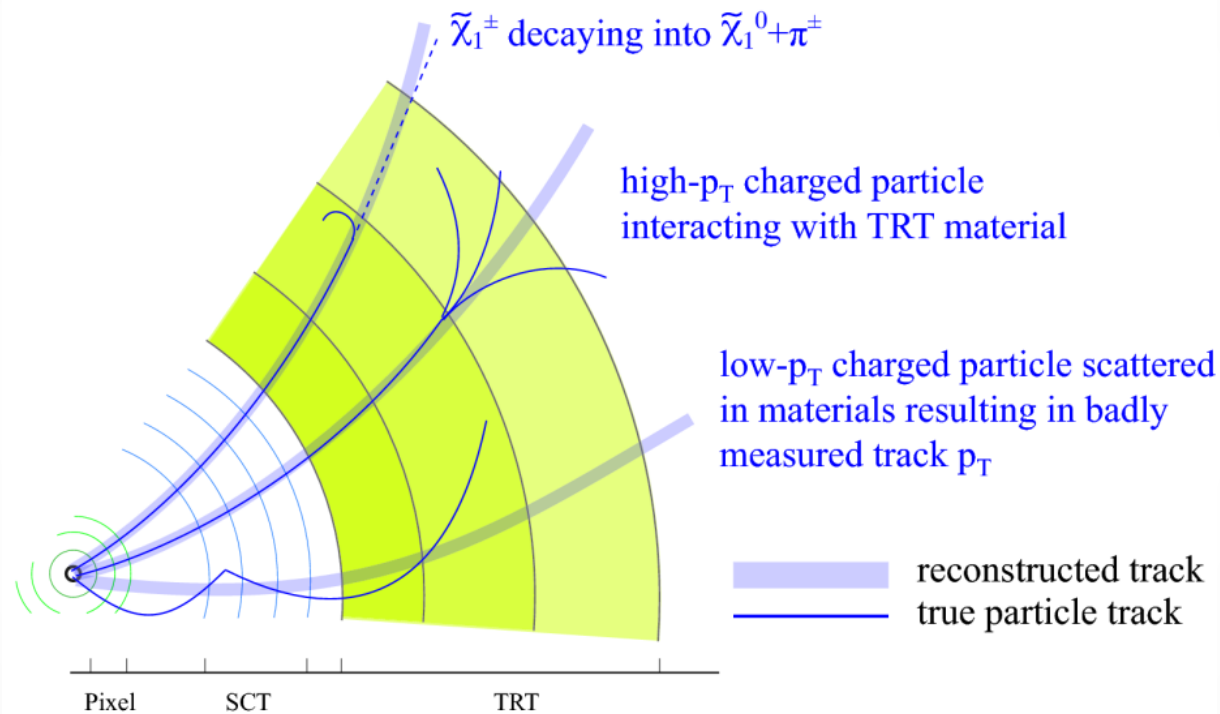


Searches for RPV-SUSY and Long-Lived particles in ATLAS



Monica D'Onofrio

University of Liverpool

GGI Meeting, Florence, 23 Oct. 2012

Searches for SUSY @ ATLAS

- Broadly and deeply cover the SUSY signature space

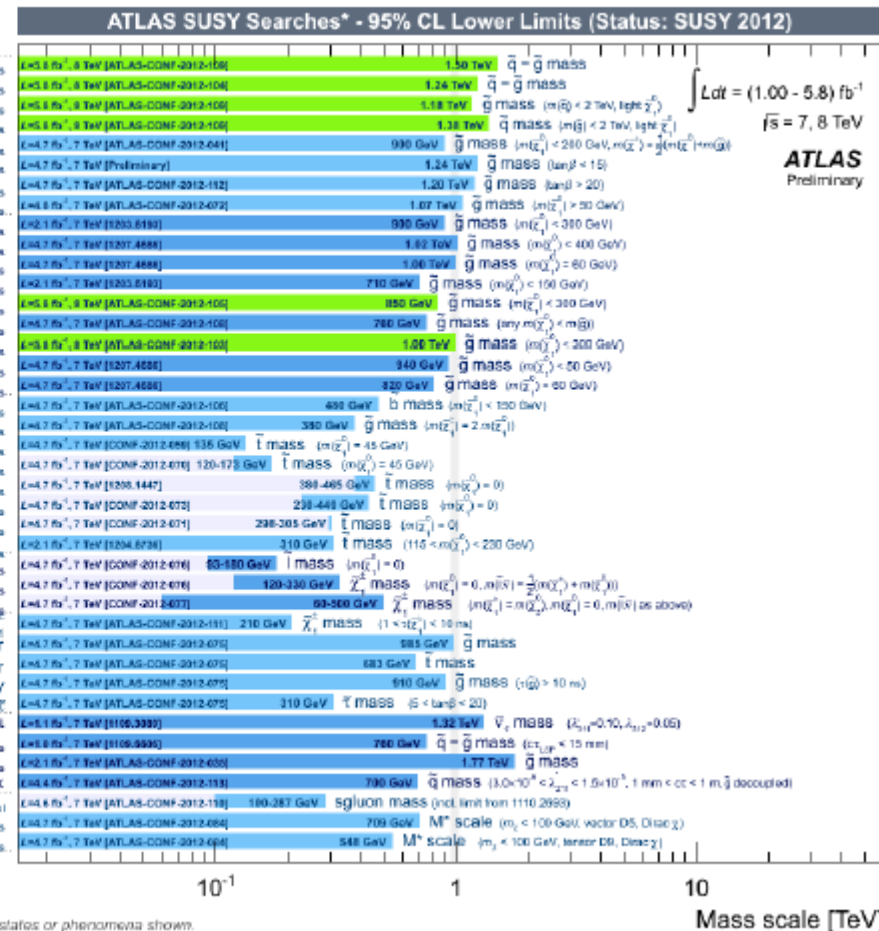
1. Inclusive search

2. Natural spectrum

3. Long-lived particles

4. Prompt RPV

5. MSSM extension



*Only a selection of the available mass limits on new states or phenomena shown.
 All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

Searches for SUSY @ ATLAS

Giacomo's talk yesterday

1. Inclusive search

2. Natural spectrum

3. Long-lived sparticles

4. Prompt RPV

5. MSSM extension

ATLAS SUSY Searches* - 95% CL Lower Limits (Status: SUSY 2012)

$L_{int} = (1.00 - 5.8) \text{ fb}^{-1}$
 $\sqrt{s} = 7, 8 \text{ TeV}$
 ATLAS Preliminary

Search	Model	Lower Limit [TeV]	Signature
MSUGRA/CMSSM	0 lep + j 's + $E_{T,miss}$	1.50 TeV	$\tilde{g} = \tilde{q}$ mass
MSUGRA/CMSSM	1 lep + j 's + $E_{T,miss}$	1.24 TeV	$\tilde{g} = \tilde{q}$ mass
Pheno model	0 lep + j 's + $E_{T,miss}$	1.18 TeV	\tilde{g} mass ($m_{\tilde{g}} < 2 \text{ TeV}, \log_{10} \mu_{eff}^2$)
Pheno model	0 lep + j 's + $E_{T,miss}$	1.38 TeV	\tilde{q} mass ($m_{\tilde{q}} < 2 \text{ TeV}, \log_{10} \mu_{eff}^2$)
Gluino med. $\tilde{\chi}^0$	$\tilde{g} \rightarrow q\bar{q}\tilde{\chi}^0$: 1 lep + j 's + $E_{T,miss}$	930 GeV	\tilde{g} MASS ($m_{\tilde{g}} < 280 \text{ GeV}, m_{\tilde{\chi}^0} = \frac{1}{2}(m_{\tilde{g}} + m_{\tilde{q}})$)
GMSB	2 lep (OS) + j 's + $E_{T,miss}$	1.24 TeV	\tilde{g} MASS ($\tan\beta < 15$)
GMSB	1-2 τ + 0-1 lep + j 's + $E_{T,miss}$	1.20 TeV	\tilde{g} MASS ($\tan\beta > 20$)
GGM	$\tau\gamma$ + $E_{T,miss}$	1.07 TeV	\tilde{g} mass ($m_{\tilde{g}} > 30 \text{ GeV}$)
$\tilde{g} \rightarrow b\bar{b}\tilde{\chi}^0$ (virtual b)	0 lep + 1/2 b - j 's + $E_{T,miss}$	930 GeV	\tilde{g} mass ($m_{\tilde{g}} < 300 \text{ GeV}$)
$\tilde{g} \rightarrow b\bar{b}\tilde{\chi}^0$ (virtual b)	0 lep + 3 b - j 's + $E_{T,miss}$	1.62 TeV	\tilde{g} mass ($m_{\tilde{g}} < 400 \text{ GeV}$)
$\tilde{g} \rightarrow b\bar{b}\tilde{\chi}^0$ (real b)	0 lep + 3 b - j 's + $E_{T,miss}$	1.06 TeV	\tilde{g} mass ($m_{\tilde{g}} < 80 \text{ GeV}$)
$\tilde{g} \rightarrow t\bar{t}\tilde{\chi}^0$ (virtual t)	1 lep + 1/2 b - j 's + $E_{T,miss}$	710 GeV	\tilde{g} mass ($m_{\tilde{g}} < 100 \text{ GeV}$)
$\tilde{g} \rightarrow t\bar{t}\tilde{\chi}^0$ (virtual t)	2 lep (SS) + j 's + $E_{T,miss}$	850 GeV	\tilde{g} mass ($m_{\tilde{g}} < 300 \text{ GeV}$)
$\tilde{g} \rightarrow t\bar{t}\tilde{\chi}^0$ (virtual t)	3 lep + j 's + $E_{T,miss}$	700 GeV	\tilde{g} mass (any $m_{\tilde{g}} < m_{\tilde{g}}$)
$\tilde{g} \rightarrow t\bar{t}\tilde{\chi}^0$ (virtual t)	0 lep + multi- j 's + $E_{T,miss}$	1.08 TeV	\tilde{g} mass ($m_{\tilde{g}} < 200 \text{ GeV}$)
$\tilde{g} \rightarrow t\bar{t}\tilde{\chi}^0$ (virtual t)	0 lep + 3 b - j 's + $E_{T,miss}$	340 GeV	\tilde{g} mass ($m_{\tilde{g}} < 50 \text{ GeV}$)
$\tilde{g} \rightarrow t\bar{t}\tilde{\chi}^0$ (real t)	0 lep + 3 b - j 's + $E_{T,miss}$	320 GeV	\tilde{g} mass ($m_{\tilde{g}} < 50 \text{ GeV}$)
$bb, b_s \rightarrow b\bar{b}\tilde{\chi}^0$	0 lep + 2 b -jets + $E_{T,miss}$	490 GeV	b MASS ($m_{\tilde{b}} < 150 \text{ GeV}$)
$bb, b_s \rightarrow b\bar{b}\tilde{\chi}^0$	3 lep + j 's + $E_{T,miss}$	390 GeV	\tilde{g} mass ($m_{\tilde{g}} = 2 \cdot m_{\tilde{b}}^2$)
\tilde{t} (very light), $\tilde{t} \rightarrow b\bar{b}\tilde{\chi}^0$	2 lep + $E_{T,miss}$	536 GeV	\tilde{t} mass ($m_{\tilde{t}}^2 = 45 \text{ GeV}^2$)
\tilde{t} (light), $\tilde{t} \rightarrow b\bar{b}\tilde{\chi}^0$	1/2 lep + b -jet + $E_{T,miss}$	120-175 GeV	\tilde{t} mass ($m_{\tilde{t}}^2 = 45 \text{ GeV}^2$)
\tilde{t} (heavy), $\tilde{t} \rightarrow b\bar{b}\tilde{\chi}^0$	0 lep + b -jet + $E_{T,miss}$	380-485 GeV	\tilde{t} MASS ($m_{\tilde{t}}^2 = 0$)
\tilde{t} (heavy), $\tilde{t} \rightarrow b\bar{b}\tilde{\chi}^0$	1 lep + b -jet + $E_{T,miss}$	258-443 GeV	\tilde{t} MASS ($m_{\tilde{t}}^2 = 0$)
\tilde{t} (heavy), $\tilde{t} \rightarrow b\bar{b}\tilde{\chi}^0$	2 lep + b -jet + $E_{T,miss}$	298-305 GeV	\tilde{t} MASS ($m_{\tilde{t}}^2 = 0$)
\tilde{t} (GMSB), $\tilde{t} \rightarrow b\bar{b}\tilde{\chi}^0$	2 lep + b -jet + $E_{T,miss}$	310 GeV	\tilde{t} mass ($115 < m_{\tilde{t}}^2 < 230 \text{ GeV}^2$)
$\tilde{t} \rightarrow b\bar{b}\tilde{\chi}^0$	2 lep + $E_{T,miss}$	85-120 GeV	\tilde{t} MASS ($m_{\tilde{t}}^2 = 0$)
$\tilde{t} \rightarrow b\bar{b}\tilde{\chi}^0$	2 lep + $E_{T,miss}$	120-330 GeV	\tilde{t} MASS ($m_{\tilde{t}}^2 = 0, m_{\tilde{g}} = \frac{1}{2}(m_{\tilde{g}} + m_{\tilde{q}}) + m_{\tilde{t}}^2$)
AMSB (direct $\tilde{\chi}^0$ pair prod.)	long-lived $\tilde{\chi}^0$	210 GeV	$\tilde{\chi}^0$ mass ($1 < \tau_{\tilde{\chi}^0} < 10 \text{ ns}$)
Stable \tilde{g} R-hadrons	Full detector	585 GeV	\tilde{g} mass
Stable \tilde{t} R-hadrons	Full detector	683 GeV	\tilde{t} mass
Metastable \tilde{g} R-hadrons	Pixel det. only	610 GeV	\tilde{g} MASS ($\tau_{\tilde{g}} > 10 \text{ ns}$)
GMSB: stable $\tilde{\tau}$		310 GeV	$\tilde{\tau}$ MASS ($\beta < \tan\beta < 20$)
RPV: high-mass eq.		1.32 TeV	\tilde{V}_τ MASS ($\tilde{\tau}_{1,2} = 0, \tilde{\tau}_{3,4} = 0.05$)
Bilinear RPV	1 lep + j 's + $E_{T,miss}$	760 GeV	$\tilde{q} = \tilde{g}$ MASS ($\tau_{\tilde{q}} < 15 \text{ ns}$)
BC1 RPV	4 lep + $E_{T,miss}$	1.77 TeV	\tilde{g} mass
RPV $\tilde{\chi}^0 \rightarrow qq_1 + \mu + \text{heavy displaced vertex}$		790 GeV	\tilde{q} mass ($3.0 \cdot 10^{-5} < \tilde{\tau}_{\tilde{q}} < 1.5 \cdot 10^{-3}, 1 \text{ mm} < cc < 1 \text{ m}, \tilde{q}$ decoupled)
Hypercolour scalar gluons	4 jets, $m_{\tilde{g}} = m_{\tilde{h}}$	180-287 GeV	sgluon mass (incl. limit from 1110 2959)
Spin dep. WIMP interaction: monojet + $E_{T,miss}$		700 GeV	M^* scale ($m_{\tilde{g}} < 100 \text{ GeV}$, vector DS, Dirac \tilde{g})
Spin indep. WIMP interaction: monojet + $E_{T,miss}$		548 GeV	M^* scale ($m_{\tilde{g}} < 100 \text{ GeV}$, tensor DS, Dirac \tilde{g})

this talk

R-parity

- Define R-parity = $(-1)^{3(B-L)+2s}$

- R = 1 for SM particles
- R = -1 for MSSM partners

If conserved (RPC), sparticles pair produced and Lightest SUSY Particle is stable (MET signatures)

If not conserved (RPV) \rightarrow different terms, couplings constraint by proton decay

$$W_{RP} = \lambda_{ijk} \hat{L}_i \hat{L}_j \hat{E}_k^C + \lambda'_{ijk} \hat{L}_i \hat{Q}_j \hat{D}_k^C + \epsilon_i \hat{L}_i \hat{H}_u + \lambda''_{ijk} \hat{U}_i^C \hat{D}_j^C \hat{D}_k^C$$

L-number violating terms

bilinear terms

B-number violating terms

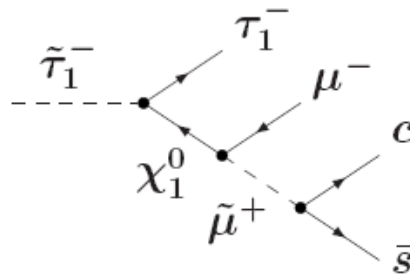
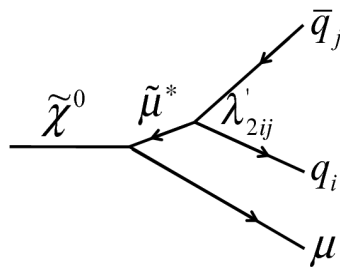
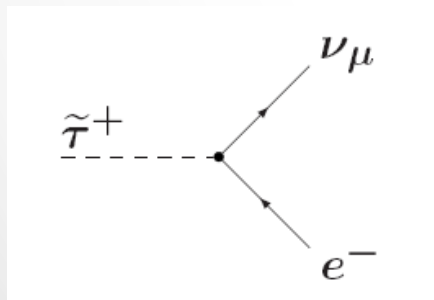
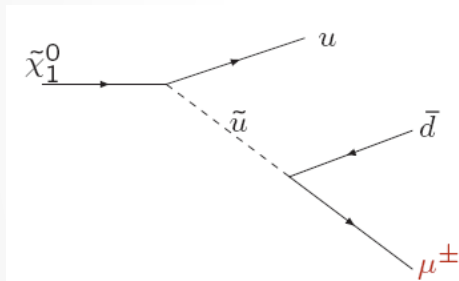
$\Delta L = 1$, 9 λ couplings, 27 λ' couplings

Pletora of new couplings, only partially constraints (m/100 GeV)

	$\lambda_{ijk} L_i L_j \bar{E}_k$	$\lambda'_{1jk} L_1 Q_j \bar{D}_k$	$\lambda'_{2jk} L_2 Q_j \bar{D}_k$	$\lambda'_{3jk} L_3 Q_j \bar{D}_k$
weakest	0.07	0.28	0.56	0.52
strongest	0.05	$5 \cdot 10^{-4}$	0.06	0.11

RPV scenarios

- Many final states to explore and not yet searched for:
 - LSP no longer stable



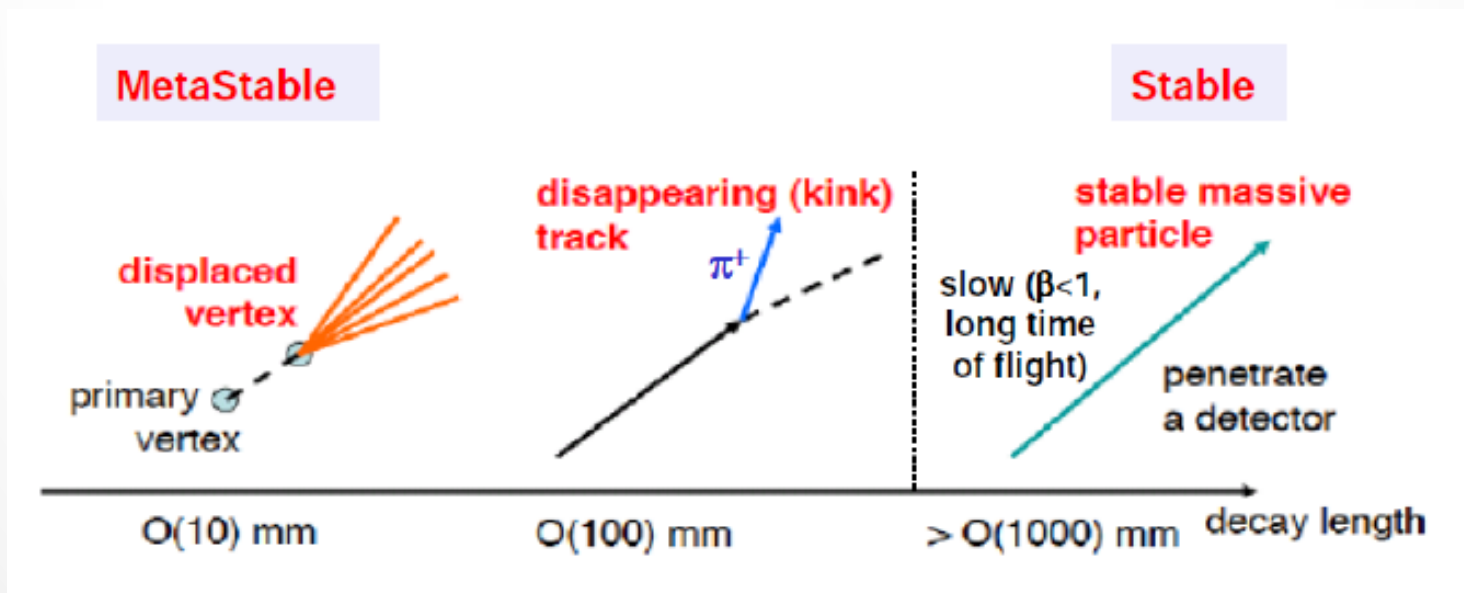
(pair production: $\tilde{q}\tilde{q}, \tilde{q}\tilde{g}, \tilde{g}\tilde{g}$)
 resonant $\tilde{\ell}$ production

$$\left(\begin{array}{c} \text{LSP} \\ \tilde{\chi}_1^0 \\ \tilde{\chi}_1^\pm \\ \tilde{\nu}_L \\ \tilde{\ell}_{L,R}^\pm \\ \tilde{\tau}_1^\pm \\ \tilde{q}_{L,R} \\ \tilde{t}_1 \\ \tilde{g} \end{array} \right) \otimes \left(\begin{array}{c} \text{Operator} \\ L_1 L_2 \bar{E}_1 \\ \vdots \\ L_2 L_3 \bar{E}_3 \\ L_e Q_1 \bar{D}_1 \\ \vdots \\ L_\mu Q_1 \bar{D}_1 \\ \vdots \\ L_\tau Q_3 \bar{D}_3 \\ \bar{U}_i \bar{D}_j \bar{D}_k \end{array} \right)$$

- Something like > 700 possibilities, with final state signatures involving leptons and/or jets
- If $\lambda, \lambda', \lambda''$ very small, can lead to long-lived LSP

Long-Lived particles in SUSY

- Several new physics models could give rise to new, massive particles with long-lifetime. Long-lived (LL) particles in RPC scenarios:
 - If $\Delta M(\text{chargino-neutralino}) \approx 100 \text{ MeV}$ (eg. in AMSB):
 - Long-lived charginos \rightarrow disappearing tracks
 - If very heavy squarks mediate gluinos decay (strong virtuality):
 - Long-lived gluinos \rightarrow R-hadrons (eg. Split SUSY)
 - If NLSP-gravitino GMSB couplings are weak: long-lived sleptons



- **‘prompt’ RPV SUSY:**
 - **Leptonic RPV (production and/or decay):**
 - search for multilepton final states
 - emu resonances and continuum
 - **Hadronic RPV:**
 - search for 2x3jet resonances
 - search for 2x2jet resonances (currently focusing on scalar gluons searches)
- **Long-lived SUSY:**
 - Displaced vertex
 - Disappearing tracks
 - R-hadrons and sleptons

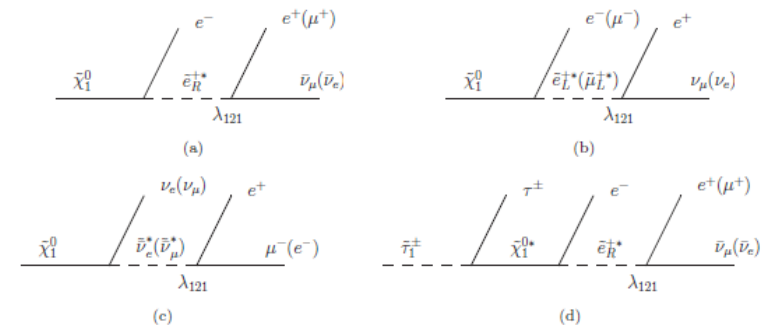
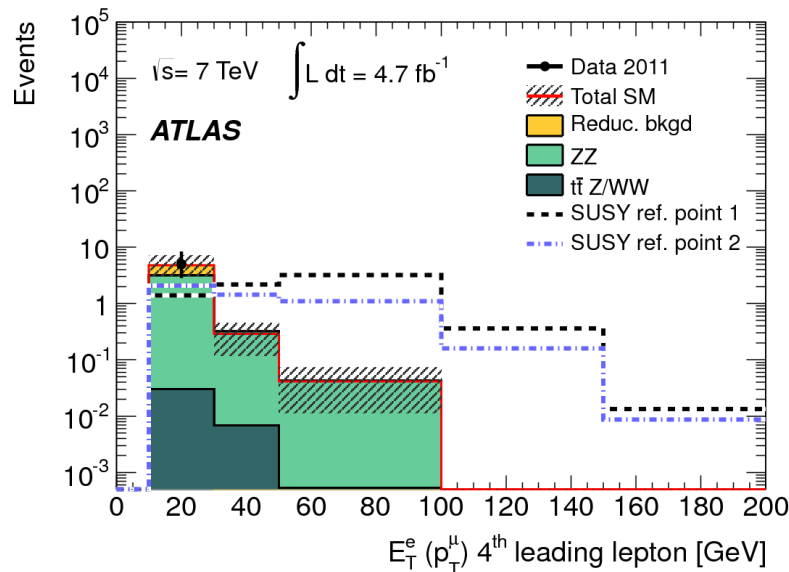
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RPV@decay: Multilepton

- Signature expected in case of RPV in the lepton sector

$$W_{RPV} = \lambda_{ijk} L_i L_j \bar{E}_k$$

- High p_T leptons, possibly non large missing E_T



- Two SR defined based on MET and total transverse energy (m_{eff})
- Validation regions used for SM background checks:
 - Reducible bkg (fakes, from data)
 - Irreducible bkg (ZZ, ttZ, ttWW): from MC

Selection	SR1	SR2	VR1	VR2	VR3
Number of leptons	≥ 4	≥ 4	3	≥ 4	≥ 4
Z-candidate	veto	veto	veto	requirement	veto
$E_T^{\text{miss}}/\text{GeV}$	> 50	-	> 50	-	< 50
$m_{\text{eff}}/\text{GeV}$	-	> 300	-	-	< 300

$$m_{\text{eff}} = E_T^{\text{miss}} + \sum_{\mu} p_T^\mu + \sum_e E_T^e + \sum_j E_T^j$$

Multilepton: background

- **Reducible background (WZ, ttW, tt, WW):**

- Events with 3 real and one fake leptons, or 2 real and 2 fake (3,4 fakes negligible)
- Estimated with a *weighting method*

$$[N_{\text{data}}(3\ell_S + \ell_L) - N_{\text{MCirr}}(3\ell_S + \ell_L)] \times F(\ell_L) - [N_{\text{data}}(2\ell_S + \ell_{L1} + \ell_{L2}) - N_{\text{MCirr}}(2\ell_S + \ell_{L1} + \ell_{L2})] \times F(\ell_{L1}) \times F(\ell_{L2})$$

$$F = \sum_{i,j} (\alpha^i \times R^{ij} \times f^{ij})$$

Fake ratio:

- ℓ_S and ℓ_L : signal and loose leptons
 - i: type of fakes (HF, conversion)
 - j: process category
 - f^{ij} : fake rate of fake type i process j
- F weighted for the fractional contribution of each process

Selection	VR1	VR2	VR3
SUSY ref. point 1	4.6 ± 0.5	1.38 ± 0.16	0.004 ± 0.006
SUSY ref. point 2	3.5 ± 0.4	2.32 ± 0.34	0.120 ± 0.029
ZZ	3.2 ± 1.9	38 ± 7	2.9 ± 1.5
ttZ	0.70 ± 0.35	0.64 ± 0.32	$(3.5 \pm 3.6) \times 10^{-3}$
ttWW	0.08 ± 0.06	$(1.3 \pm 1.0) \times 10^{-3}$	$(2.2 \pm 2.1) \times 10^{-4}$
WZ (†)	34.6 ± 6	–	–
ttW (†)	2.6 ± 0.8	–	–
Σ Irreducible	41 ± 8	38 ± 7	2.9 ± 1.5
Reducible	95 ± 32	-0.25 ± 0.96	1.3 ± 1.3
Σ SM	136 ± 33	38 ± 7	4.2 ± 1.8
Data	152	40	2
p_0 -value (σ)	0.33 (0.45)	0.42 (0.21)	0.80 (–0.85)

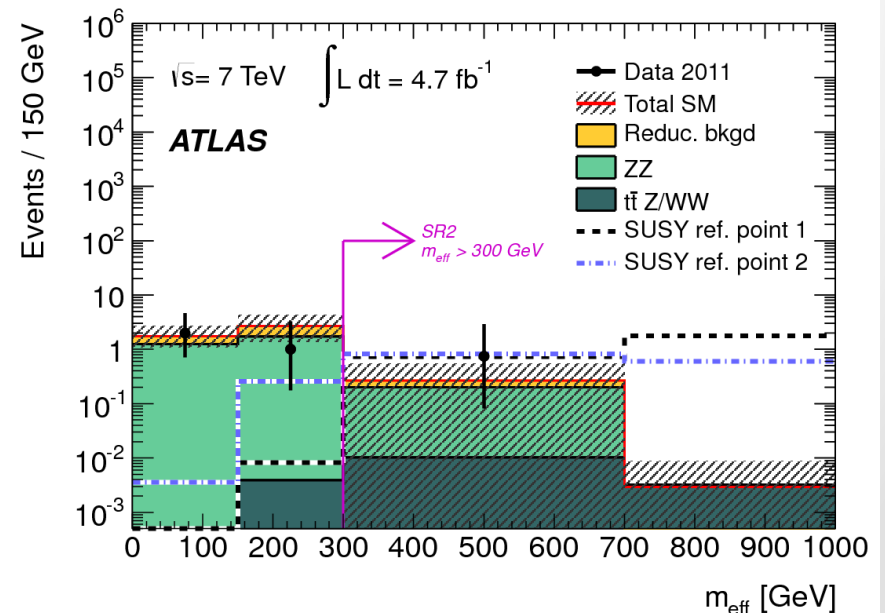
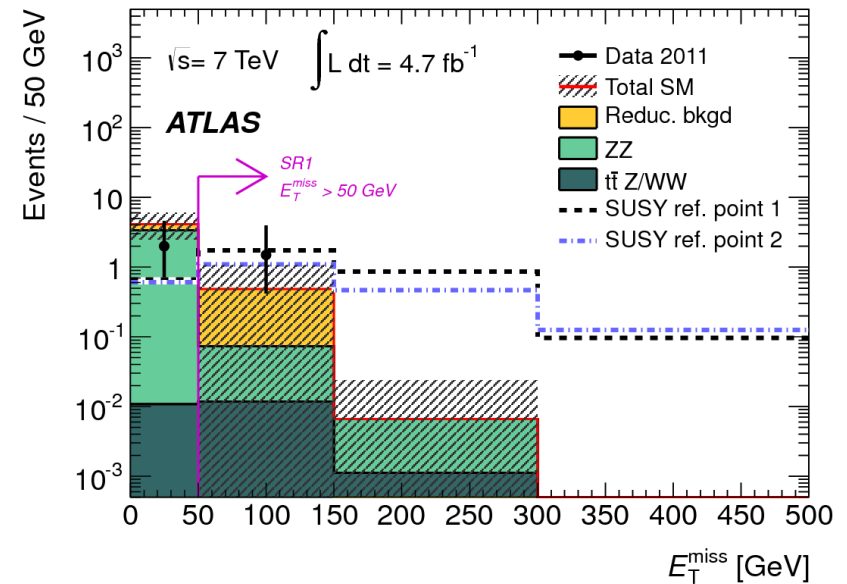
- **Irreducible background:**

- ZZ, ttZ, ttWW

Multilepton: results

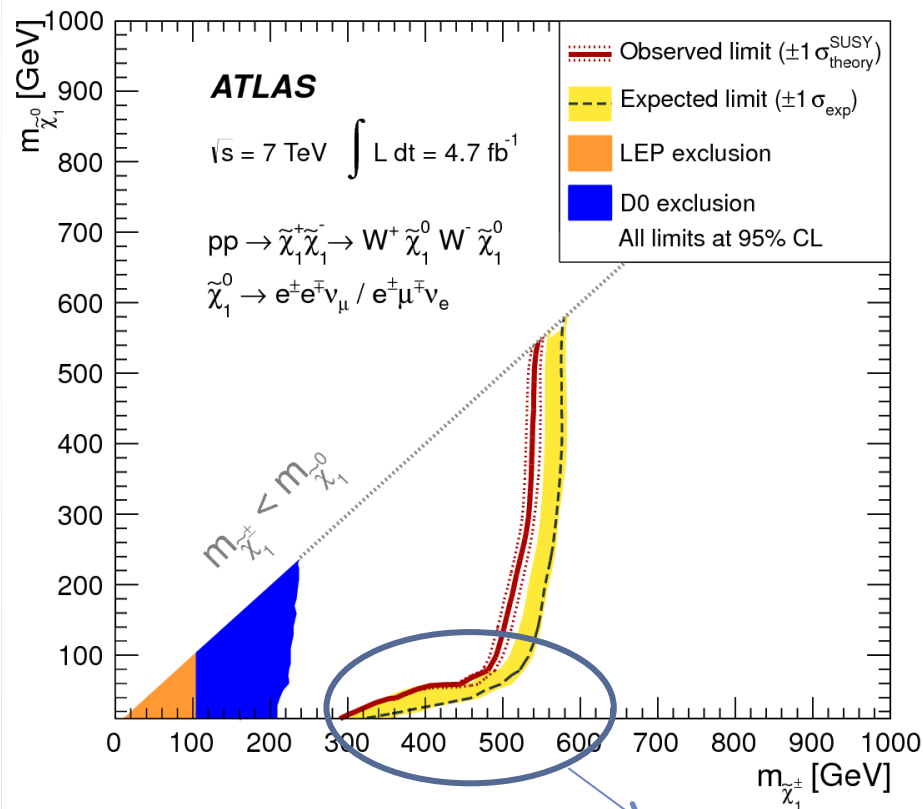
- Major uncertainties:
 - Irreducible bkg: generator uncertainties on ZZ MC modeling
 - Reducible bkg: dominated by unc. on contributions from each process

Selection	SR1	SR2
SUSY ref. point 1	6.5 ± 0.6	7.1 ± 0.7
SUSY ref. point 2	4.2 ± 0.6	4.5 ± 0.6
ZZ	0.14 ± 0.11	0.51 ± 0.30
$t\bar{t}Z$	0.023 ± 0.014	0.029 ± 0.016
$t\bar{t}WW$	0.0044 ± 0.0035	0.005 ± 0.004
Σ Irreducible	0.17 ± 0.12	0.54 ± 0.31
Reducible	0.8 ± 0.8	0.18 ± 0.26
Σ SM	1.0 ± 0.8	0.7 ± 0.4
Data	3	2
p_0 -value (σ)	0.05 (1.7)	0.07 (1.5)
σ_{vis} obs (exp)	1.3 (0.8)	1.1 (0.7)

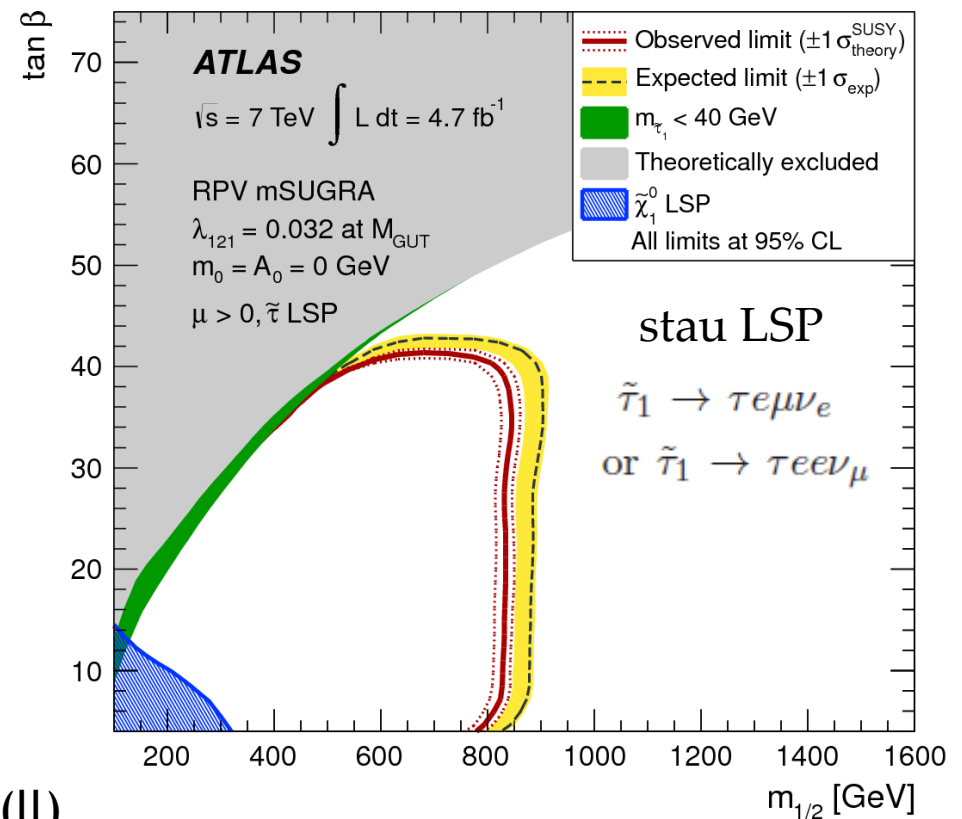


Multilepton: interpretation

- Interpretation in simplified models for chargino pair production as well as mSUGRA/cMSSM breaking models

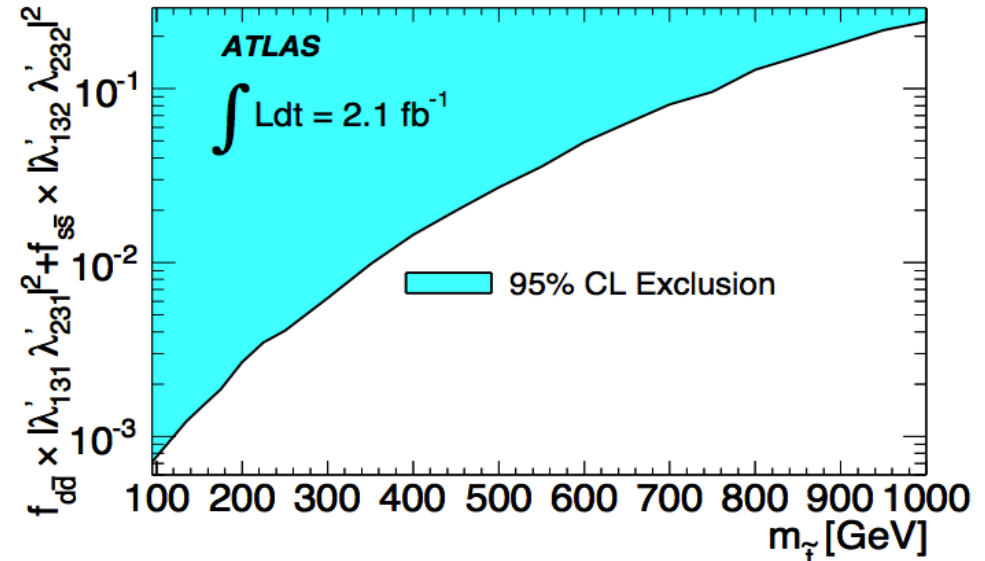
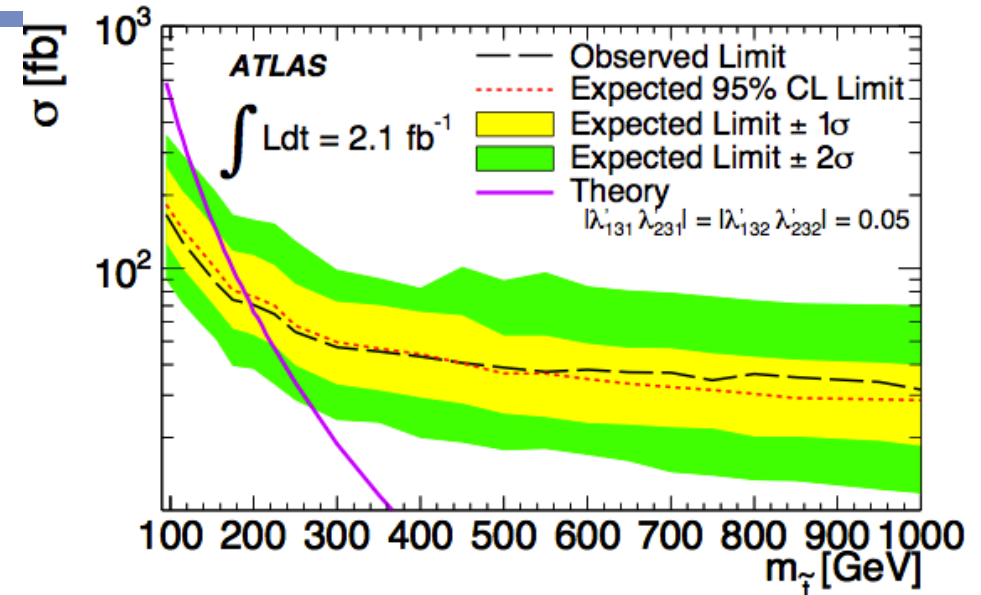
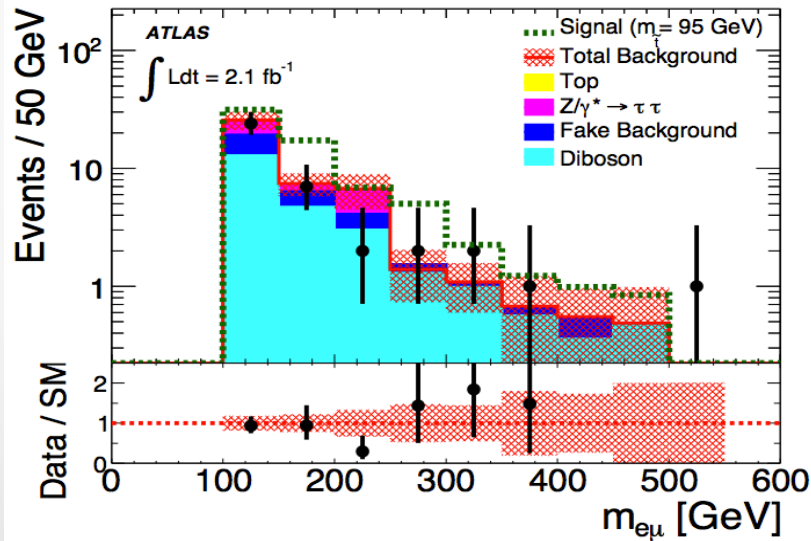
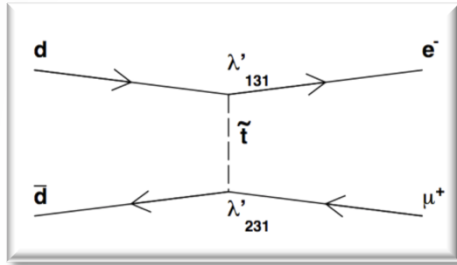


Low $\Delta R(\text{leptons})$ and low $m(\text{ll})$



RPV@Production

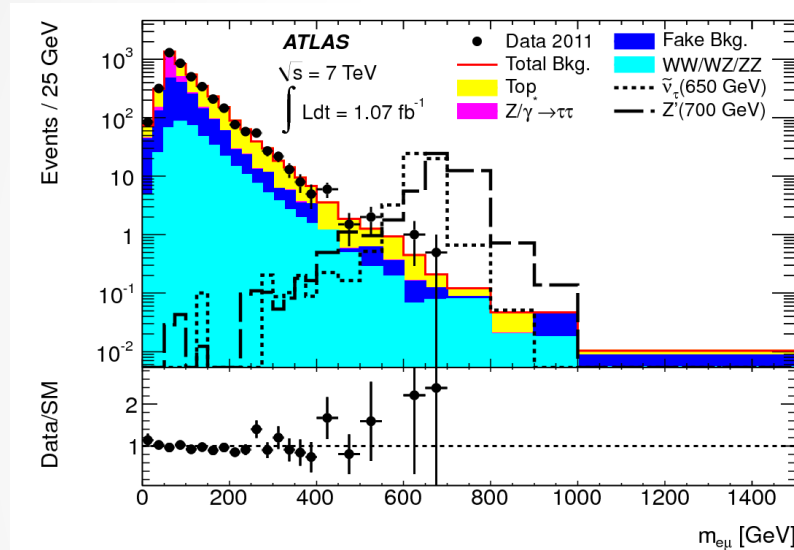
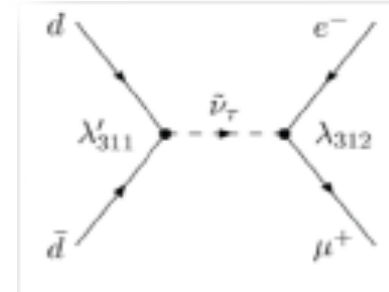
- Look for excess in the tail of different flavor di-leptons
- In this case: $e\mu$ invariant mass
 - No jets, low E_{miss}^T



e-mu resonances: sneutrinos

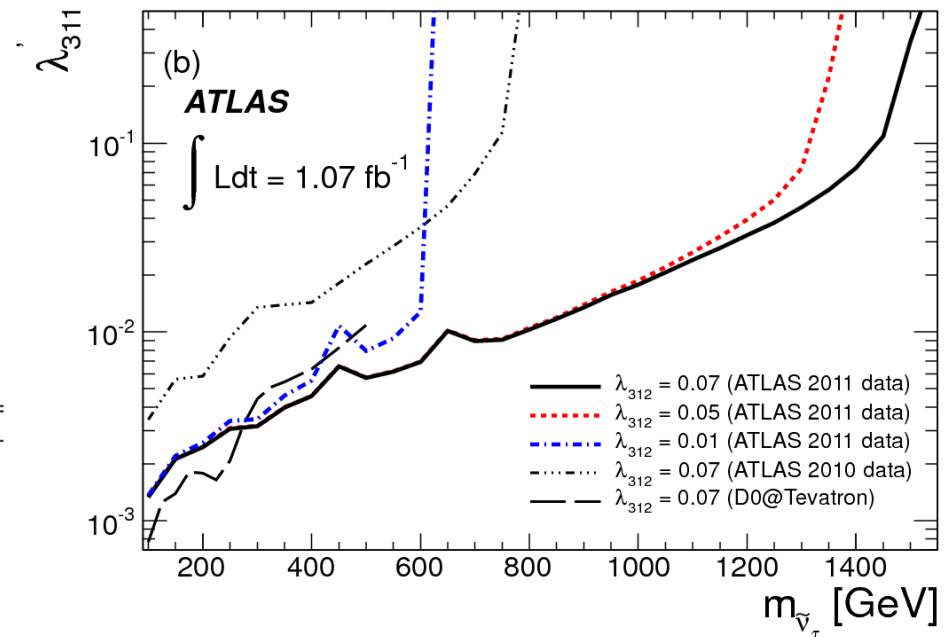
RPV couplings at production and decay

- SM background dominated by fakes and ttbar



Scan in
 $e\mu$ mass

$m_{e\mu}$	Data	SM prediction
> 200 GeV	286	288 ± 22
> 250 GeV	152	136 ± 11
> 300 GeV	70	67 ± 6
> 350 GeV	35	34.0 ± 3.0
> 400 GeV	22	17.7 ± 1.7
> 450 GeV	10	10.5 ± 1.2
> 500 GeV	7	6.8 ± 0.9
> 550 GeV	3	4.3 ± 0.6
> 600 GeV	3	2.4 ± 0.4
> 650 GeV	1	1.49 ± 0.31
> 700 GeV	0	1.07 ± 0.25

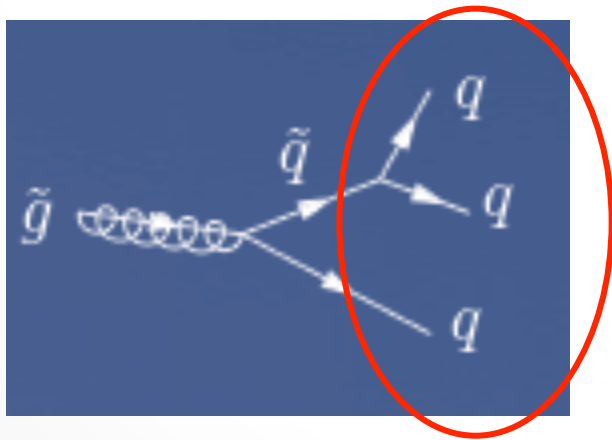


- **‘prompt’ RPV SUSY:**
 - **Leptonic RPV:**
 - search for multilepton final states
 - emu resonances and continuum
 - **Hadronic RPV:**
 - search for 2x3jet resonances
 - search for 2x2jet resonances (currently focusing on scalar gluons searches)
- **Long-lived SUSY:**
 - Displaced vertex
 - Disappearing tracks
 - R-hadrons and sleptons

'Hadronic' RPV: 2x3jets

- UDD couplings least constrained via collider experiments
 - Very hard given the all-hadronic decays and overwhelming multijet bkg
- Several interesting possibilities:
 - Here, consider **2 x 3jets** for gluino pair production focusing on **uds λ'' couplings**

$$W_{Rp} = \frac{1}{2} \lambda''_{ijk} \bar{U}_i \bar{D}_j \bar{D}_k$$



2 analyses:

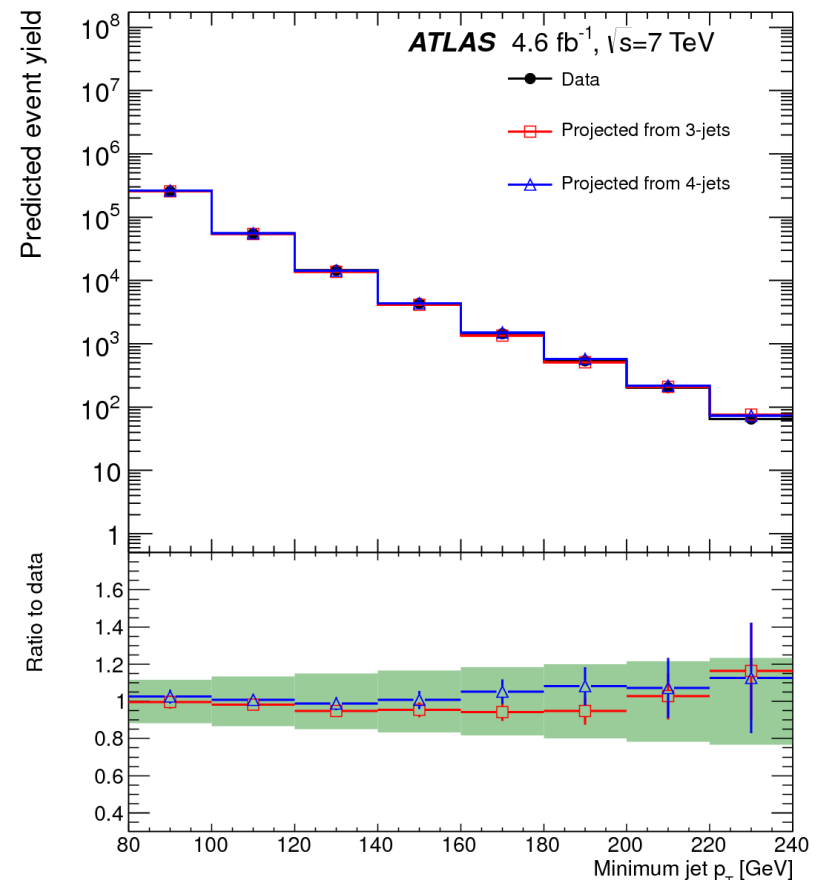
- Resolved
- Boosted

2x3jets: resolved

- signal discriminated from multijet background by exploiting the large p_T of jets produced in gluino decays.
- p_T of 6th leading jets used to discriminate signal from bkg:
 - Three selections used: **> 80 GeV, 120 GeV and 160 GeV**
- **Multijet bkg** estimated using normalization procedure from low to high N jets

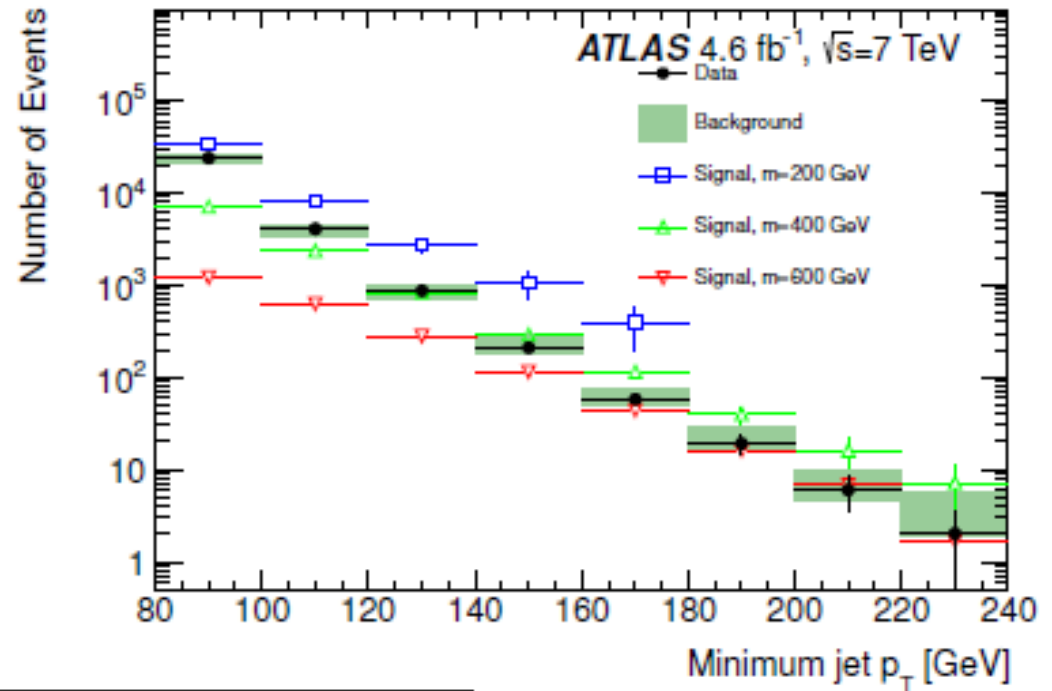
$$N_{data}^{n-jet} = N_{data}^{m-jet} \frac{N_{MC}^{n-jet}}{N_{MC}^{m-jet}}$$

- Tested from 3j / 4j \rightarrow 5j (signal depleted)
- Final estimate using 3 jets \rightarrow \geq 6 jets



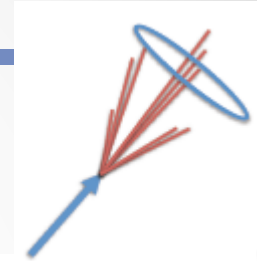
2x3jets: resolved (II)

- Good agreement between data and SM expectation
- Systematic uncertainties mostly from JES and JER, followed by PDF
 - Systematics on color flow and QCD radiation not considered due to the poor knowledge in case of RPV-vertex



Model ($m_{\tilde{g}}$)	$p_{T,\min}^{6\text{th-jet}}$	Data	Background	Signal bias [%]	Signal
100 GeV	80 GeV	23600	23500 ± 2800	8.5	99200 ± 20000
200 GeV	120 GeV	856	851 ± 140	3.7	2700 ± 500
300 GeV	120 GeV	856	851 ± 140	1.0	1460 ± 240
400 GeV	160 GeV	57	62 ± 13	0.8	110 ± 13
500 GeV	160 GeV	57	62 ± 13	0.3	67 ± 9
600 GeV	160 GeV	57	62 ± 13	0.1	43 ± 7
800 GeV	160 GeV	57	62 ± 13	0.0	20 ± 3

2x3jets: boosted jets



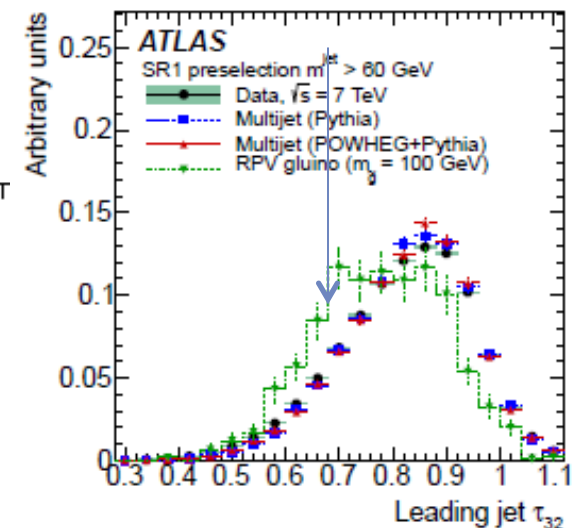
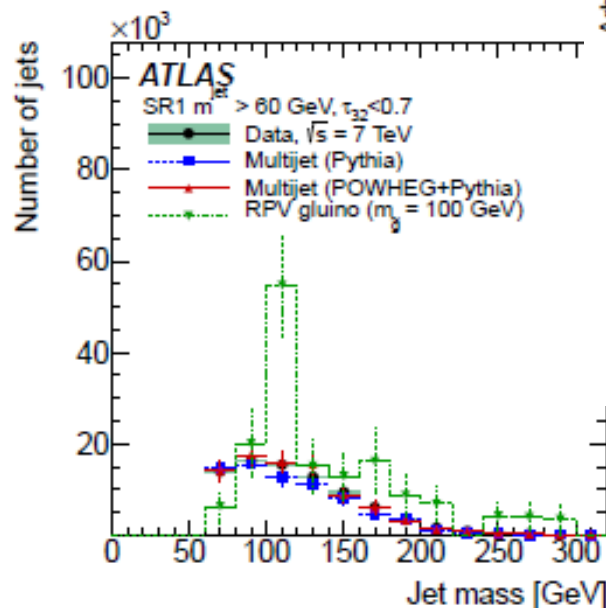
- Jet shape variable sensitive to the N-body structure from jets containing boosted-gluino decay products.

N-subjettiness
$$\tau_N = \frac{1}{d_0} \sum_k p_{T_k} \times \min(\delta R_{1k}, \delta R_{2k}, \dots, \delta R_{Nk}) , \quad \text{with} \quad d_0 \equiv \sum_k p_{T_k} \times R$$

- $\tau_{32} = \tau_3 / \tau_2$
 \rightarrow if $\approx 1(0)$, jet well described by 2(3)-subjets

After that, the trimming procedure is applied:

- Soft radiation removed
- Invariant mass of the jet considered (if $\tau_{32} < 0.7$)



2x3jets: boosted (II)

- Select events with two massive large-R jets and $\tau_{32} < 0.7$
- High normal jet multiplicity also required to reduce multijet bkg
 - Use ABCD method

Selection	Baseline Selection	SR1	SR2
Small- R ($R = 0.4$) jet p_T^{jet}	$p_T^{\text{jet}} > 30$ GeV	$p_T^{\text{jet}} > 30$ GeV	$p_T^{\text{jet}} > 30$ GeV
Large- R ($R = 1.0$) jet p_T^{jet}	$p_T^{\text{jet}} > 200$ GeV	$p_T^{\text{jet}} > 200$ GeV	$p_T^{\text{jet}} > 350$ GeV
Scalar sum $\sum_{i=1}^{N_{\text{jet}}^{R4}} p_T^{\text{jet}}$	(—)	600 GeV	(—)
Small- R jet multiplicity	(—)	$N_{\text{jet}}^{R4} \geq 4$	$N_{\text{jet}}^{R4} \geq 4$
Large- R jet multiplicity	$N_{\text{jet}} \geq 2$	$N_{\text{jet}} \geq 2$	$N_{\text{jet}} \geq 2$
Large- R jet mass	(—)	$m_{J_1, J_2}^{\text{jet}} > 60$ GeV	$m_{J_1, J_2}^{\text{jet}} > 140$ GeV
Large- R jet τ_{32}	(—)	$\tau_{32} < 0.7$	$\tau_{32} < 0.7$

Region	Jet (J_1) selections	Jet (J_2) selections	Description
CR-A	$m^{\text{jet}} < M_{\text{threshold}}$	$m^{\text{jet}} < M_{\text{threshold}}$	Low-mass jets, to validate τ_{32} shape
CR-B	$m^{\text{jet}} > M_{\text{threshold}}$ $\tau_{32} < 0.7$	$m^{\text{jet}} < M_{\text{threshold}}$	Signal-like leading jet, to validate m^{jet}
CR-C	$m^{\text{jet}} < M_{\text{threshold}}$	$m^{\text{jet}} > M_{\text{threshold}}$ $\tau_{32} < 0.7$	Signal-like subleading jet, to validate m^{jet}

$$N_{SR} = N_{CR-C} \times \left(\frac{N_{CR-B}}{N_{CR-A}} \right) \times \alpha$$

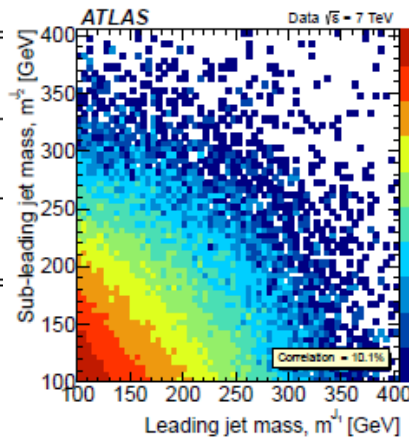
$$\alpha = \left(\frac{N_{SR} / N_{CR-C}}{N_{CR-B} / N_{CR-A}} \right) \Big|_{\text{MC}}$$

α = correlation factor, evaluated from MC

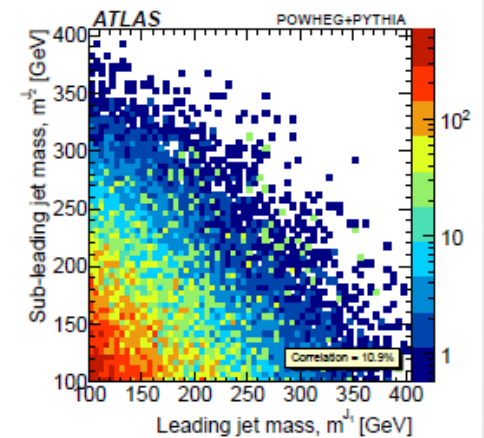
- Monica D'Onofrio, GGI, Florence

Data

MC (POWHEG+PYTHIA)



(c) Data, $m^{\text{jet}} > 100$ GeV

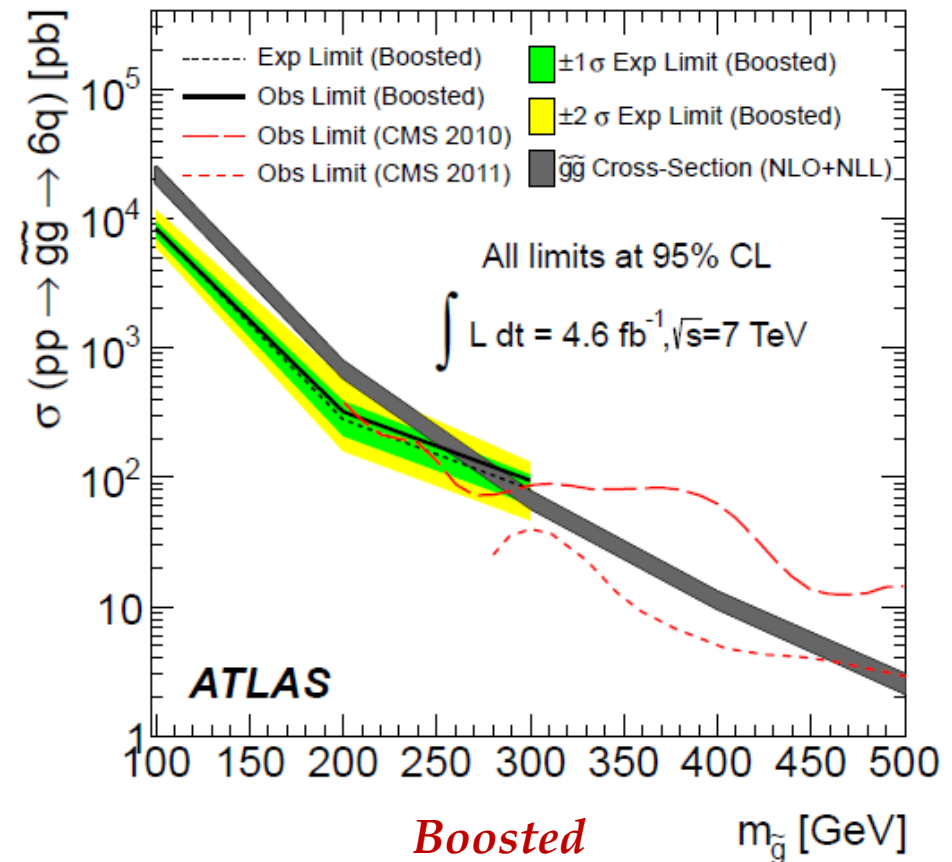
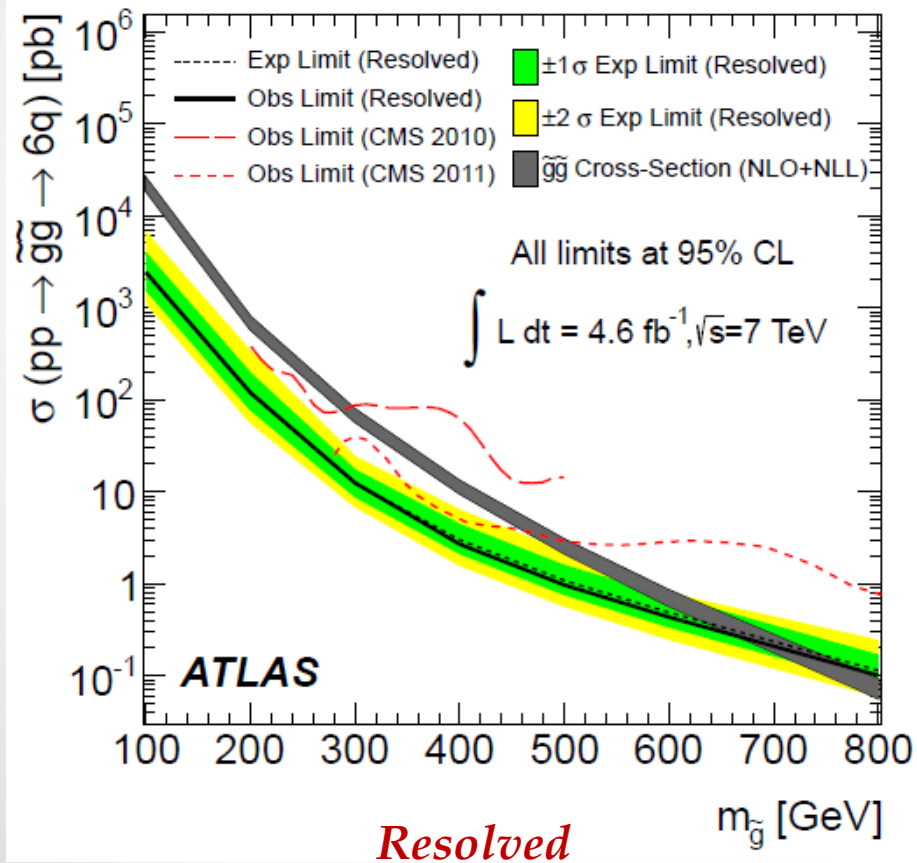


(d) POWHEG+PYTHIA, $m^{\text{jet}} > 100$ GeV

Agreement with data is tested looking at correlations of 1st and 2nd leading jet M

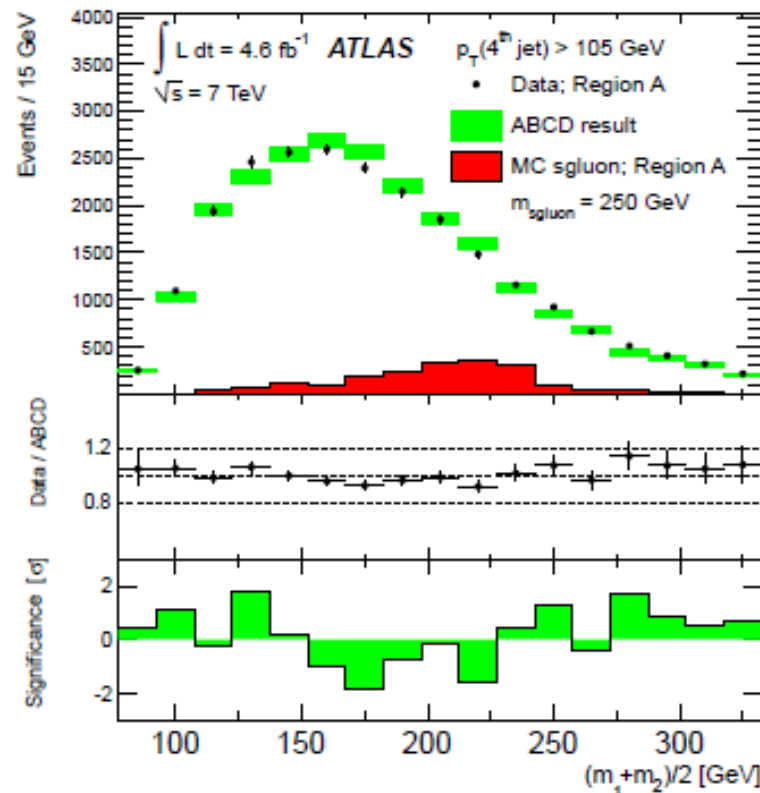
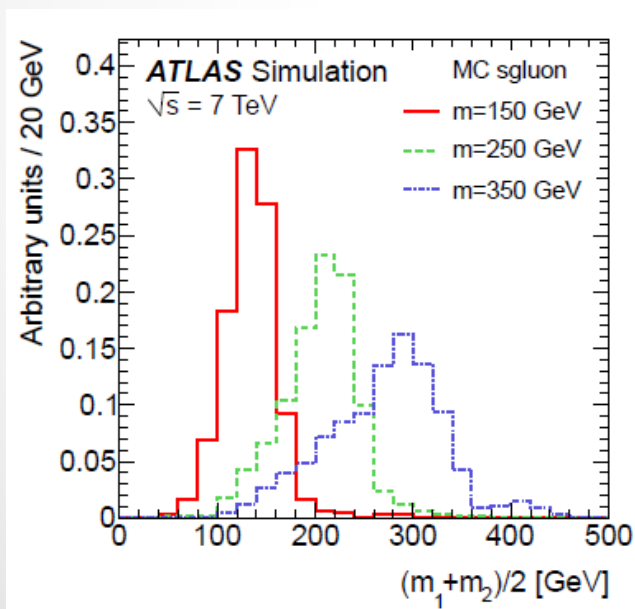
2x3jets: results

- Boosted analysis sensitive for $m(\text{gluinos})$ up to 300 GeV
- Resolved analysis excludes $m(\text{gluinos}) < 670$ GeV



2x2jets

- Searches for 2jets resonances potentially sensitive also to RPV sparticle production ($\rightarrow jj$ final states)
 - In this case, search for pair-produced massive coloured scalars in four-jet
 - Background estimate using **ABCD method**:
 - Regions defined by m_1, m_2 and scattering angle, $\cos(\theta^*)$

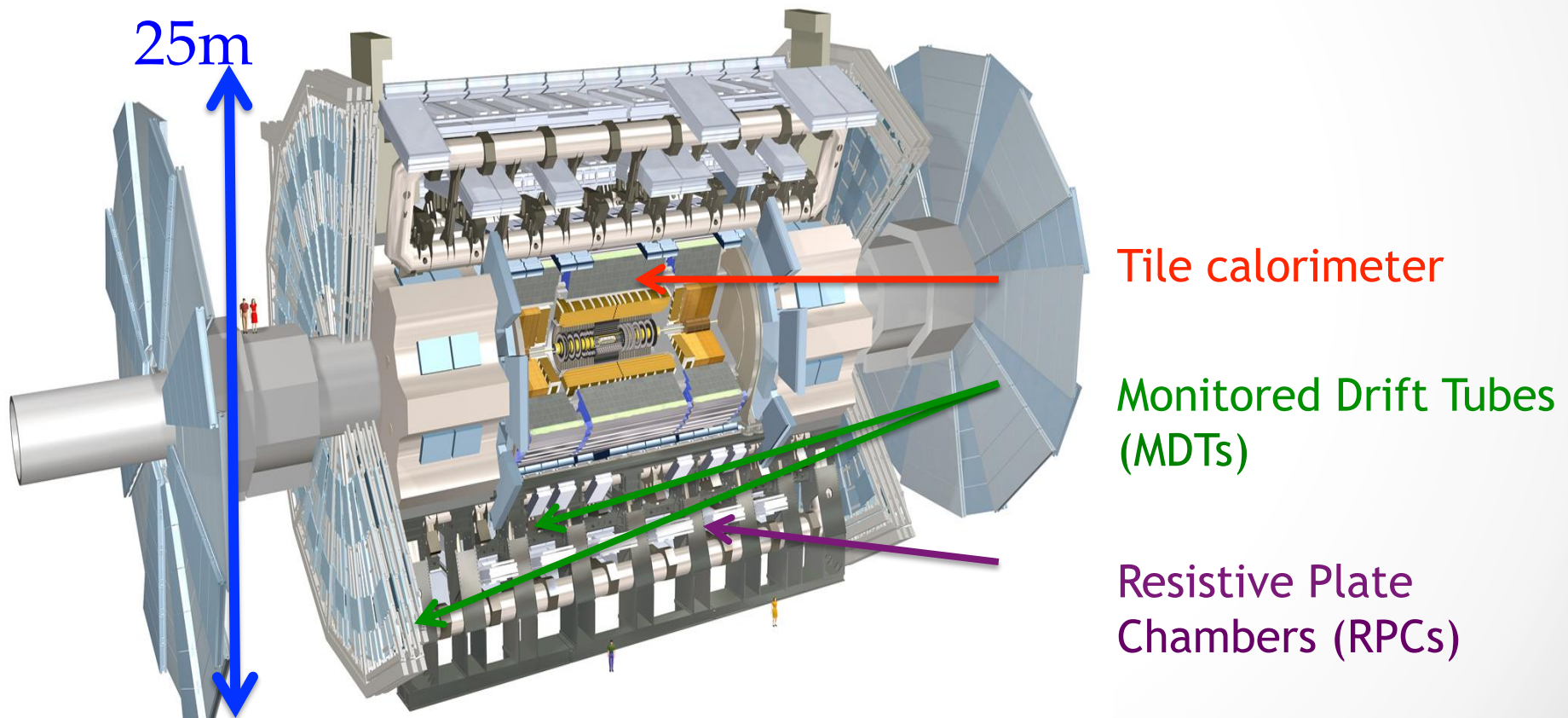


Exclude sgluons
between 150 and
290 GeV

- ‘prompt’ RPV SUSY:
 - Leptonic RPV:
 - search for multilepton final states
 - emu resonances and continuum
 - Hadronic RPV:
 - search for 2x3jet resonances
 - search for 2x2jet resonances (currently focusing on scalar gluons searches)
- **Long-lived SUSY:**
 - Displaced vertex
 - Disappearing tracks
 - R-hadrons and sleptons

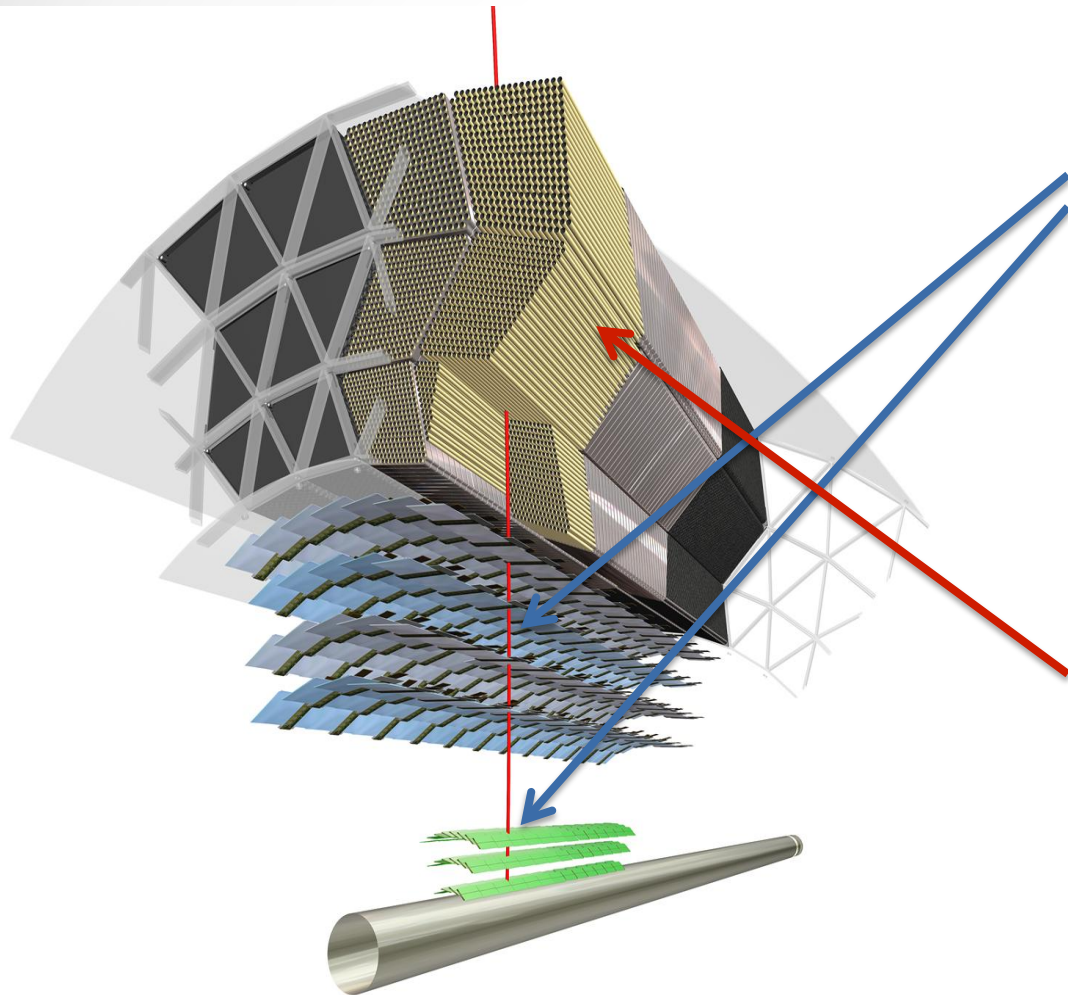
To make it possible: ATLAS detector

Several sub-detectors with excellent time resolution, including (but not only):



Can measure time-of-flight

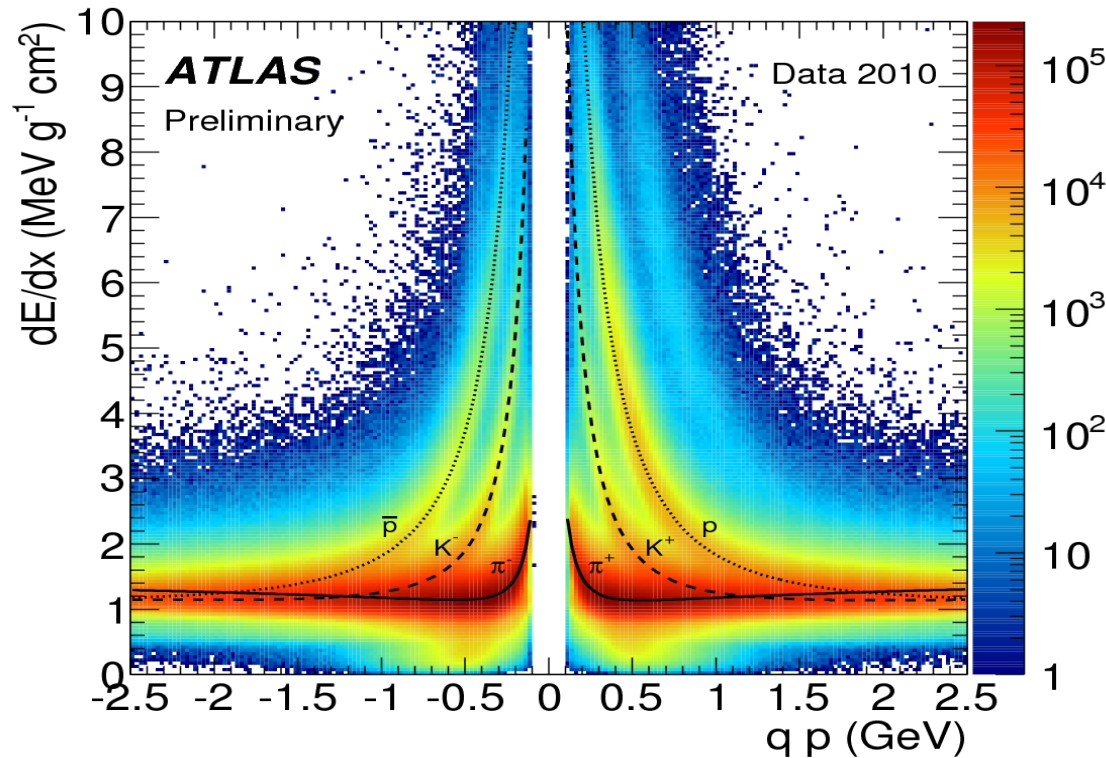
The ATLAS Inner Detector (ID)



- Pixel detector
- Semiconductor Tracker (SCT)
 - Good for finding vertices
- Transition Radiation Tracker (TRT) → a continuous tracker
 - Can detect kinked or disappearing tracks

All within a 2T solenoidal B-field

dE/dx measurements

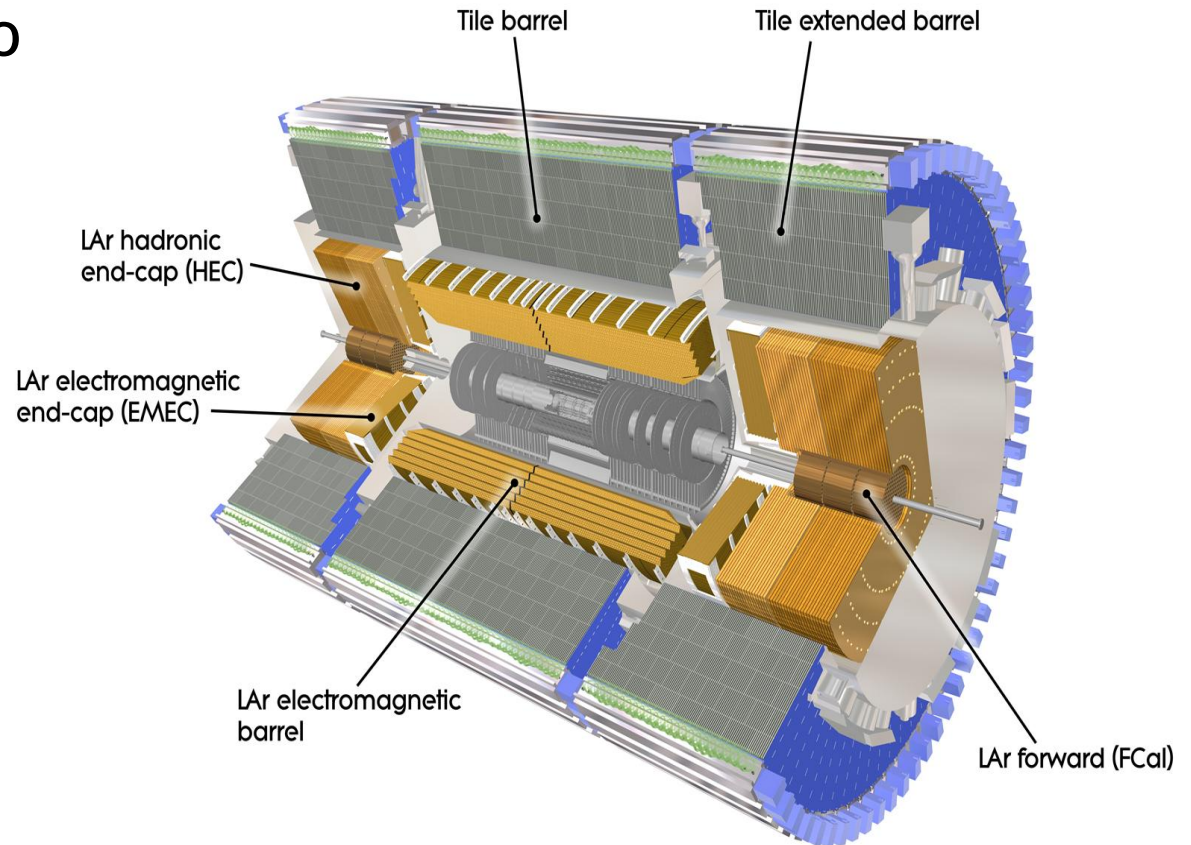
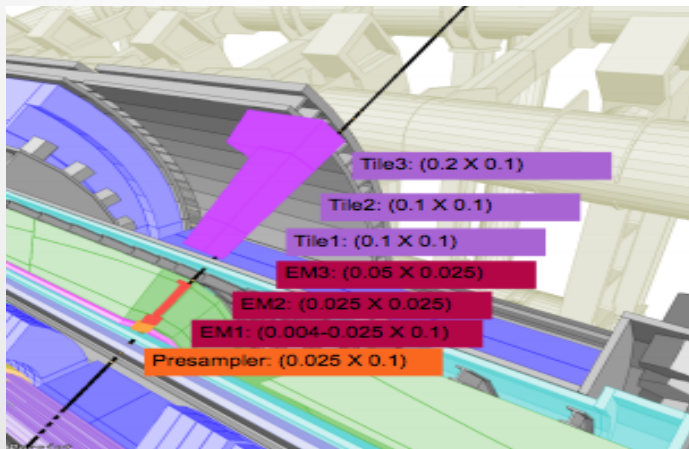


Pixel detector can measure ionization energy loss dE/dx via charge deposited (calculated from Time-over-Threshold)

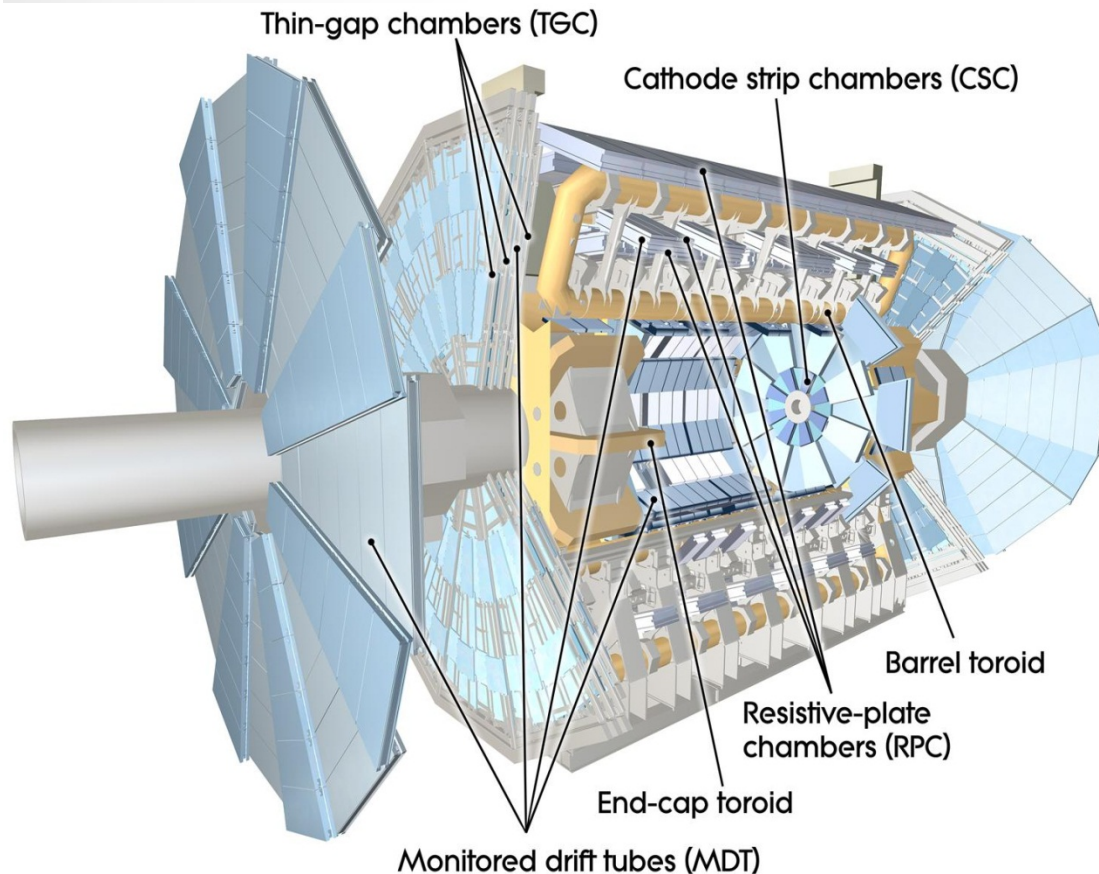
TRT can also measure dE/dx via Time-over-Threshold

The ATLAS Calorimeters

- Both LAr and Tile calorimeters can also measure dE/dx by summing energy deposits over path length.



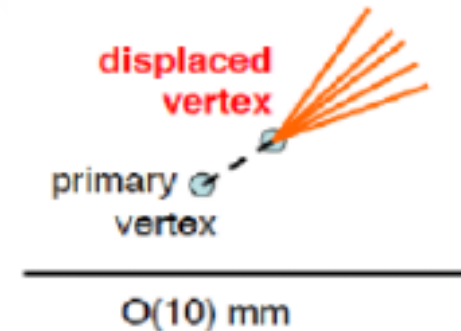
The ATLAS Muon Spectrometer



- Precision muon chambers can reconstruct “standalone” tracks
 - i.e. can find particles that did not leave tracks in Inner Detector (e.g. decay products of LLPs)
- MDTs can also measure dE/dx (similar principle to TRT).

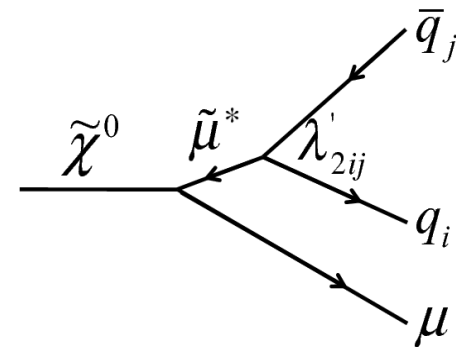
Displaced vertices in ID

- Particles with lifetimes of order a few picoseconds could decay within the inner tracking detector, giving rise to **displaced vertices**



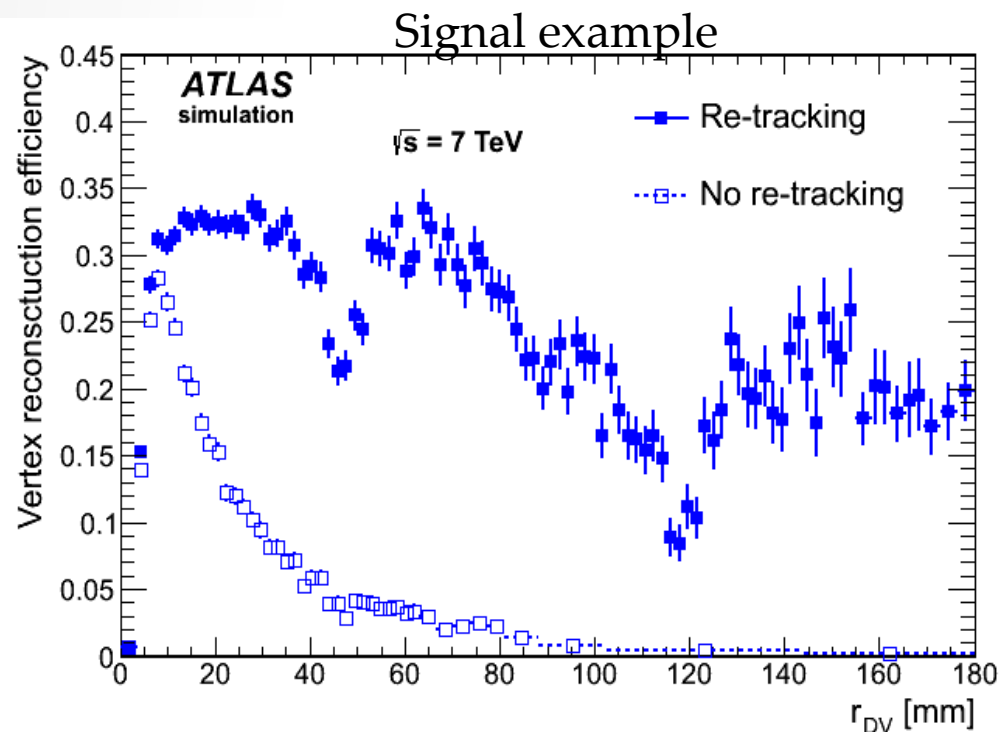
→ *Here*: consider **muon+jets** topology, i.e. from RPV neutralinos decay with a non-zero (but small) λ'_{211} coupling

- Presence of a muon useful for triggering and background rejection
- High track multiplicity helps vertex reconstruction



track and vertex reconstruction

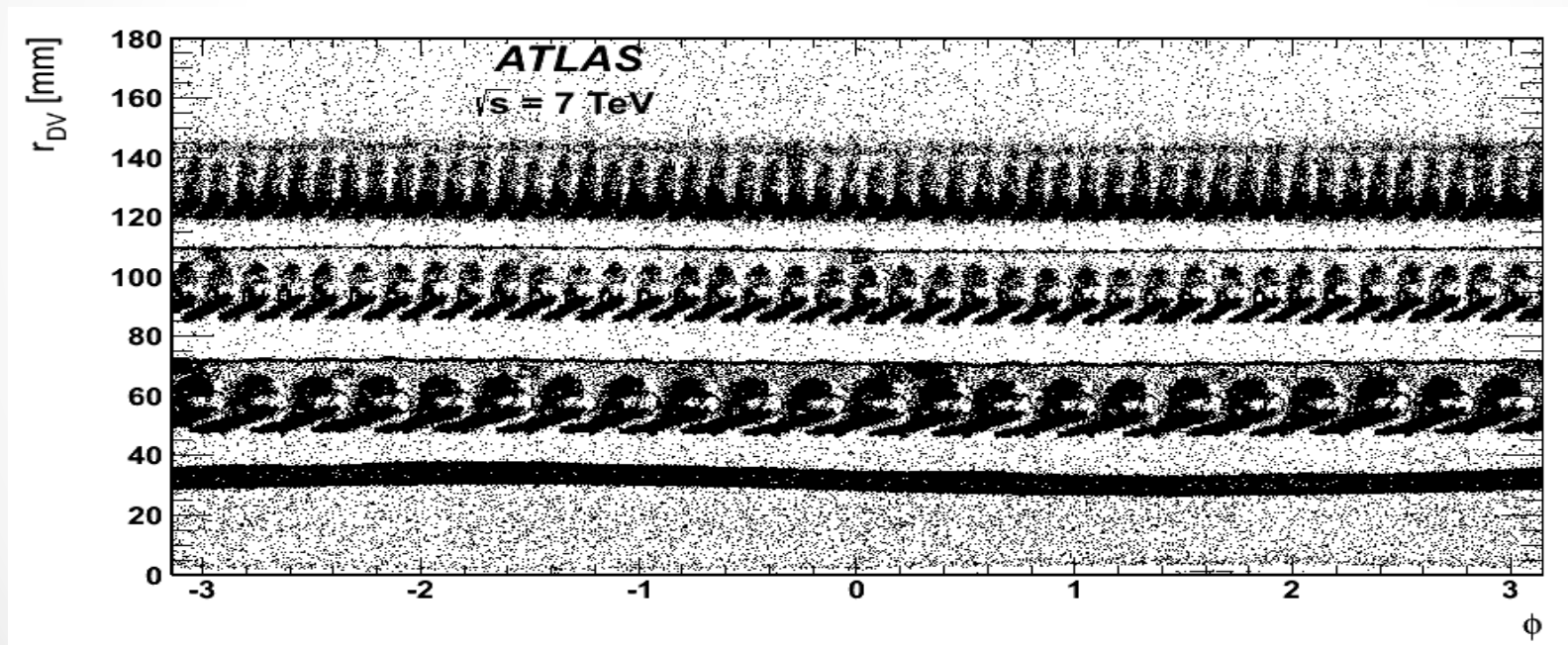
- Standard ATLAS tracking is highly optimized for tracks coming from the primary interaction point (IP)
- To increase efficiency for secondary tracks
 - re-run Silicon-seeded tracking algorithm, with looser cuts on transverse impact parameter, using “left-over” (shared) hits from Standard tracking.



- Vertex-finding algorithm based on an iterative disambiguation process
- then splits/merges/refits vertices until no tracks are shared between vertices.

Displaced vertices - selection

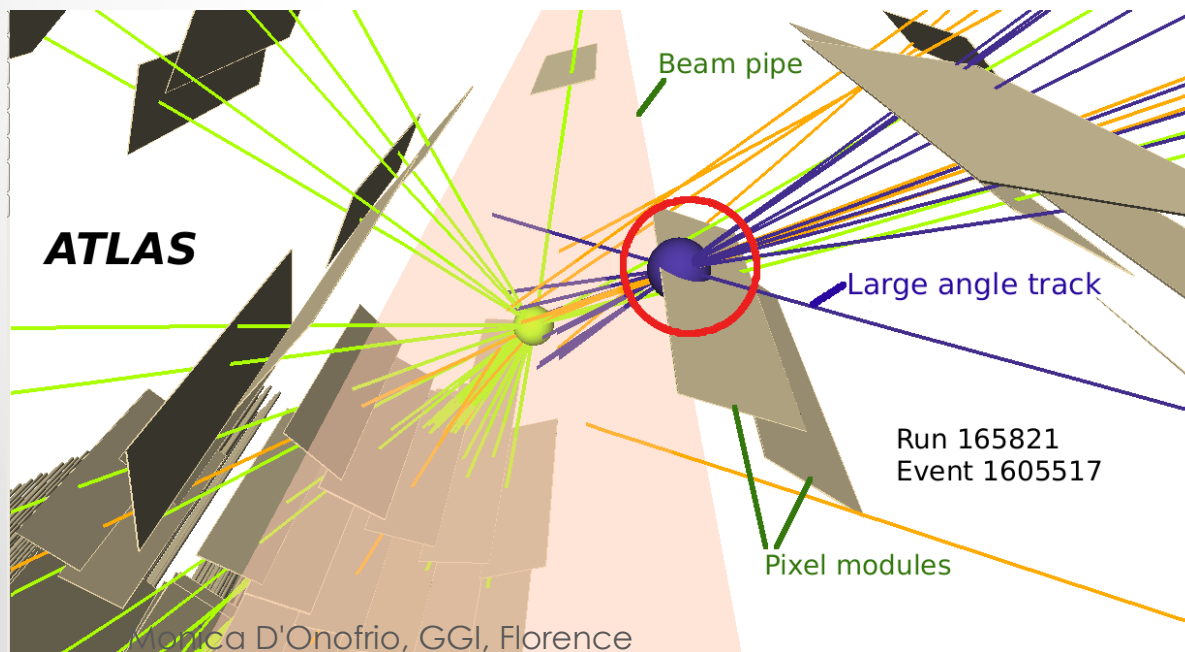
- Events triggered by high- p_T muon trigger, with no ID track requirement.
- Use tracks with $|d_0| > 2\text{mm}$, $p_T > 1\text{ GeV}$ as input to vertexing
- Look in fiducial volume roughly corresponding to Pixel barrel
- Require at least 5 tracks in vertex, and mass $> 10\text{ GeV}$.
- Require high- p_T muon passing within 0.5mm of reco vertex



→ Veto vertices reconstructed in regions
with high material density

Displaced vertices - backgrounds

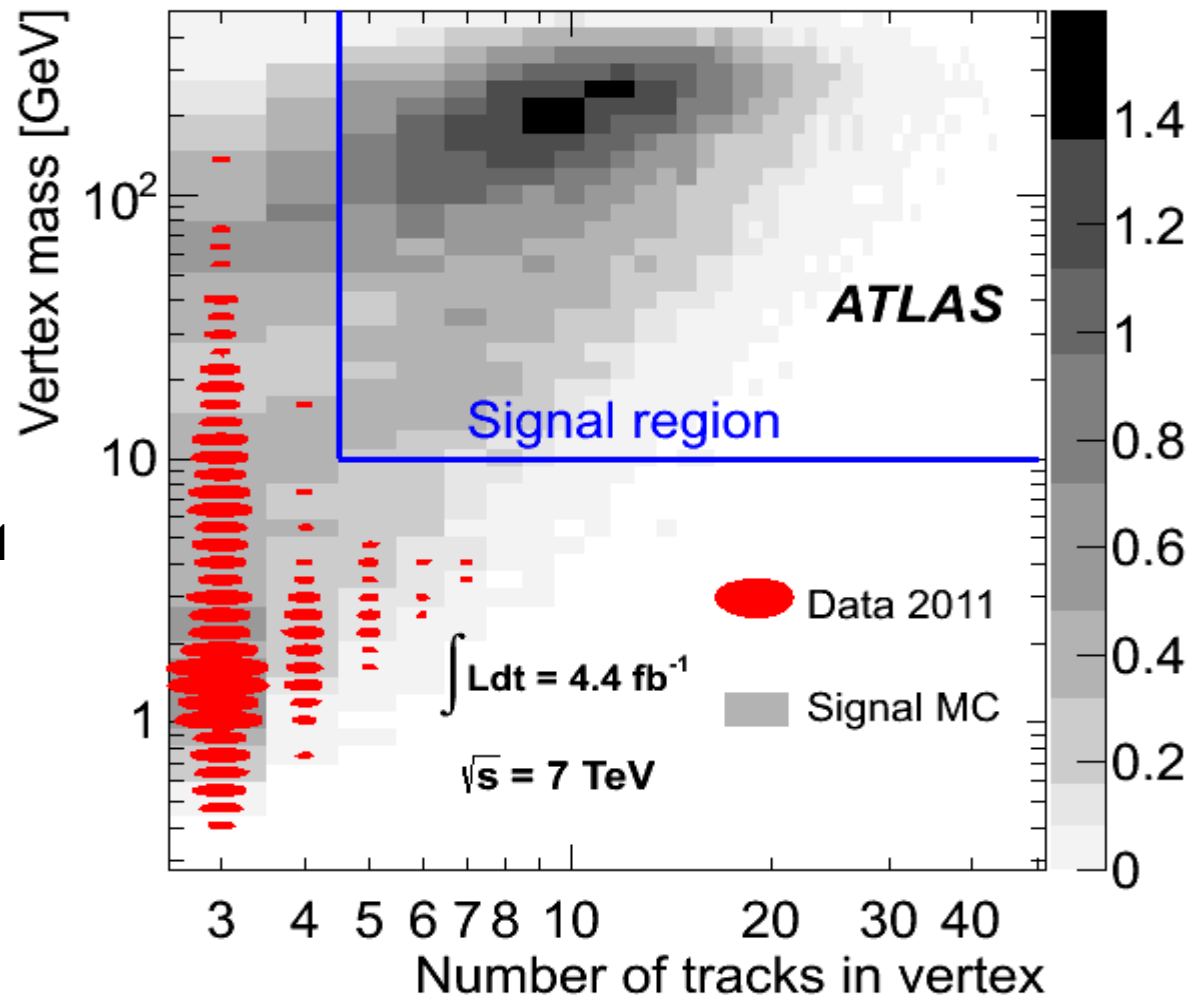
- Two sources of background vertices considered:
 - Purely random combinations of tracks inside the beampipe (where vacuum is good, but track density is high).
 - High-mass tail of distribution of real vertices from hadronic interactions with gas molecules.
 - Particularly if vertex is crossed by random (real or fake) track at large angle.



- Use different data-driven method for each source: total estimate is $(4 \pm 60) \cdot 10^{-3}$ vertices

Displaced vertices - results

- **Zero** vertices passing selection requirements observed in 4.4 fb^{-1} data sample.

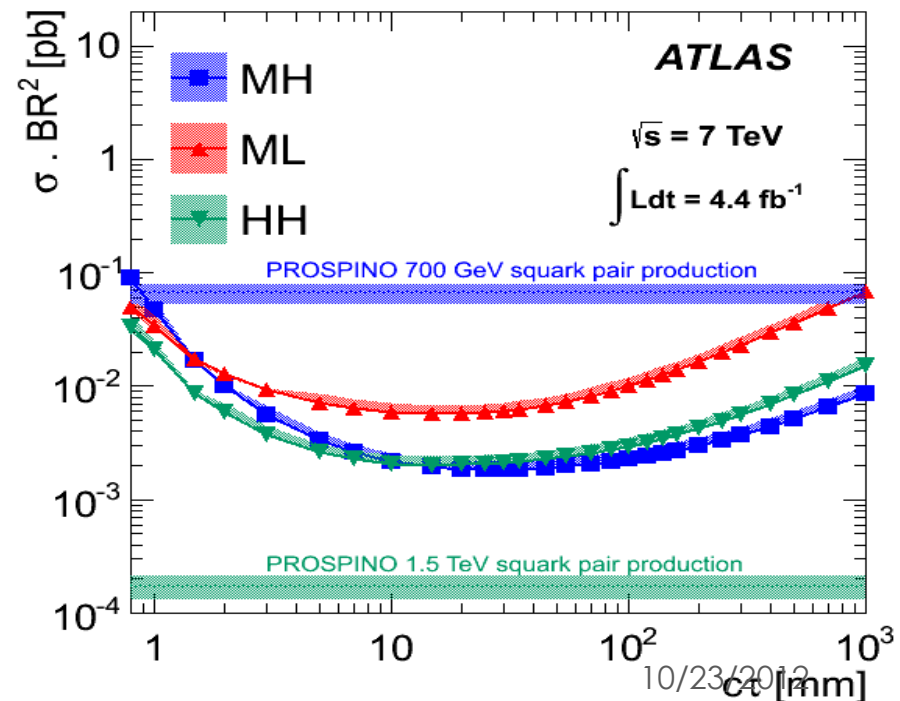
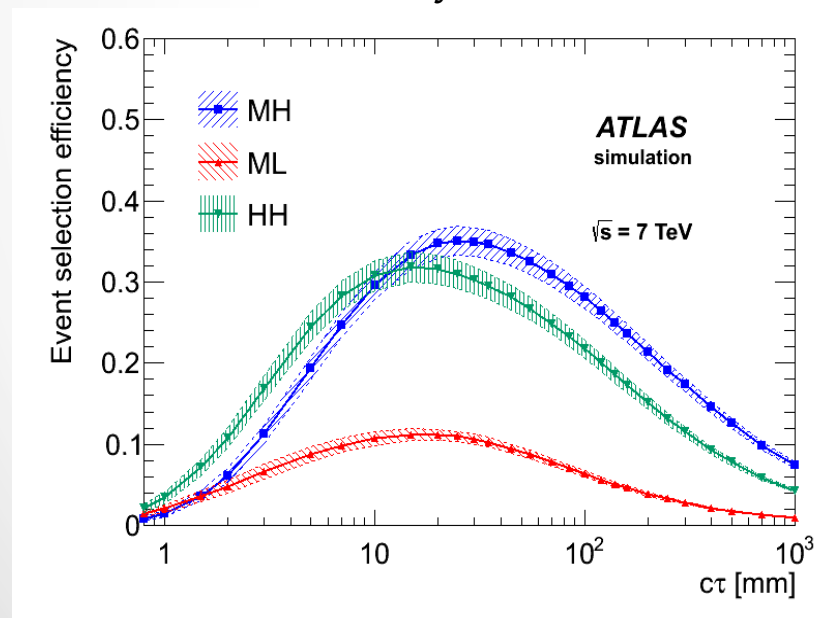


Displaced vertices - interpretation

- Exclusion limits set on squark pair production
 - Assume squark decays directly to LL neutralino, which decays to muon plus jets.
- Three combinations of squark and neutralino masses considered, to study the effect of LLP mass and boost on reconstruction efficiency

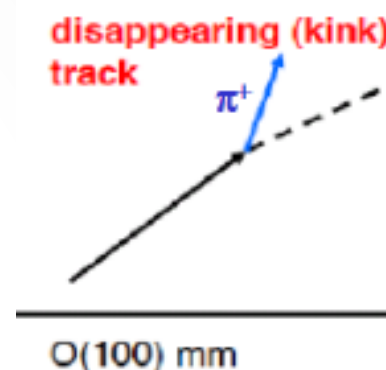
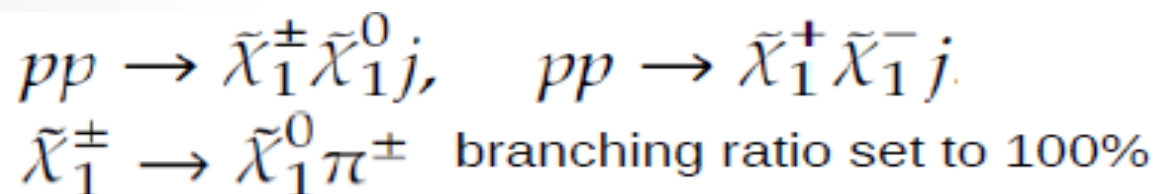
Sample	$m_{\tilde{q}}$ [GeV]	$m_{\tilde{\chi}_1^0}$ [GeV]	$c\tau$ [mm]
MH	700	494	78
ML	700	108	101
HH	1500	494	82

efficiency-vs- $c\tau$



Disappearing tracks

- In AMSB model, could have long-lived charginos decaying to neutralino and soft pion
- Look for production processes as:



Use jet from ISR to trigger on events

- Resulting final state:
 - High p_T jet + large missing ET
 - High- p_T disappearing track (or “kinked” track):
 - Isolated tracks that stop in outer TRT

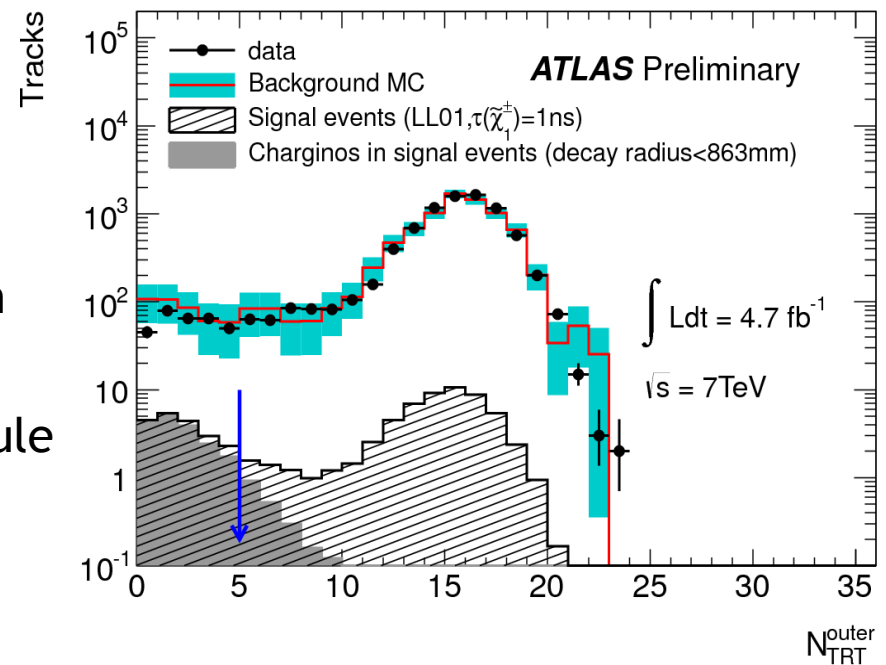
Disappearing tracks - selection

- Event selection:

- missing $E_T > 90\text{GeV}$ and ≥ 1 jet with $p_T > 90\text{GeV}$
- Jets well separated from missing E_T direction in φ ($\Delta\varphi > 1.5$)
- Veto events with reconstructed e or μ candidates.

- Disappearing track candidate selection:

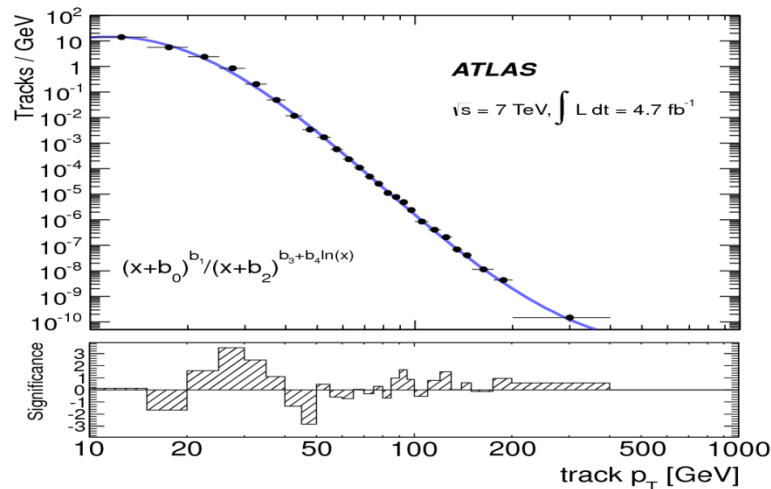
- Track must be isolated
- have $p_T > 10\text{GeV}$,
- at least 1 Pixel hit and 6 SCT hits,
- originate from primary vertex, and point to TRT barrel (but not region around $|\eta|=0$)
- Fewer than 5 hits in TRT outer module



Disappearing tracks - backgrounds

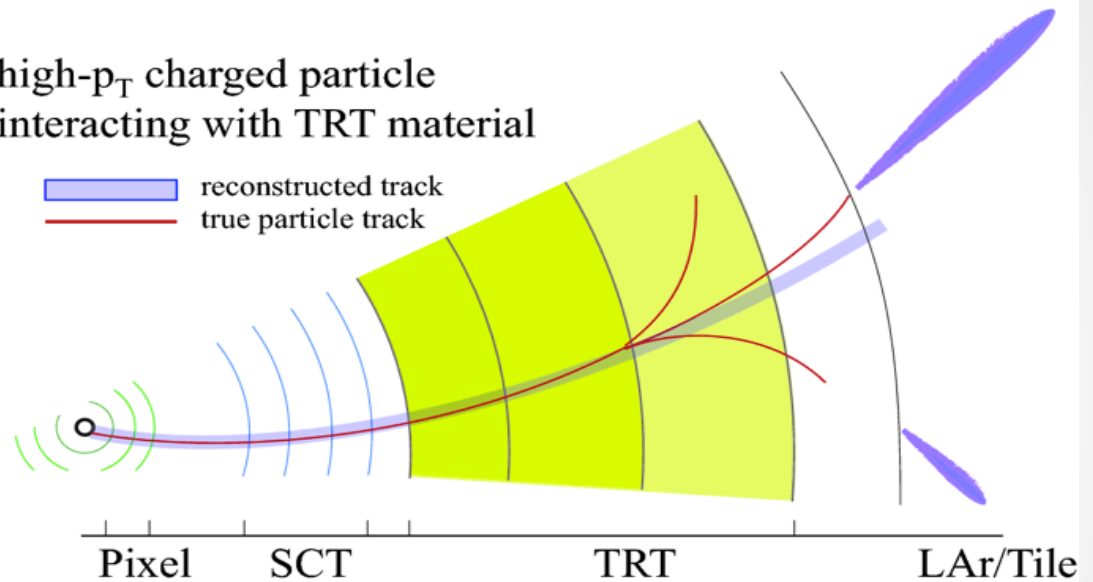
- Potential background sources are:

Charged hadrons
interacting with detector
material (80%)



high- p_T charged particle
interacting with TRT material

reconstructed track
true particle track

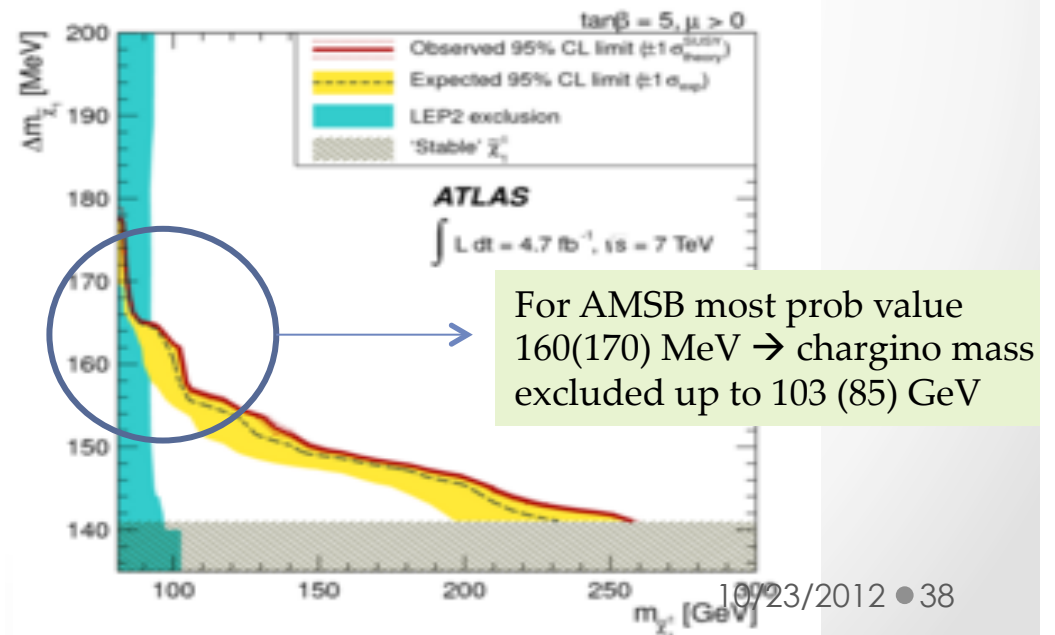
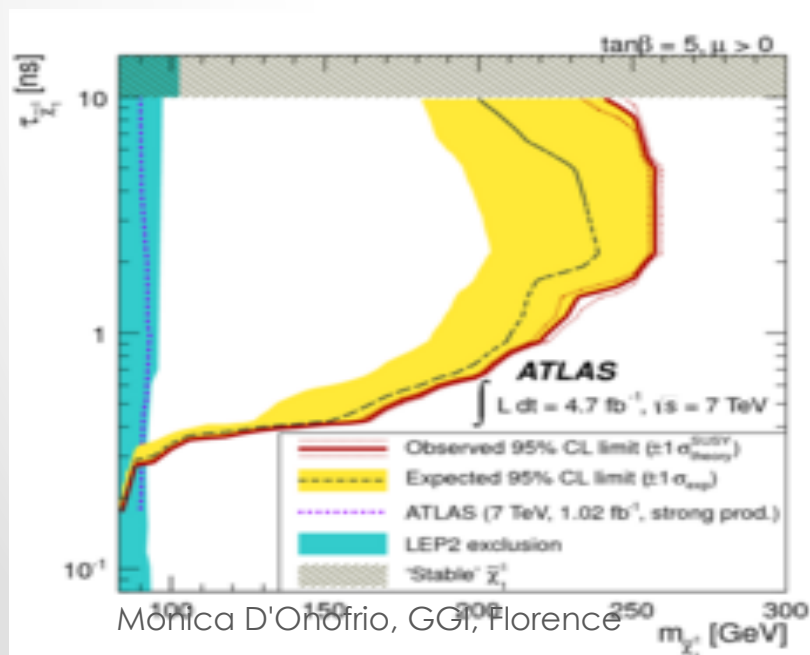
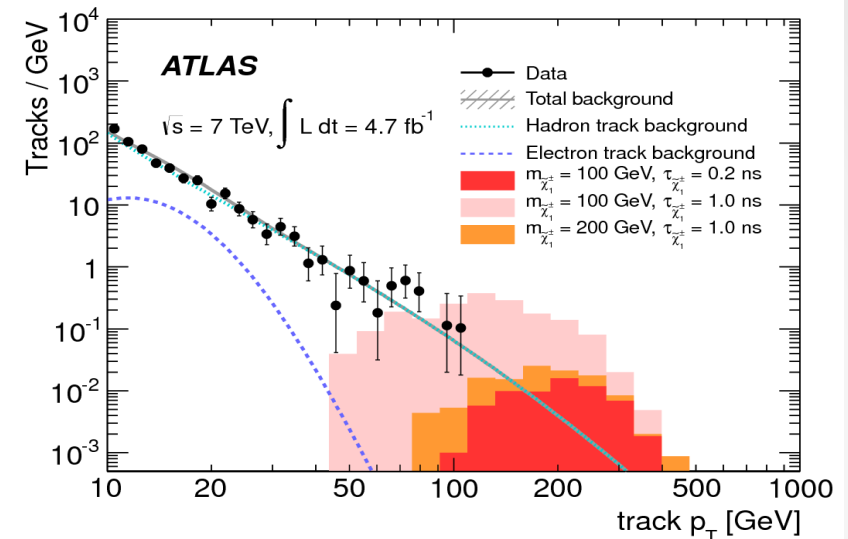


Electrons surviving lepton veto,
undergoing bremsstrahlung

Obtain p_T spectrum for both sources of
background using data control samples

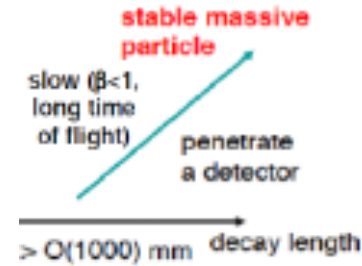
Disappearing tracks - results

- Use signal+background likelihood fit to track p_T spectrum, to test different signal hypotheses
- No significant excess found \rightarrow Set limits on chargino mass and lifetime and chargino-neutral ΔM :

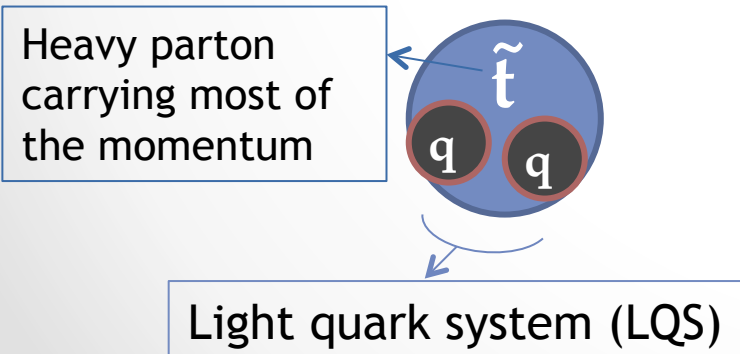


Stable Massive Particles (SMPs)

- SMPs predicted in SUSY several other BSM scenarios
- Stable $\rightarrow c\tau \geq$ size of detector
- Produced with $\beta < 1$
- Within SUSY:
 - \tilde{t} and $\tilde{\chi}^+$
 - \tilde{q}/\tilde{g} (bound states)



Colored sparticles can hadronise into long-lived bound hadronic states
R-hadrons



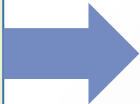
SMP	LSP	Scenario	Conditions
$\tilde{\tau}_1$	$\tilde{\chi}_1^0$	MSSM	$\tilde{\tau}_1$ mass (determined by $m_{\tilde{L},R}^2, \mu, \tan \beta$, and A_τ) close to $\tilde{\chi}_1^0$ mass.
	\tilde{G}	GMSB	Large N , small M , and/or large $\tan \beta$.
	\tilde{g}	\tilde{g} MSSB	No detailed phenomenology studies, see [23].
	\tilde{g}	SUGRA	Supergravity with a gravitino LSP, see [24].
$\tilde{\tau}_1$	$\tilde{\tau}_1$	MSSM	Small $m_{\tilde{L},R}$ and/or large $\tan \beta$ and/or very large A_τ .
	$\tilde{\tau}_1$	AMSB	Small m_0 , large $\tan \beta$.
	\tilde{g}	\tilde{g} MSSB	Generic in minimal models.
\tilde{e}_1	\tilde{G}	GMSB	$\tilde{\tau}_1$ NLSP (see above). \tilde{e}_1 and $\tilde{\mu}_1$ co-NLSP and also SMP for small $\tan \beta$ and μ .
	$\tilde{\tau}_1$	\tilde{g} MSSB	\tilde{e}_1 and $\tilde{\mu}_1$ co-LSP and also SMP when stau mixing small.
$\tilde{\chi}_1^+$	$\tilde{\chi}_1^0$	MSSM	$m_{\tilde{\chi}_1^+} - m_{\tilde{\chi}_1^0} \lesssim m_{\pi^+}$. Very large $M_{1,2} \gtrsim 2 \text{ TeV} \gg \mu $ (Higgsino region) or non-universal gaugino masses $M_1 \gtrsim 4M_2$, with the latter condition relaxed to $M_1 \gtrsim M_2$ for $M_2 \ll \mu $. Natural in O-II models, where simultaneously also the \tilde{g} can be long-lived near $\delta_{GS} = -3$.
	$\tilde{\chi}_1^0$	AMSB	$M_1 > M_2$ natural. m_0 not too small. See MSSM above.
\tilde{g}	$\tilde{\chi}_1^0$	MSSM	Very large $m_{\tilde{g}}^2 \gg M_3$, e.g. split SUSY.
	\tilde{G}	GMSB	SUSY GUT extensions [25–27].
\tilde{g}	\tilde{g}	MSSM	Very small $M_3 \ll M_{1,2}$, O-II models near $\delta_{GS} = -3$.
	\tilde{G}	GMSB	SUSY GUT extensions [25–29].
\tilde{t}_1	$\tilde{\chi}_1^0$	MSSM	Non-universal squark and gaugino masses. Small $m_{\tilde{q}}^2$ and M_3 , small $\tan \beta$, large A_t .
\tilde{b}_1	$\tilde{\chi}_1^0$	MSSM	Small $m_{\tilde{q}}^2$ and M_3 , large $\tan \beta$ and/or large $A_b \gg A_t$.

Table 1: Brief overview of possible SUSY SMP states considered in the literature. Classified by SMP, LSP, scenario, and typical conditions for this case to materialise in the given scenario.

SMPs: how-to

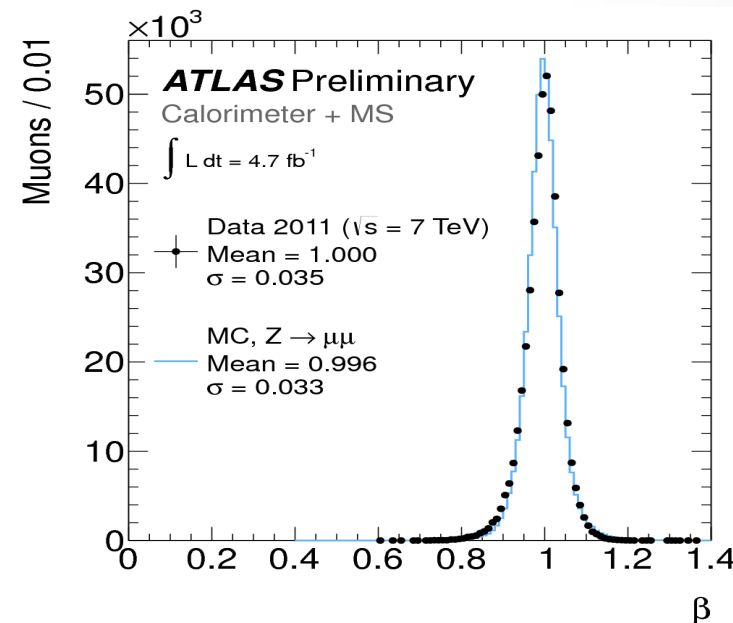
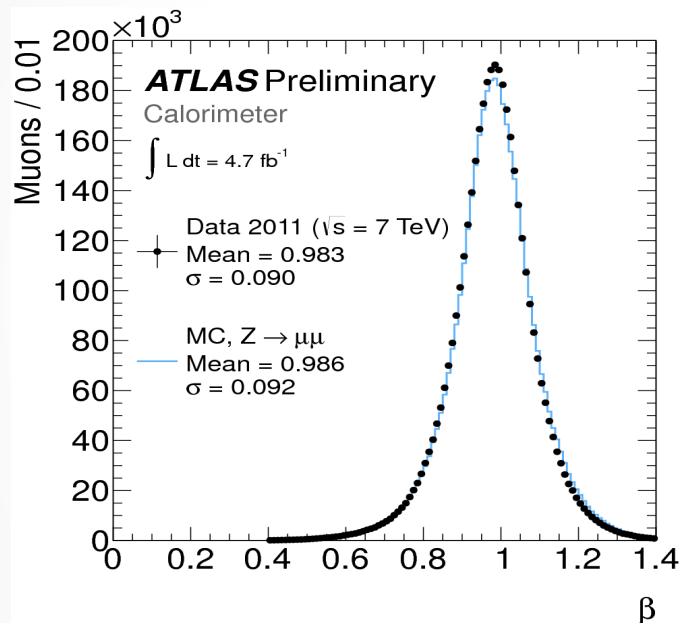
- Particles with $c\tau \geq$ size of detector:
 - If neutral and weakly interacting \rightarrow missing E_T
 - If charged (at any point!) or strongly interacting \rightarrow can detect them directly
 - If massive, could be produced with low velocities: $\beta < 1$
- Mass of SMPs:

- measure charged particle momentum p in ID and MS
- measure energy loss dE/dx in several subdetectors
 - For these analyses, use Pixel
 - dE/dx is related to relativistic boost factor $\beta\gamma$
- measure time-of-flight in several sub-detectors
 - For these analyses, use Tile, RPC, MDT
 - Can measure velocity β


$$p = \beta\gamma m$$

β measurements and bkg

- Use $Z \rightarrow \mu\mu$ events to calibrate β measurements
- If β measurements from different systems are > 0.2 and internally consistent \rightarrow combined in a weighted average.

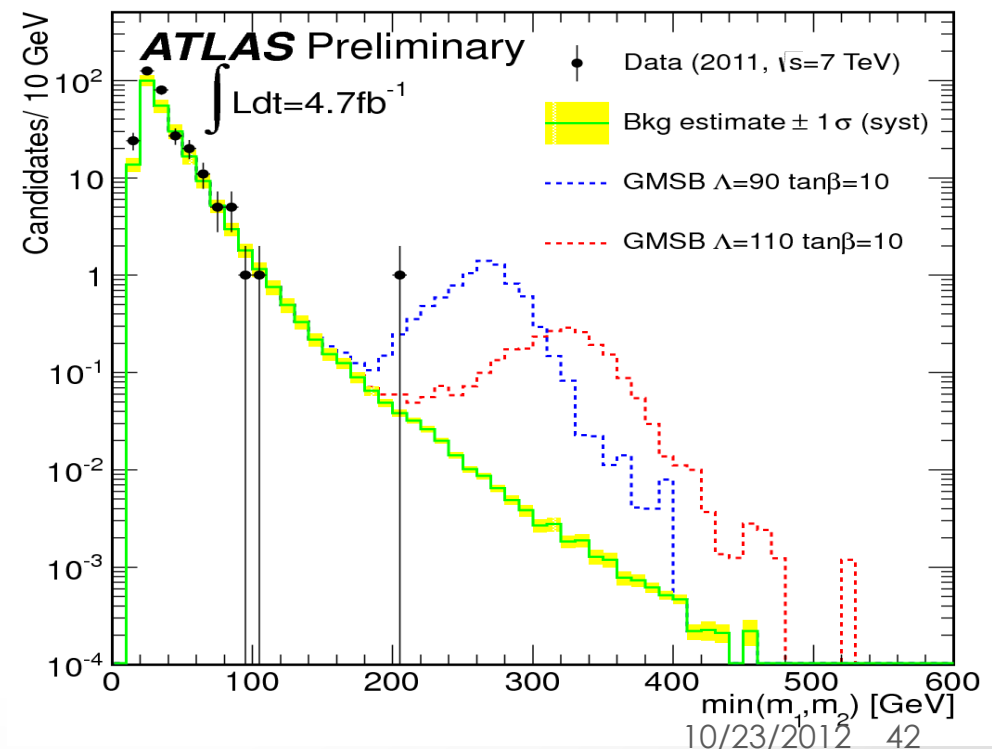


Main background are high- p_T muons with mismeasured β .

- Exploit fact that mis-measurements of β or β_y in different subdetectors are uncorrelated.
- Use data-driven method based on randomly sampling β or β_y values from CR distributions.

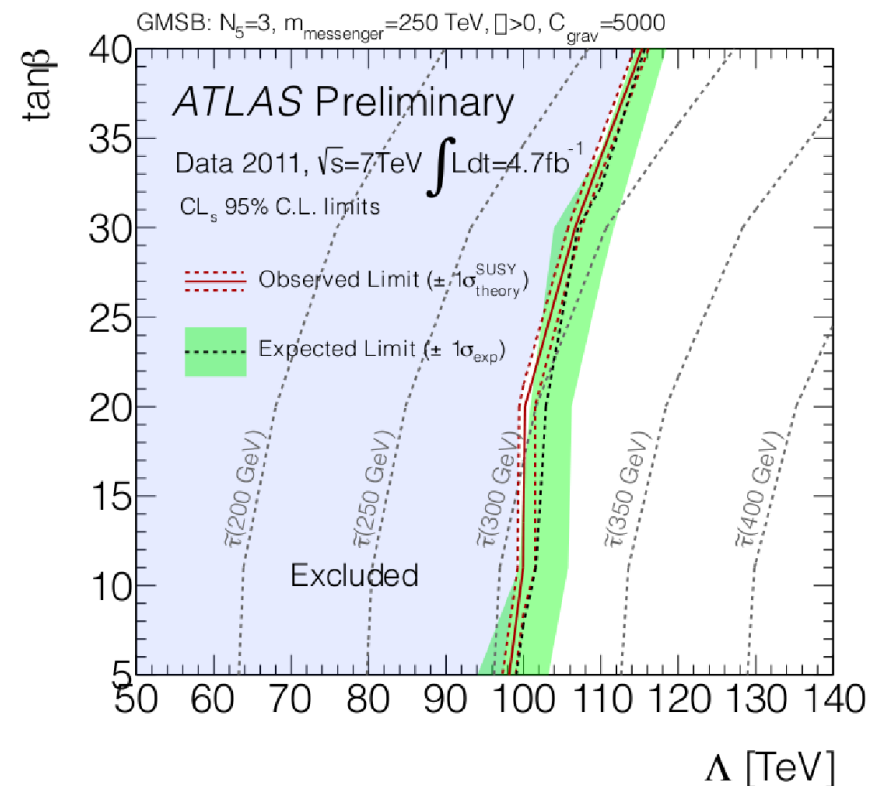
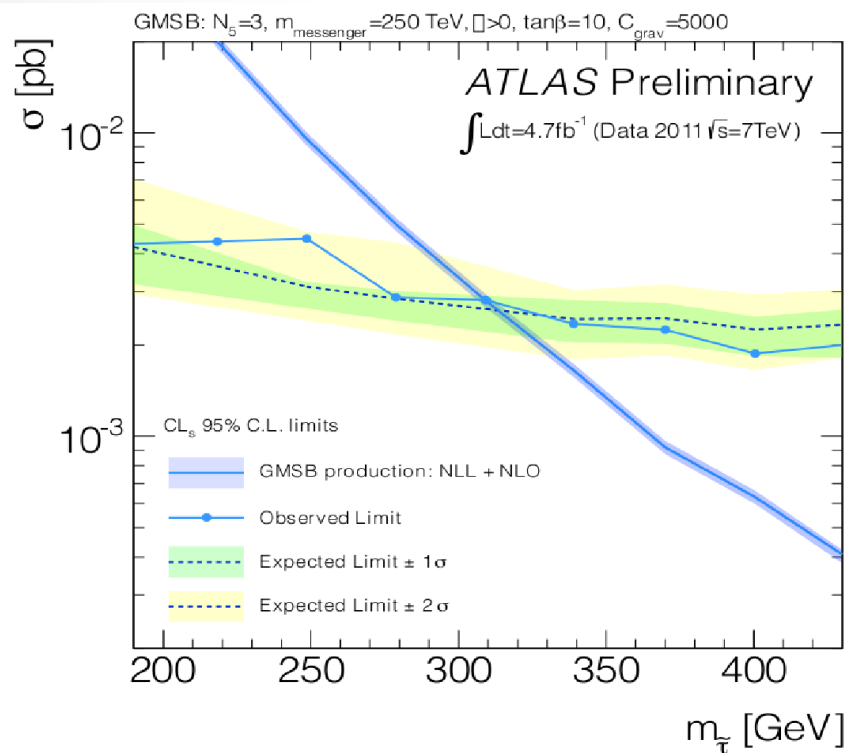
Long-lived sleptons - selection

- “Heavy muons”-like, releasing energy throughout detector
 - Use single μ trigger (70-85% efficiency for GMSB slepton events):
 - Efficiency for particles arriving late at MS estimated from MC signal
 - 2 μ candidates per event, $p_{T \text{ loose(tight)}} > 50(70)$ GeV.
-
- Either two loose candidates or one must pass tight selection.
 - Bkg estimate: from iterative procedure based on muon- β PDF:
 - Assume bkg is due to measurement resolution
 - Apply cut on candidate mass $m_\beta = p/\gamma\beta$ (depending on stau mass hypothesis: 122 - 465 GeV range)



Slepton search - results

- No excess above background expectation
- Constraints on stau mass in GMSB scenario
- Constraints in $\tan\beta$ vs Λ in GMSB scenario
 - Direct sleptons production dominate around limit

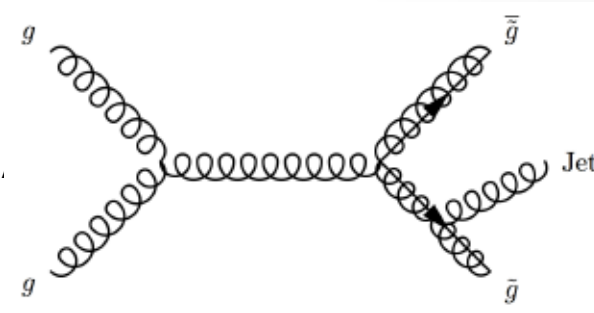


R-hadrons - selection

- Undergo interactions with detector material
 - can also change charge as it moves through detector
- If β is too low, particle might be associated with following bunch crossing by the time it gets to MS.
- Due to both these effects, efficiency for single muon trigger can be quite low.

→ use also missing E_T trigger

- Exploit possible additional jets from ISR.



- **Three different analyses:**

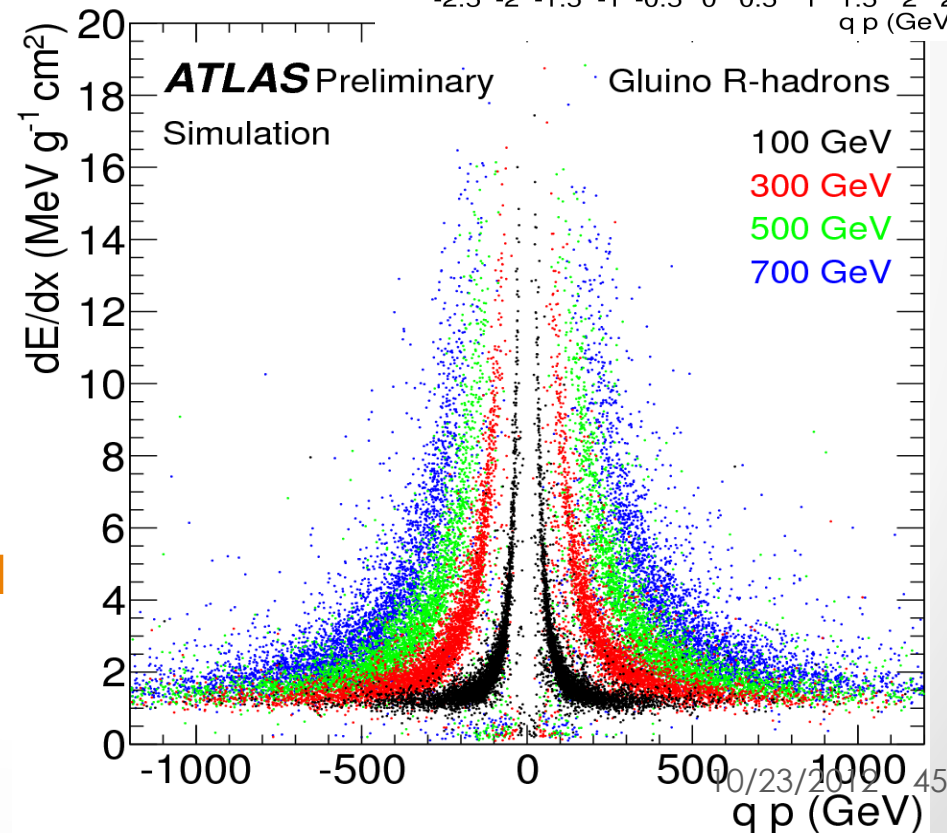
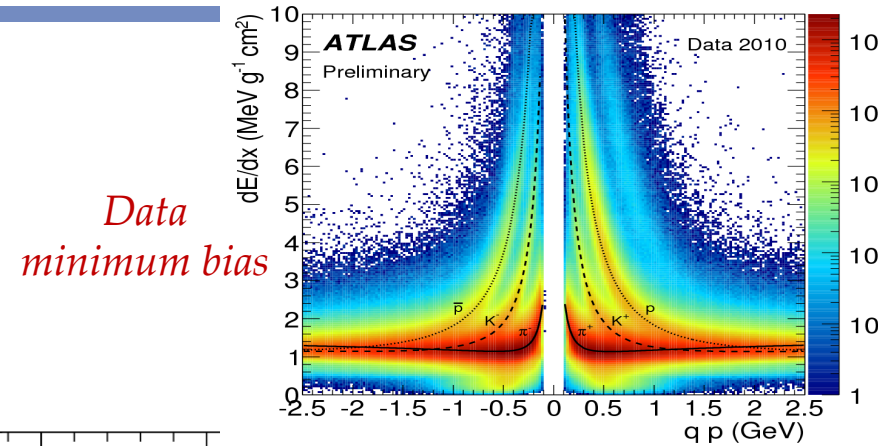
- “Full detector”
- “ID only”
- “MS agnostic”



Cover the lack of knowledge of R-hadron interactions with the detector and the lifetimes for which they would not reach the calorimeters

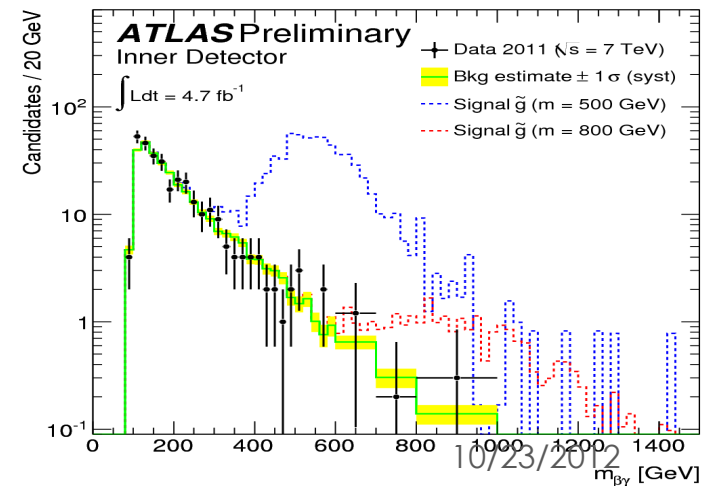
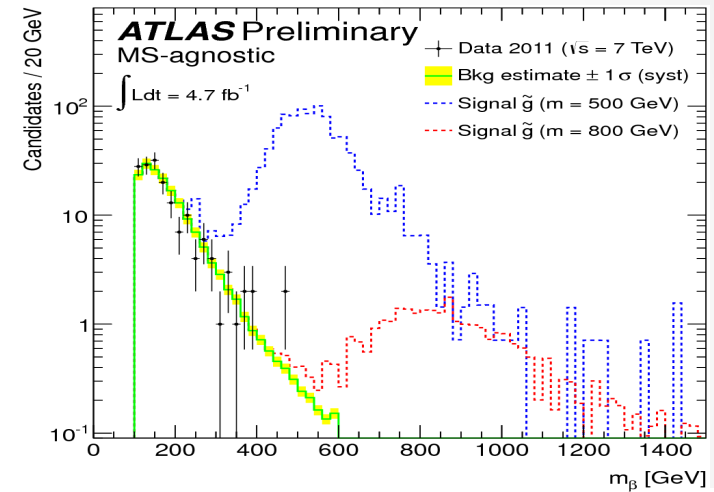
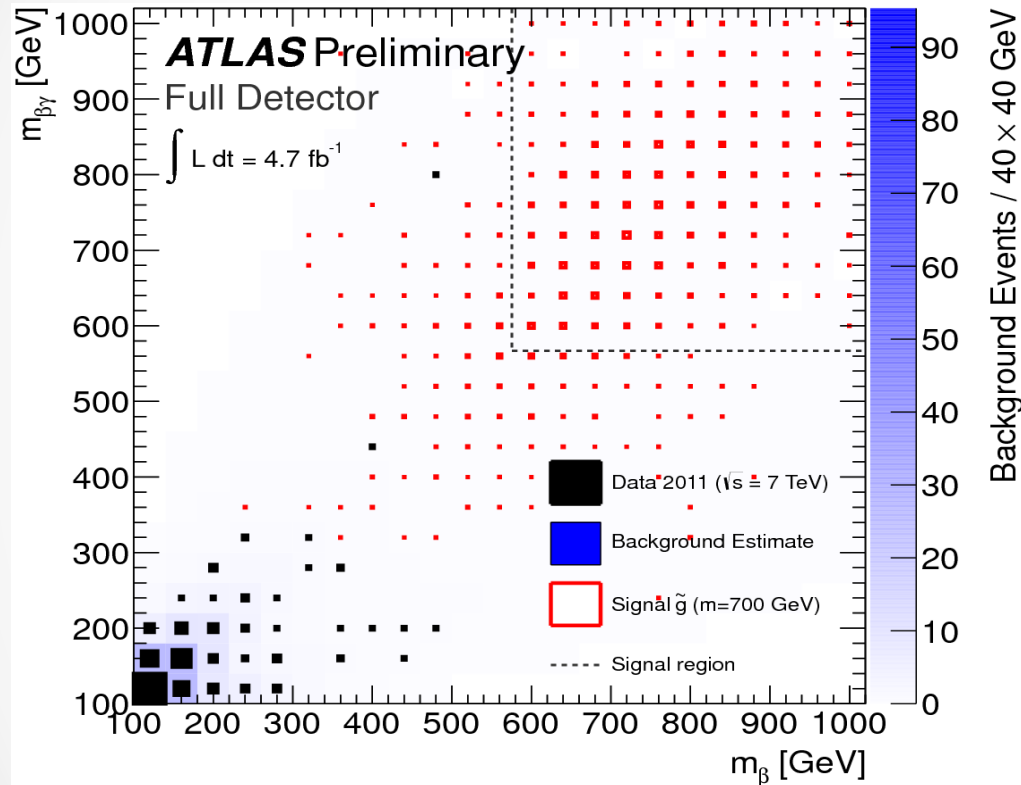
R-hadrons - selection (2)

- All three analyses require good quality, isolated, high-momentum ID track.
- “MS agnostic”:
 - missing E_T triggers
 - calorimeter-only timing measurement
 - require ID trk $p_T > 140$ GeV
- “ID only”:
 - Offline missing E_T cut
 - Tighter cuts on isolation and number of silicon hits



R-hadron searches - results

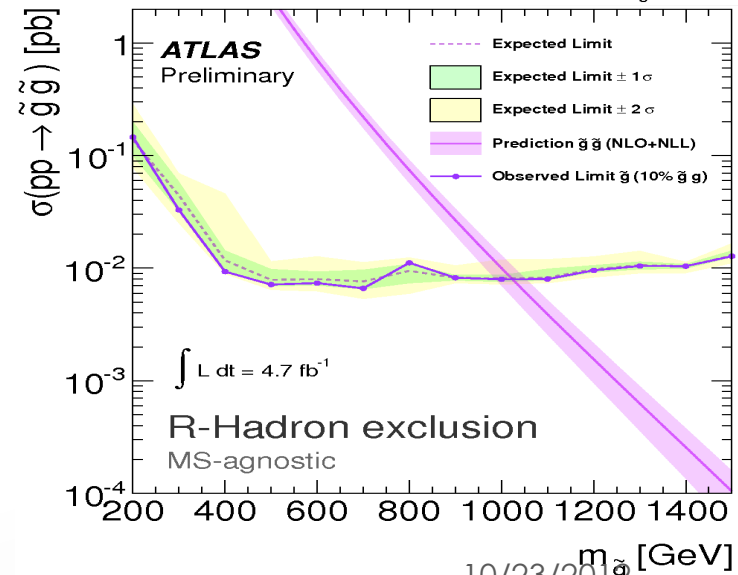
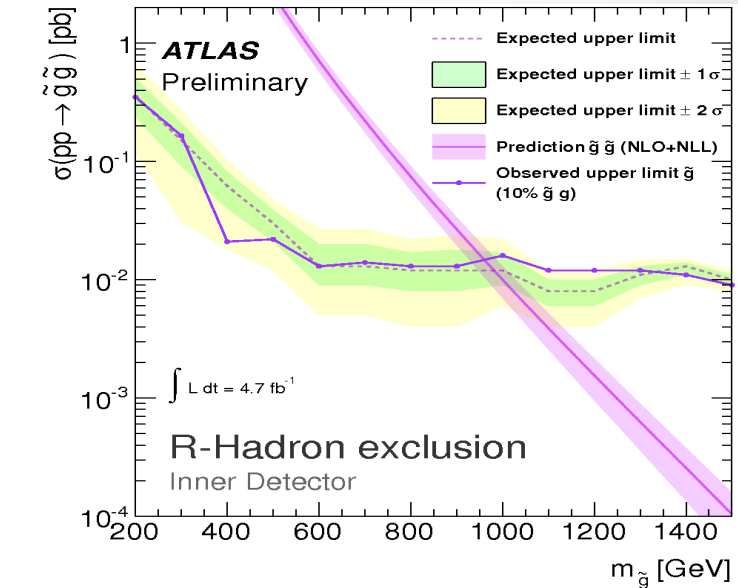
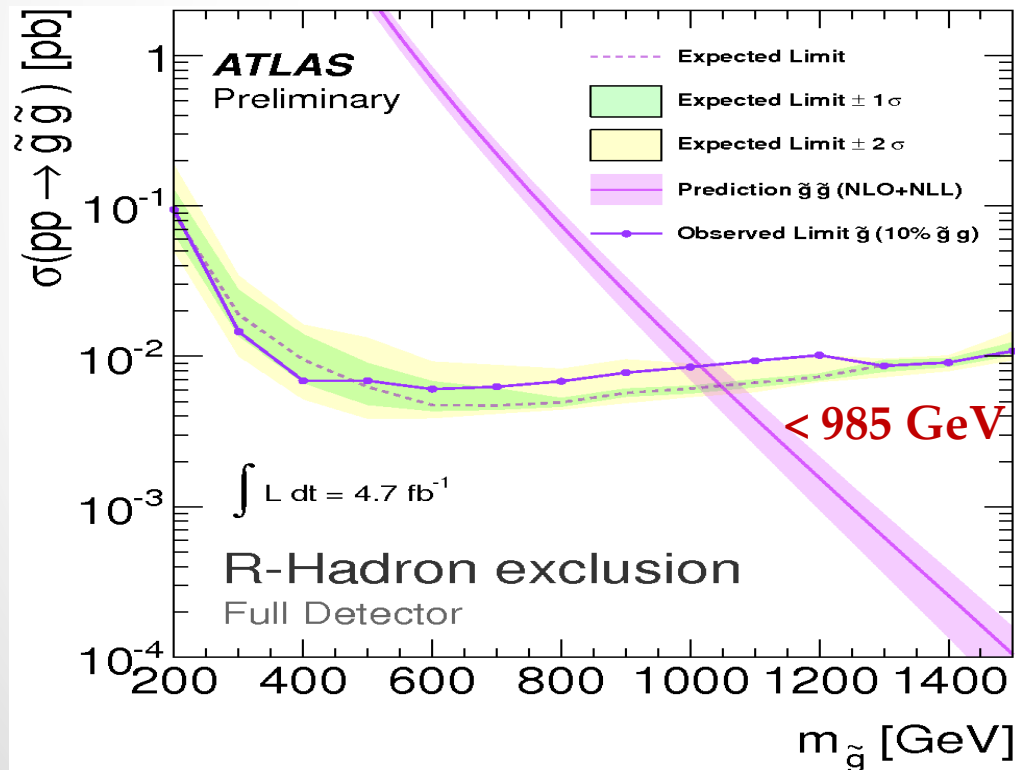
- MS-agnostic and FD: final selection based on requirements on $\beta\gamma$ and β



- ID-only: selection based on exceeding dE/dx thresholds

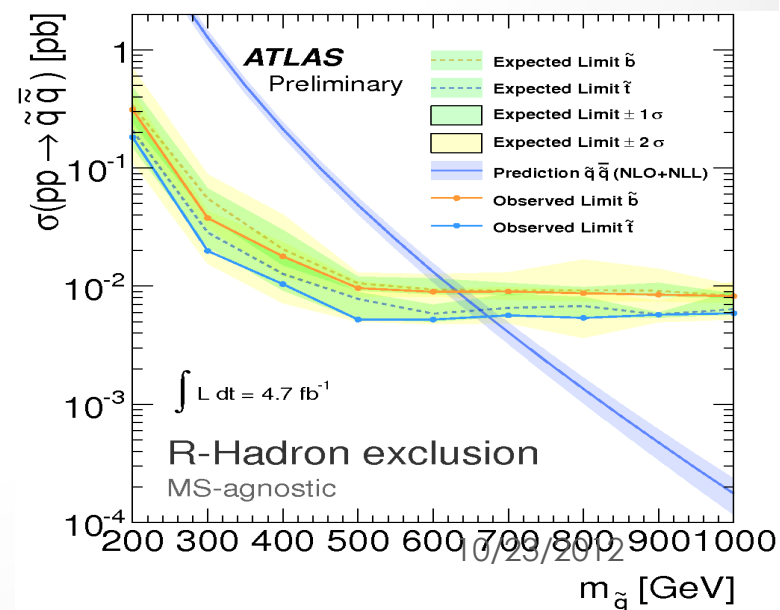
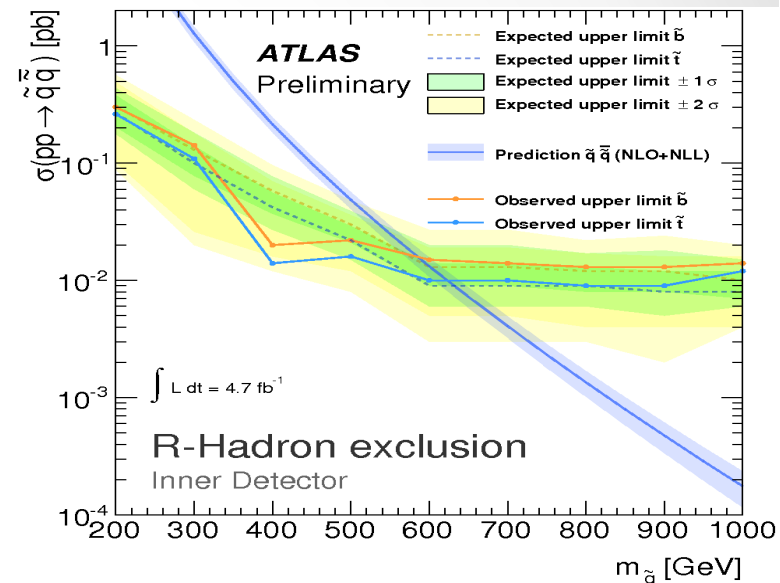
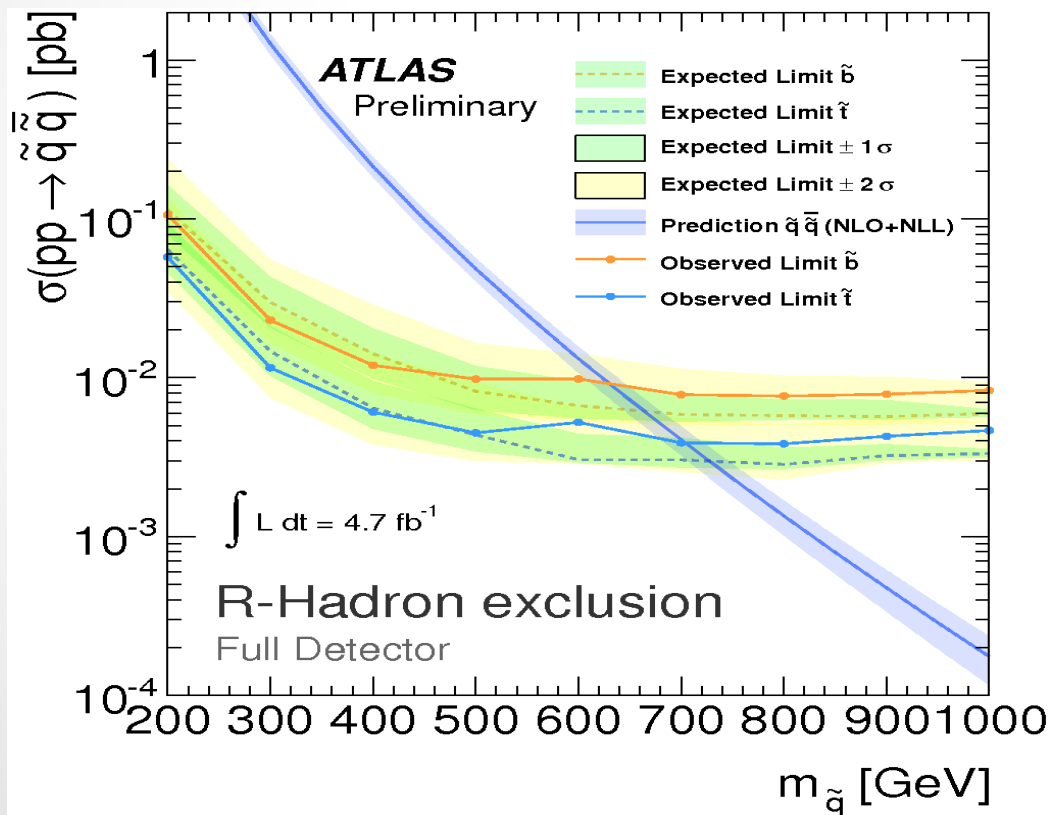
R-hadron searches - results

- No excess above background expectation seen in any of the three analyses
- Limits set by counting N events passing model-dependent mass cuts
- **gluino R-hadrons:**



R-hadron searches - results

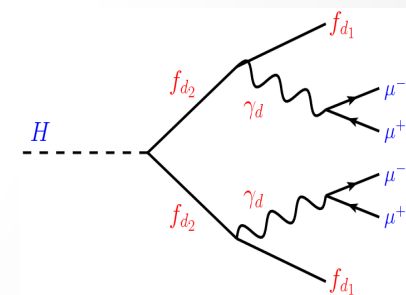
- squark (stop & sbottom) R-hadrons:



Other searches for LLP particles

Several New Physics models (well beyond SUSY) could give rise to new massive particles, with long lifetimes:

- Hidden sector interacts with SM via (heavy) Communicator particle(s)
 - **light Higgs-to-LLP search:**
 - look for displaced vertices at larger radii, near outer radius of hadronic calorimeter, or in the MS
 - Use specially developed trigger algorithm, and specialized tracking and vertexing, to reconstruct vertices in MS
 - **Higgs decay to hidden-sector fermions** → decay to a (potentially long-lived) neutral hidden-sector particle Υ_d and a stable hidden sector fermion
 - those escape detection
 - Decay of Υ_d could give rise to collimated pairs of leptons
- **Magnetic monopoles:**
 - Experimental signature → large, localized energy deposit in EM calorimeter, associated with region of high ionization in TRT.



Conclusions

- Several SUSY models lead to hard to constrain or ‘unconventional’ signatures:
 - RPV multilepton or multijet scenarios
 - Long-lived particles
- Wide range of analyses, looking for many different signatures, sometimes using the detector in interesting and “non-standard” ways.
- No sign of New Physics so far....
- **BUT:**
 - All these analyses are being updated (and improved) with 2012 data.
 - More to come

Doing our best to cover as much parameter space as possible, as well as and get maximum possible value out of our fantastic ATLAS detector