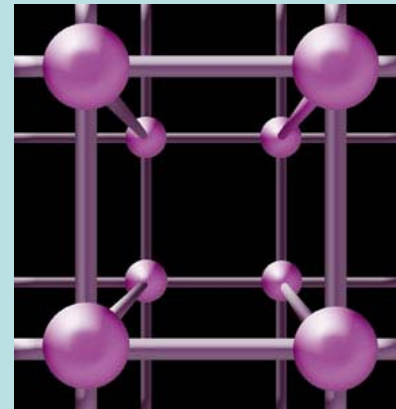


Conformality and The Lattice



Florence

September, 2012

Lattice Simulations

1. Collaboration with George Fleming and Ethan Neil

- 1) arXiv:0712.0609, PRL 100, 171607, 2008
- 2) arXiv:0901.3766 PR D79, 076010, 2009



2. LSD collaboration J. C. Osborn, R. Babich, R. C. Brower, M.I. Buchoff, M. A. Clark, C. Rebbi, D. Schaich, M. Cheng, J. Kiskis, R. Soltz, P. M. Vranas, T. Appelquist, G. T. Fleming, E. T. Neil, M. Lin, G. Voronov, J. Wasem

- 1) arXiv:0910.2224 PRL 104, 071601 (2010) 2→6 Flavors
- 2) arXiv:1009.5967 PRL106:231601,2010 Parity Doubling and S
- 3) arXiv: 1201.3977 PR D85, 074505 (2012) W-W scattering
- 4) arXiv: 1204.6000 --> PRL 10 Flavors
- 5) Underway: SU(2)



3. GF, EN, Meifeng Lin, David Schaich arXiv: 1106.2148 PRD D84, 054501, 2011

Disclaimer:

I am not a lattice gauge theorist.

OUTLINE

1. Introduction - Conformality
2. Degree-of Freedom Inequality Constraint
~ c-theorems , a-theorems
3. Conformal Window via the Lattice:
SU(3) with $N_f = 8, 12$
4. Creeping toward the conformal window
SU(3) with $N_f = 2, 6, 8, 10,$
5. A light dilaton in walking gauge theories?

New Strong Dynamics at TeV Energies?

- New, SM-singlet Sector

Conformal behavior

- Electroweak breaking

Near-conformal IR: walking technicolor (1980's + ...)

Composite Higgs models

Conformal Symmetry / Scale Invariance / Dilatation Symmetry

No fixed scales (masses)

Expected possibly at high energies or
temperatures

Examples:

$$\sigma (e^+e^- \rightarrow \text{Hadrons}) \sim 1/ E_{\text{cm}}^2$$

$$F(T) \sim T^4$$

Not Quite

Quantum field theories *typically* depend on a renormalization scale Λ .

$$\alpha \equiv g^2/4\pi = \alpha(q^2 / \Lambda^2)$$

For a Yang-Mills theory

$$q^2 \rightarrow \infty \\ \sim 1 / \log(q^2 / \Lambda^2)$$

Asymptotic
freedom

QCD Infrared Features

1. Quark and gluon confinement

$$\Lambda \sim 200 \text{ MeV}$$

2. Spontaneous chiral symmetry
(nearly massless quarks)

$$\text{breaking } SU(2)_L \times SU(2)_R \rightarrow SU(2)$$

3 Pions = PN Goldstone B's

$$F = 93 \text{ MeV}$$

$$\langle \bar{\psi}\psi \rangle \sim 4 \pi F^3$$

Beyond QCD

- SU(N) Gauge Theories with N_f Massless Fermions (fundamental and other reps) **Screening**

Asymptotically-free $N_f < N_{af}$

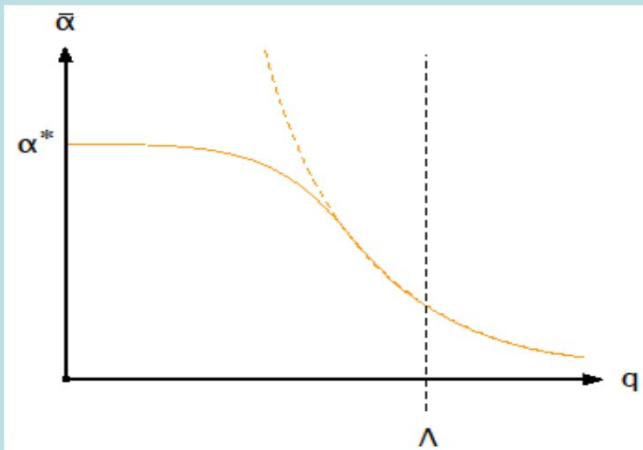
Vary N_f and study how the **infrared** behavior changes

(Chiral symmetry breaking and confinement versus *infrared* conformal behavior.)

SUSY Theories?

Conformal Window

$$N_{af} > N_f > N_{fc} \quad (\text{Fermion screening})$$



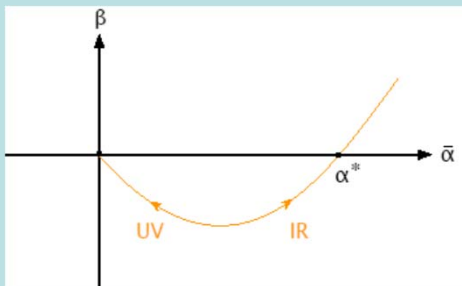
Perturbative for $N_f \rightarrow N_{af}$

Gross and Wilczek, antiquity

William Caswell, 1974 \dashrightarrow

Banks and Zaks, 1982

.....



α^* increases as N_f decreases.

$N_f \rightarrow N_{fc}$ (Non-perturbative)

SU(3) Perturbation Theory

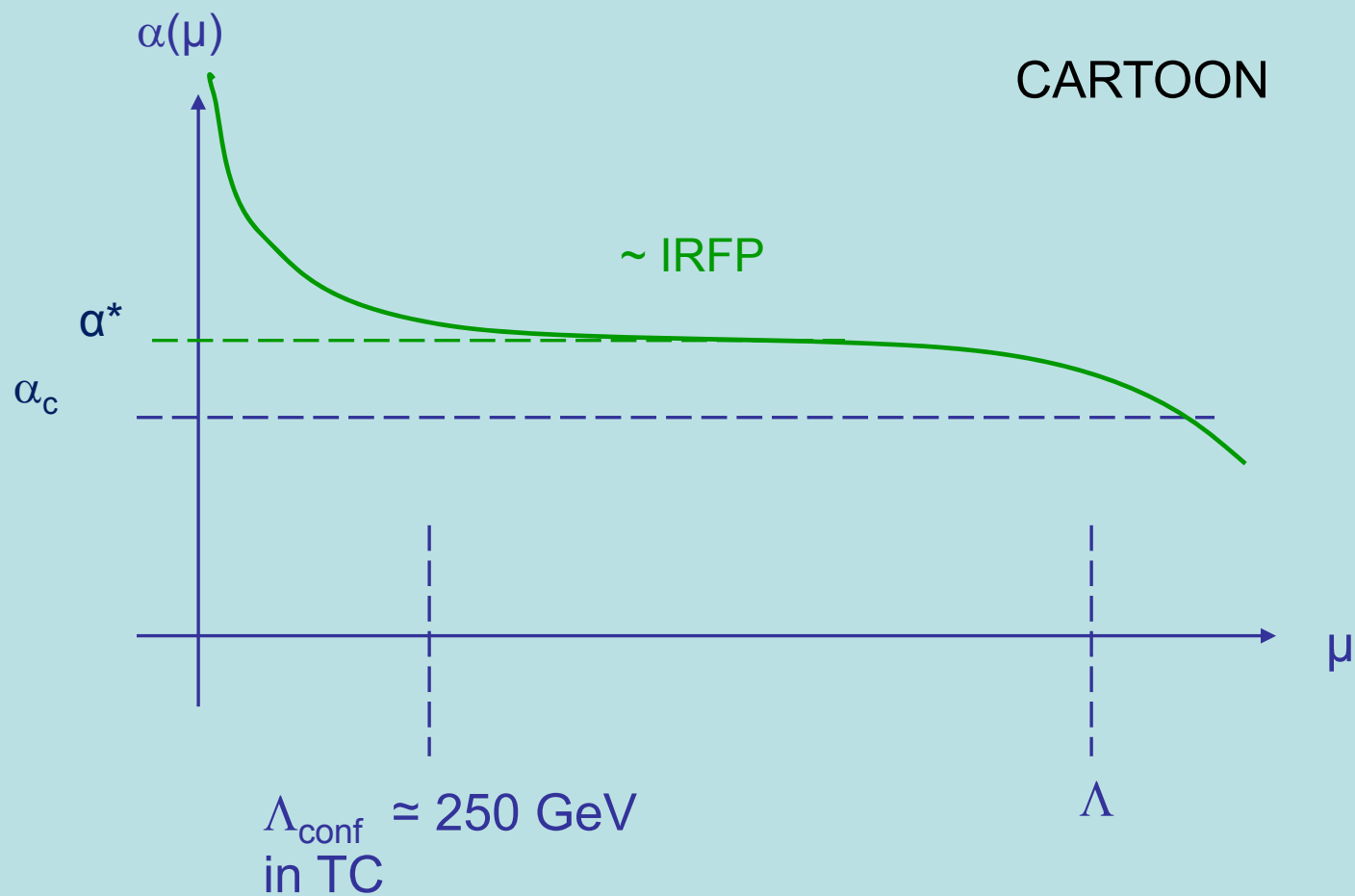
$$\mu \frac{\partial}{\partial \mu} \bar{g}^2(\mu) = \beta(\bar{g}^2(\mu)) = -b_1 \bar{g}^4(\mu) - b_2 \bar{g}^6(\mu) - b_3 \bar{g}^8(\mu) + \dots$$

$$b_1 = \frac{2}{(4\pi)^2} \left(11 - \frac{2}{3} N_f \right), \quad b_2 = \frac{2}{(4\pi)^4} \left(102 - \frac{38}{3} N_f \right)$$

$$b_3^{\overline{MS}} = \frac{2}{(4\pi)^6} \left[\frac{2857}{2} - \frac{5033}{18} N_f + \frac{325}{54} N_f^2 \right]$$

NEAR-CONFORMAL BEHAVIOR

For N_f just below N_{fc} (Walking)



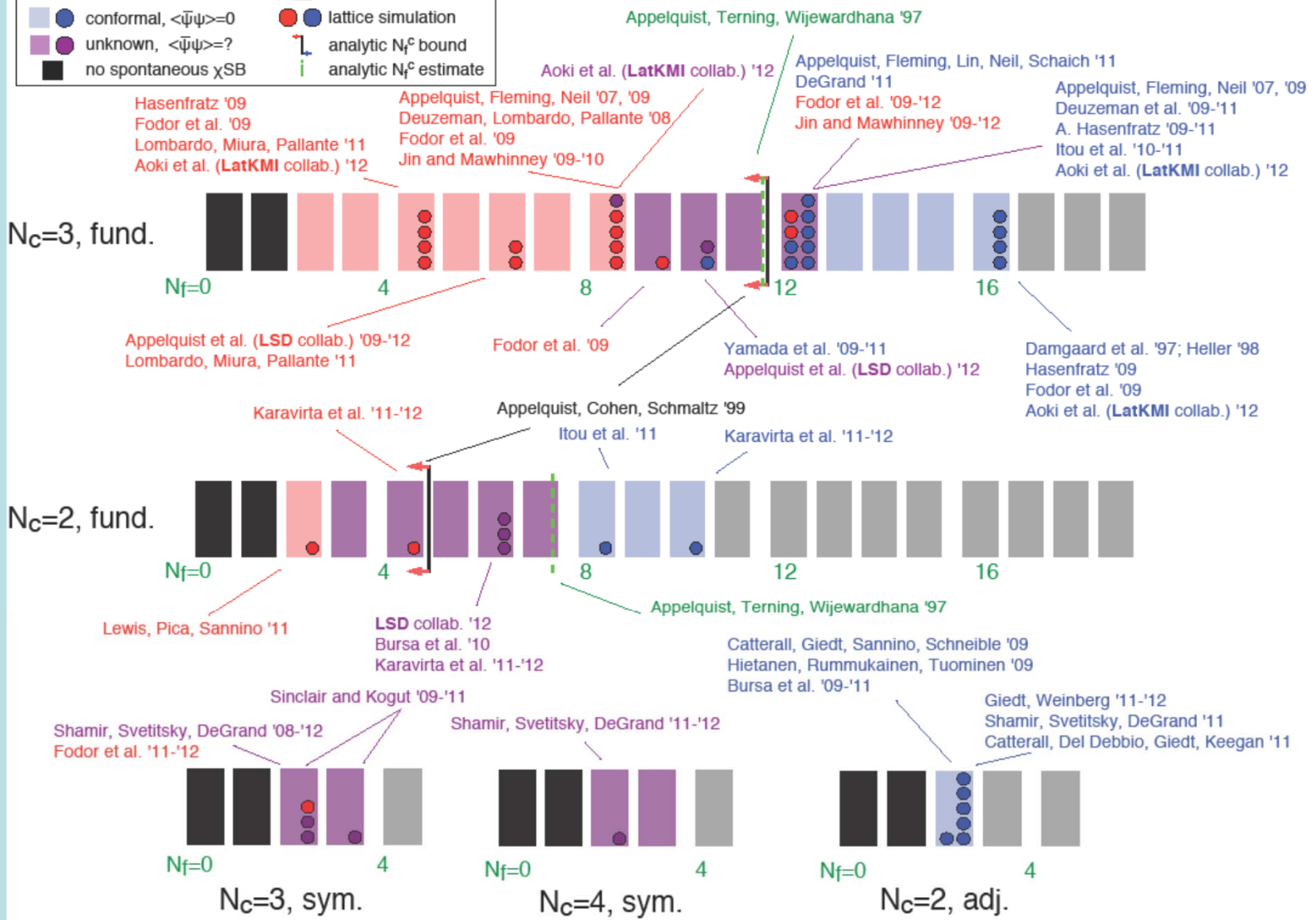
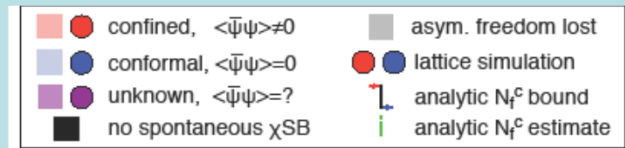
Questions

1. Value of N_{fc} ? (Order of transition?)
2. Inside the conformal window ($N_f > N_{fc}$)
3. Below and near the transition ($N_f \lesssim N_{fc}$)

Approximate IRFP (Walking)?

Condensate enhancement? Parity Doubling? EW
precision studies (S parameter)?

Dilaton ? (Light 0^{++} State)



I will not
discuss: **Holographic Methods
(gauge-gravity duality)**

Holographic Walking Technicolor from D-branes
Lilia Anguelova, Peter Suranyi, L.C. Rohana Wijewardhana
arXiv:1105.4185 Nucl. Phys. B852 (2011) 39-60

Holographic Electroweak Symmetry Breaking from D-branes,
C. Carone, J. Erlich and M. Sher,
arXiv:0704.3084, Phys. Rev. D76 (2007) 0150

Holographic Technicolor Models and their S-parameter,
O. Mintakevich and J. Sonnenschein,
arXiv:0905.3284 JHEP 0907 (2009) 032,

A Negative S Parameter from Holographic Technicolor,
J. Hirn and V. Sanz,
Phys. Rev. Lett. 97 (2006) 121803, hep-ph/0606086;

Dynamical Electroweak Symmetry Breaking from Deformed AdS: Vector Mesons
and Eective Couplings M. Fabbrichesi, M. Piai and L. Vecchi, ,
Phys. Rev. D78 (2008) 045009, arXiv:0804.0124

+

N_{fc} in SU(N) Gauge Theory

- Degree-of-Freedom Inequality (Cohen, Schmaltz, TA 1999). Fundamental rep:

← Conjecture

$$\begin{aligned} N_{fc} &< 4N[1 - 1/18N^2 + \dots] & N = 3, 4, 5, \dots \\ &< 5 & N = 2 \end{aligned}$$

- Gap-Equation Studies, Instantons (1996):

← Suspect methods

$$N_{fc} \cong 4N$$

Pioneers:
Maskawa & Nakajima,
Kugo & Fukuda 1970's

Thermal Inequality

Introduction and Discussion

Related to c-theorems, a-theorems

Komargodski and Schwimmer 1107.3987

Luty, Polchinski, and Rattazzi 1204.5221

+

2+1 QED and Abelian Higgs model
(including recent lattice studies)

Inequality Constraint on Strongly Coupled Field Theories Based on the Thermodynamic Free Energy

A. Cohen
M. Schmaltz
T. Appelquist
hep-th/9901109

- Basic Idea

$$F(T) \sim T^4 f(T)$$

←
“Counts” Thermodynamic
Degrees of Freedom

$$f(T \rightarrow 0) \leq f(T \rightarrow \infty)$$

NATURE
“ABHORS”
MASSLESS
PARTICLES

FINITE TEMPERATURE

$$Z = \text{Tr } e^{-\beta H} \quad \beta = T^{-1}$$

$$A_\mu(\beta, \vec{\alpha}) = A_\mu(0, \vec{\alpha})$$

$$\psi(\beta, \vec{\alpha}) = -\psi(0, \vec{\alpha})$$

$$F(T) = -\ln Z / \beta V$$

$$P(T) = -F(T)$$

$$s = \partial P / \partial T > 0$$

$$C_v = T \partial s / \partial T > 0$$

$$F = \rho - Ts$$

FREE FIELD THEORY

Path Integral \rightarrow Gaussian

$$F(T) = \frac{-\pi^2 T^4}{90} [N_s + 2 N_v + \frac{7}{8} \cdot 4 N_f]$$

Massless

+ Additive, T-Independent
Quartically Divergent
Zero-Point Energy
(Non-SUSY)

Perturbative Corrections

More Generally, Focus on Theories for Which

$$F(T \rightarrow \infty) \rightarrow \frac{-\pi^2 T^4}{90} f(\infty)$$

Finite Constant Counts Fundamental
Degrees of Freedom For A.F. Theory

$$F(T \rightarrow 0) \rightarrow \frac{-\pi^2 T^4}{90} f(0)$$

Finite Constant Counts Massless
L.E. Degrees of Freedom for I.F. Theory

Will apply to theories with non-trivial
IR fixed points

THE INEQUALITY

$$f(0) \leq f(\infty)$$

Natural Route to Proof:

Define: $f(T) \equiv \frac{-\pi^2 T^4}{90} F(T)$

Prove: $f(T)$ is Monotonic in T

A “c-theorem”

Problem: It ain't so.

SUSY SU(N) N_f Matter Multiplets

Finite T Destroys SUSY

$$f(\infty) = \left(1 + \frac{7}{8}\right) [2(N^2 - 1) + 4 N_f N]$$

$$f_{\text{Free Mag}}(0) = \underset{\substack{\text{Duality} \\ N_f \geq N + 2}}{\left(1 + \frac{7}{8}\right)} \left\{ 2[(N_f - N)^2 - 1] \right. \\ \left. + 4 N_f (N_f - N) \right. \\ \left. + 2 N_f^2 \right\}$$

Theorem \Rightarrow

$$N_f \leq \frac{3}{2} N$$

SU(N) (Vector) Gauge Theory

N_f Dirac Fermions

$$f(\infty) = 2(N^2 - 1) + \frac{7}{8} 4N_f N$$

$$f_{\text{NG}}(0) = N_f^2 - 1$$

IR Free

Inequality \Rightarrow

$$N_f \leq 4N \left\{ 1 - \frac{1}{18N^2} + \dots \right\}$$

For N-G Phase

Interesting ?

SU(2) with N_f Dirac Fermions

$$f(\infty) = 6 + 7 N_f$$

Broken Vacuum: $SU(2 N_f) \rightarrow SP(2 N_f)$

$$f(0) = [4 N_f^2 - 1] - [N_f (2N_f + 1)]$$

$$f(0) < f(\infty) \quad \rightarrow \quad N_f < 4.8$$

$$T \frac{\partial}{\partial T} f(T)$$

$$(1) \quad T \frac{\partial}{\partial T} f(g(T)) = \beta(g) f'(g)$$

(One Coupling)

$$(2) \quad T \frac{\partial}{\partial T} f(T) = \frac{90}{\pi^2} \left\{ \frac{s}{T^4} - 4 \frac{P}{T^4} \right\}$$

$$= \frac{90}{\pi^2} \frac{\rho - 3P}{T^4}$$

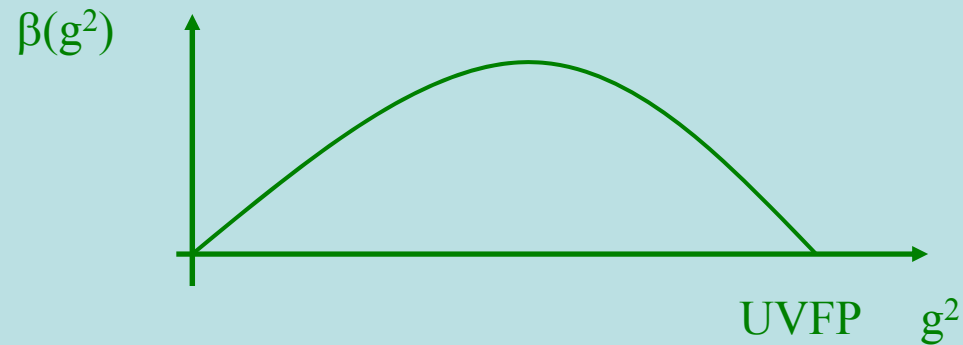
$$= \frac{90}{\pi^2} \frac{\Theta^\mu{}_\mu}{T^4}$$

$$= \frac{90}{\pi^2} \frac{\beta(g)}{g} \langle F_{\mu\nu}^2 \rangle / T^4$$

**Classically scale-invariant
Gauge theory.**

SUGGESTS HOW TO FIND A COUNTER-EXAMPLE:

Find a gauge theory for which



❖ NO 4D EXAMPLE

❖ THERE EXIST 3D EXAMPLES

RESTRICT TO AF THEORIES

SU(N) Theory, N_f Fermions

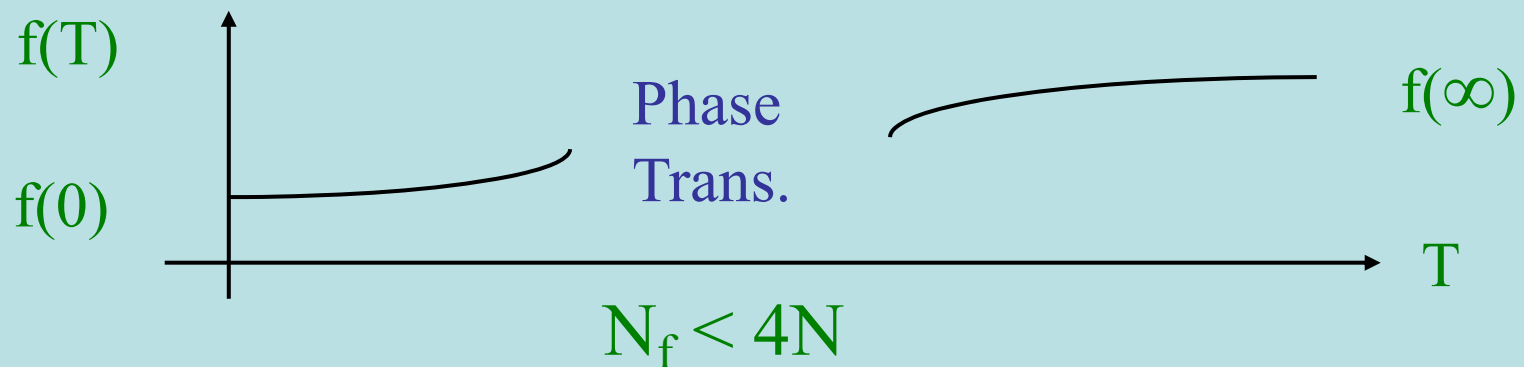
$$T \rightarrow \infty: \quad f(T) \rightarrow \quad 2(N^2 - 1) + \frac{7}{2} N_f N$$

$$- 10 \frac{g^2(T)}{16\pi^2} (N^2 - 1) \left(N + \frac{5}{4} N_f \right)$$

$$+ 0(g^3)$$

$$T \rightarrow 0: \quad f_{\text{NG}}(T \rightarrow 0) = N_f^2 - 1$$

$$+ \frac{N_f^2 - 1}{36} \frac{(\pi T)^4}{(\pi F_\pi)^4} \ln \frac{F\pi}{T} + \dots$$



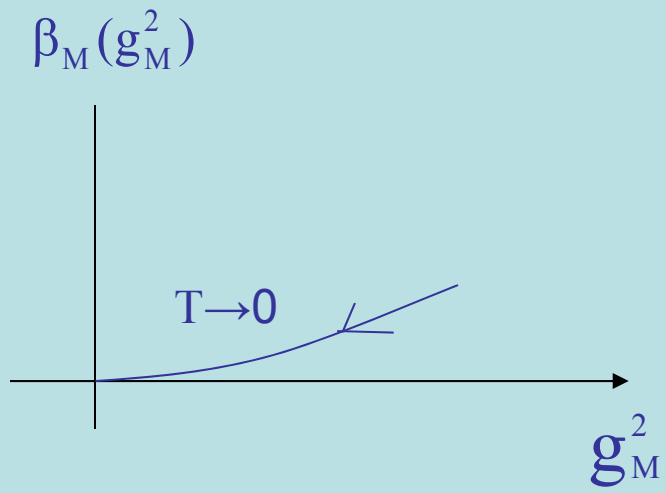
MONOTONICITY COUNTER – EXAMPLE(S)

$$\text{BUT } f(0) \leq f(\infty)$$

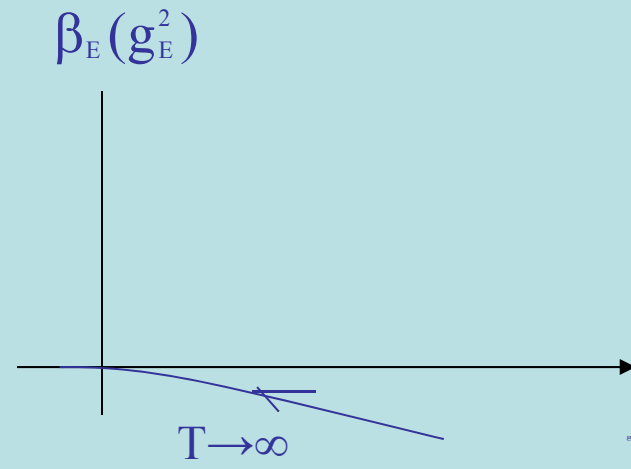
$$\text{SUSY SU}(N) \quad N_f < 3/2 N$$

$$T \rightarrow \infty: \quad f(T) \rightarrow \left(1 + \frac{7}{8}\right) [2(N^2 - 1) + 4N_f N] \\ - (N^2 - 1)(N + 3N_f) \frac{45g^2(T)}{32\pi^2} + \dots$$

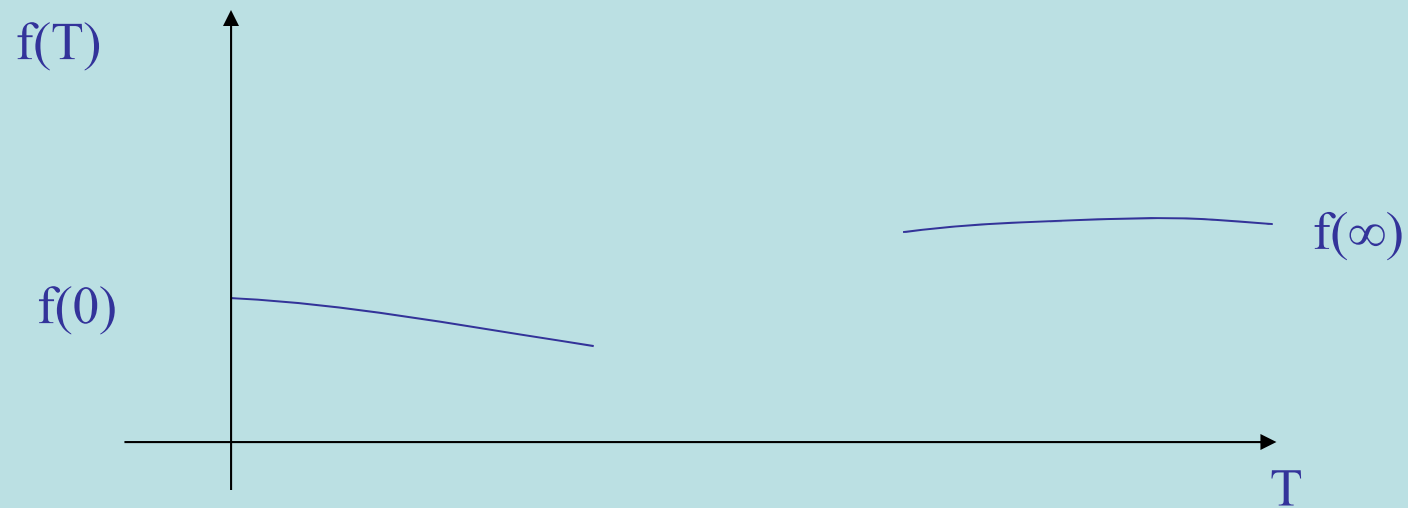
$$T \rightarrow 0: \quad f_{\text{Free Mag}}(T) \rightarrow \left(1 + \frac{7}{8}\right) [2((N_f - N)^2 - 1) \\ + 4N_f(N_f - N) + 2N_f^2] \\ - ((N_f - N)^2 - 1)(4N_f - N) \frac{45g^2(T)}{32\pi^2} \\ - 3N_f^2 N \frac{45y^2(T)}{32\pi^2}$$



$$N_f < \frac{3}{2}N$$



$$N_f < 3N$$



SUMMARY

1. Correct for all AF theories where it can be checked.
2. Interesting constraint on $SU(N)$ theory with N_f fermions.
Saturated?
3. Proof ?
Homework assignment.
4. Counter examples with UV fixed points.

PHASE STRUCTURE OF NON-COMPACT QED3

With L.C.R. Wijewardhana
hep-ph/0403250

- An Old Problem (1988-present)
“Simple”, strongly-coupled QFT
- Planar condensed matter systems
Anti-ferromagnet
High- T_c Cuprates

- The Model

Couple A_μ to N Massless Fermions

- The Question

Spontaneous (dynamical) Fermion Mass

Generation vs. N .

(Chiral symmetry breaking)

- Condensed Matter Correspondence

Antiferromagnetic (Neel) ordering ($N = 2$)

But, it's compact QED3

POLYAKOV INSTANTONS?

THE ABELIAN HIGGS MODEL

$$\begin{aligned} \mathcal{L} = & -\frac{1}{4} (F_{\mu\nu})^2 + \sum_{j=1}^N \bar{\psi}_j (i\partial - eA) \psi_j \\ & + \frac{1}{2} D_\mu \phi^* D^\mu \phi - \lambda (\phi^* \phi - Nv)^2 \end{aligned}$$

Remove Higgs

$$\lambda \rightarrow \infty \quad v \text{ fixed}$$

- Dimensionful Coupling

$$\alpha = e^2 N / 8 \quad (\text{UV cutoff})$$

- Gauge Boson Mass

$$M = e \sqrt{Nv} = \sqrt{8\alpha v}$$

FERMIONS

- $\psi_i =$ Four-Component Fermions

$$\gamma^0 = \begin{pmatrix} \sigma_3 & 0 \\ 0 & -\sigma_3 \end{pmatrix} \quad \gamma^1 = \begin{pmatrix} i\sigma_1 & 0 \\ 0 & -i\sigma_1 \end{pmatrix} \quad \gamma^2 = \begin{pmatrix} i\sigma_2 & 0 \\ 0 & -i\sigma_2 \end{pmatrix}$$

- Two further anticommuters

$$\gamma^3 = i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \gamma^5 = i \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

- U(2) “chiral” symmetry for each spinor

$$\mathbf{1}, -i\gamma^3, \gamma^5, \frac{1}{2}[\gamma^3, \gamma^5] = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

\Rightarrow U(2N) chiral symmetry

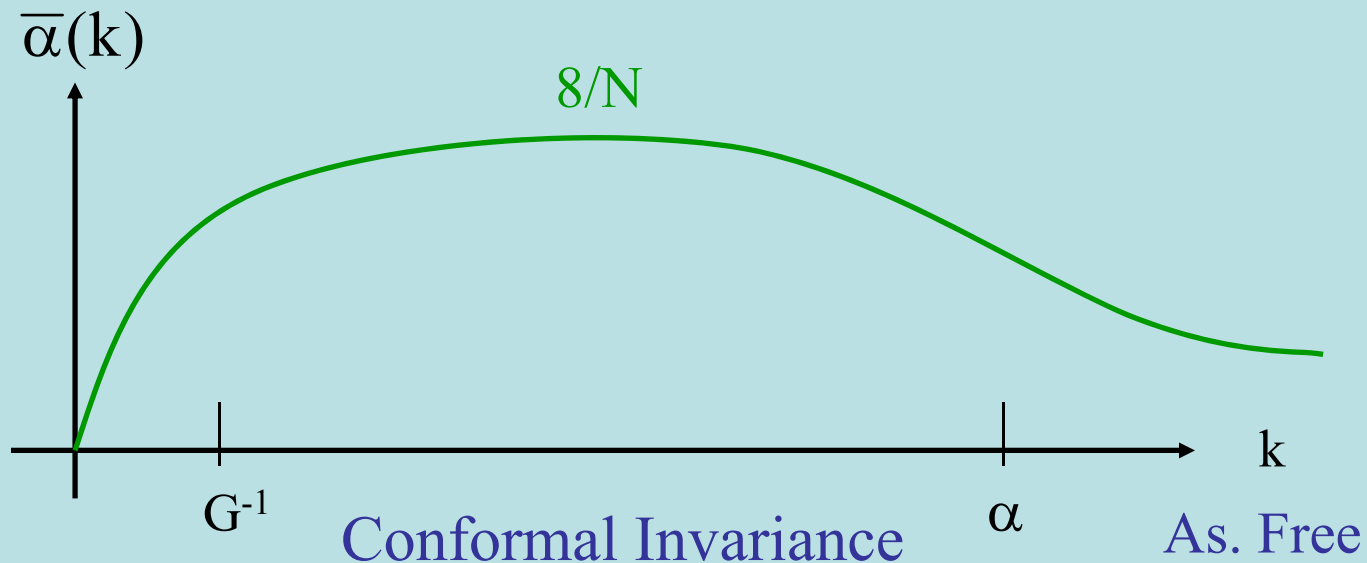
LARGE N

Dimensionless running coupling

$$\bar{\alpha} \equiv \frac{e^2 k}{k^2 + M^2 + \alpha k} = \frac{8}{N} \frac{k}{k^2 / \alpha + G^{-1} + k}$$



$$G = \alpha/M^2$$



Finite N?

Dynamical Fermion Mass?

- Chiral Symmetry Breaking $(U(2N) \rightarrow U(N) \times U(N))$

$$m \bar{\psi}_j \psi_j = m \psi_j^\dagger \begin{pmatrix} \sigma_3 & 0 \\ 0 & -\sigma_3 \end{pmatrix} \psi_j$$

Invariant under parity

$$P \psi_j(t, \vec{x}) P^{-1} = \begin{pmatrix} 0 & \sigma_1 \\ \sigma_1 & 0 \end{pmatrix} \psi_j(t, \vec{x}_P)$$

- Alternate (Parity Violating) Mass

$$m \bar{\psi}_j \frac{[\gamma^3, \gamma^5]}{2} \psi_j = m \psi_j^\dagger \begin{pmatrix} \sigma_3 & 0 \\ 0 & \sigma_3 \end{pmatrix} \psi_j$$

$U(2N)$ Invariant \longleftrightarrow Chern-Simons Term

Vafa-Witten: Not spontaneously

Consensus: Yes for $N \leq N_c(M/\alpha)$

QED3:

$M \rightarrow 0$, α fixed ($G^{-1} = M^2/\alpha \rightarrow 0$): $N_c(0)$

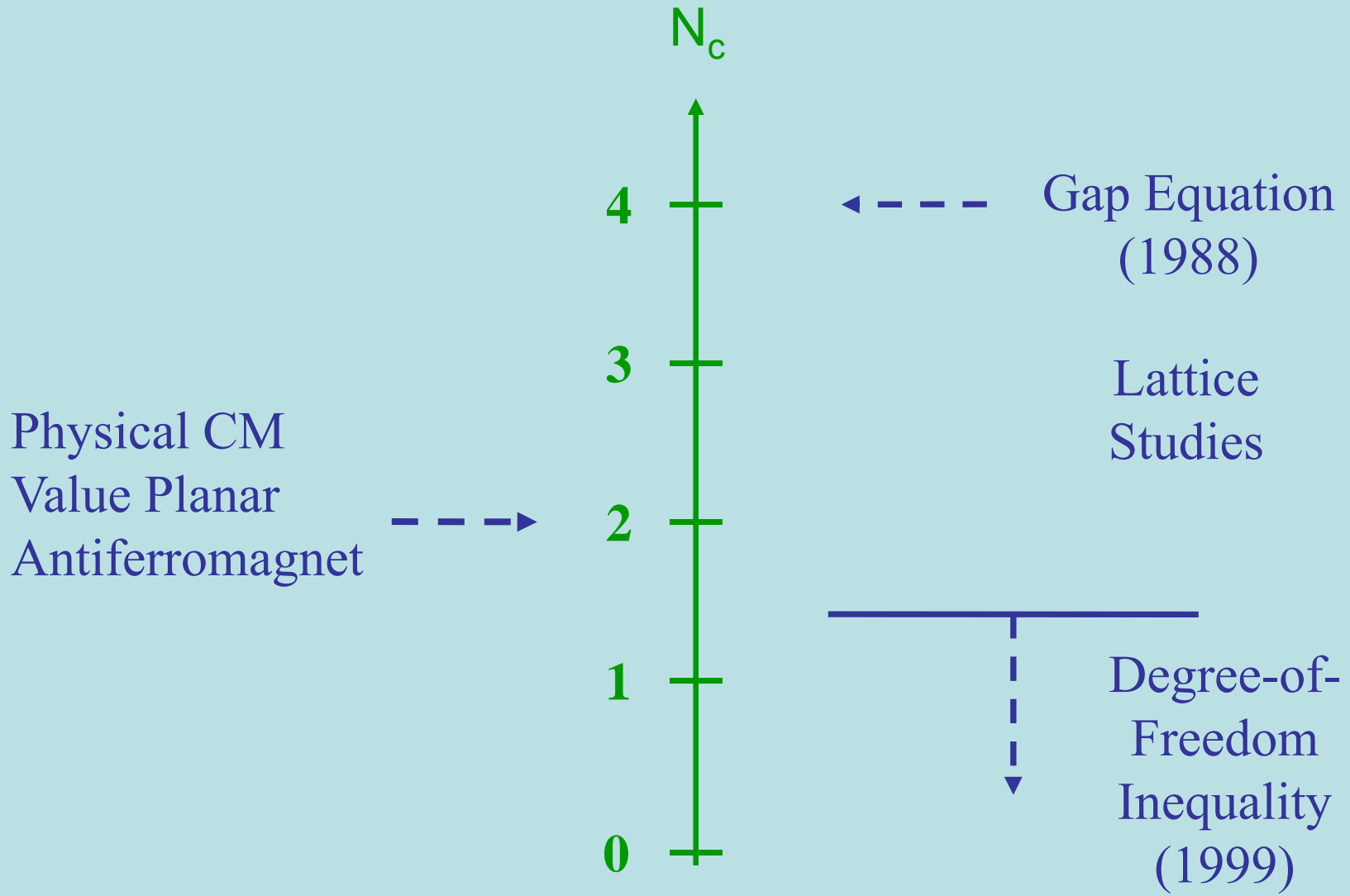
$N > N_c(0)$: IRFP

$N \lesssim N_c(0)$: Approximate IRFP (walking)

Thirring Model:

$G^{-1} = M^2/\alpha$ fixed, $\alpha \rightarrow \infty$ ($M/\alpha \rightarrow 0$): $N_c(0)$

QED3:



Degree-of-Freedom Inequality

Thermodynamic Free Energy

$$f(T) \equiv \frac{F(T)}{T^3} \Omega_{d=3} \quad \left[\begin{array}{l} = 1 \text{ for a} \\ \text{free massless} \\ \text{boson} \end{array} \right]$$

Conjecture for AF Theories

$$f(0) \leq f(\infty)$$

CONCERN

Mermin-Wagner-Coleman Theorem

No NGB's in 2D

Finite-T 3D \rightarrow 2D in IR

Rosenstein, Warr, Park (1990)

Symmetry unbroken at Finite-T

$$m_{\text{NGB}}^2 \sim T^2 e^{-\alpha/T}$$

VANISHES MORE RAPIDLY THAN T AS $T \rightarrow 0$
NO EFFECT ON $f(0)$

QED3 Ultraviolet

$$f(\infty) = 1 + \frac{3}{4}(4N)$$

QED3 Infrared (Broken Phase)

$$f(0) = 1 + 2N^2$$

(Broken Chiral
Symmetry)

\Rightarrow

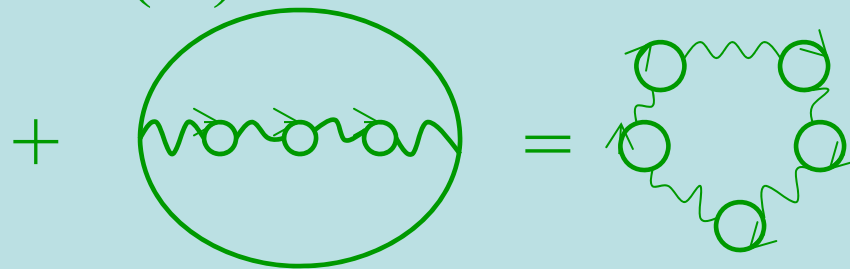
$$N_c \leq 3/2$$

LARGE-N QED3

$$f(\infty) = \frac{3}{4}4N + 1$$

Weak
IRFP

$$f(0) \cong f(\infty)$$



Expect Negative

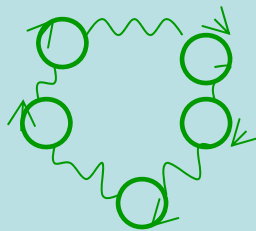
LARGE-N THIRRING MODEL

$\alpha \rightarrow \infty$, FIXED $G^{-1} \equiv M^2/a$

$$f(0) = 3N + 1$$

Weak
UVFP

$$f(\infty) = f(0) +$$



Negative? Then $f(0) > f(\infty)$

Recent Lattice work

Kogut and Strouthos

- **Chiral Symmetry breaking in Three Dimensional QED.**
[Costas Strouthos \(Cyprus U.\)](#) , [John B. Kogut \(DOE, Wash., D.C. & Maryland U.\)](#) . Aug 2008.
25th International Conference on Low Temperature Physics (LT25), Amsterdam, The Netherlands, 6-13 Aug 2008.
J.Phys.Conf.Ser.150:052247,2009.
arXiv:0808.2714
-
- **The Phases of Non-Compact QED(3).**
[Costas Strouthos \(Lausanne\)](#) , [John B. Kogut \(DOE, Wash., D.C. & Maryland U.\)](#) . Apr 2008.
Presented at 25th International Symposium on Lattice Field Theory, Regensburg, Germany, 30 Jul - 4 Aug 2007.
PoS LAT2007:278,2007.
arXiv:0804.0300