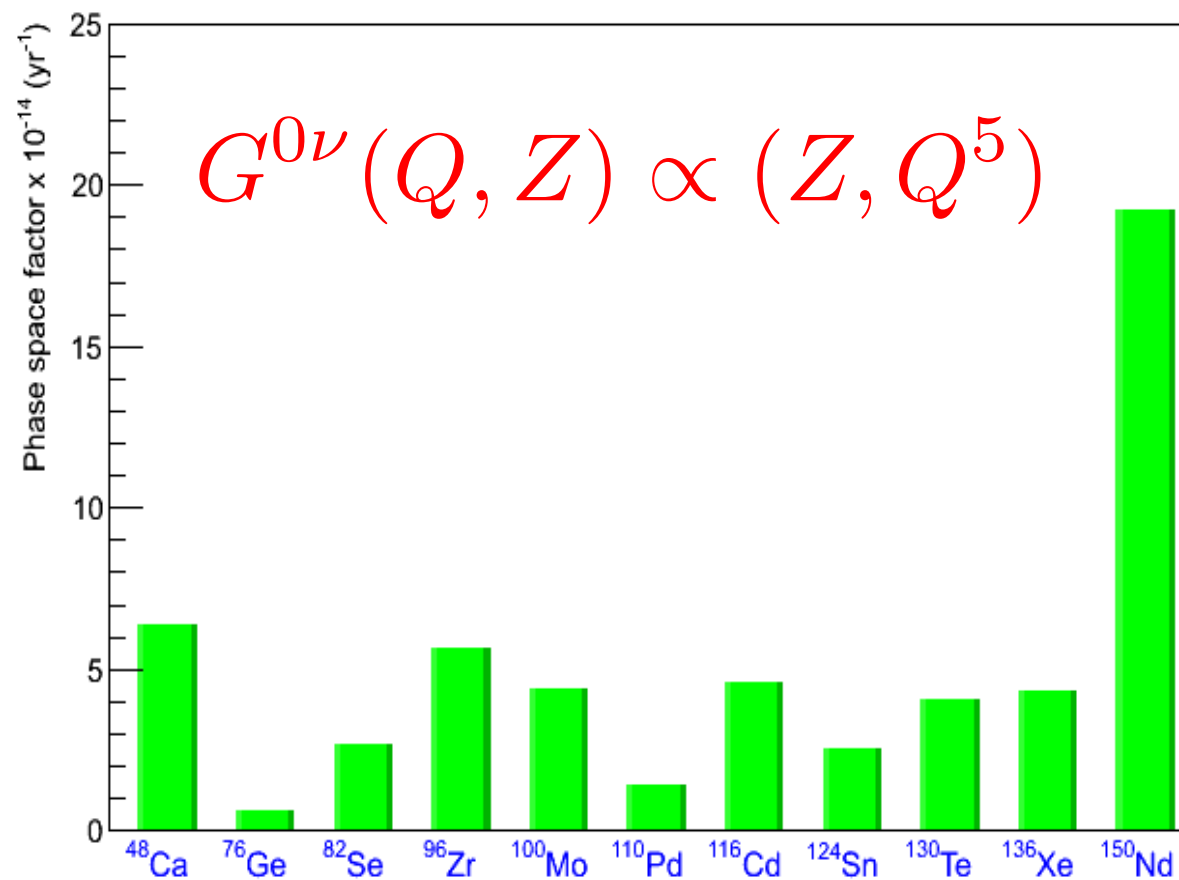


expected:
“peak” at the Q-value of the decay

$$Q = E_{e1} + E_{e2} - 2m_e$$

Phase space

$$\Gamma^{0\nu} = \frac{1}{T_{1/2}^{0\nu}} = G^{0\nu}(Q, Z) |M^{0\nu}|^2 \frac{|m_{\beta\beta}|^2}{m_e^2}$$



F. Piquemal, Neutrino2012, Kyoto

Transition	G [10^{-14} yr^{-1}]	Q [keV]
⁴⁸ Ca \rightarrow ⁴⁸ Ti	6.35	4373.7
⁷⁶ Ge \rightarrow ⁷⁶ Se	0.63	2039.1
⁸² Se \rightarrow ⁸² Kr	2.70	2995.5
¹⁰⁰ Mo \rightarrow ¹⁰⁰ Ru	4.36	3035
¹¹⁶ Cd \rightarrow ¹¹⁶ Sn	4.62	2809
¹³⁰ Te \rightarrow ¹³⁰ Xe	4.09	2530.3
¹³⁶ Xe \rightarrow ¹³⁶ Ba	4.31	2461.9
¹⁵⁰ Nd \rightarrow ¹⁵⁰ Sm	19.2	3367.3

Matrix elements

$$\Gamma^{0\nu} = \frac{1}{T_{1/2}^{0\nu}} = G^{0\nu}(Q, Z) |M^{0\nu}|^2 \frac{|m_{\beta\beta}|^2}{m_e^2}$$

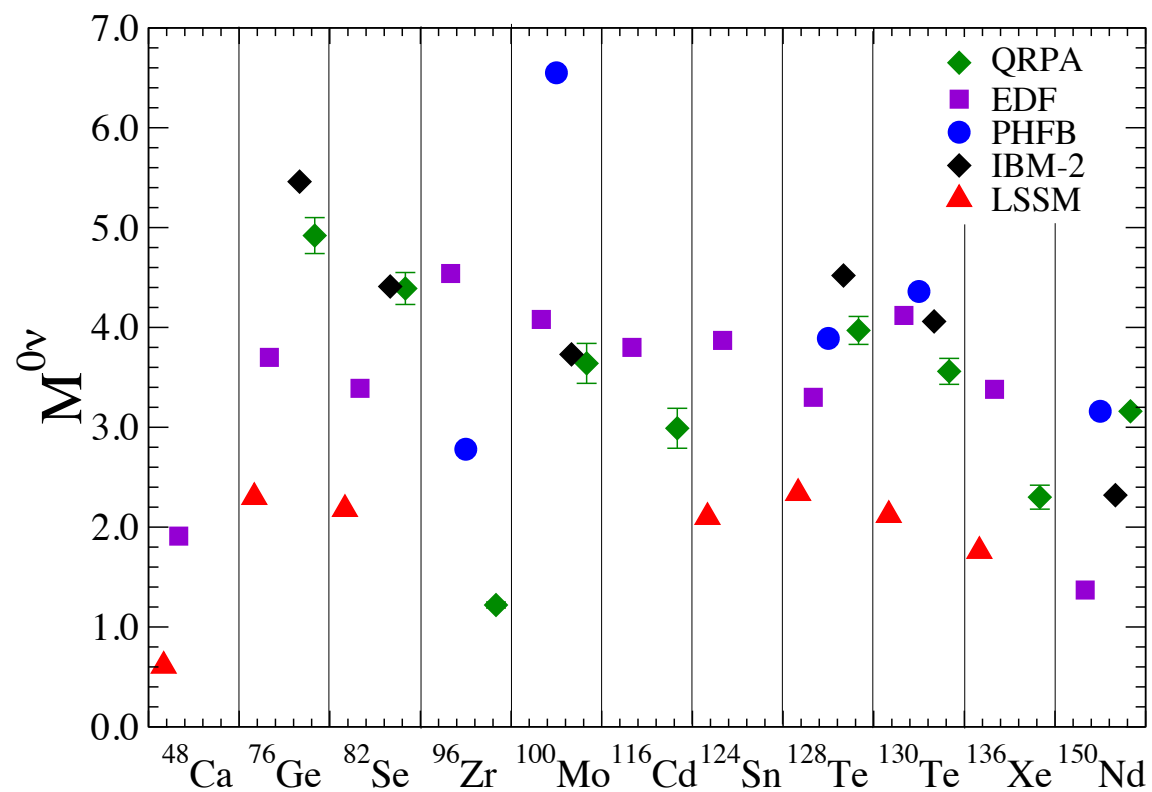


Fig. 3. Values of the NME calculated with the methods in Tab. 2 ⁷⁴.

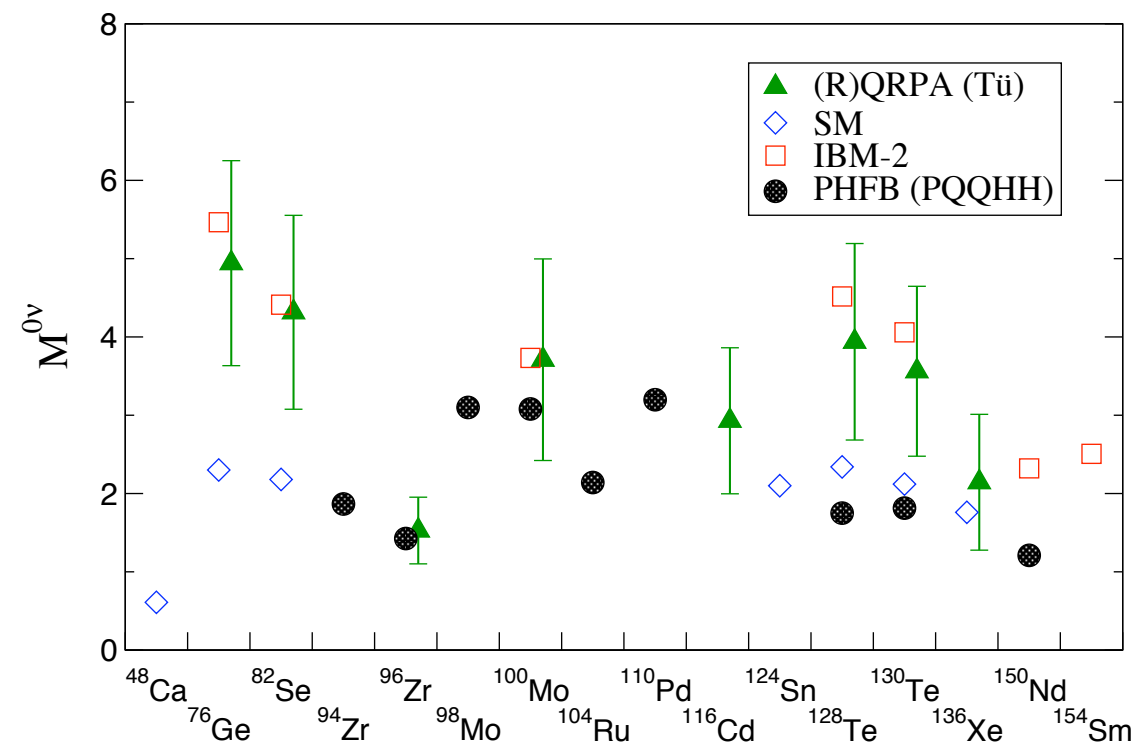


FIG. 7: (Color online) Neutrinoless double beta decay transition matrix elements for the different approaches: QRPA [5, 6], the SM [8–10], the projected HFB method [14] and the IBM [15]. The error bars for the QRPA are calculated as the highest and the lowest values for three different single nucleon basis sets, two different axial charges $g_A = 1.25$ and the quenched value $g_A = 1.00$ and two different treatments of short range correlations (Jastrow-like [25] and the Unitary Correlator Operator Method (UCOM) [26]). The radius parameter is as in this whole work $r_0 = 1.2$ fm.

Matrix elements: vary by a factor of 2- 3 for a given A

Effective Majorana neutrino mass

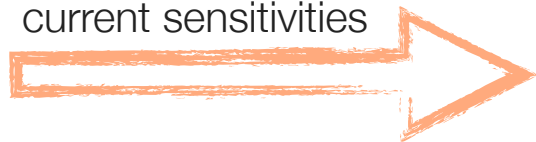
- $|m_{\beta\beta}|$ is a mixture of m_1, m_2, m_3 , proportional to the U_{ei}^2 , where U_{ei} are complex entries

$$|m_{\beta\beta}| = |m_1|U_{e1}|^2 + m_2|U_{e2}|^2 e^{i(\alpha_1 - \alpha_2)} + m_3|U_{e3}|^2 e^{i(-\alpha_1 - 2\delta)}|$$

- where U = neutrino mixing matrix, $c_{ij} = \cos\theta_{ij}$, $s_{ij} = \sin\theta_{ij}$, α_1, α_2 = Majorana phases

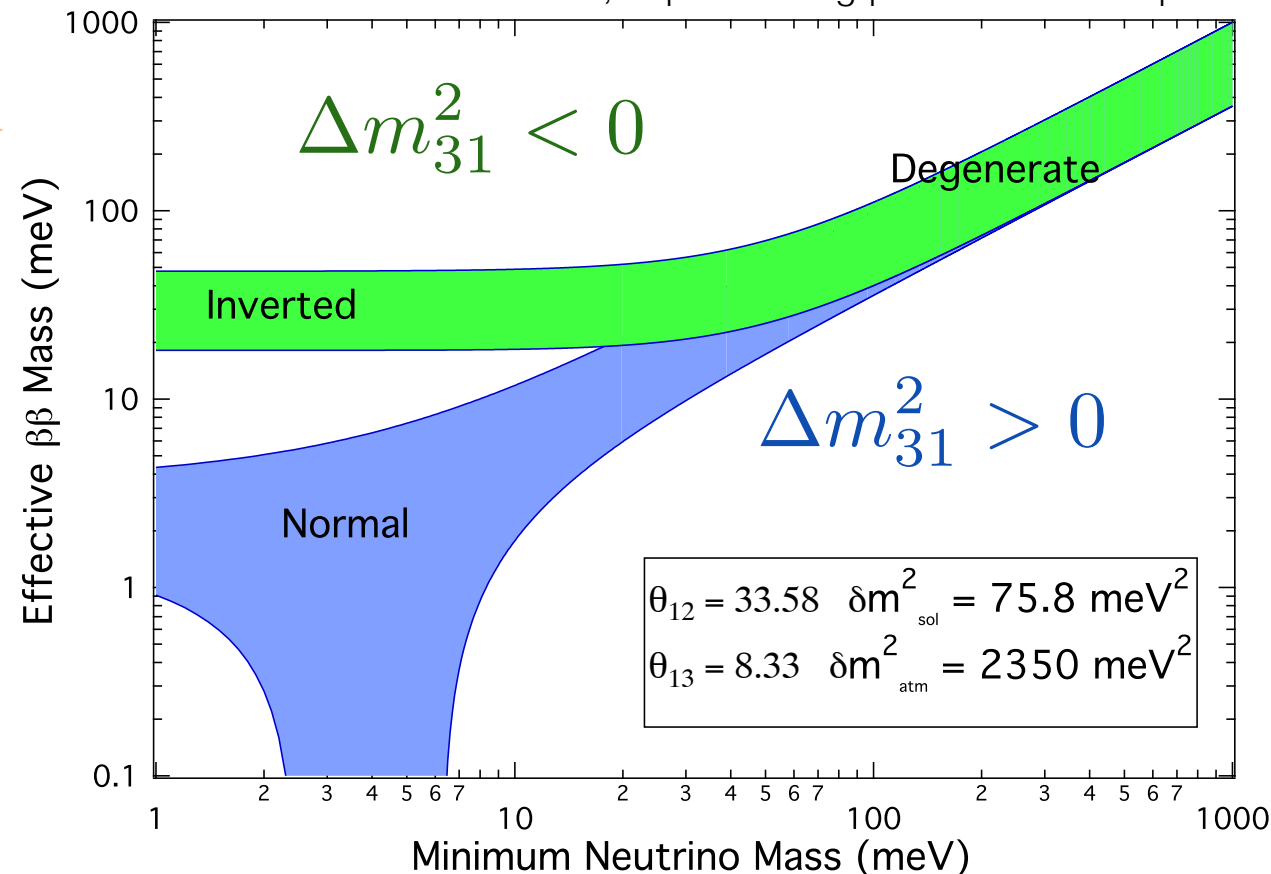
S. Elliott, <http://arxiv.org/pdf/1203.1070v1.pdf>

current sensitivities



$$T_{1/2} \sim 10^{27} \text{ yr}$$

$$T_{1/2} \sim 10^{29} \text{ yr}$$



Remark: here the exchange of a light neutrinos is considered; many other contributions are possible (Majoron, heavy Majorana neutrino exchange, right-handed currents, SUSY, etc)

For a recent review, see:
<http://xxx.lanl.gov/pdf/1205.0649.pdf>

Experimental sensitivity

- Experiments observe:

$$N_{\beta\beta}^{0\nu} = \frac{a \cdot M \cdot N_A}{A} \frac{\ln 2}{T_{1/2}^{0\nu}} \cdot \epsilon \cdot t$$

- with a non-zero number of background events:

$$N_{bg} = M \cdot t \cdot B \cdot \Delta E$$

- The experimental sensitivity is thus:

$$T_{1/2}^{0\nu}(n_\sigma) = \frac{N_A \ln 2}{\sqrt{2} n_\sigma} \frac{a \cdot \epsilon}{A} \sqrt{\frac{M \cdot t}{B \cdot \Delta E}}$$

a = enrichment

ϵ = detector efficiency

M = total mass

t = measuring time

ΔE = energy resolution

B = background index

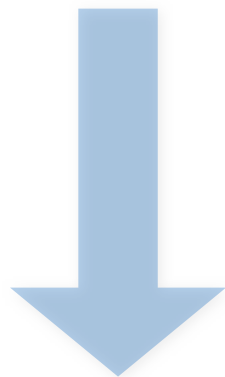
n_σ = confidence level in units of sigma

Experimental requirements

- Experiments thus measure the half life of the decay, $T_{1/2}$

$$T_{1/2}^{0\nu} \propto a \cdot \epsilon \cdot \sqrt{\frac{M \cdot t}{B \cdot \Delta E}}$$

$$\langle m_{\beta\beta} \rangle \propto \frac{1}{\sqrt{T_{1/2}^{0\nu}}}$$



Minimal requirements:

large detector masses
enriched materials
ultra-low background noise
excellent energy resolution



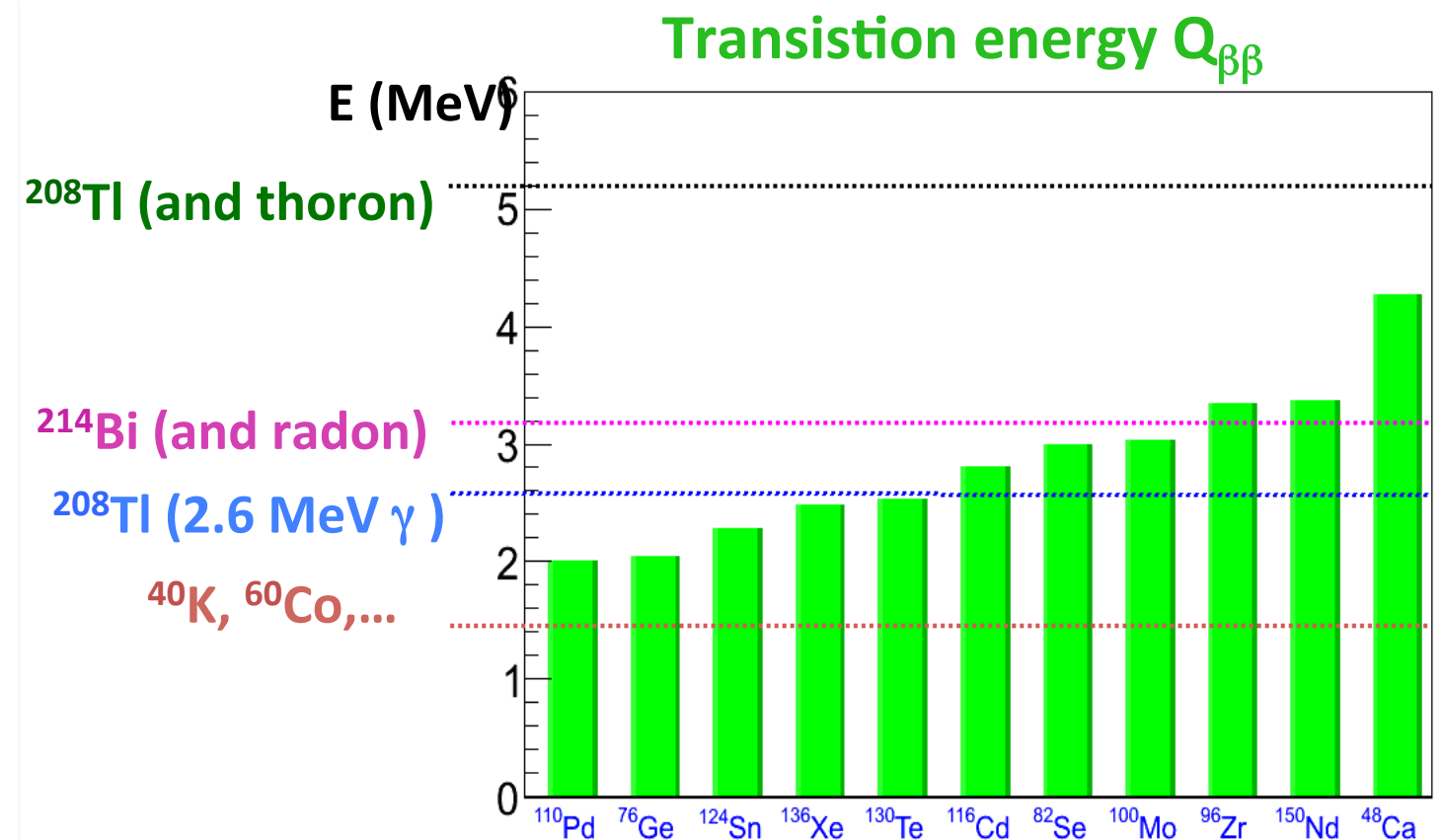
Additional tools to distinguish signal
from background:

angular distribution
decay to excited states (gamma-rays)
identification of daughter nucleus

Backgrounds for double beta experiments

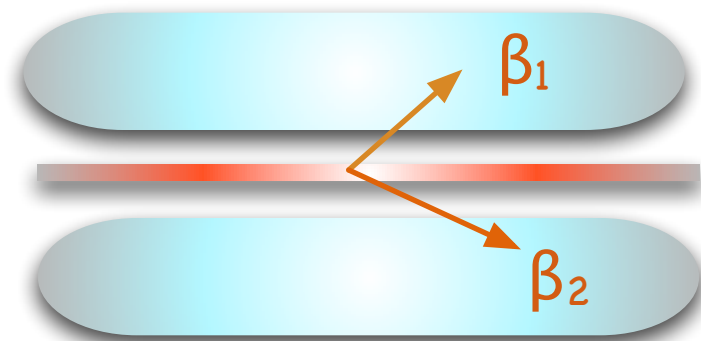
- ⊛ primordial radionuclides (^{238}U , ^{232}Th , ^{40}K) in the detector materials, in the shielding and the concrete/rock (alpha, beta, gamma and neutrons)
- ⊛ cosmic activation of detector materials (^{60}Co , ^{54}Mn , ^{65}Zn ,...)
- ⊛ cosmic rays - muons - and secondary particles
- ⊛ radon in air, radon emanation of materials,....
- ⊛ anthropogens (^{85}Kr , ^{137}Cs , ^{207}Bi ,...)

2 $\nu\beta\beta$ -events: irreducible background
 an excellent energy resolution of the
 detector is crucial



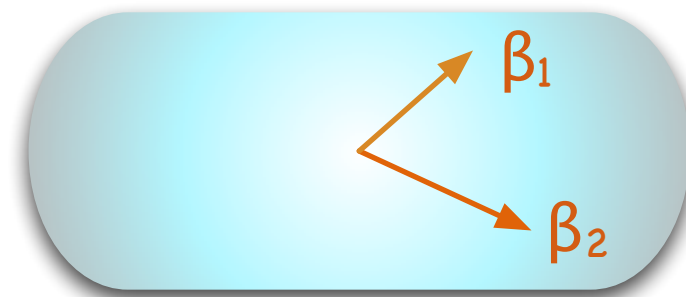
Experiments: Main Approaches

Source \neq Detector

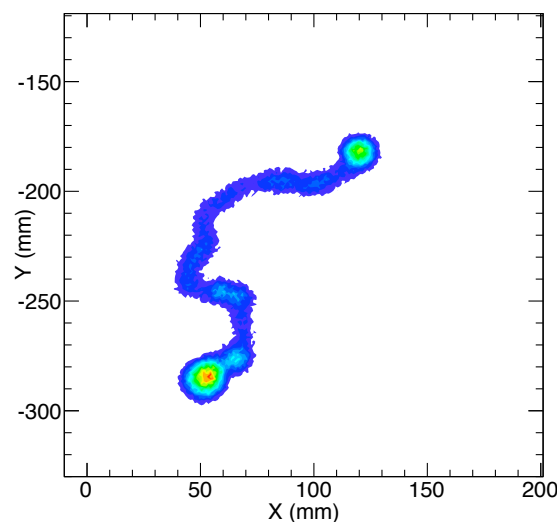


Source as thin foil
Electrons detected with: scintillator, TPC, drift chamber, semiconductor detectors
Event topology
Low energy resolution and detection efficiency

Source = Detector (calorimeters)



The sum of the energy of the two electrons is measured
Signature: peak at the Q-value of the decay
Scintillators, semiconductors, bolometers
High resolution + detection efficiency
No event topology (unless pixellized)



Source = Detector = Tracker

Source is - for example - the (high-pressure) gas of a TPC
Charge and light detected with electron multipliers and/or photosensors
Good energy and position resolution, high efficiency
Event topology very helpful in reducing the background and *in identifying the potential signal*

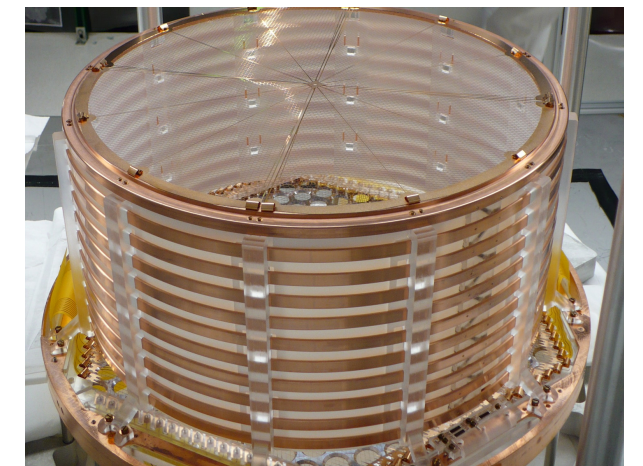
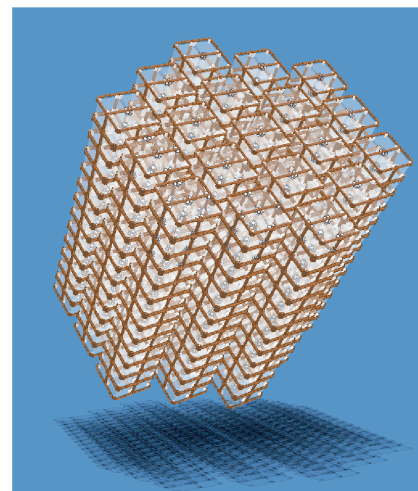
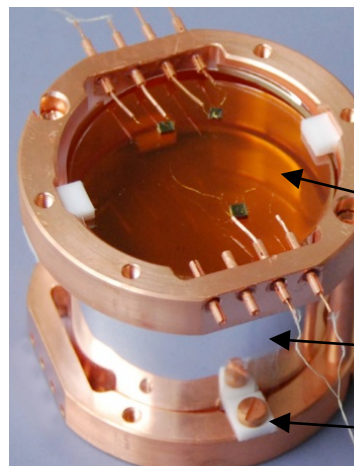
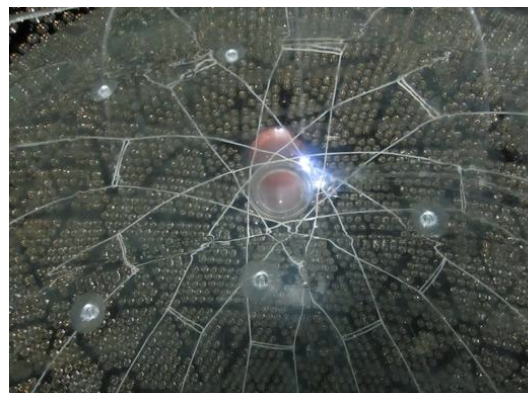
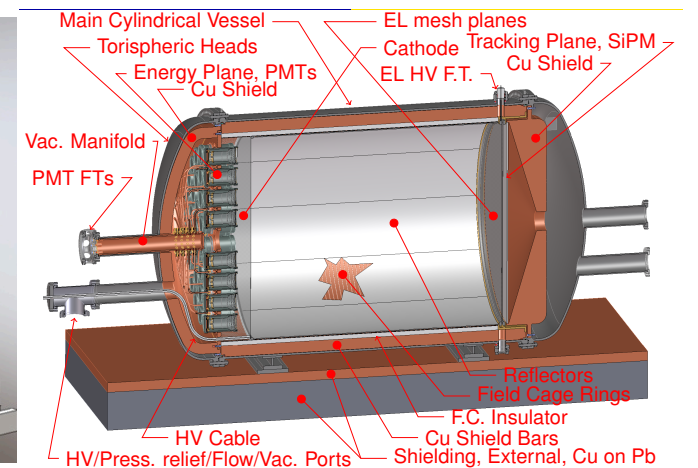
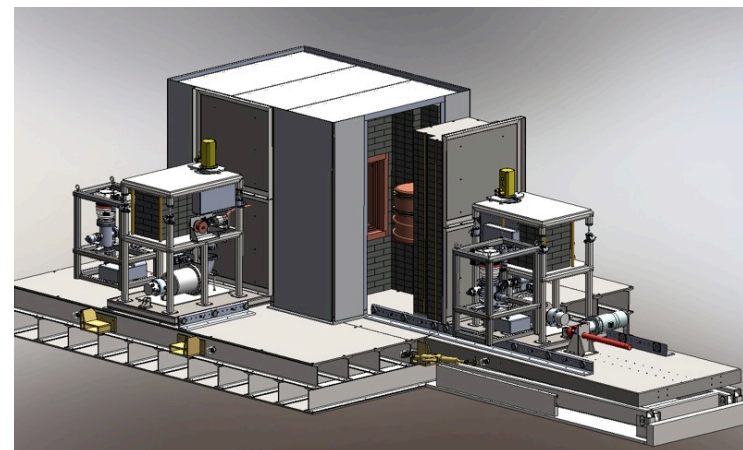
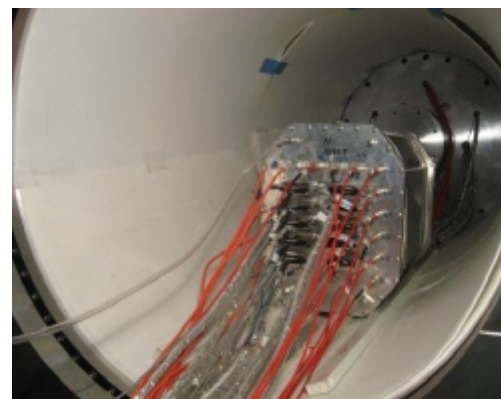
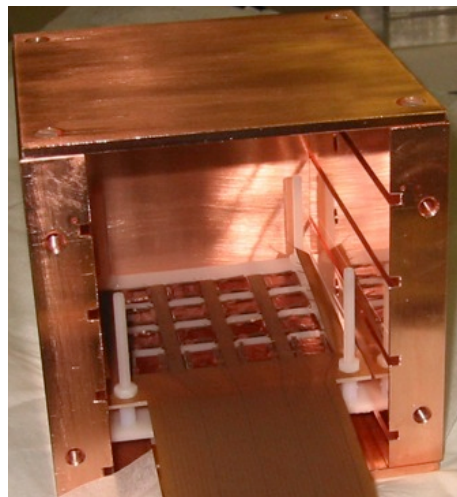
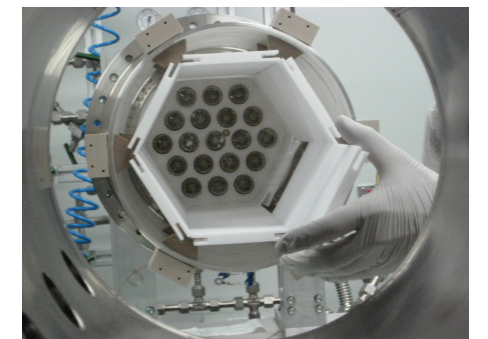
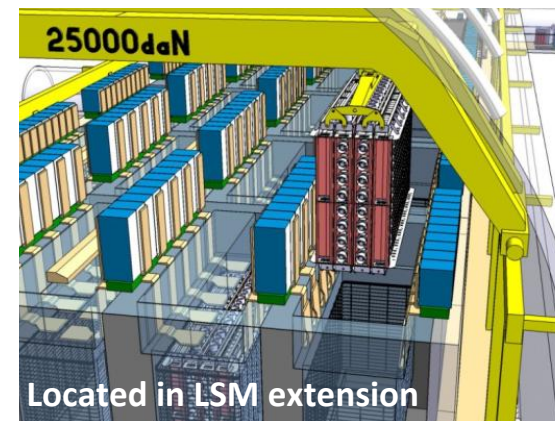
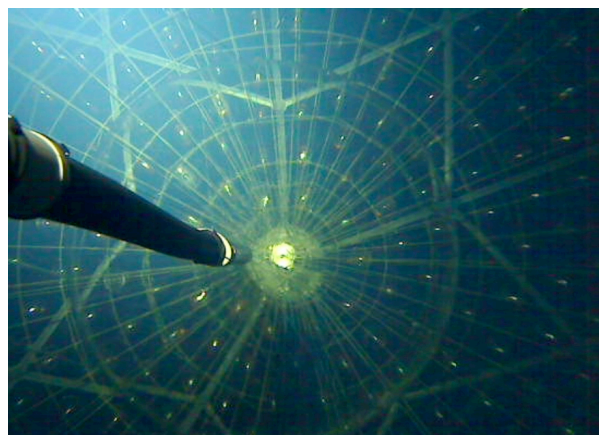
Existing experimental limits on $T_{1/2}$ and the effective Majorana neutrino mass

Current best sensitivities are around a few 100 meV

Table 1. A list of recent $0\nu\beta\beta$ experiments and their 90% confidence level (except as noted) limits on $T_{1/2}^{0\nu}$. The $\langle m_{\beta\beta} \rangle$ limits are those quoted by the authors using the $M_{0\nu}$ of their choice.

Isotope	Technique	$T_{1/2}^{0\nu}$	$\langle m_{\beta\beta} \rangle$ (eV)	Reference
^{48}Ca	CaF_2 scint. crystals	$> 1.4 \times 10^{22}$ y	$< 7.2\text{-}44.7$	14
^{76}Ge	^{enr}Ge det.	$> 1.9 \times 10^{25}$ y	< 0.35	15
^{76}Ge	^{enr}Ge det.	$(1.19_{-0.50}^{+2.99}) \times 10^{25}$ y (3σ)	0.24-0.58	16
^{76}Ge	^{enr}Ge det.	$> 1.57 \times 10^{25}$ y	$< (0.33\text{-}1.35)$	17
^{82}Se	Thin metal foils and tracking	$> 3.6 \times 10^{23}$ y	$< (0.89\text{-}2.54)$	18
^{96}Zr	Thin metal foils and tracking	$> 9.2 \times 10^{21}$ y	$< (7.2\text{-}19.5)$	19
^{100}Mo	Thin metal foils and tracking	$> 1.1 \times 10^{24}$ y	$< (0.45\text{-}0.93)$	18
^{116}Cd	$^{116}\text{CdWO}_4$ scint. crystals	$> 1.7 \times 10^{23}$ y	< 1.7	20
^{128}Te	geochemical	$> 7.7 \times 10^{24}$ y	$< (1.1\text{-}1.5)$	21
^{130}Te	TeO_2 bolometers	$> 2.8 \times 10^{24}$ y	$< (0.3\text{-}0.7)$	22
^{136}Xe	Xe dissolved in liq. scint.	$> 5.7 \times 10^{24}$ y	$< (0.3\text{-}0.6)$	23
^{150}Ne	Thin metal foil within TPC	$> 1.8 \times 10^{22}$ y	N.A.	24

Current, near-future, future experiments



Current, near-future, future experiments

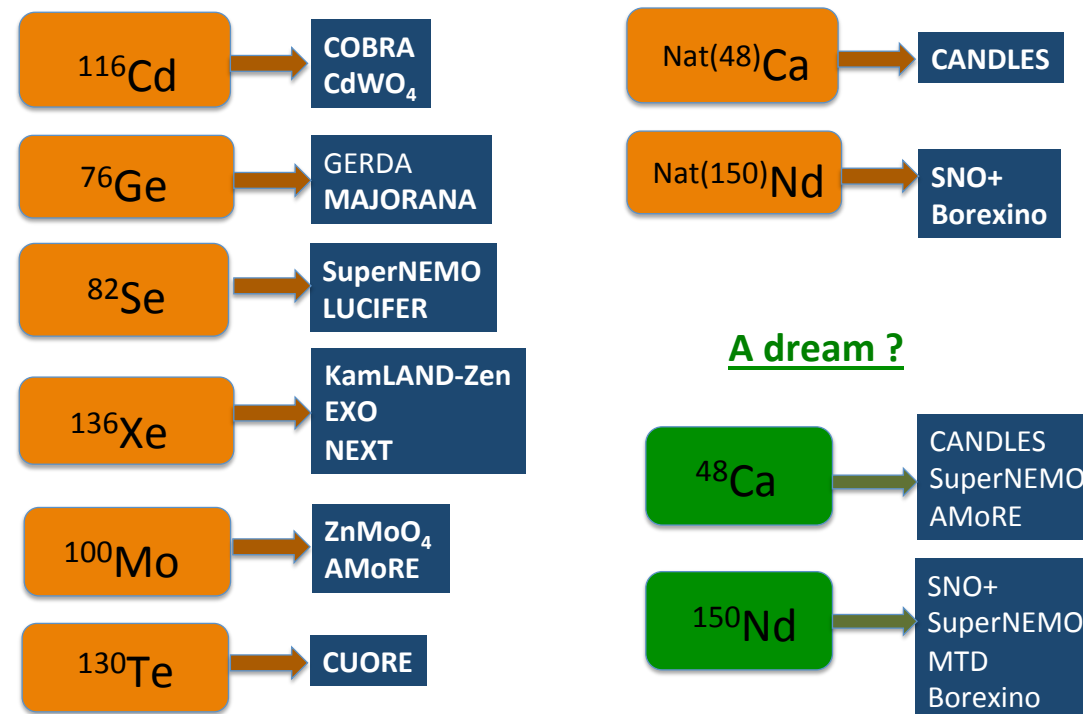


Existing and proposed experiments

Table 2. A summary list of the $0\nu\beta\beta$ proposals and experiments.

Experiment	Isotope	Mass	Technique	Present Status	Location
AMoRE ^{89,90}	¹⁰⁰ Mo	50 kg	CaMoO ₄ scint. bolometer crystals	Development	Yangyang
CANDLES ⁹¹	⁴⁸ Ca	0.35 kg	CaF ₂ scint. crystals	Prototype	Kamioka
CARVEL ⁹²	⁴⁸ Ca	1 ton	CaF ₂ scint. crystals	Development	Solotvina
COBRA ⁹³	¹¹⁶ Cd	183 kg	^{enr} Cd CZT semicond. det.	Prototype	Gran Sasso
CUORE-0 ⁶⁹	¹³⁰ Te	11 kg	TeO ₂ bolometers	Construction - 2012	Gran Sasso
CUORE ⁶⁹	¹³⁰ Te	203 kg	TeO ₂ bolometers	Construction - 2013	Gran Sasso
DCBA ⁹⁴	¹⁵⁰ Ne	20 kg	^{enr} Nd foils and tracking	Development	Kamioka
EXO-200 ⁵⁷	¹³⁶ Xe	160 kg	Liq. ^{enr} Xe TPC/scint.	Operating - 2011	WIPP
EXO ⁷⁰	¹³⁶ Xe	1-10 t	Liq. ^{enr} Xe TPC/scint.	Proposal	SURF
GERDA ⁷¹	⁷⁶ Ge	≈35 kg	^{enr} Ge semicond. det.	Operating - 2011	Gran Sasso
GSO ⁹⁵	¹⁶⁰ Gd	2 ton	Gd ₂ SiO ₅ :Ce crys. scint. in liq. scint.	Development	
KamLAND-Zen ⁹⁶	¹³⁶ Xe	400 kg	^{enr} Xe dissolved in liq. scint.	Operating - 2011	Kamioka
LUCIFER ^{97,98}	⁸² Se	18 kg	ZnSe scint. bolometer crystals	Development	Gran Sasso
MAJORANA ^{77,78,79}	⁷⁶ Ge	26 kg	^{enr} Ge semicond. det.	Construction - 2013	SURF
MOON ⁹⁹	¹⁰⁰ Mo	1 t	^{enr} Mofolios/scint.	Development	
SuperNEMO-Dem ⁸⁷	⁸² Se	7 kg	^{enr} Se foils/tracking	Construction - 2014	Fréjus
SuperNEMO ⁸⁷	⁸² Se	100 kg	^{enr} Se foils/tracking	Proposal - 2019	Fréjus
NEXT ^{82,83}	¹³⁶ Xe	100 kg	gas TPC	Development - 2014	Canfranc
SNO+ ^{84,85}	¹⁵⁰ Nd	55 kg	Nd loaded liq. scint.	Construction - 2013	SNOLab

Steve Elliott: <http://arxiv.org/pdf/1203.1070v1.pdf>

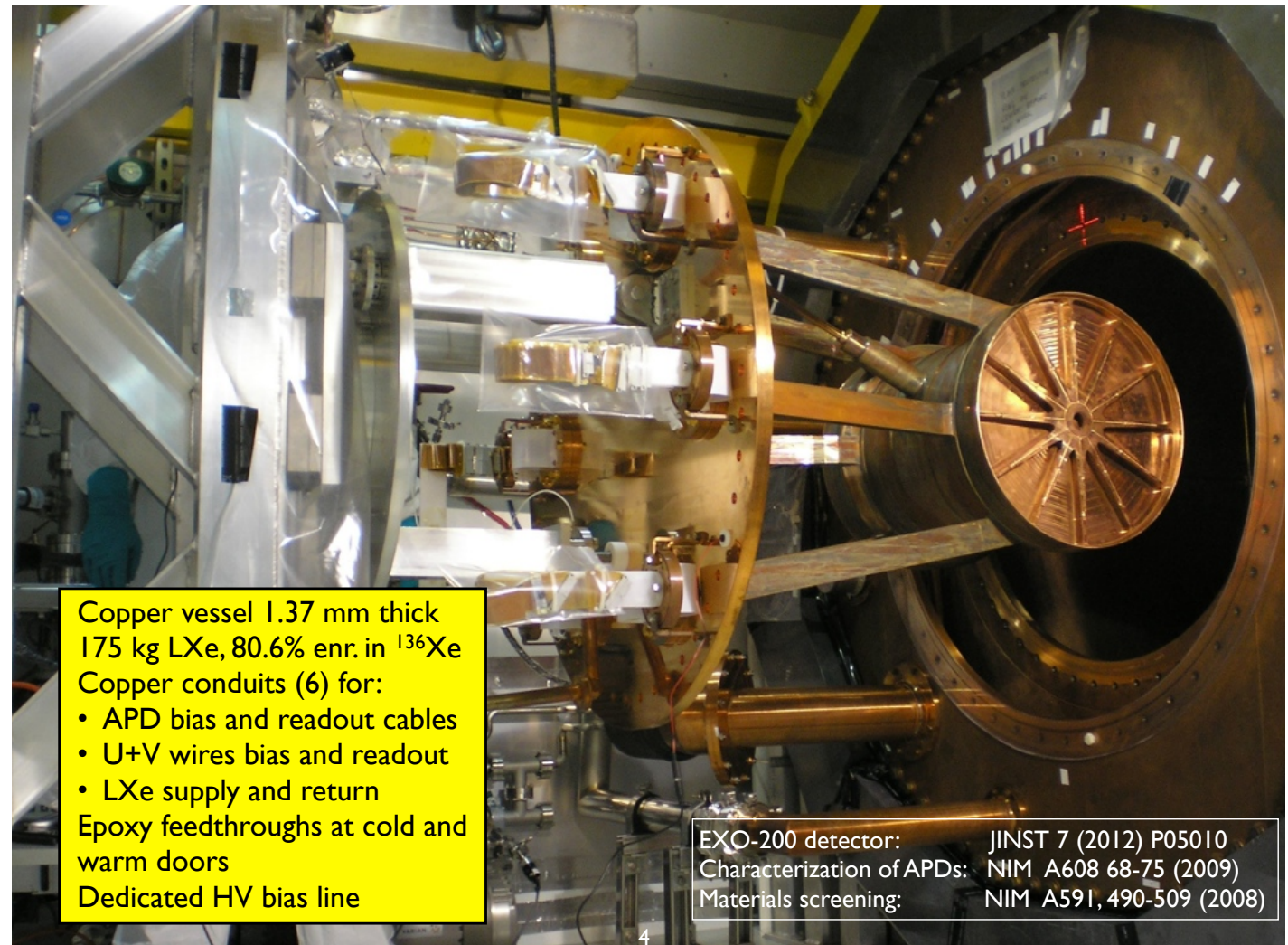
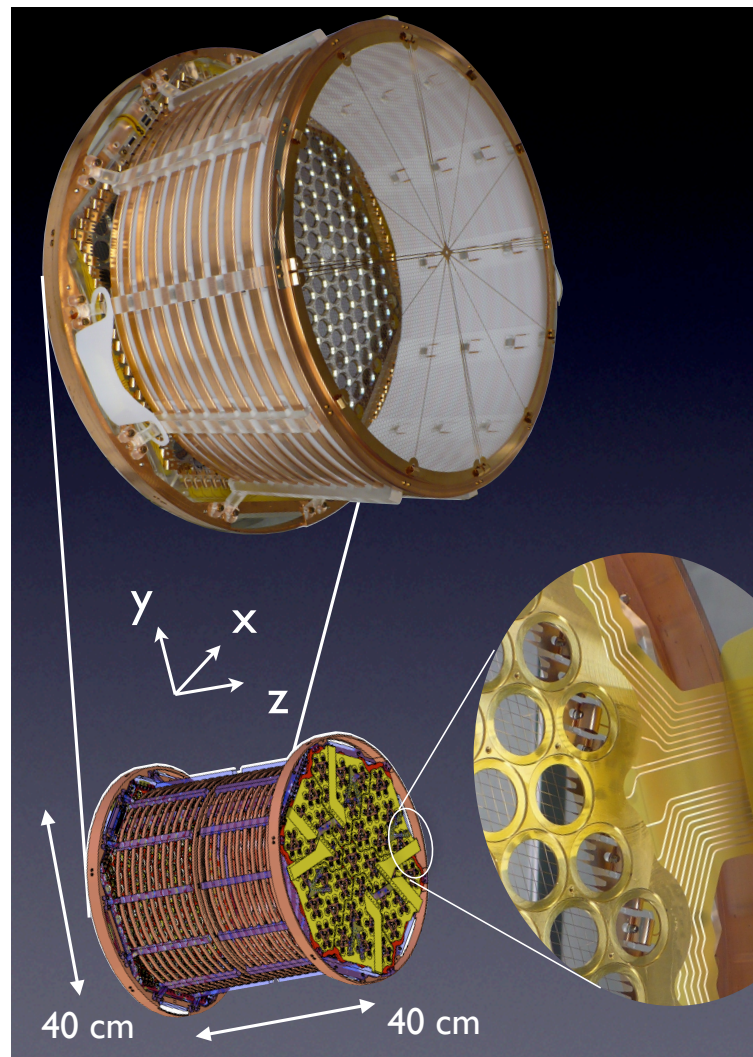


F. Piquemal, talk at Neutrino2012, Kyoto

Recent results

EXO-200

- Liquid xenon TPC: 175 kg LXe, 80.6% enriched in ^{136}Xe
- Charge and light readout (triplet wire channels and large area avalanche photodiodes)
- Drift field: 376 V/cm



EXO-200

- So far, 2 data taking phases
- First measurement of ^{136}Xe 2-neutrino half life; limit on the 0-neutrino mode

	Run I	Run 2 (this analysis)
Period	May 21, 11 – Jul 9, 11	Sep 22, 11 – Apr 15, 12
Live Time	752.7 hr	2,896.6 hr
Exposure (^{136}Xe)	4.4 kg-yr	26.3 kg-yr
Publ.	PRL 107 (2011) 212501	arXiv:1205:5608

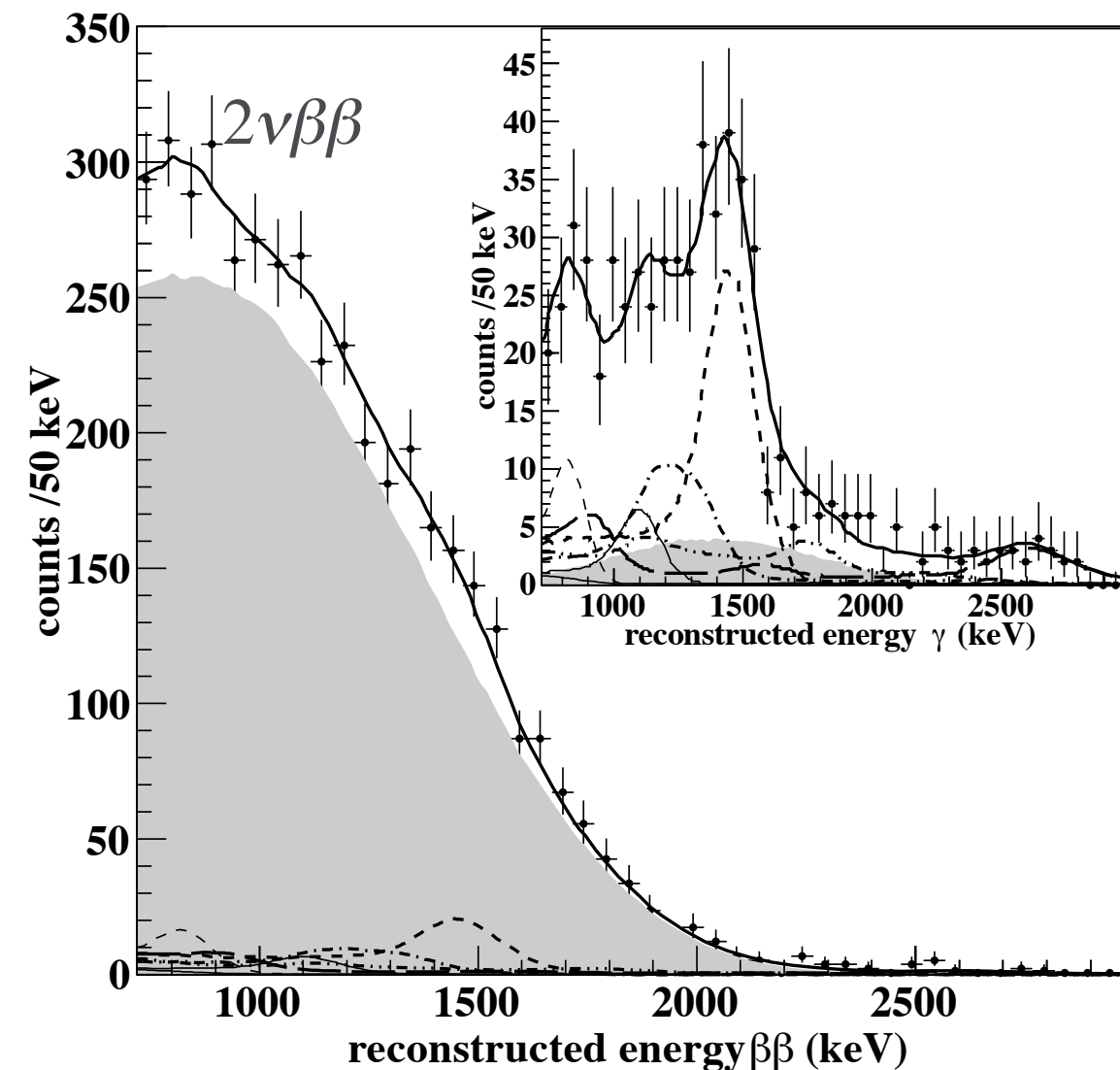
Run I Results:

$$T_{1/2}^{2\nu\beta\beta} (^{136}\text{Xe}) = (2.11 \pm 0.04 \text{ stat} \pm 0.21 \text{ sys}) \cdot 10^{21} \text{ yr}$$

In disagreement with previously reported limits by

R. Bernabei et al. Phys. Lett. B 546 (2002) 23, and
Yu. M. Gavriljuk et al. , Phys. Atom Nucl. 69 (2006)

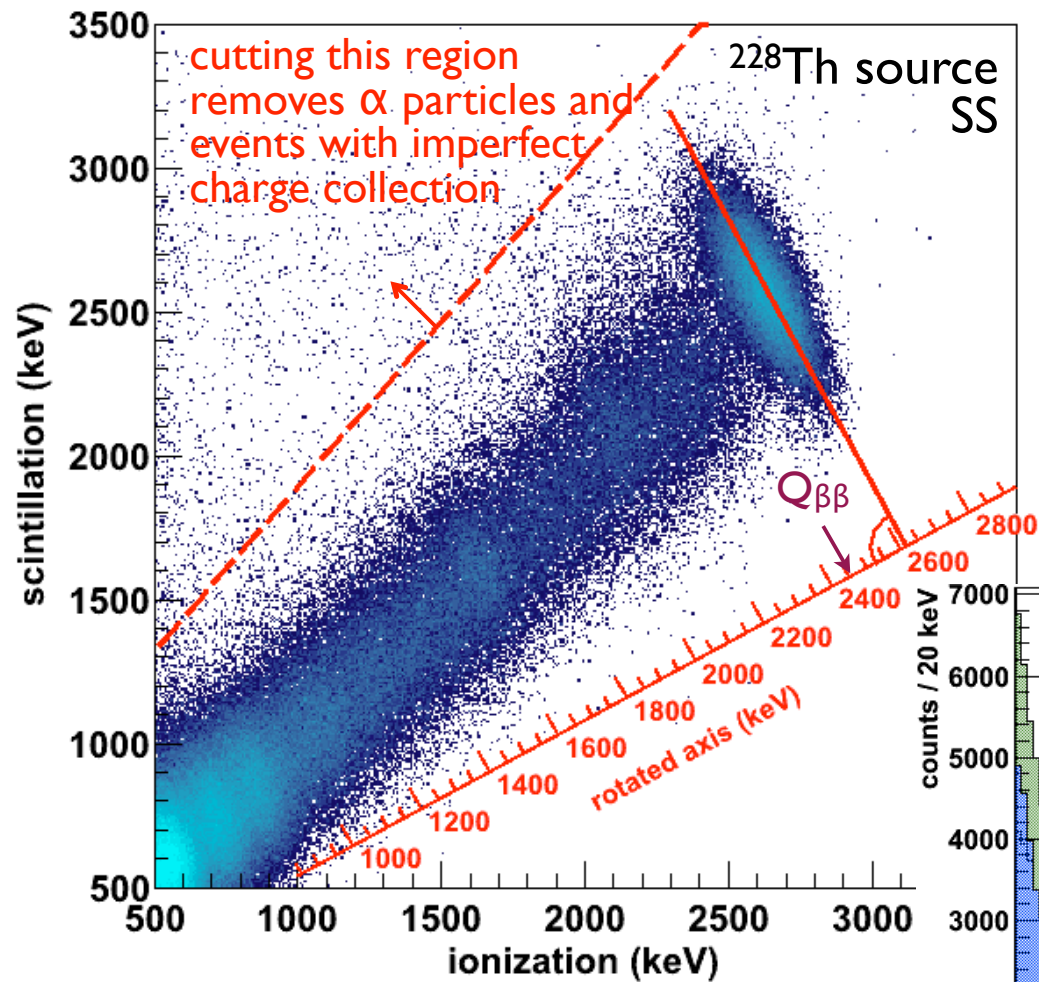
This was also a measurement of a nuclear matrix element of 0.019 MeV^{-1} , the smallest measured among the $2\nu\beta\beta$ emitters



EXO-200: resolution and calibration

- Good energy resolution by linear combination of scintillation and charge signals

At $Q_{\beta\beta}$ (2458 keV):
 $\sigma/E = 1.67\%$ (SS)
 $\sigma/E = 1.84\%$ (MS)

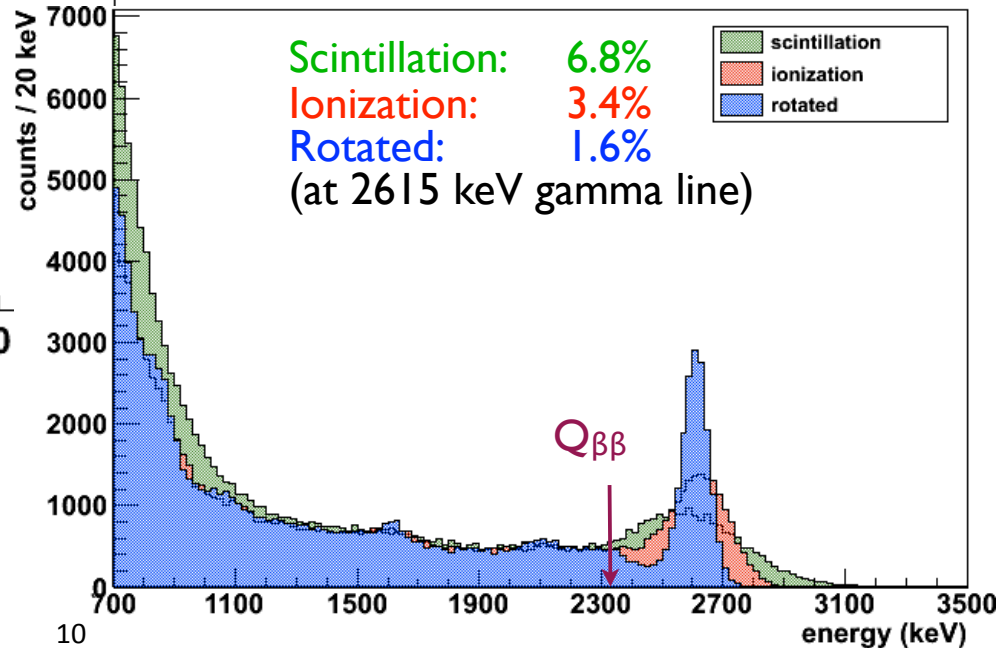


Rotation angle chosen to optimize energy resolution at 2615 keV

Properties of xenon cause increased scintillation to be associated with decreased ionization (and vice-versa)

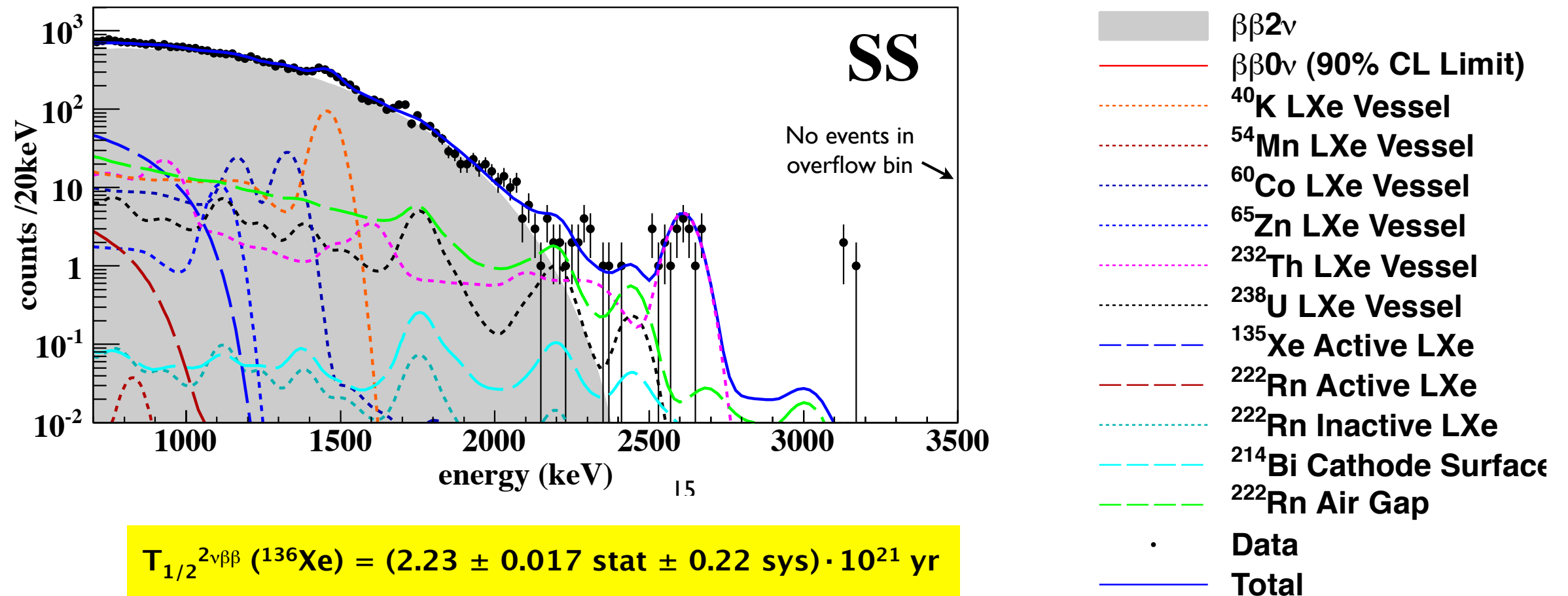
E. Conti et al. Phys. Rev. B 68 (2003) 054201

Use projection onto a rotated axis to determine event energy



EXO-200: low-background spectrum

- Observed 22'000 2-neutrino events in 32.5 kg yr exposure
- Background PDFs fitted along with 2-neutrino and 0-neutrino PDFs



$$T_{1/2}^{2\nu\beta\beta} (^{136}\text{Xe}) = (2.23 \pm 0.017 \text{ stat} \pm 0.22 \text{ sys}) \cdot 10^{21} \text{ yr}$$

In agreement with previously reported value by

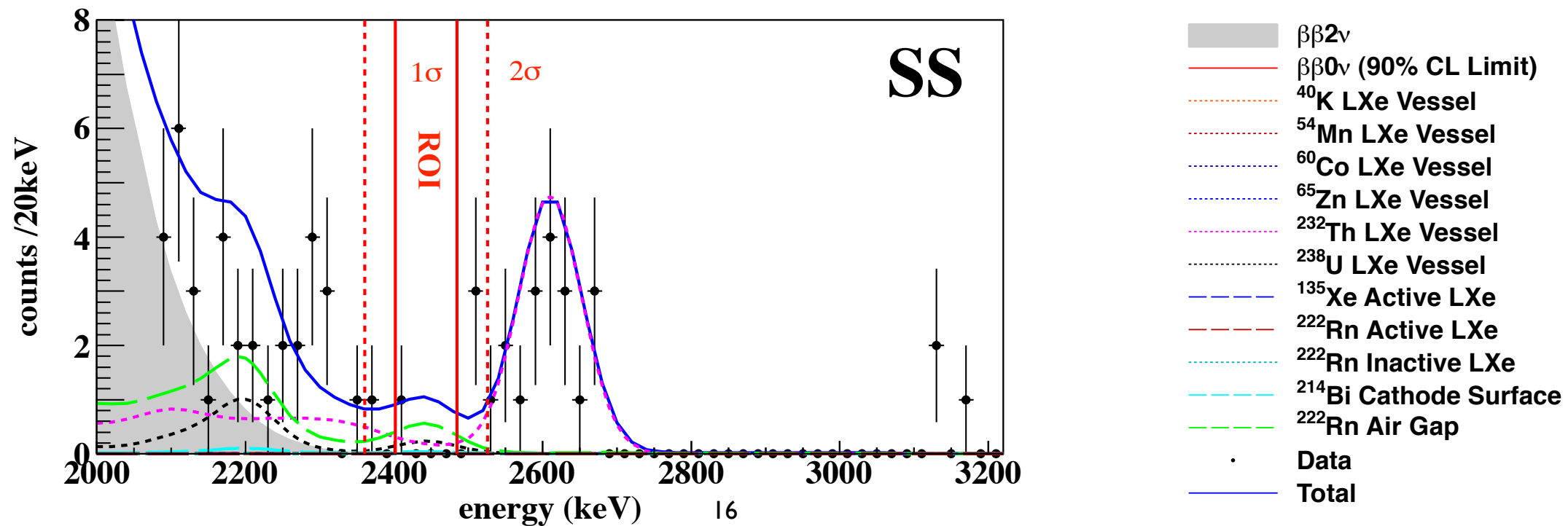
EXO-200 Phys.Rev.Lett. 107 (2011) 212501

and

KamLAND-ZEN Phys.Rev.C85:045504,2012)

EXO-200: low-background spectrum

- No 0-neutrino signal observed => lower limit on $T_{1/2}$



	Expected events from fit			
	$\pm 1 \sigma$		$\pm 2 \sigma$	
^{222}Rn in cryostat air-gap	1.9	± 0.2	2.9	± 0.3
^{238}U in LXe Vessel	0.9	± 0.2	1.3	± 0.3
^{232}Th in LXe Vessel	0.9	± 0.1	2.9	± 0.3
^{214}Bi on Cathode	0.2	± 0.01	0.3	± 0.02
All Others	~ 0.2		~ 0.2	
Total	4.1	± 0.3	7.5	± 0.5
Observed	1		5	
Background index b ($\text{kg}^{-1}\text{yr}^{-1}\text{keV}^{-1}$)	$1.5 \cdot 10^{-3}$	± 0.1	$1.4 \cdot 10^{-3}$	± 0.1

From profile likelihood:

$$T_{1/2}^{0\nu\beta\beta} > 1.6 \cdot 10^{25} \text{ yr}$$

$$\langle m_{\beta\beta} \rangle < 140\text{--}380 \text{ meV}$$

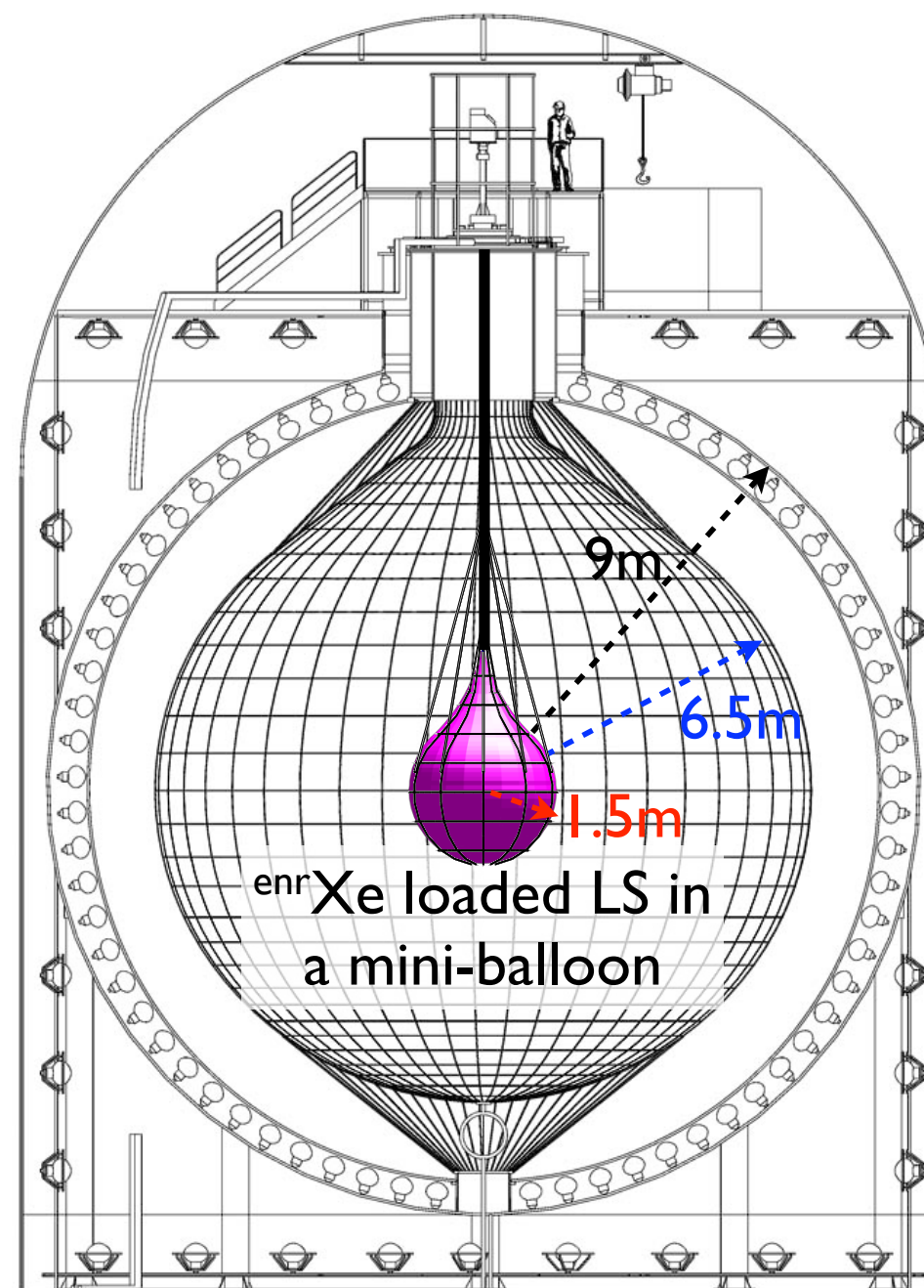
(90% C.L.)

arXiv:1205.5608 – Subm. to PRL

KAMLAND-Zen

- Scintillator loaded with xenon
- 320 kg 90% enriched ^{136}Xe so far (more than 600 kg in the Kamioka mine)
- Advantages: huge and clean (U: $3.5\text{e-}18$ g/g, Th: $5.2\text{e-}17$ g/g) running detector
- Xe-LS can be purified, and is highly scalable
- No escape or invisible energy from gammas and beta: good background identification
- Disadvantage: relatively poor energy resolution
- no beta/gamma discrimination
- limited LS composition

Zero Neutrino
double beta decay search



KamLAND-Zen: installation

Installation in a class 10~100 clean room built at the top of KamLAND



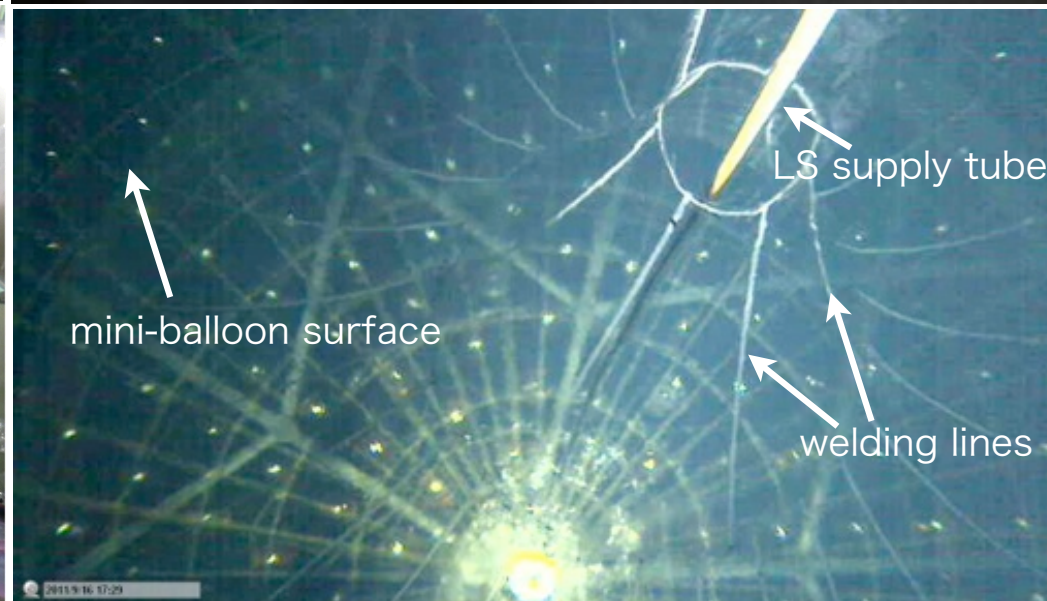
balloon and corrugated tube deployment



balloon went through the black sheet



installation completed



mini-balloon surface

LS supply tube

welding lines

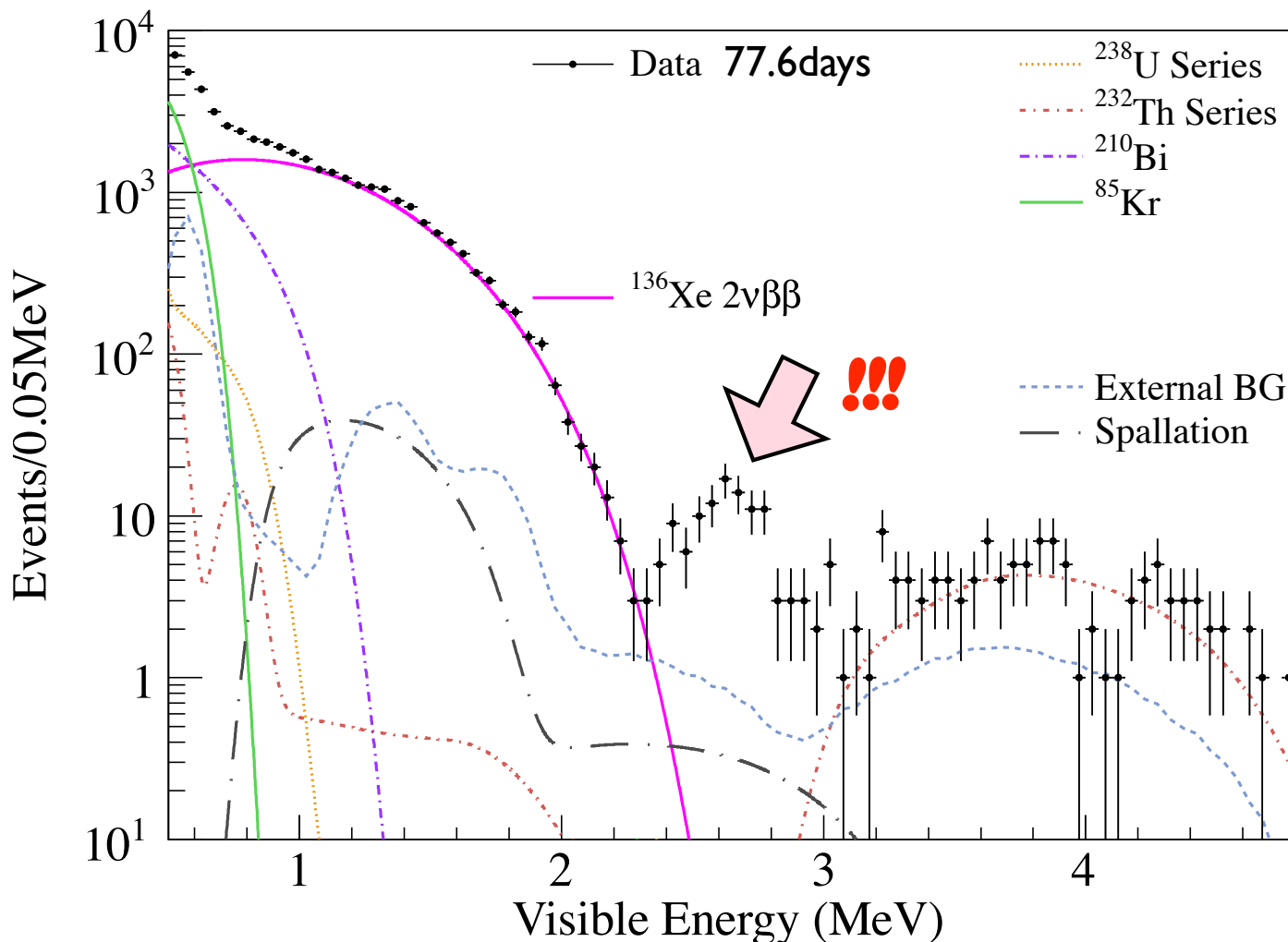
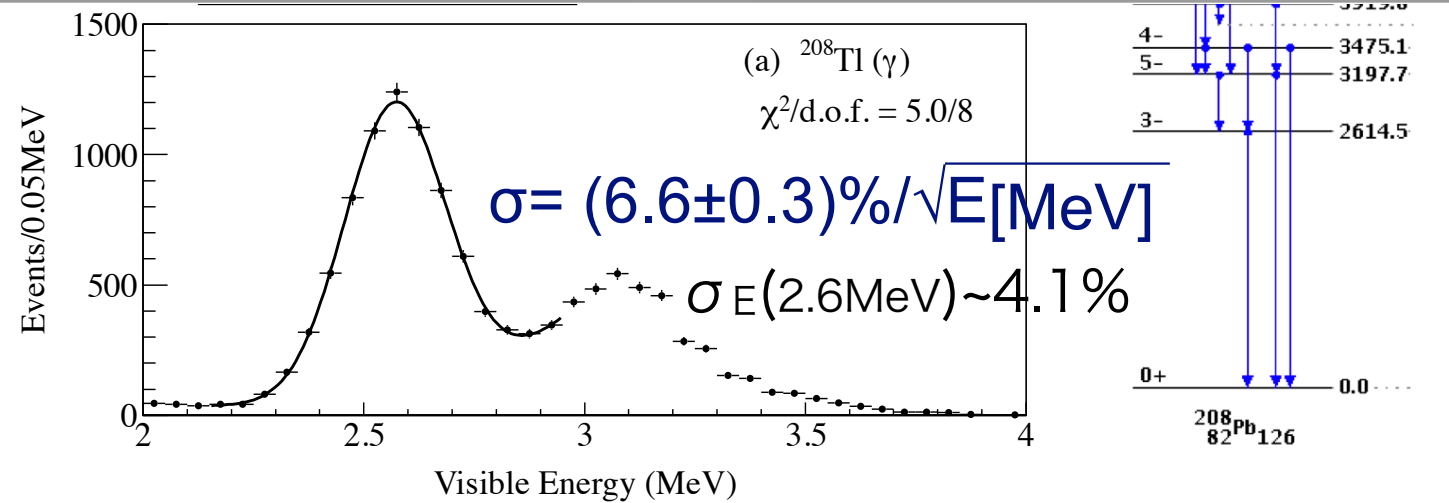
mini-balloon inflated with dummy LS and then replaced with Xe-loaded LS
density tuning finished and tubes to be extracted₅

KamLAND-Zen



KamLAND-Zen: energy calibration and low-background spectrum

- Resolution at 2.6 MeV: $\sigma \sim 4.1\%$



KamLAND-Zen (2012)

Xe loaded liquid scintillator

$$T^{2\nu}_{1/2} = 2.38 \pm 0.02(\text{stat}) \pm 0.14(\text{syst}) \times 10^{21} \text{ years}$$

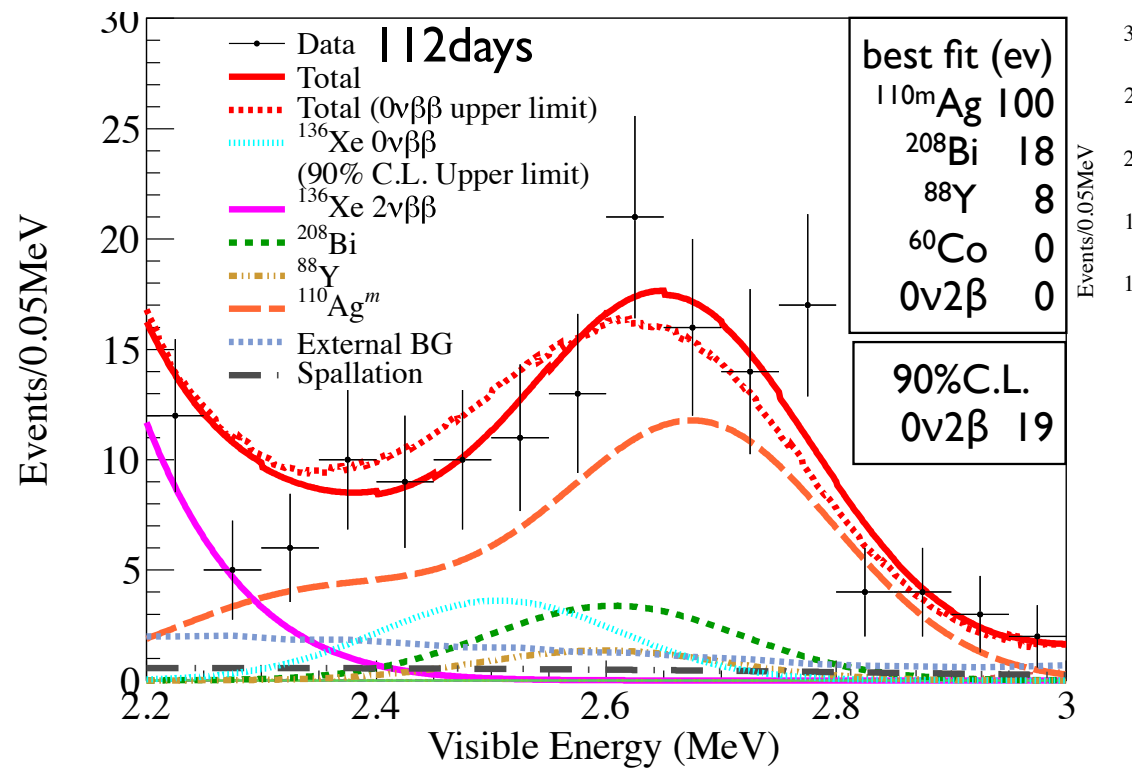
Phys.Rev.C85,045504(2012)

update $T^{2\nu}_{1/2} = 2.30 \pm 0.02(\text{stat}) \pm 0.12(\text{syst}) \times 10^{21} \text{ years}$

arXiv:1205.6372

KamLAND-Zen: low-background spectrum

- Peak around the Q-value; however, peak position is different



(χ^2 at 2.2~3.0MeV)

	χ^2 112days	
simul. fit	11.6	
$0\nu+^{110m}\text{Ag}$	13.1	
$0\nu+^{208}\text{Bi}$	22.7	△
$0\nu+^{88}\text{Y}$	22.2	△
$0\nu+^{60}\text{Co}$	82.9	×
0ν only	85.0	×

BG is likely to be ^{110m}Ag

simultaneous fit and 90% CL upper limit for $0\nu 2\beta$

$$T_{1/2}^{0\nu} > 5.7 \times 10^{24} \text{ years at 90\% C.L. (78days)}$$

factor 5 improvement from DAMA

$$T_{1/2}^{0\nu} > 6.2 \times 10^{24} \text{ years (KL-Zen 112days)}$$

(ref. current best is 1.6×10^{24} years from EXO-200)

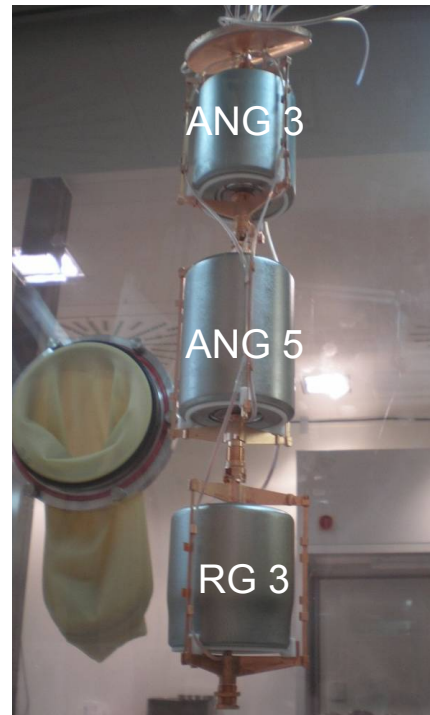


(R)QRPA (CCM SRC)
Phys.Rev.C79,055501(2009)

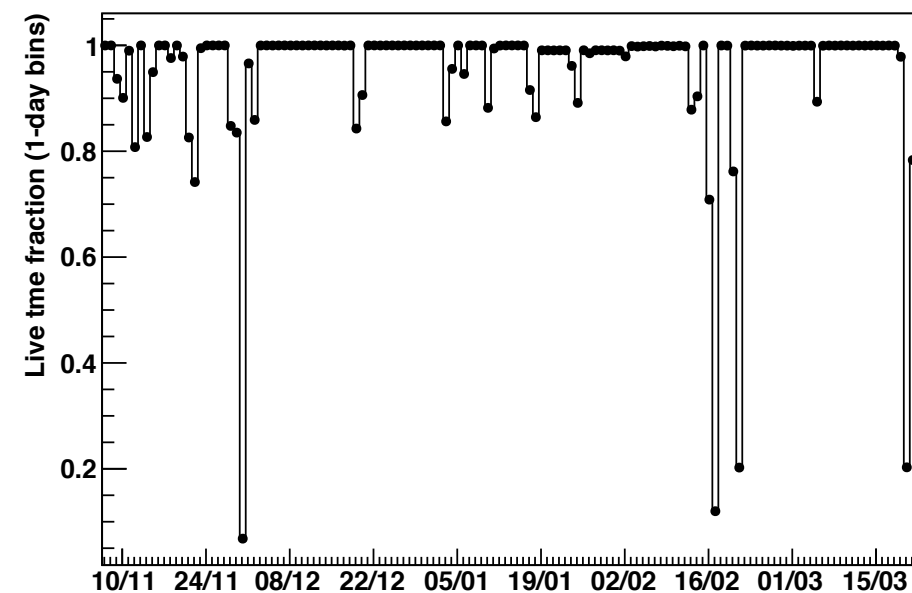
$$\langle m_{\beta\beta} \rangle < 0.26 \sim 0.54 \text{ eV @90\% C.L.}_{15}$$

GERDA

- HPGe detectors in liquid argon (U/Th in LAr $< 7 \times 10^{-4} \mu\text{Bq/kg}$)

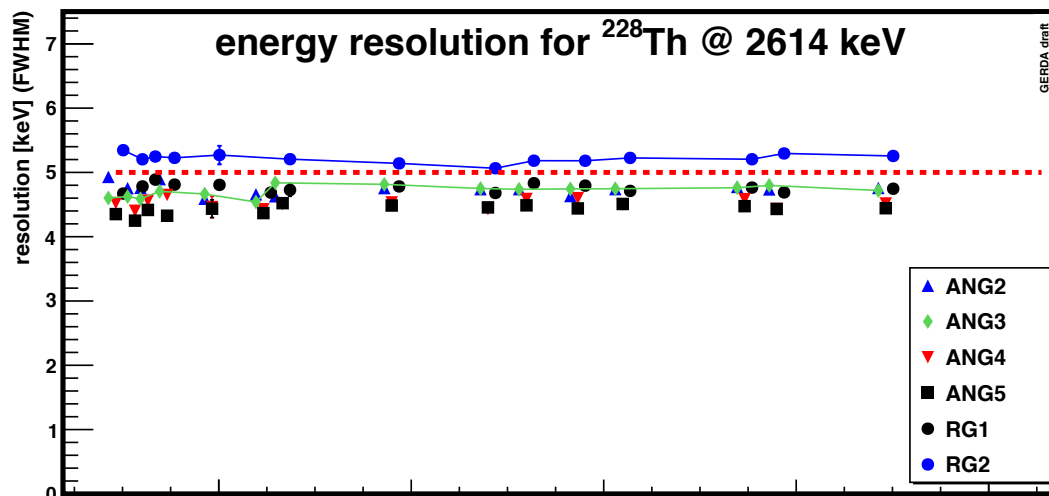
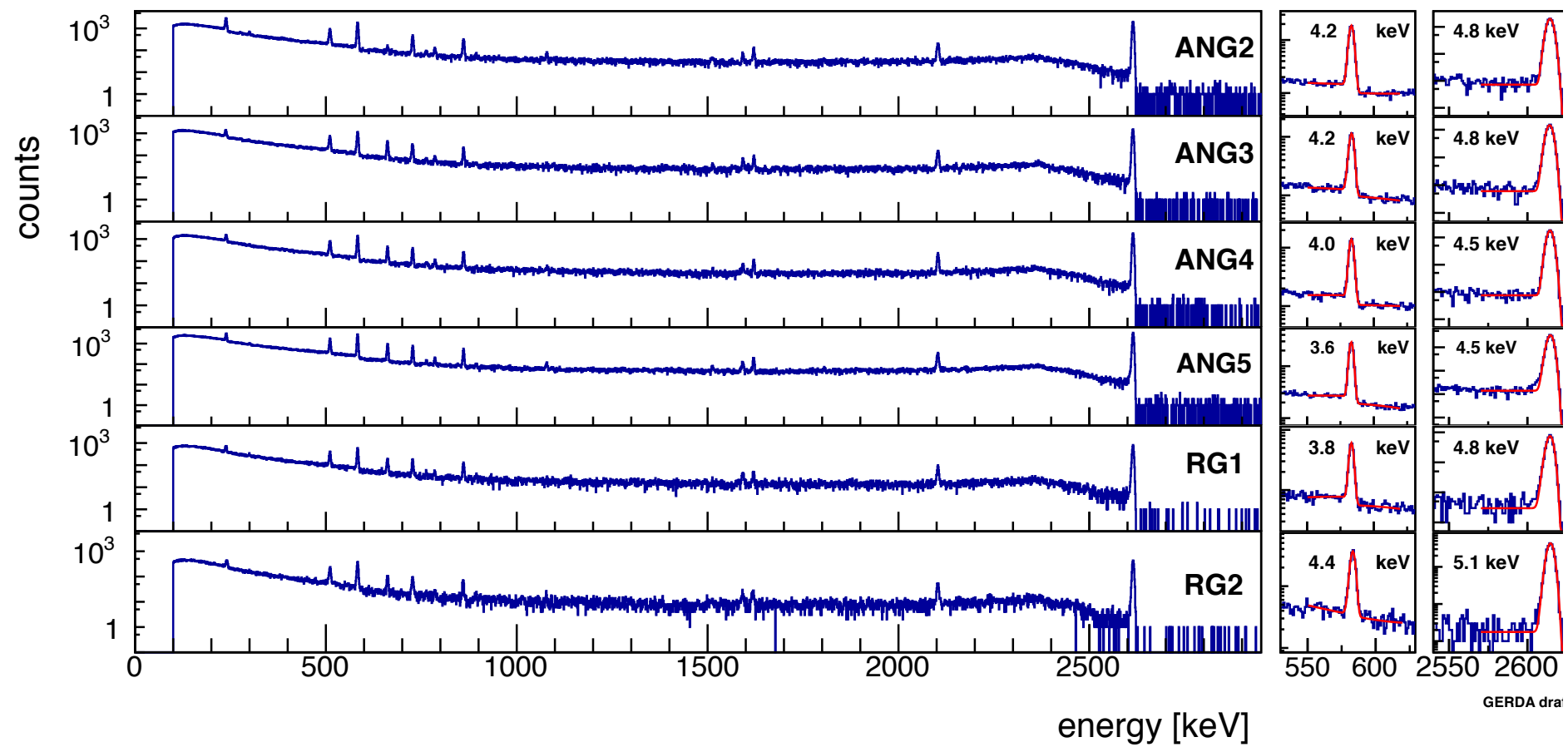


- Physics run started on November 9, 2011



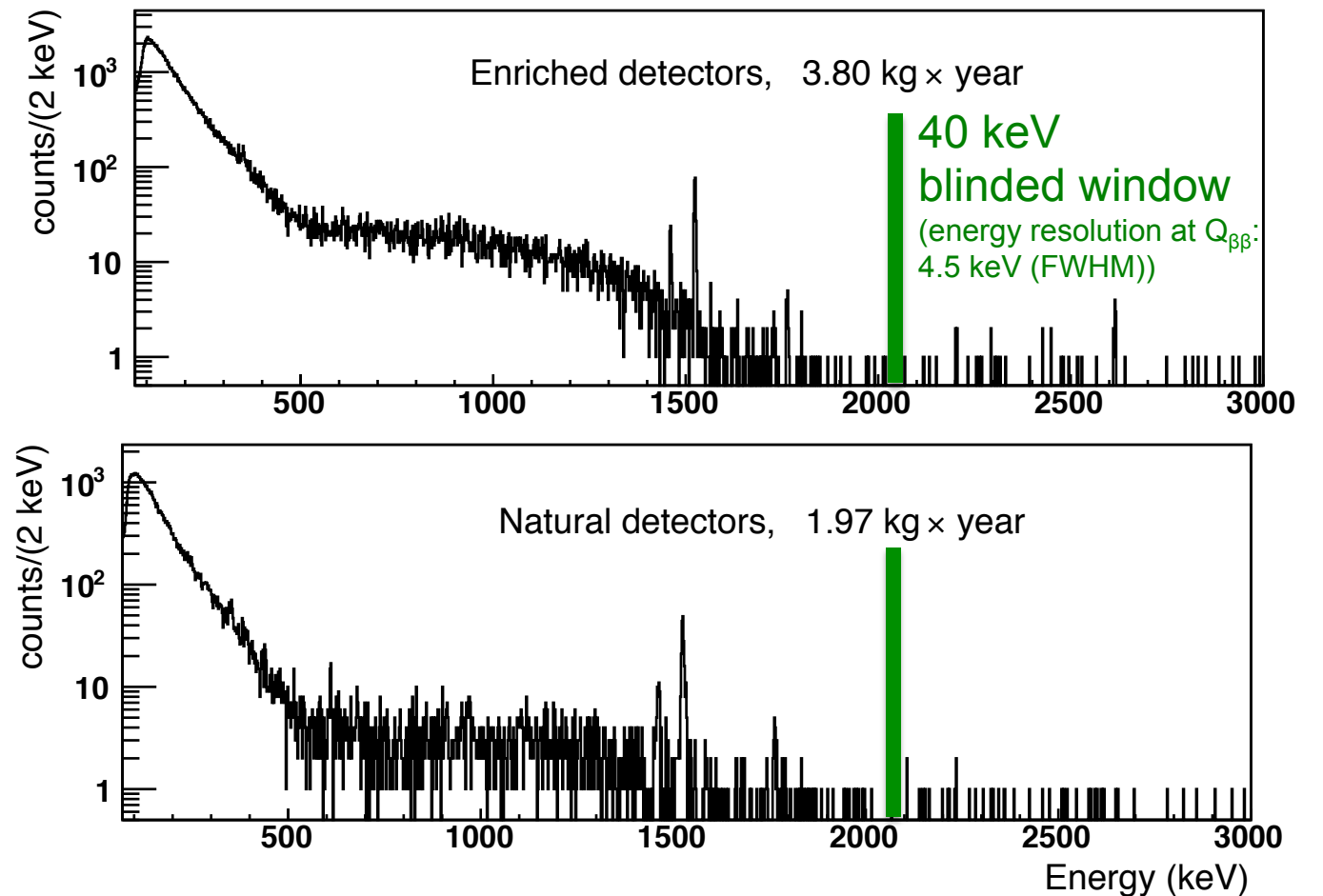
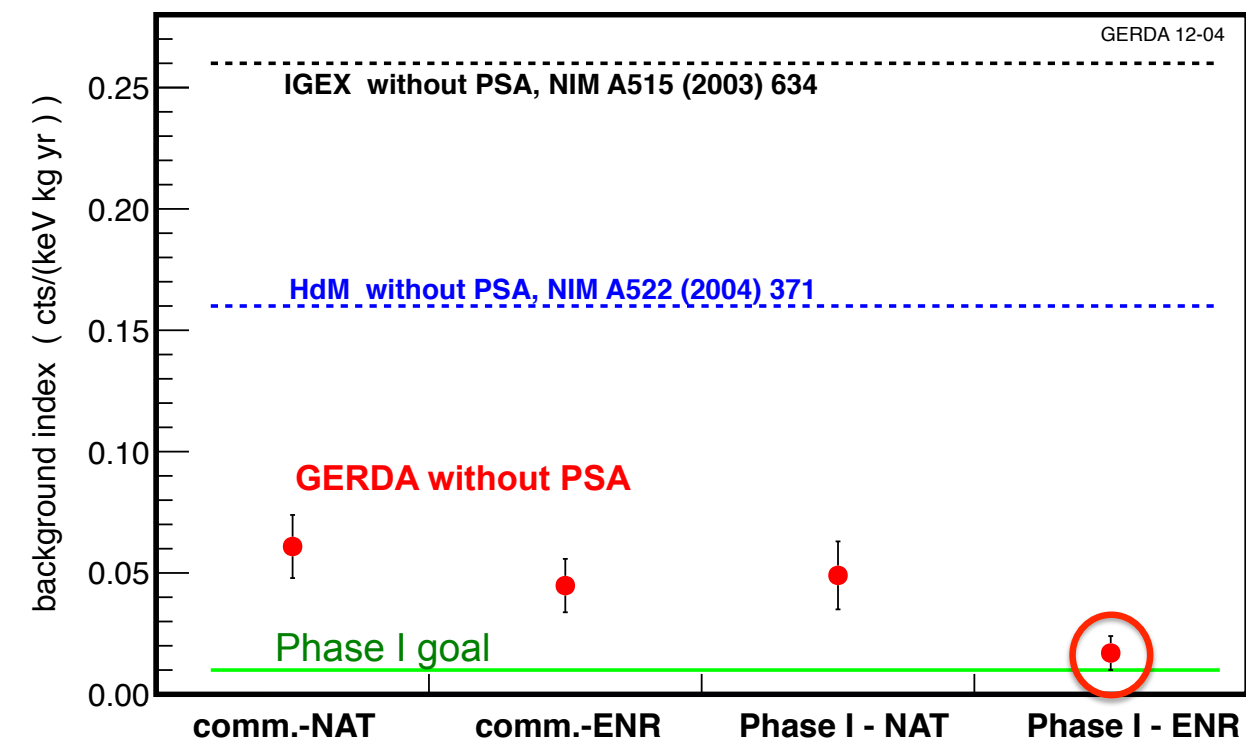
GERDA Calibration

- Energy resolution: $\sim 4.5 - 5$ keV (FWHM) at 2.6 MeV



GERDA low-background spectrum

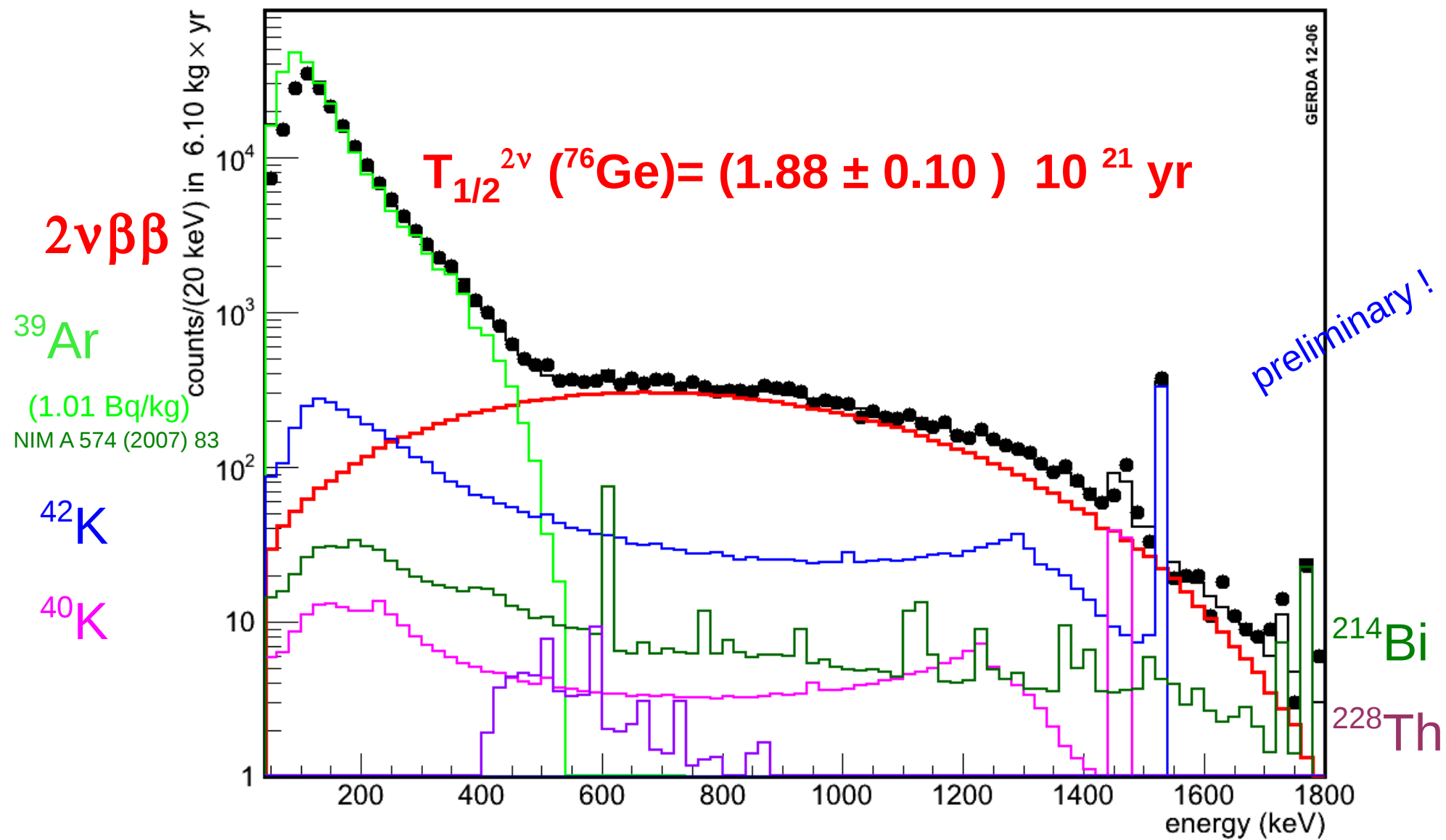
- Background goal of $\sim 10^{-2}$ events/(kg yr keV) was reached
- Phase II (BEGe) detectors in production and testing
- LAr instrumentation (PMTs or SiPM & scintillating fibers) in development
- End of phase I and start of phase II: spring 2013



GERDA low-background spectrum

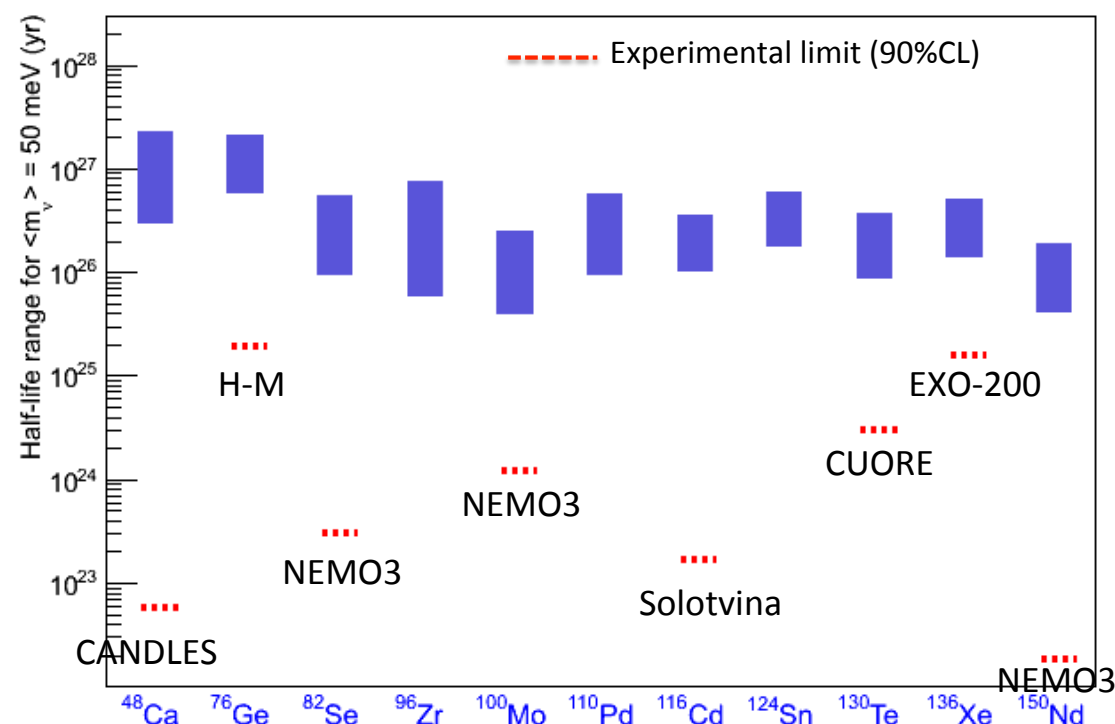
- Analysis of 2-neutrino decay mode is in progress

exposure : 6.1 kg yr



Summary

- Two-neutrino decay mode was measured for the first time in ^{136}Xe
- Xenon experiments provide competitive limits to germanium for the neutrinoless mode
- Several experiments are taking data, new results are expected soon
- Experiments under construction (or phase II of existing experiments) should achieve a sensitivity of 50 - 100 meV
- To go beyond, much lower backgrounds and larger masses are needed
- Tracking will be important to confirm a potential signal



Let us hope that...

- this prediction is true - it could be probed with future double beta experiments!

Tsutomu Yanagida
(Kavli IPMU)

Conclusion

The seesaw with Occam's razor

Frampton, Glashow, Yanagida

CP violation in neutrino oscillation

↔ Universe's baryon asymmetry

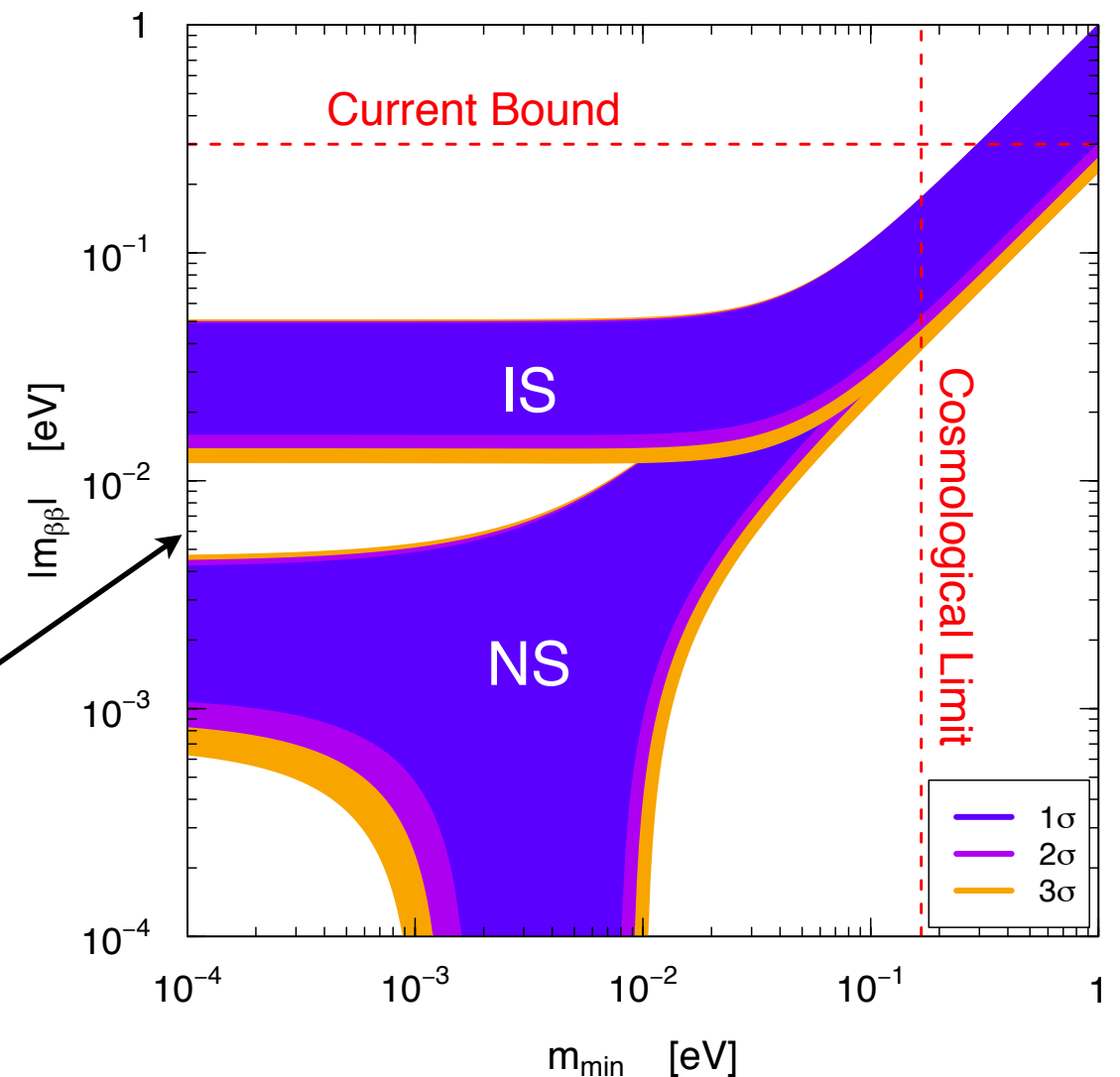
The normal hierarchy is excluded and
it is consistent with the inverted hierarchy !!!

$$|\delta_{CP}| = \frac{\pi}{2} \pm 0.02$$

It predicts

$$m_{ee} = (47 \pm 1) \text{ meV}$$

Bilenky, Giunti: <http://xxx.lanl.gov/abs/1203.5250>



End

Double beta decay

- If simple β^- or β^+ -decay is forbidden on energetic grounds a nucleus can decay through a double beta mode:



- The probability for a decay is very small, the mean lifetime of a nucleus is much larger than the age of the universe ($\tau_U \sim 1.4 \times 10^{10}$ a)

$$\tau_{2\nu} \approx 10^{20} \text{ a}$$

- This is indeed a very rare process (as for instance proton decay, which was not yet observed)
- Nonetheless - if one uses a large amount of nuclei, the process can be observed experimentally

