# Interquark Potentials in Charmonium at Finite Temperature<sup>a</sup>

<sup>*a*</sup> Wynne Evans, Jon-Ivar Skullerud, CA

Interquark Potentials in Charmonium at Finite Temperature<sup>a</sup>

#### and related topics<sup>b</sup>

<sup>*a*</sup> Wynne Evans, Jon-Ivar Skullerud, CA

<sup>b</sup> Gert Aarts, Tim Harris, Seyong Kim, Maria Paola Lombardo, Mehmet Oktay, Sinead Ryan, Don Sinclair, Jon-Ivar Skullerud, CA Interquark potential in charmonium at finite temperature

- Schrödinger Equation Approach
  - Nambu-Bethe-Salpeter wavefunction
- Conventional Exponential Fits
- Maximum Entropy Method

Interquark potential in charmonium at finite temperature

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Related topics = bottomonium at finite temperature

- Spectral Functions
- Non-zero momenta
- "Ground state" masswidth

#### **Particle Data Book**



#### $\sim 1,500~\mathrm{pages}$

#### zero pages on Quark-Gluon Plasma...

## Static Quark Potential (T = 0)

#### UKQCD Collaboration [pre-history]



N-N potential



HAL QCD Collaboration, Aoki, Doi, Hatsuda, Ikeda, Inoue, Ishii, Murano, Nemura, Sasaki

## **Inter-quark potential from the Lattice**

Kawanai, Sasaki, arXiv:1111.025 Ikeda, Iida, arXiv:1102.2097

- finite-quark mass
- quenched

We extend this by using:

- 2 (+1) flavour
- finite temperature
- anisotropic lattices

Hatsuda, PoS CD09 (2009) 068 use the Schrödinger equation to "reverse engineer" the potential, V(r), given the Nambu-Bethe-Salpeter wavefunction,  $\psi(r)$ :

input input

$$\begin{pmatrix} \frac{p^2}{2M} + V(r) \\ \downarrow \end{pmatrix} \begin{pmatrix} \downarrow & \downarrow & \downarrow \\ \psi(r) = E & \psi(r) \\ \downarrow \end{pmatrix}$$

output

 $\psi(r)$  is determined from a lattice simulation from correlators of *non-local* (point-split) operators,  $J(x; \vec{r}) = q(x) \Gamma U(x, x + \vec{r}) \overline{q}(x + \vec{r})$ 

$$\begin{array}{lll} C(\vec{r},t) & = & \displaystyle \sum_{\vec{x}} < J(0;\vec{r}) \; J(x;\vec{r}) > \\ & \longrightarrow & |\psi(r)|^2 \; e^{-Et} \end{array}$$

#### Dublin-Maynooth $N_f = 2$ configurations

| $N_s$ | $N_{\tau}$ | T(MeV) | $T/T_c$ | $N_{\rm cfg}$ |
|-------|------------|--------|---------|---------------|
| 12    | 80         | 90     | 0.42    | 250           |
| 12    | 32         | 230    | 1.05    | 1000          |
| 12    | 28         | 263    | 1.20    | 1000          |
| 12    | 24         | 306    | 1.40    | 500           |
| 12    | 20         | 368    | 1.68    | 1000          |
| 12    | 18         | 408    | 1.86    | 1000          |
| 12    | 16         | 458    | 2.09    | 1000          |

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anisotropic lattice with  $\xi = a_s/a_\tau \approx 6$ 

 $a_s = 0.167 \; {\rm fm}$ 

Vector and Pseudoscalar Channels Charm treated relativistically

























$$V_{q\overline{q}} = \frac{1}{4} [V_{\mathsf{PS}}(r) + 3V_V(r)]$$

with the Schrödinger Eq'n used to define V(r):

$$V_{\Gamma}(r) = E + \frac{1}{\psi(r)} \frac{\nabla^2}{2\mu} \psi(r)$$

where  $\mu$  is the reduced mass:

$$\mu = \frac{1}{2}m_Q$$
 where  $m_Q \approx M_H/2$ 









# $V_{q-\overline{q}}$ Potential (exp fitting)



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#### MEM

Do bound hadronic states persist into the "quark-gluon" plasma phase? How can we extract transport coefficients?

Spectral functions can answer this!

$$C(t, \vec{p}) = \int \rho(\omega, \vec{p}) K(t, \omega) d\omega$$

$$\uparrow \qquad \downarrow \qquad \swarrow$$
Euclidean
Spectral
(Lattice)
Correlator
Function
Kernel

where the (lattice) Kernel is:

$$K(t,\omega) = \frac{\cosh[\omega(t-N_t/2)]}{\sinh[\omega/(2T)]}$$
$$\sim \exp[-\omega t]$$








## **Spectral Functions via MEM**

- Extraction of a spectral density from a lattice correlator is an ill-posed problem:
  - Given C(t) derive  $\rho(\omega)$
  - More  $\omega$  data points then t data points!
- Requires the use of Bayesian analysis -Maximum Entropy Method (MEM)
  - Hatsuda, Asakawa et al
  - Commonly used in other areas...
- Need to check MEM output w.r.t. choice of:
  - Default model
  - Statistics
  - Energy range
  - Euclidean time range











Range in  $\omega$  spans the ground state mass from exp fit











### Wavefunctions: MEM v exp (PS)



### Wavefunctions: MEM v exp (PS)



### Wavefunctions: MEM v exp (PS)



## **Spectral Functions: Excited State (PS)**











Successfully calculated the inter-quark potential in charmonium at finite temperature.

Future Plans:

- Increase from  $12^3$  to  $24^3$  and  $32^3$  volumes with  $N_f = 2 + 1$
- Will study P-wave states

#### Related topics = bottomonium at finite temperature













Pb-Pb collisions

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anisotropic lattice with  $\xi = a_s/a_\tau \approx 6$  ( $a_s = 0.167$  fm) Bottom quark = NRQCD Have < one part per mille statistical error in correlators  $p_{\text{latt}} = (1, 0, 0), \dots (2, 2, 0)$  i.e.  $p = 0.634, \dots 1.73$  GeV

$$G_P(\tau) \sim \int \frac{d^3 p}{(2\pi)^3} p^2 \exp(-2E\tau) \sim \tau^{-5/2}$$





Confirmation that (high temp) P-wave state is "free"

| state                      | $a_{\tau}\Delta E$ | Mass (MeV) | Expt (MeV)      |  |
|----------------------------|--------------------|------------|-----------------|--|
| $1^1 S_0(\eta_b)$          | 0.118(1)           | 9438(8)    | 9390.9(2.8)     |  |
| $2^{1}S_{0}(\eta_{b}(2S))$ | 0.197(2)           | 10019(15)  | -               |  |
| $1^3S_1(\Upsilon)$         | 0.121(1)           | 9460*      | 9460.30(26)     |  |
| $2^{3}S_{1}(\Upsilon')$    | 0.198(2)           | 10026(15)  | 10023.26(31)    |  |
| $1^1 P_1(h_b)$             | 0.178(2)           | 9879(15)   |                 |  |
| $1^{3}P_{0}(\chi_{b0})$    | 0.175(4)           | 9857(29)   | 9859.44(42)(31) |  |
| $1^{3}P_{1}(\chi_{b1})$    | 0.176(3)           | 9864(22)   | 9892.78(26)(31) |  |
| $1^{3}P_{2}(\chi_{b2})$    | 0.182(3)           | 9908(22)   | 9912.21(26)(31) |  |

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| $1^1 P_1(h_b)$             | 0.178(2)           | 9879(15)   | <b>9898.3</b> $\pm 1.1^{+1.0}_{-1.1}$ | prediction |
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**Belle Collaboration** 

### Zero temperature spectral functions, p = 0



#### Non-zero temperature spectral functions, p = 0












Pb-Pb collisions

p-p collisions

Non-zero temperature Mass, p = 0





From Brambilla et al thermal contribution to the width is

$$\frac{\Gamma}{T} = \frac{1156}{81} \alpha_s^3 \simeq 14.27 \alpha_s^3,$$

(at leading order in weak coupling and large mass expansion).

Our results  $\longrightarrow \Gamma/T \sim 1$  so  $\alpha_s \sim 0.4$ .

Also from Brambilla et al thermal contribution to the mass is

$$\delta E_{\rm thermal} = \frac{17\pi}{9} \alpha_s \frac{T^2}{M} \simeq 5.93 \alpha_s \frac{T^2}{M}$$

(see dashed line)

Non-zero temperature Mass, p = 0



Preliminary...

$$p_{\text{latt}} = (1, 0, 0), \dots (2, 2, 0)$$
 i.e.  $p = 0.634, \dots 1.73$  GeV

#### **Variation with** *T*, fixed $p \neq 0$



# Variation with $\boldsymbol{p}\text{, fixed }T$



# Single ratio



#### **Double ratio**



$$\frac{G(\tau; p, T)}{G(\tau; p = 0, T)} / \frac{G(\tau; p, T_0)}{G(\tau; p = 0, T_0)}$$

## **Double ratio**



t

#### Non-zero momentum, spectral function, T = 0



ω

### Non-zero momentum, spectral function, $T \neq 0$

Upsilon



## Mass as function of speed



### Width as function of speed



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#### the end

# Slides to help me answer difficult questions