QCD Critical Point : Inching Towards Continuum

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Introduction

Lattice QCD Results

Searching Experimentally

Summary

* *Work done with Saumen Datta & Sourendu Gupta*

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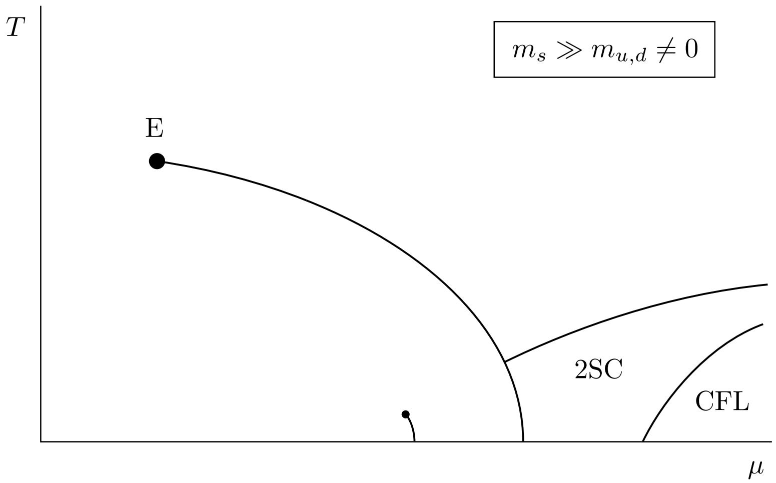

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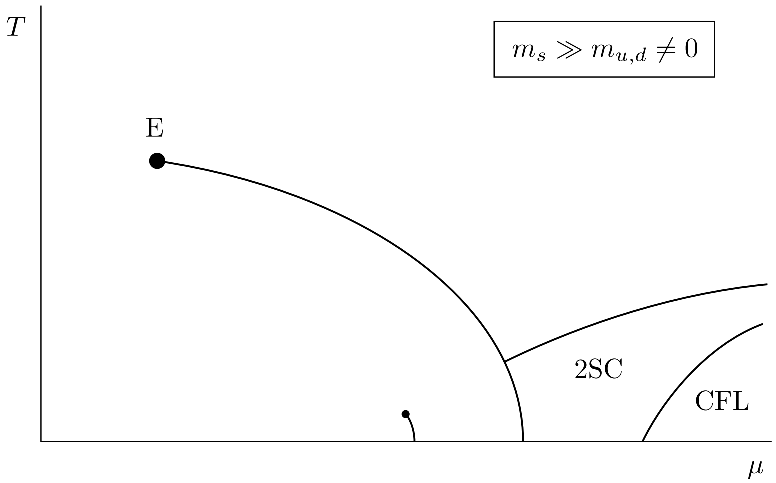

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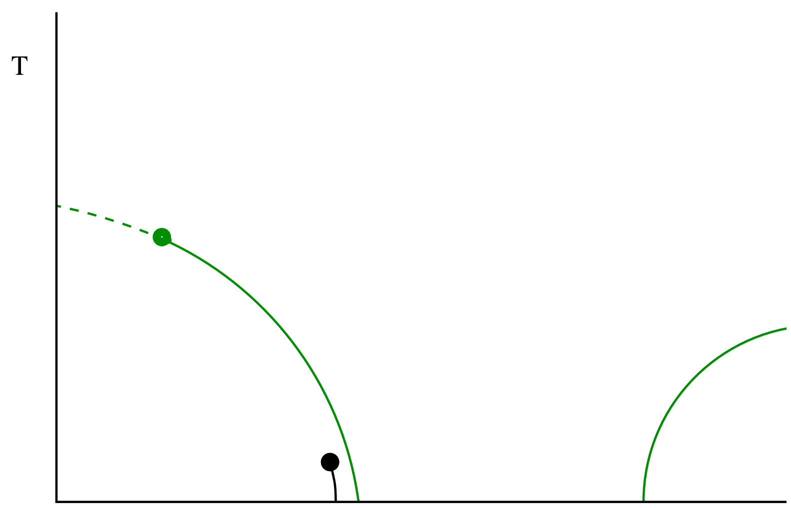

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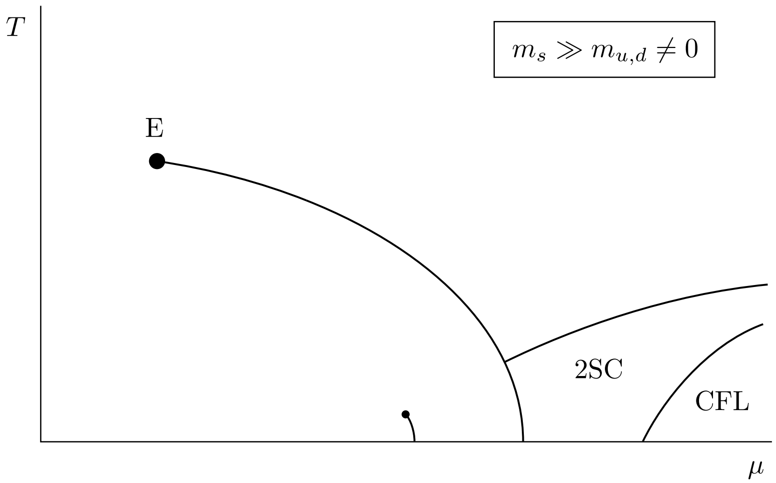

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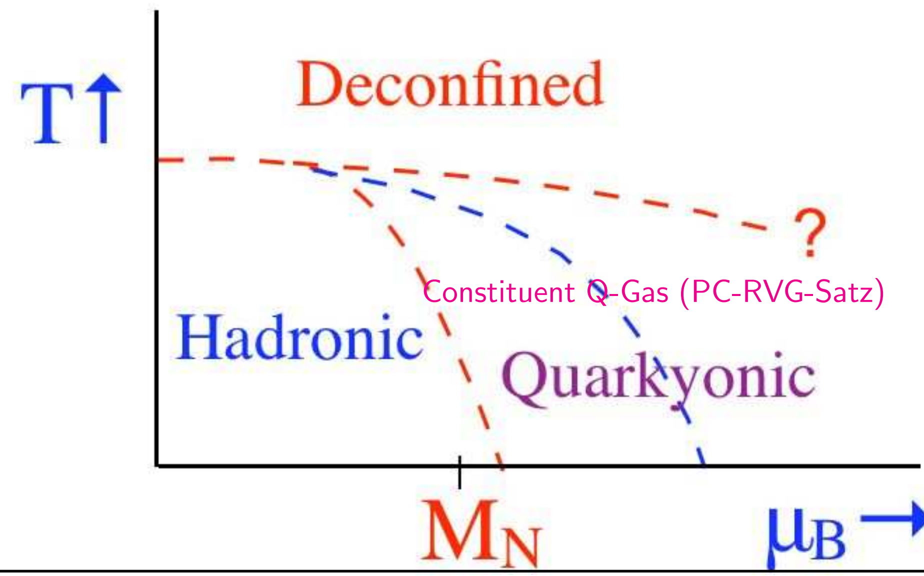

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- And, of course, between the various ongoing (RHIC/STAR) and the proposed/designed experiments (FAIR/CBM).

♡ Maybe it is the synergy between one or more of them that will eventually lead us to the holy grail.

Lattice QCD Results

- Lattice QCD – Most Reliable and Completely parameter-free way to extract non-perturbative physics relevant to Heavy Ion Colliders.

- The Transition Temperature T_c, the Equation of State (used now in 'elliptic flow' analysis), and the Wróblewski Parameter λ_s etc. (Wuppertal-Budapest, HotQCD, GG '02)

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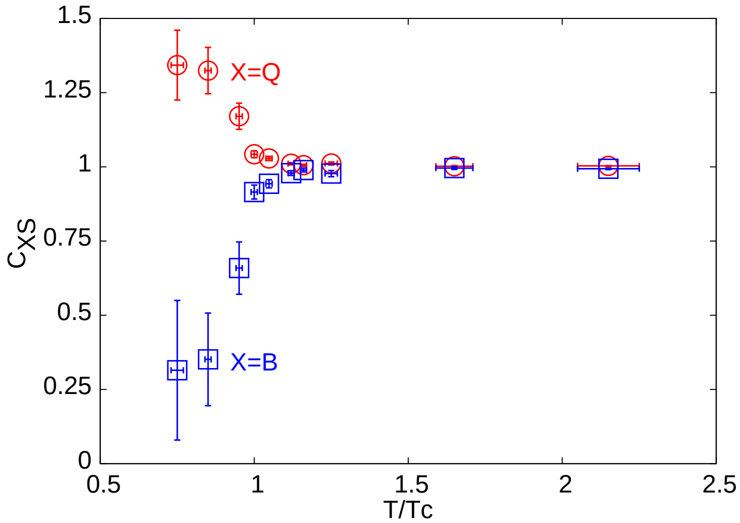
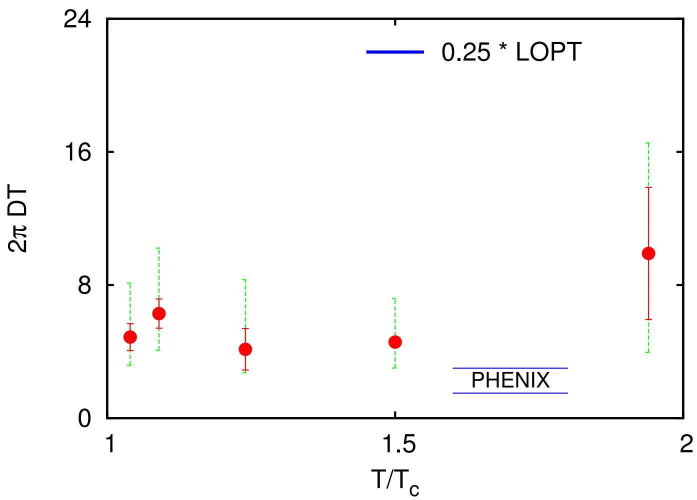
- Flavour Correlations (C_{BS}) and Charm Diffusion Coefficient D are some more such examples for RHIC Physics. (Gavai-Gupta, PRD 2006 & Banerjee et al. PRD 2012)





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- Domain Wall or Overlap Fermions better. BUT Computationally expensive.
- Introduction of μ a la Bloch & Wettig (PRL 2006 & PRD2007)
- Unfortunately breaks chiral symmetry ! (Banerjee, Gavai & Sharma PRD 2008; PoS (Lattice 2008); PRD 2009)
- Good News : Problem Solved !
  Overlap Lattice Action with exact chiral invariance at nonzero μ and any a now exists (Gavai & Sharma , arXiv : 1111.5944; PLB in press, Narayanan-Sharma JHEP '11).

- Using chiral projectors for overlap fermions as, $\psi_L = [1 - \gamma_5(1 - aD_{ov})]\psi/2$ & $\psi_R = [1 + \gamma_5(1 - aD_{ov})]\psi/2$, leaving the antiquark field decomposition as in the continuum, the overlap action for nonzero μ is

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\begin{aligned}
S^F &= \sum_n [\bar{\psi}_{n,L}(aD_{ov} + a\mu\gamma^4)\psi_{n,L} + \bar{\psi}_{n,R}(aD_{ov} + a\mu\gamma^4)\psi_{n,R}] \\
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$$

- Easy to check that under the chiral transformations, $\delta\psi = i\alpha\gamma_5(1 - aD_{ov})\psi$ and $\delta\bar\psi = i\alpha\bar\psi\gamma_5$, it is invariant or all values of $a\mu$ and a.

- Order parameter exists for all μ and T. It is $\langle\bar\psi\psi\rangle = \lim_{am\to 0}\lim_{V\to\infty}\langle \mathrm{Tr}\ \frac{(1-aD_{ov}/2)}{[aD_{ov}+(am+a\mu\gamma^4)(1-aD_{ov}/2)]}\rangle$.

The $\mu \neq 0$ problem : The Measure

Simulations can be done IF Det $M > 0$. However, det M is a complex number for any $\mu \neq 0$: The Phase/sign problem

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Several Approaches proposed in the past two decades :

- Two parameter Re-weighting (Z. Fodor & S. Katz, JHEP 0203 (2002) 014).
- Imaginary Chemical Potential (Ph. de Forcrand & O. Philipsen, NP B642 (2002) 290; M.-P. Lombardo & M. D'Elia PR D67 (2003) 014505).
- Taylor Expansion (C. Allton et al., PR D66 (2002) 074507 & D68 (2003) 014507; R.V. Gavai and S. Gupta, PR D68 (2003) 034506).
- Canonical Ensemble (K. -F. Liu, IJMP B16 (2002) 2017, S. Kratochvila and P. de Forcrand, Pos LAT2005 (2006) 167.)
- Complex Langevin (G. Aarts and I. O. Stamatescu, arXiv:0809.5227 and its references for earlier work).

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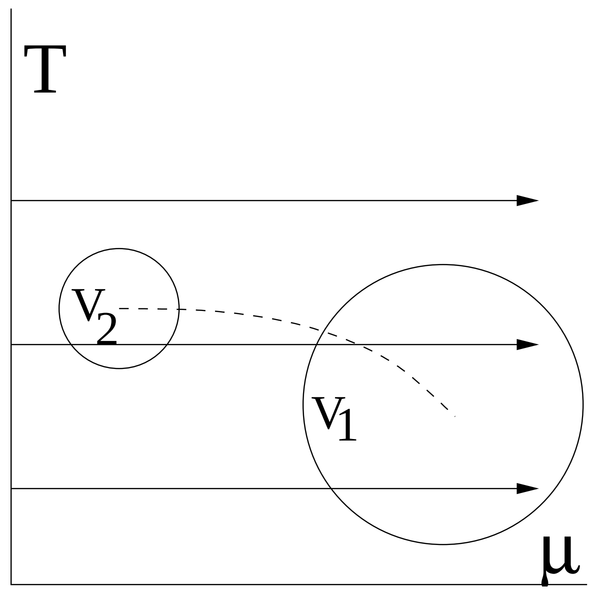

We studied volume dependence at several T to i) bracket the critical region and then ii) tracked its change as a function of volume.

Details of Expansion

Standard definitions yield various number densities and susceptibilities :

$$n_i = \frac{T}{V}\frac{\partial \ln \mathcal{Z}}{\partial \mu_i} \text{ and } \qquad \chi_{ij} = \frac{T}{V}\frac{\partial^2 \ln \mathcal{Z}}{\partial \mu_i \partial \mu_j} \quad .$$

These are also useful by themselves both theoretically and for Heavy Ion Physics (Flavour correlations, $\lambda_s \ldots$)

Denoting higher order susceptibilities by χ_{n_u,n_d}, the pressure P has the expansion in μ:

$$\frac{\Delta P}{T^4} \equiv \frac{P(\mu,T)}{T^4} - \frac{P(0,T)}{T^4} = \sum_{n_u,n_d} \chi_{n_u,n_d} \frac{1}{n_u!}\left(\frac{\mu_u}{T}\right)^{n_u} \frac{1}{n_d!}\left(\frac{\mu_d}{T}\right)^{n_d}$$

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- From this expansion, a series for baryonic susceptibility can be constructed. Its radius of convergence gives estimate of the location of nearest critical point.

- Successive estimates for the radius of convergence obtained from these using $\sqrt{\frac{n(n+1)\chi_B^{(n+1)}}{\chi_B^{(n+3)}T^2}}$ or $\left(n!\frac{\chi_B^{(2)}}{\chi_B^{(n+2)}T^2}\right)^{1/n}$. We use both and terms up to 8th order in μ.

- All coefficients of the series must be POSITIVE for the critical point to be at real μ, and thus physical.

- We (Gavai-Gupta '05, '09) use up to 8^{th} order. B-RBC so far has up to 6^{th} order.

- 10th & even 12th order may be possible : Ideas to extend to higher orders are emerging (Gavai-Sharma PRD 2012 & PRD 2010) which save up to 60 % computer time.

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All susceptibilities can be written as traces of products of M^{-1} and various derivatives of M.

At leading order,

$$\chi_{20} = \left(\frac{T}{V}\right) [\langle \mathcal{O}_2 + \mathcal{O}_{11} \rangle], \qquad \chi_{11} = \left(\frac{T}{V}\right) [\langle \mathcal{O}_{11} \rangle]$$

Here $\mathcal{O}_2 = \mathrm{Tr}\ M^{-1}M'' - \mathrm{Tr}\ M^{-1}M'M^{-1}M'$, and $\mathcal{O}_{11} = (\mathrm{Tr}\ M^{-1}M')^2$, and the traces are estimated by a stochastic method (Gottlieb et al., PRL '87):

$\mathrm{Tr}\ A = \sum_{i=1}^{N_v} R_i^\dagger A R_i / 2N_v$, and $(\mathrm{Tr}\ A)^2 = 2\sum_{i>j=1}^{L} (\mathrm{Tr}\ A)_i (\mathrm{Tr}\ A)_j / L(L-1)$, where R_i is a complex vector from a set of N_v subdivided in L independent sets.

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Higher order NLS are more involved. E.g., at the 8th order, terms involve operators up to $\mathcal{O}_8$ which in turn have terms up to 8 quark propagators and combinations of M' and M''.

In fact, the entire evaluation of the χ_{80} needs 20 inversions of Dirac matrix.

This can be reduced to 8 inversions using an action linear in μ (Gavai-Sharma PRD 2012 & PRD 2010), leading still to results in agreement with that exponential in μ.

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Our Simulations & Results

- Staggered fermions with $N_f = 2$ of $m/T_c = 0.1$; R-algorithm used.
- $m_\pi/m_\rho = 0.31 \pm 0.01$ (MILC); Kept the same as $a \to 0$ (on all N_t).
- Earlier Lattice : 4 $\times N_s^3$, $N_s =$ 8, 10, 12, 16, 24 (Gavai-Gupta, PRD 2005)
  Finer Lattice : 6 $\times N_s^3$, $N_s =$ 12, 18, 24 (Gavai-Gupta, PRD 2009).

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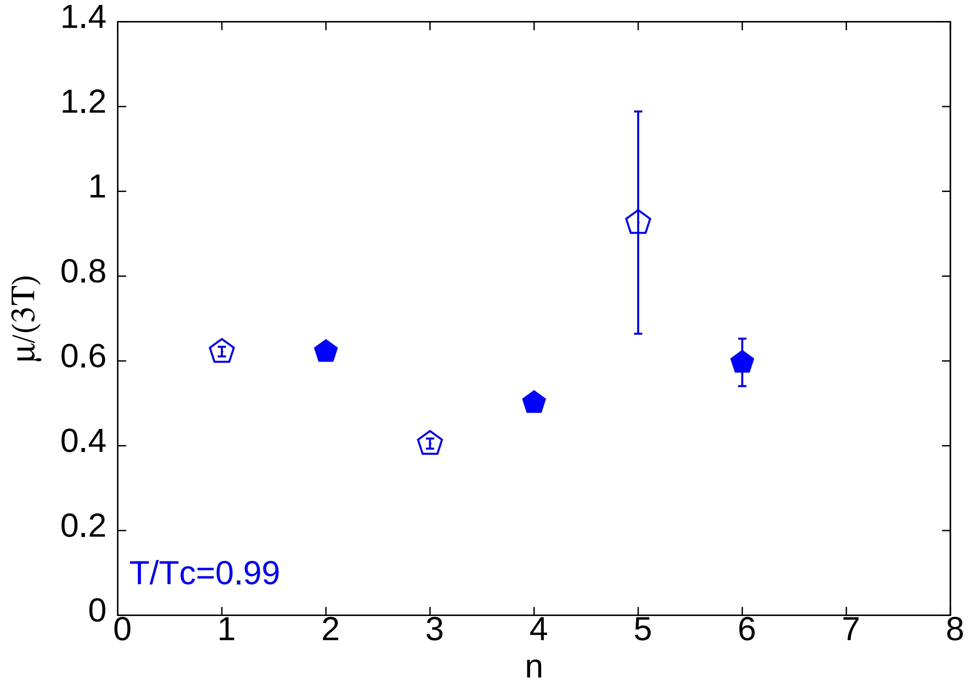
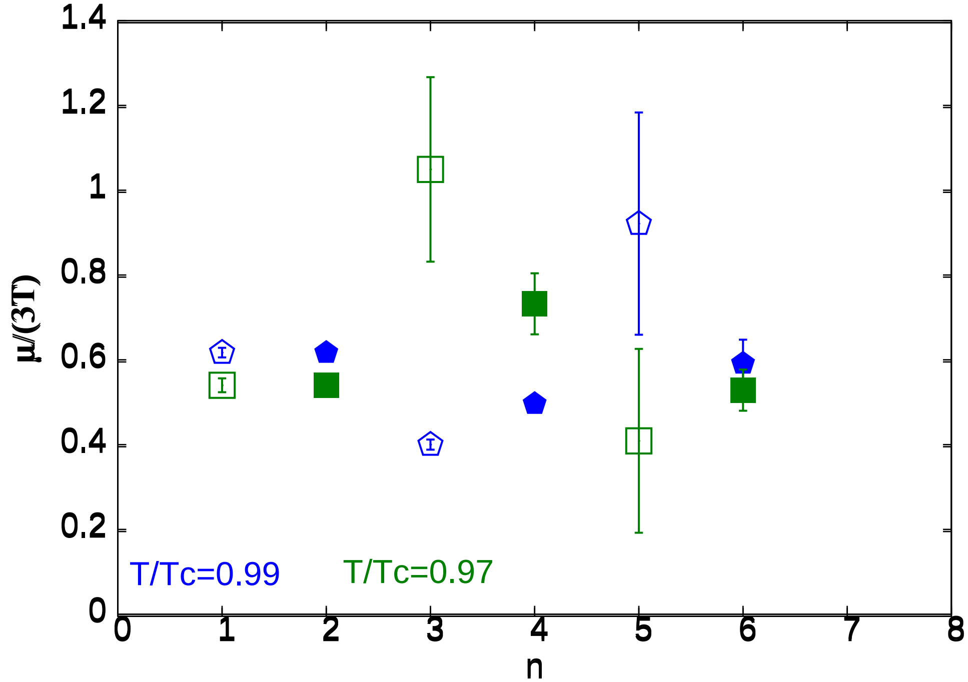
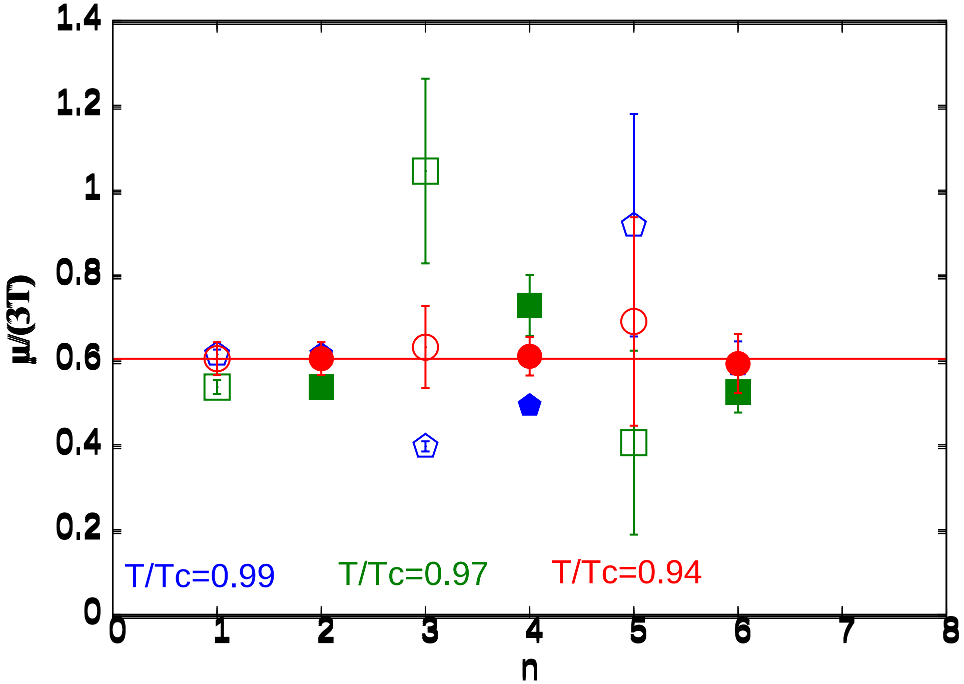
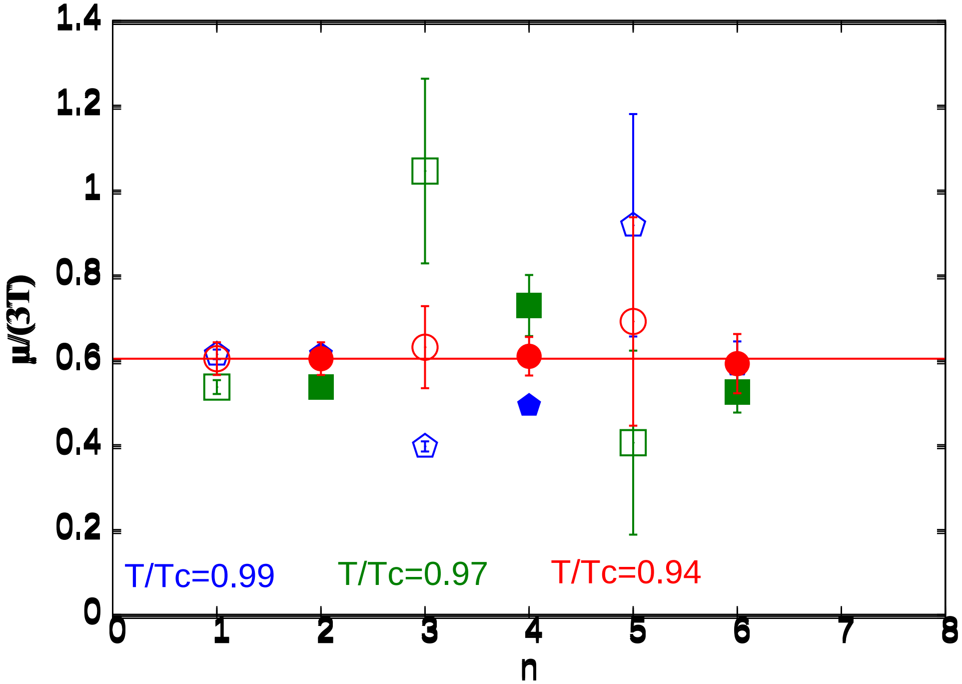
  Finer Lattice : 6 $\times N_s^3$, $N_s =$ 12, 18, 24 (Gavai-Gupta, PRD 2009).
- Even finer Lattice : 8 $\times 32^3$ — This Talk (Datta-RVG-Gupta, '12)
  Aspect ratio, N_s/N_t, maintained four to reduce finite volume effects.
- Simulations made at $T/T_c =$ 0.90, 0.92, 0.94, 0.96, 0.98, 1.00, 1.02, 1.12, 1.5 and 2.01. Typical stat. 100-200 in max autocorrelation units.
- T_c — defined by the peak of Polyakov loop susceptibility.



μ/(3T)
n
T/Tc=0.99


μ/(3T)
n
T/Tc=0.99
T/Tc=0.97


1.4
1.2
1
0.8
0.6
0.4
0.2
0
μ/(3T)
0
1
2
3
4
5
6
7
8
n
T/Tc=0.99
T/Tc=0.97
T/Tc=0.94



- $\frac{T^E}{T_c} = 0.94 \pm 0.01$, and $\frac{\mu_B^E}{T^E} = 1.8 \pm 0.1$ for finer lattice: Our earlier coarser lattice result was $\mu_B^E/T^E = 1.3 \pm 0.3$. Infinite volume result: $\downarrow$ to 1.1(1)

- Critical point at $\mu_B/T \sim 1 - 2$.

Cross Check on μ^E/T^E

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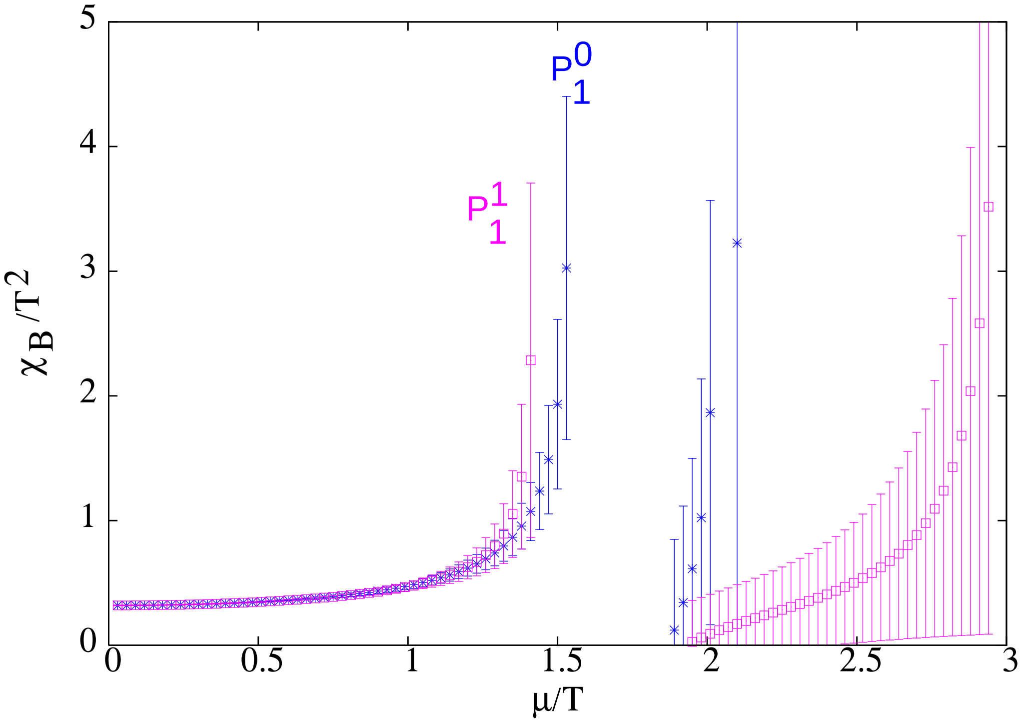
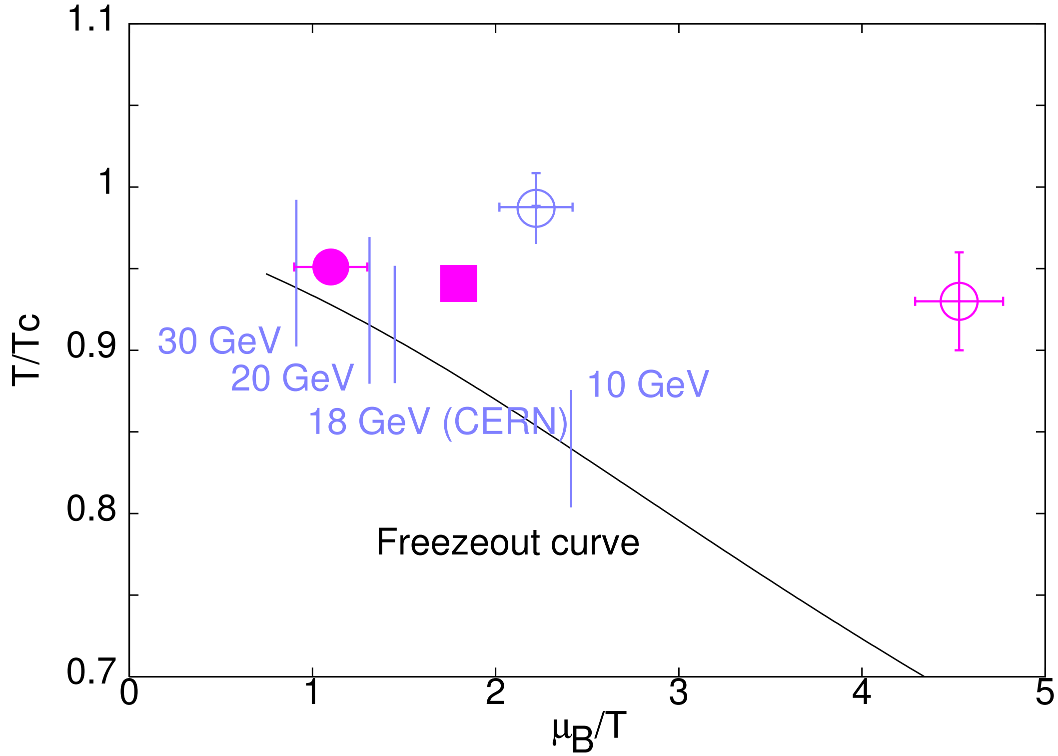


♡ Consistent Window with our other estimates.

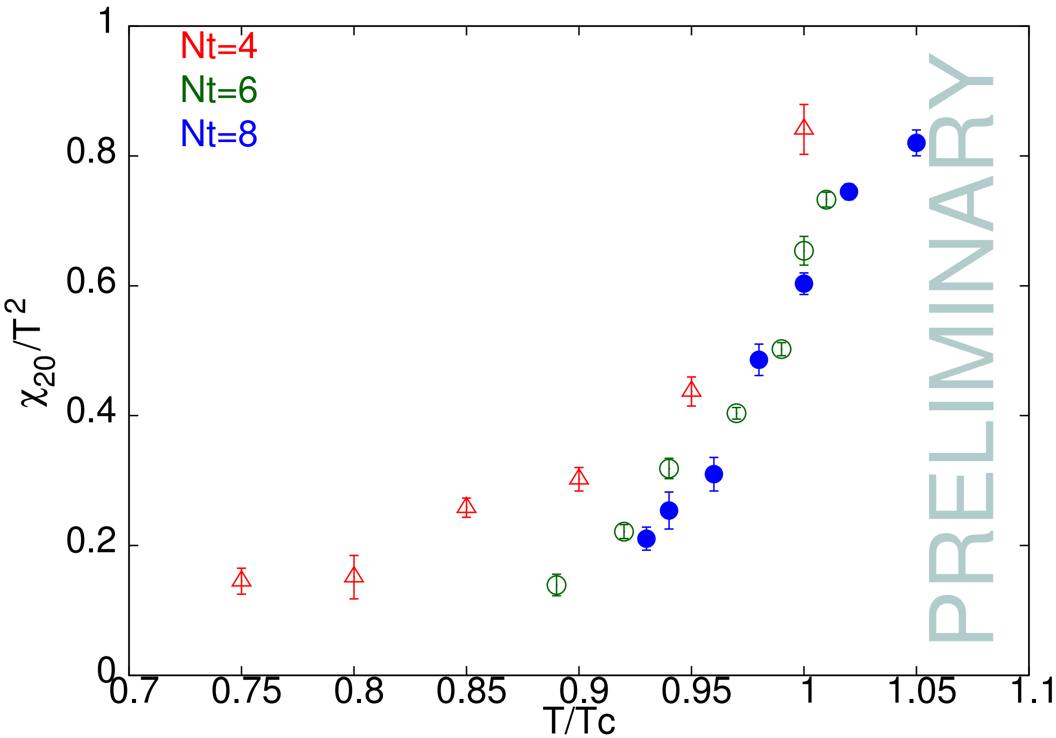
Critical Point : Story thus far



♠ $N_f = 2$ (magenta) and 2+1 (blue) (Fodor-Katz, JHEP '04).

♡ $N_t = 4$ Circles (GG '05 & Fodor-Katz JHEP '02), $N_t = 6$ Box (GG '09).

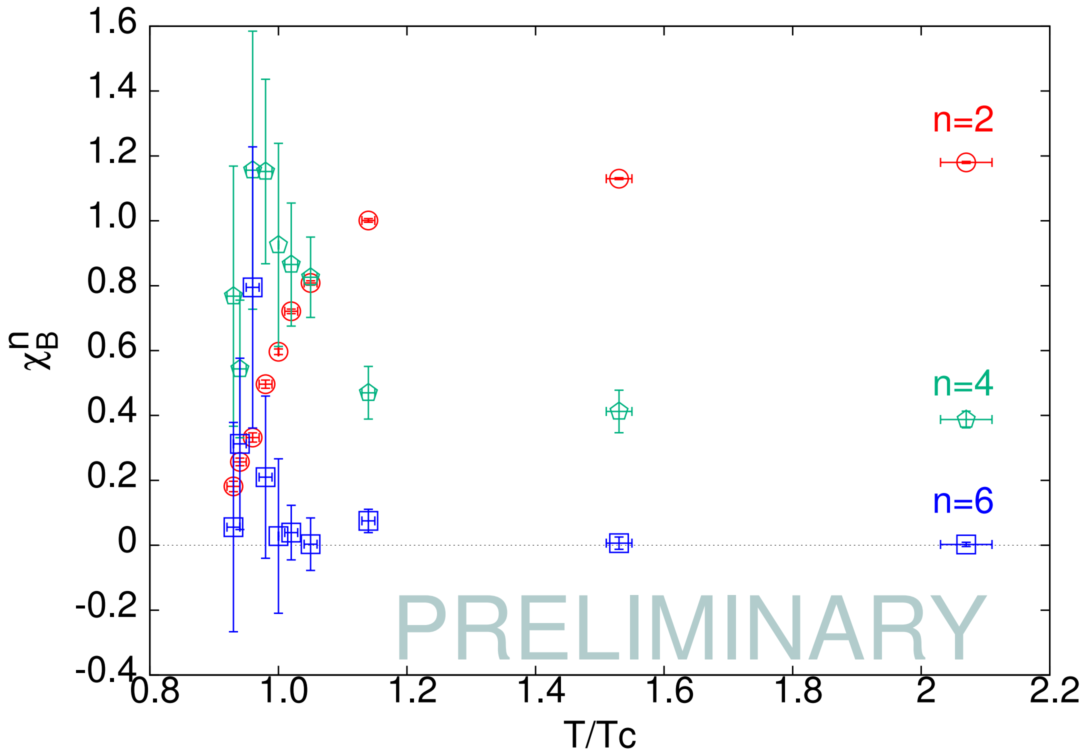
χ_2 for $N_t = 8$, 6, and 4 lattices



♠ $N_t = 8$ (Datta-Gavai-Gupta, QM12) and 6 (GG, PRD '09) results agree.

♡ $\beta_c(N_t = 8)$ agrees with Gottlieb et al. PR D47,1993.

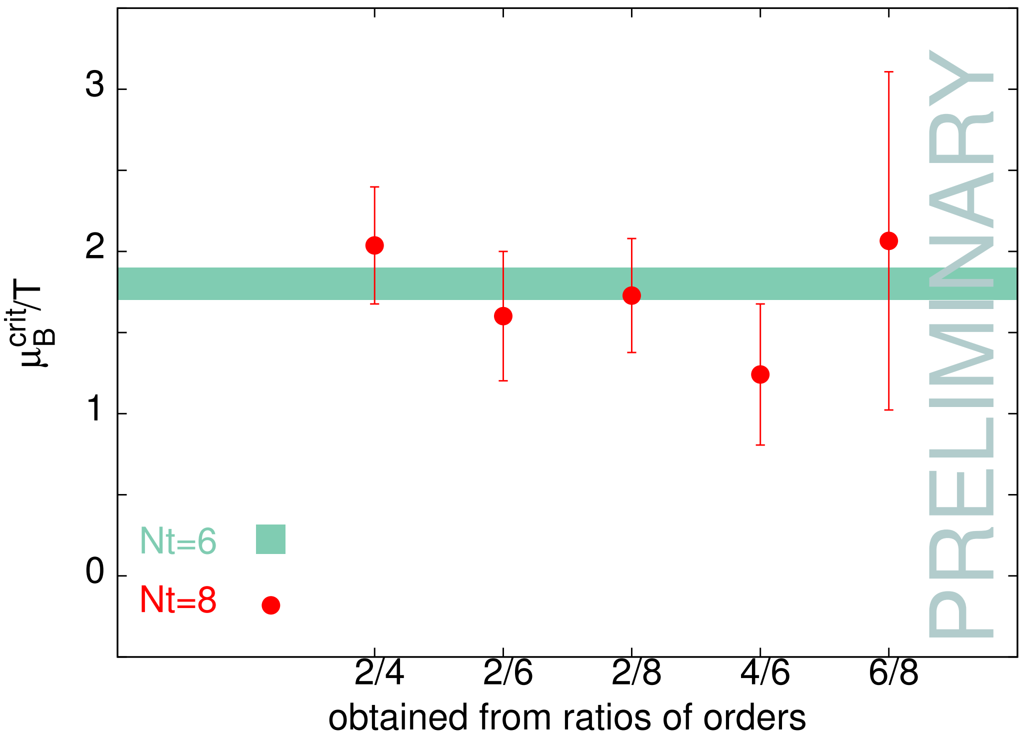
χ_B^n for $N_t = 8$ lattice



♠ 100 configurations & 1000 vectors at each point employed.

♡ More statistics coming in critical region. Window of positivity in anticipated region.

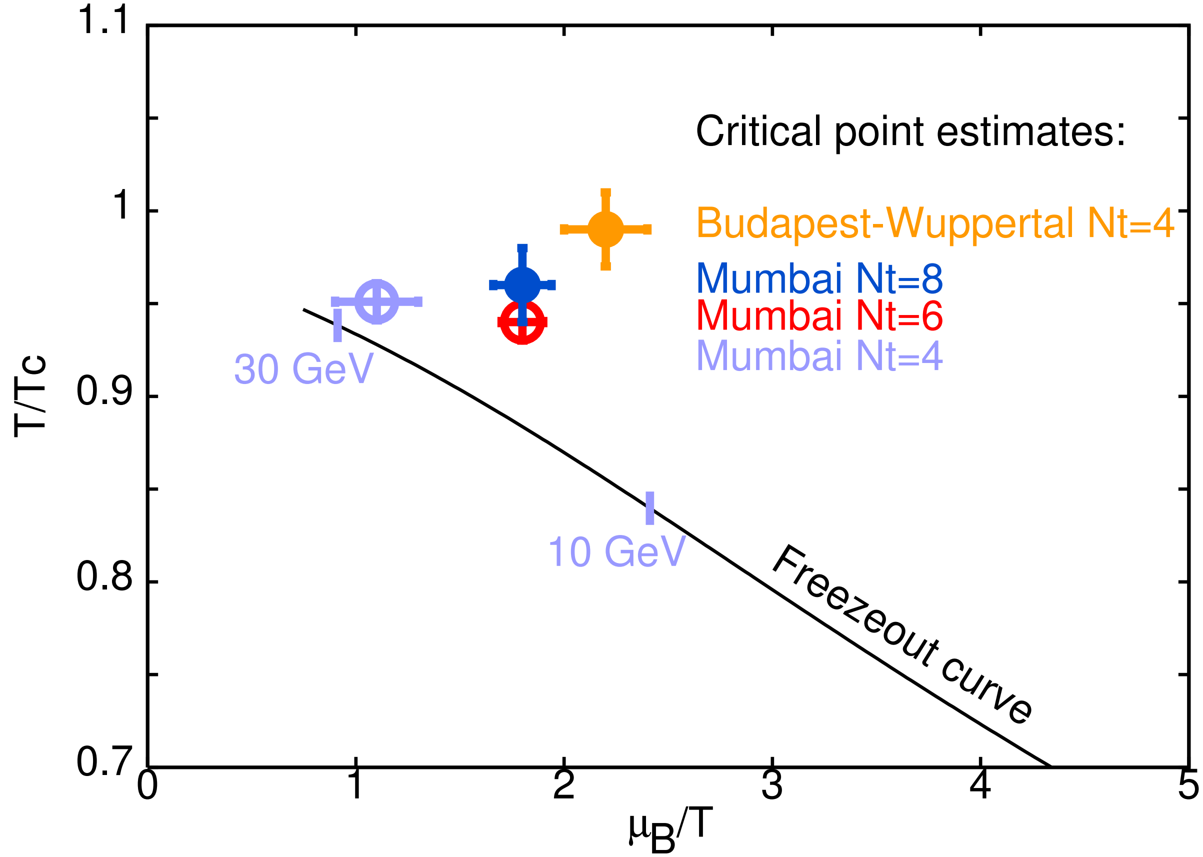
Radius of Convergence result



♠ At our (T_E, μ_E) for $N_t = 6$, the ratios display constancy for $N_t = 8$ as well.

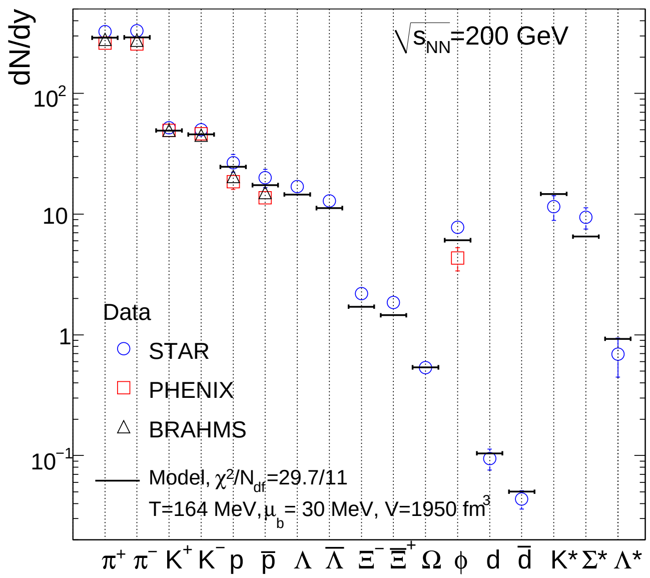
♡ Currently : Similar results at neighbouring $T/T_c \Longrightarrow$ a larger ΔT at same μ_B^E.

Critical Point : Inching Towards Continuum



Lattice predictions along the freezeout curve

- Hadron yields well described using Statistical Models, leading to a freezeout curve in the T-μ_B plane. (Andronic, Braun-Munzinger & Stachel, PLB 2009 ; Oeschler, Cleymans, Redlich & Wheaton, 2009)



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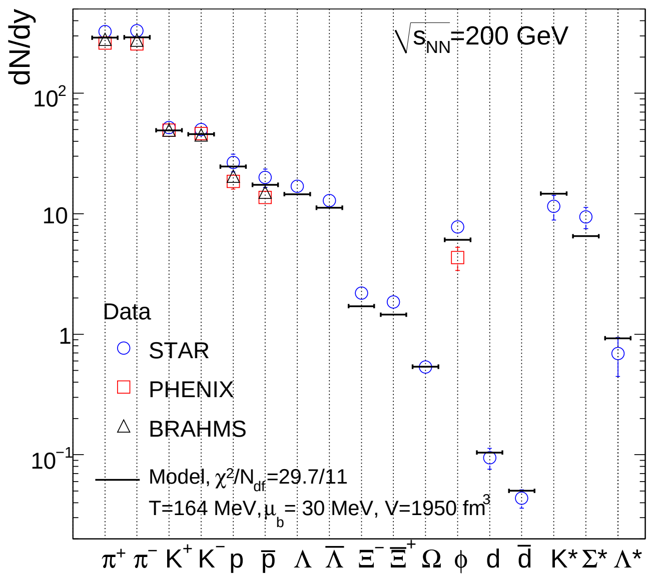

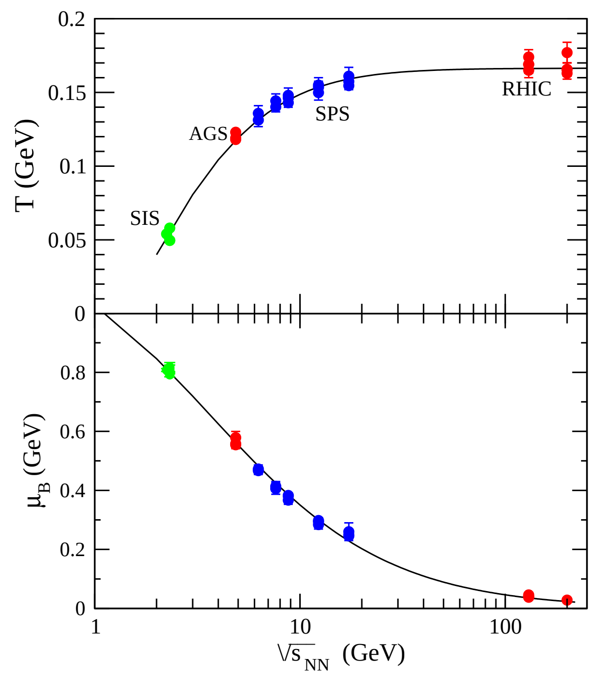

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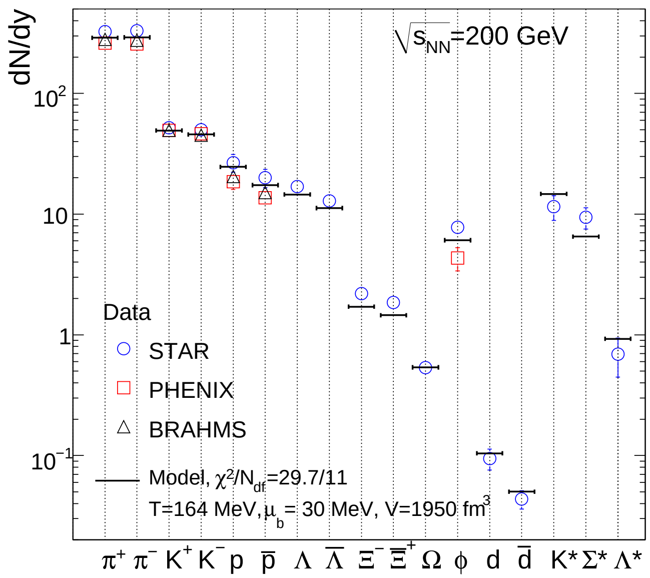
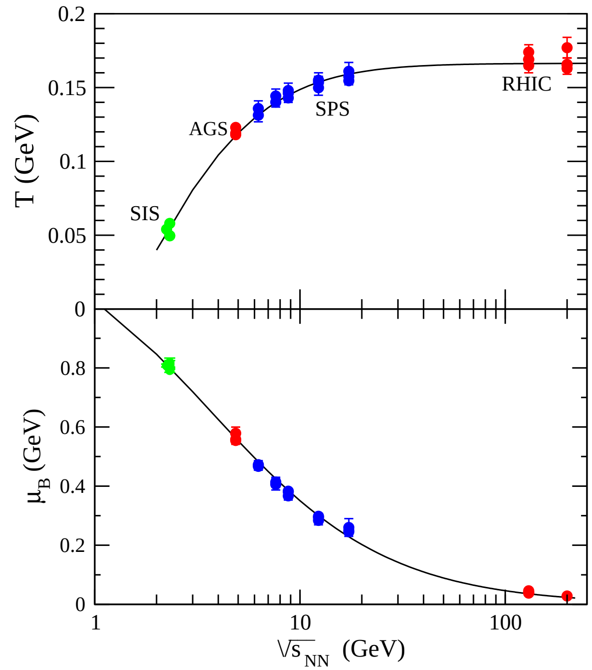
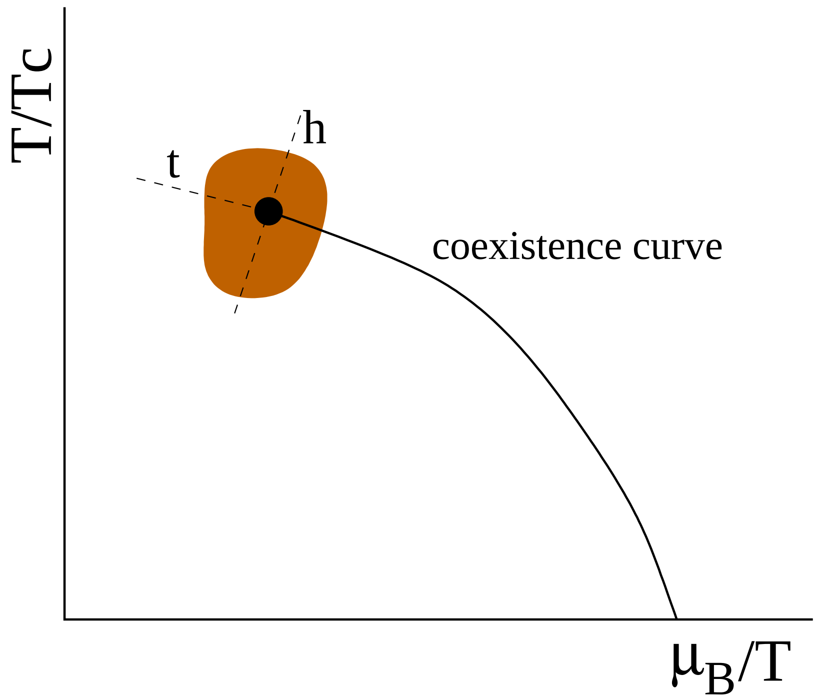
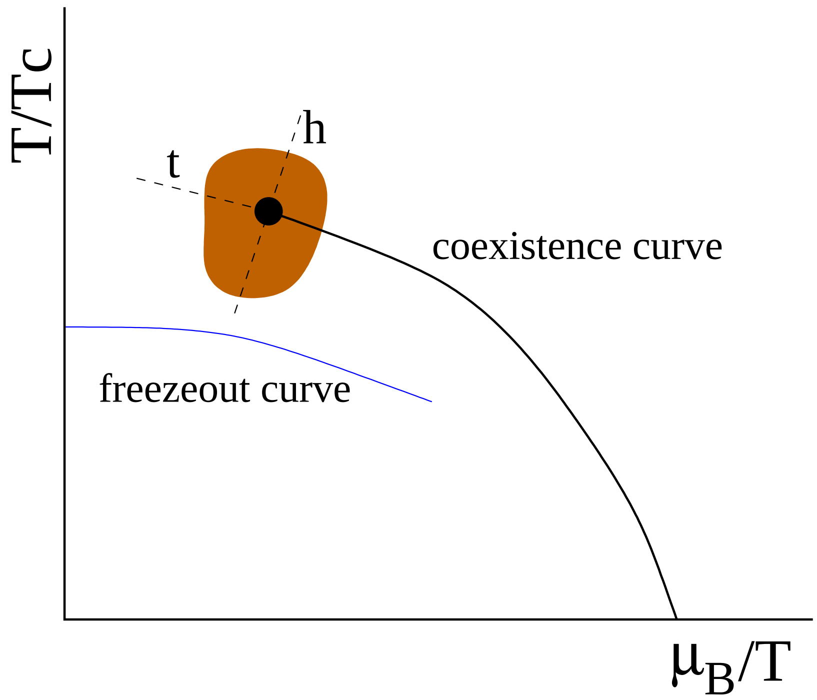
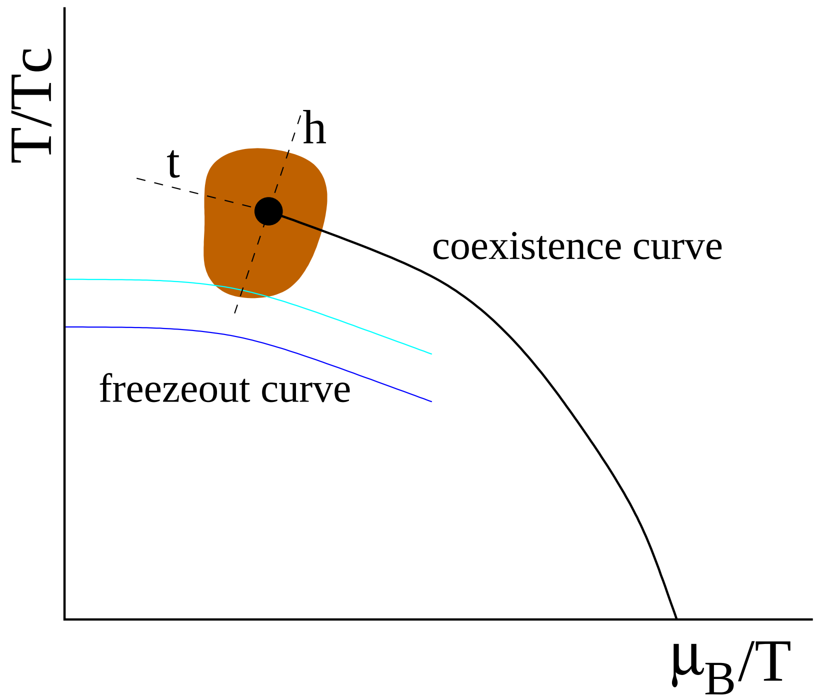




- Our Key Proposal : Use the freezeout curve from hadron abundances to *predict* fluctuations using lattice QCD along it. (Gavai-Gupta, TIFR/TH/10-01, arXiv 1001.3796)


T/Tc
h
t
coexistence curve
μB/T


T/Tc
h
t
coexistence curve
freezeout curve
μ_B/T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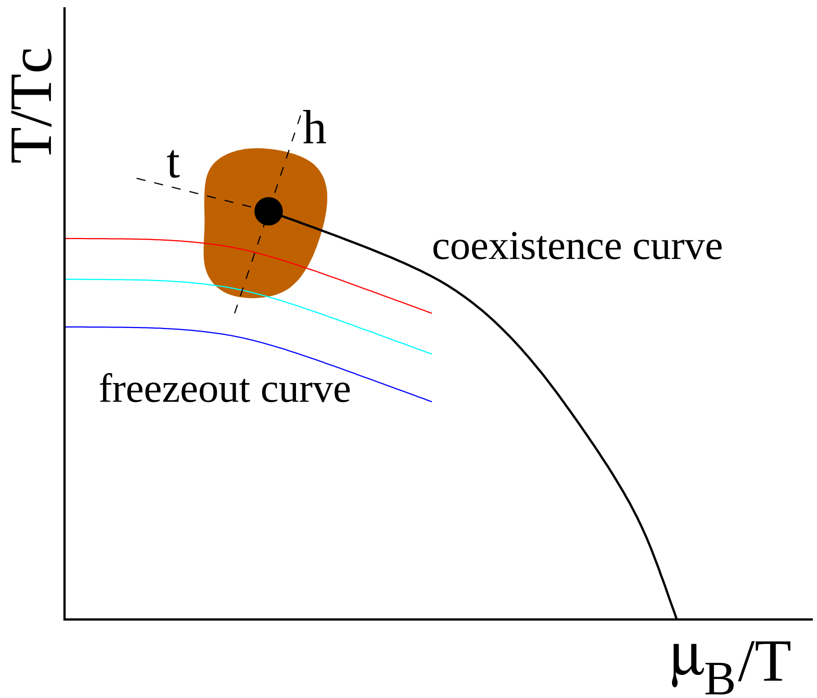
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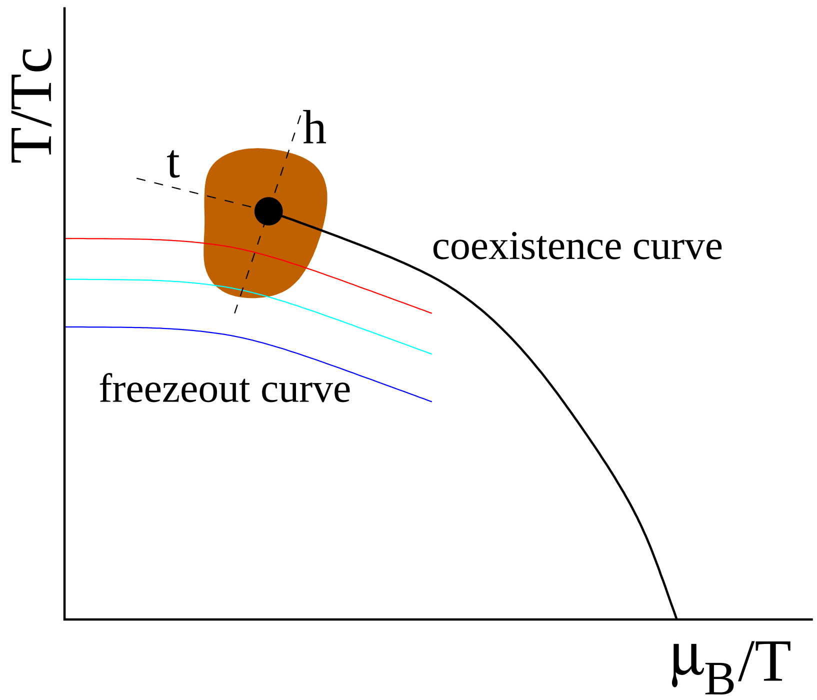
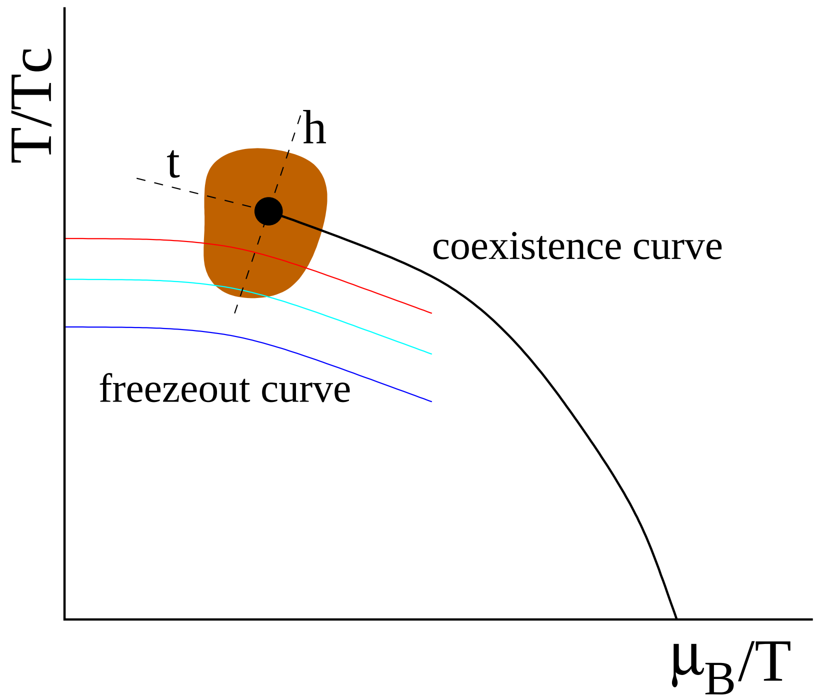


- Use the freezeout curve to relate $(T, \mu_B)$to $\sqrt{s}$ and employ lattice QCD predictions along it. (Gavai-Gupta, TIFR/TH/10-01, arXiv 1001.3796)



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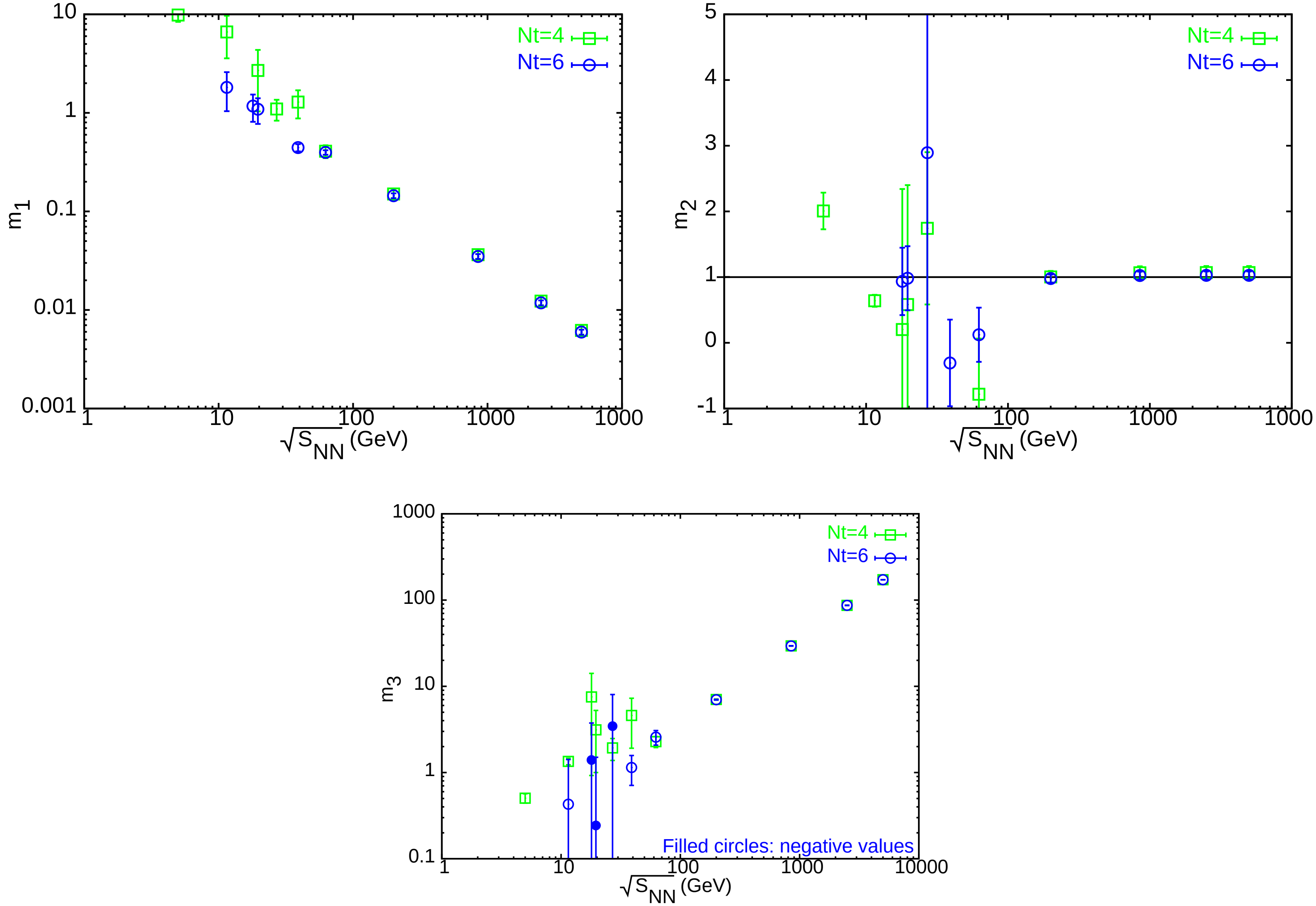
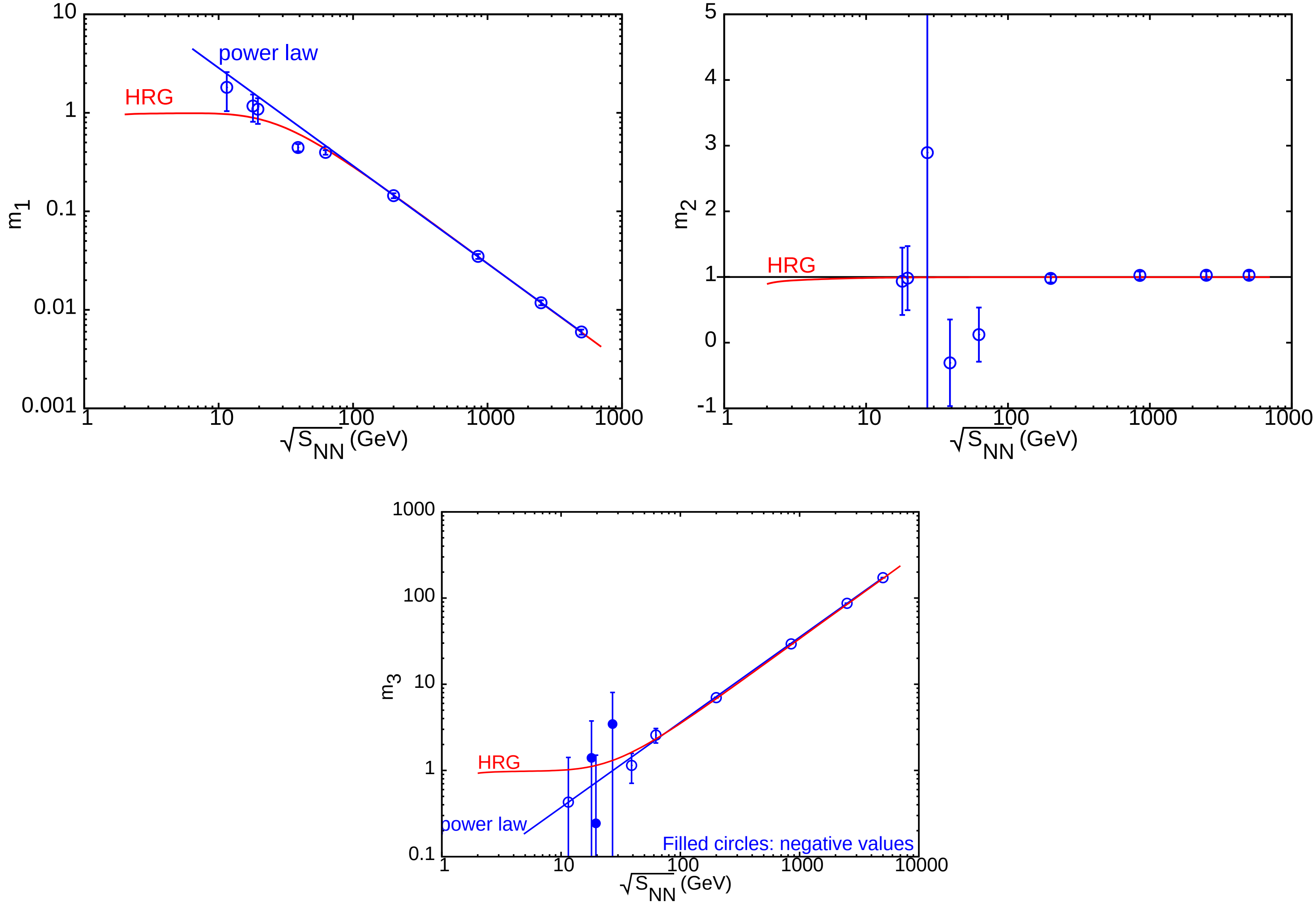
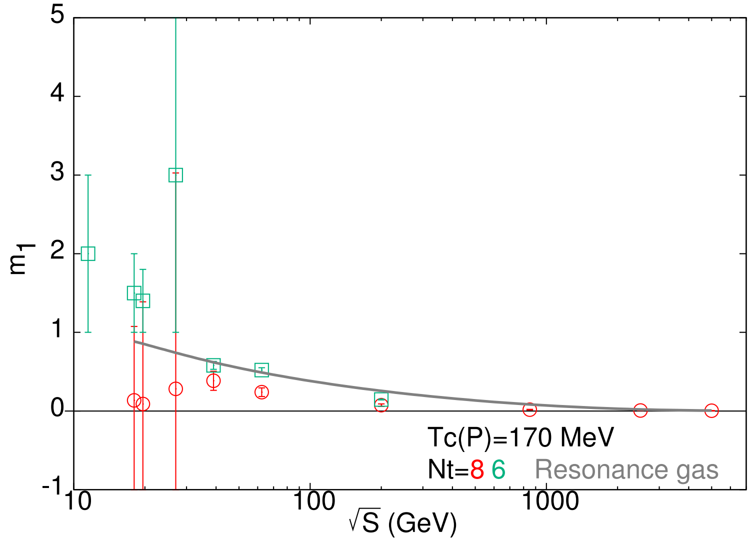
- Define $m_1 = \frac{T\chi^{(3)}(T,\mu_B)}{\chi^{(2)}(T,\mu_B)}$, $m_3 = \frac{T\chi^{(4)}(T,\mu_B)}{\chi^{(3)}(T,\mu_B)}$, and $m_2 = m_1 m_3$ and use the Padè method to construct them.

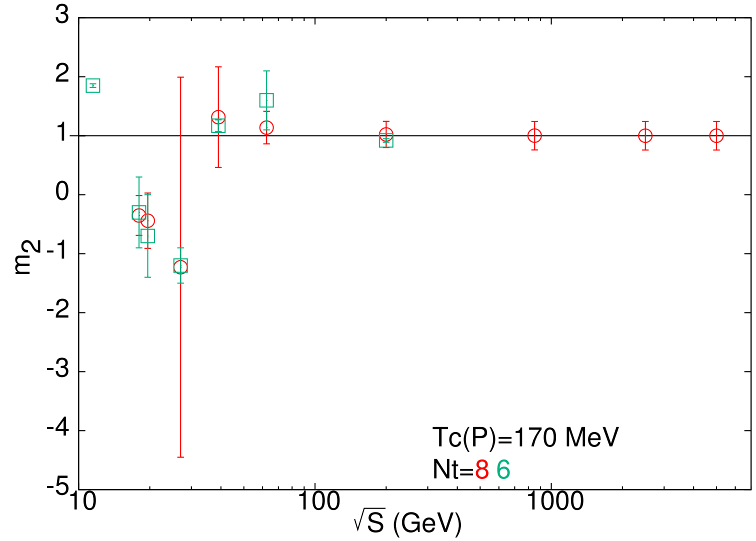


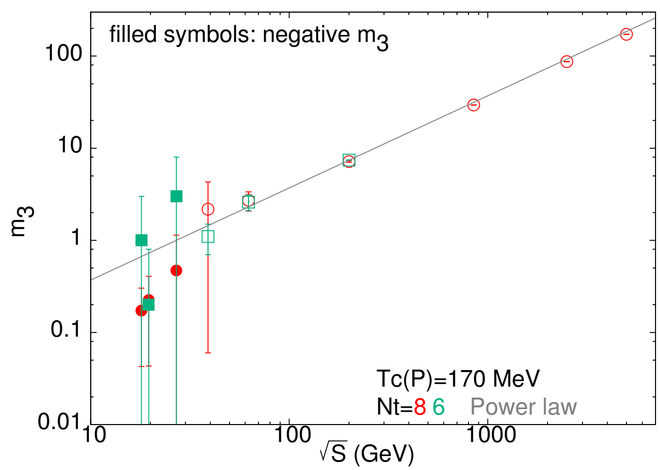
♠ Gavai & Gupta, arXiv: 1001.3796.



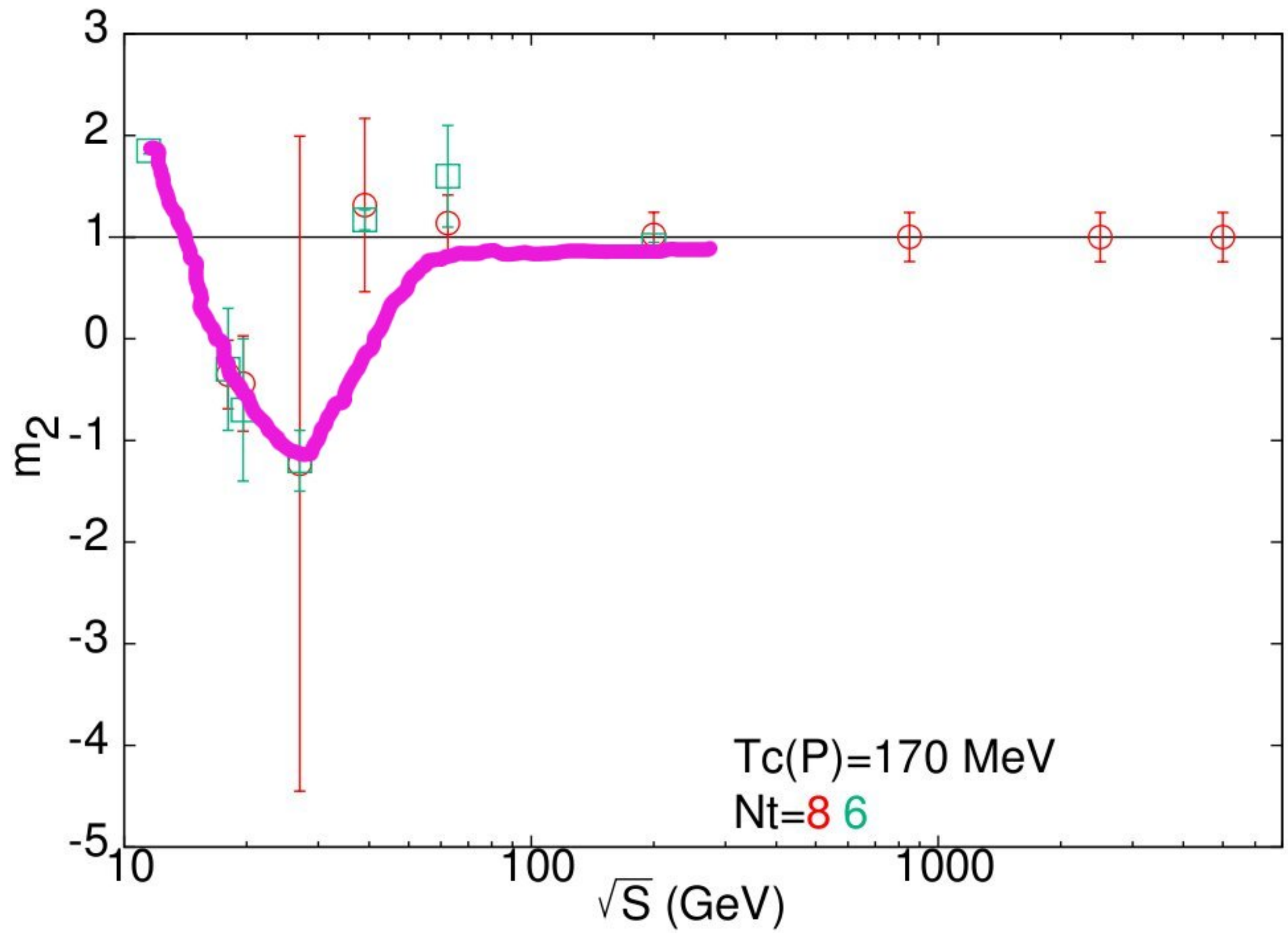
♠ Used $T_c(\mu = 0) = 170$ MeV (Gavai & Gupta, arXiv: 1001.3796).







♠ Marginal change if $T_c = 175$ MeV (Datta, Gavai & Gupta, QM '12).



Gavai-Gupta, '10 & Datta-Gavai-Gupta, QM '12

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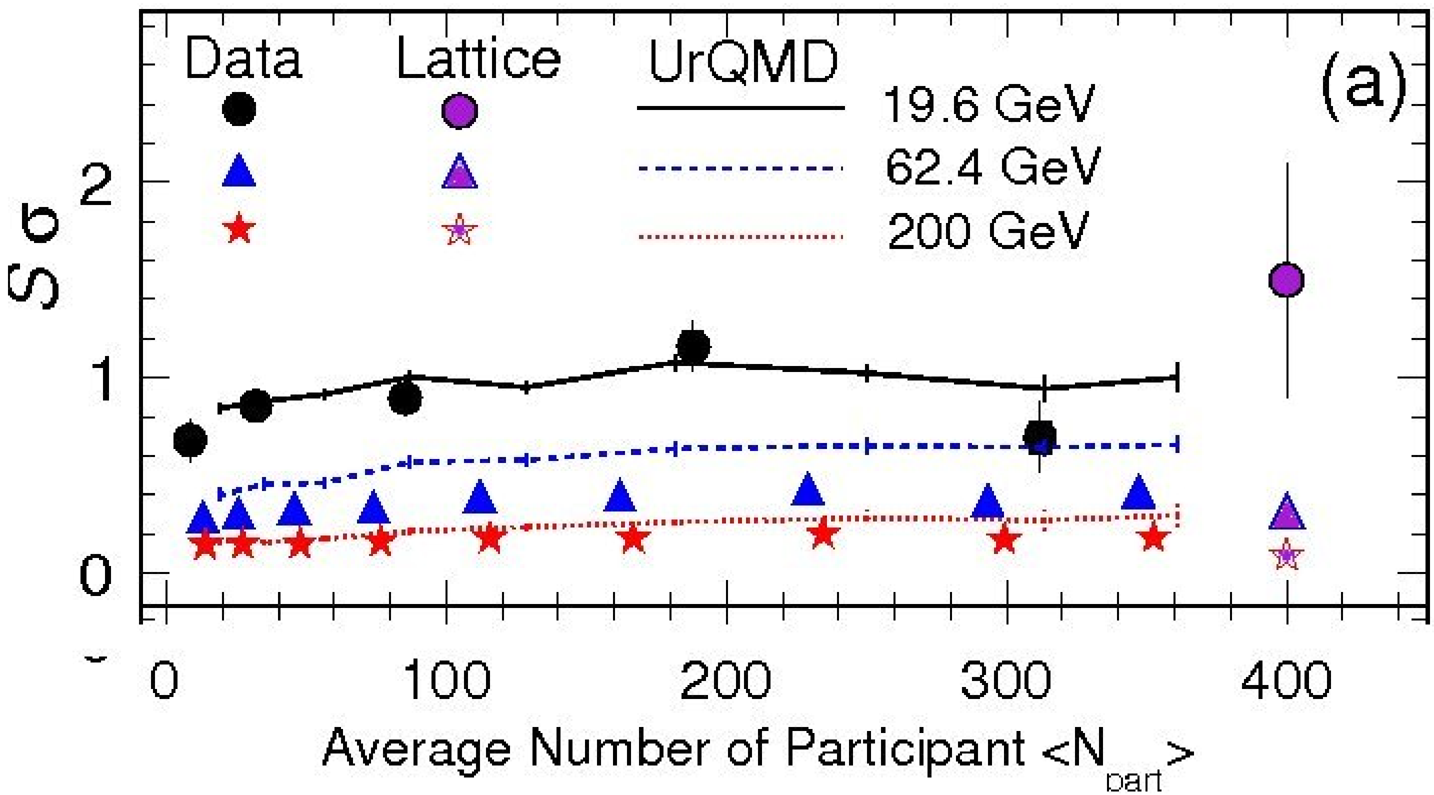
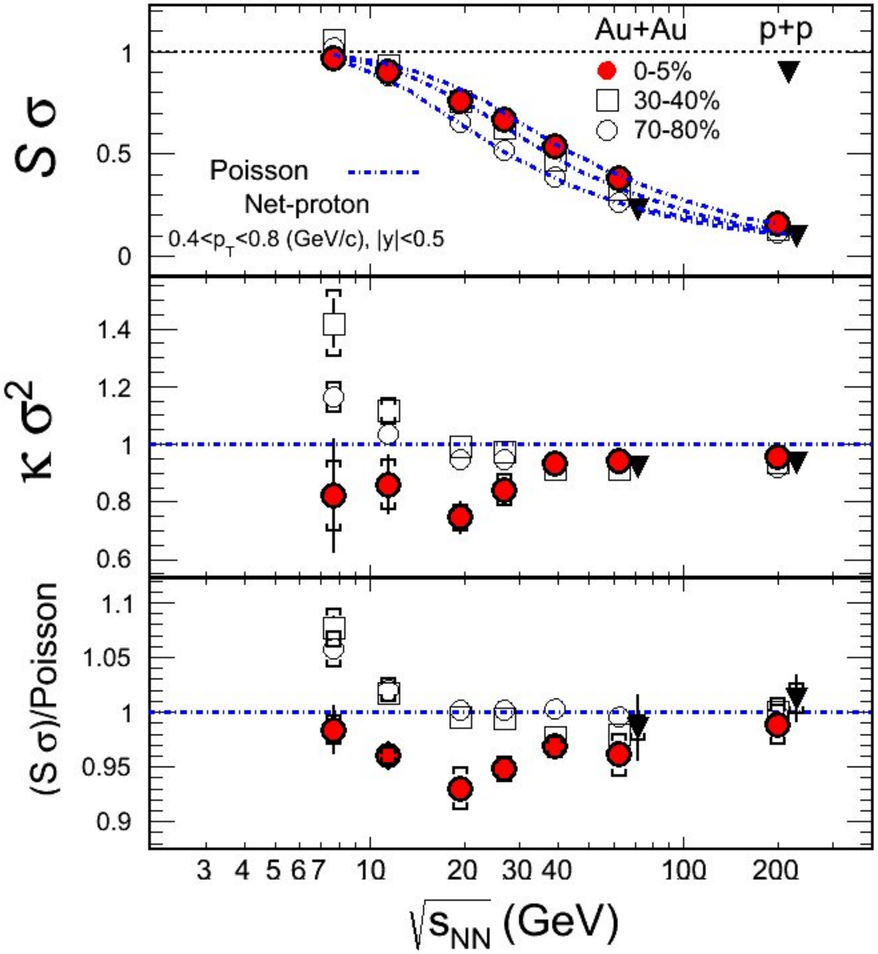
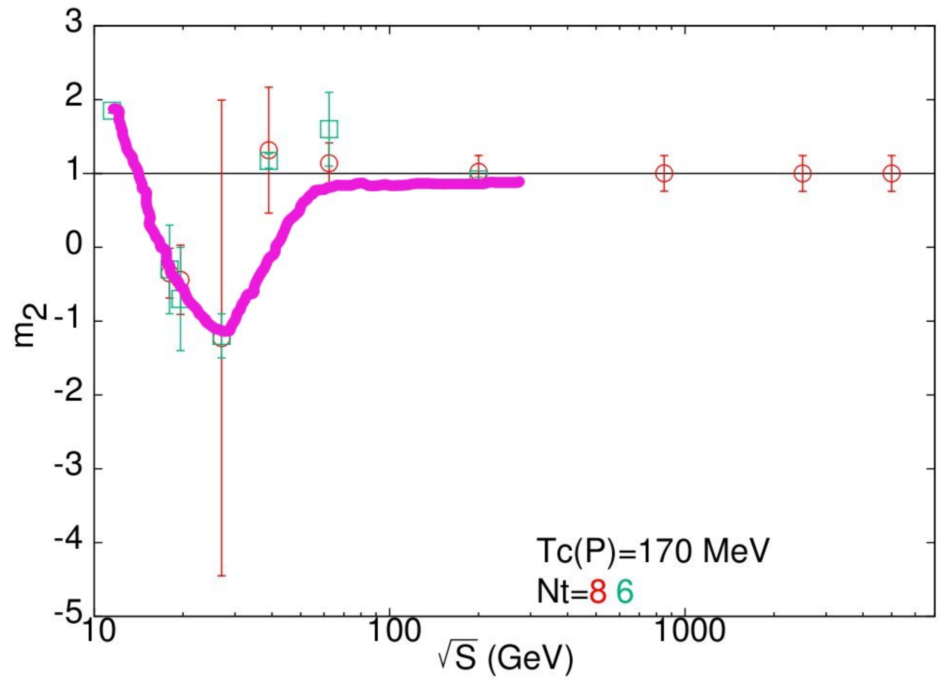
- Leads to a ratio χ_Q:χ_I:χ_B = 1:0:4
- Assuming protons, neutrons, pions to dominate, both χ_Q and χ_B can be shown to be proton number fluctuations only.



Aggarwal et al., STAR Collaboration, arXiv : 1004.4959

• Reasonable agreement with our lattice results. Where is the critical point ?





Xiaofeng Luo, QM'12
From STAR Collaboration

Gavai-Gupta, '10
Datta-Gavai-Gupta, QM '12

Summary

- Phase diagram in $T - \mu$ has begun to emerge: Different methods, $\rightsquigarrow$ similar qualitative picture. Critical Point at $\mu_B/T \sim 1 - 2$.

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- Our results for $N_t = 8$ first to begin the inching towards continuum limit.

- Critical Point leads to structures in m_i on the Freeze-Out Curve.

- STAR results appear to agree with our Lattice QCD predictions.