

Mass hierarchy and electroweak symmetry breaking

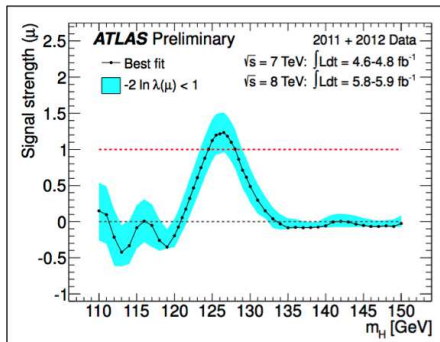
I. Antoniadis



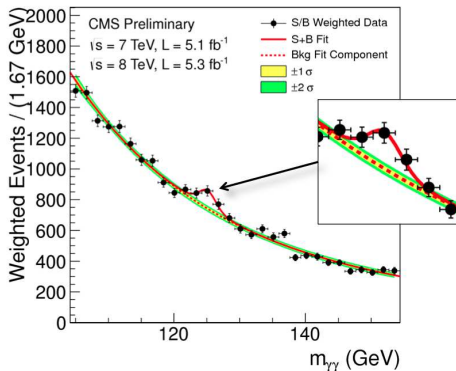
Beyond the Standard Model after the first run of LHC
GGI, Florence, 9-12 July 2013

- Low energy SUSY and 125 GeV Higgs
- Live with the hierarchy
- Extra $U(1)$'s
- Low scale strings and extra dimensions

Higgs boson discovery



$$m_H = 125.5 \pm 0.2 \text{ (stat.)} \pm 0.5 \text{ (syst.)}$$



$$m_H = 125.7 \pm 0.3 \pm 0.3 \text{ GeV}$$

Beyond the Standard Model of Particle Physics: driven by the mass hierarchy problem

Standard picture: low energy supersymmetry

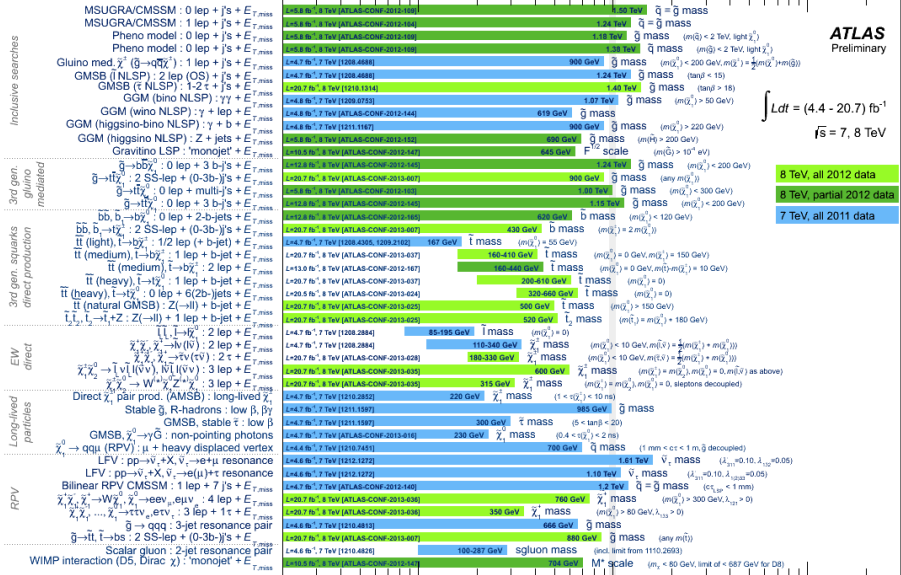
Natural framework: Heterotic string (or high-scale M/F) theory

Advantages:

- natural elementary scalars
- gauge coupling unification
- LSP: natural dark matter candidate
- radiative EWSB

Problems:

- too many parameters: soft breaking terms
- MSSM : already a % - ‰ fine-tuning 'little' hierarchy problem



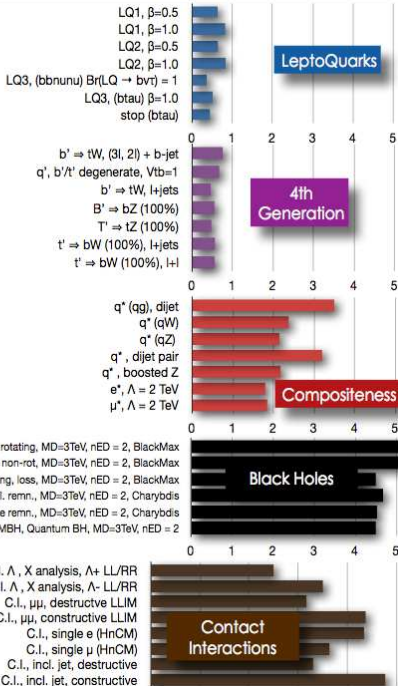
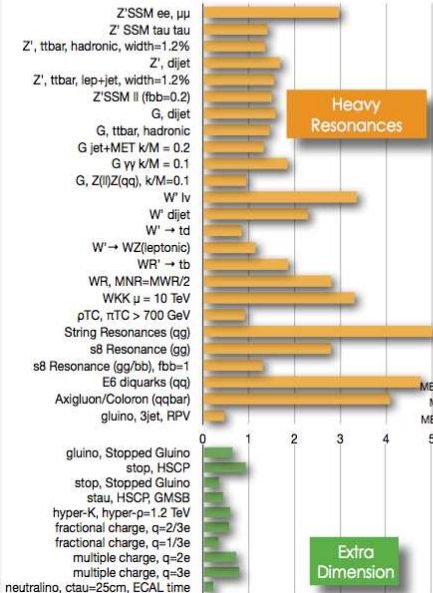
$\int L dt = (4.4 - 20.7) \text{ fb}^{-1}$
 $\sqrt{s} = 7, 8 \text{ TeV}$

8 TeV, all 2012 data
 8 TeV, partial 2012 data
 7 TeV, all 2011 data

*Only a selection of the available mass limits on new states or phenomena shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty

CMS EXOTICA

95% CL EXCLUSION LIMITS (TeV)

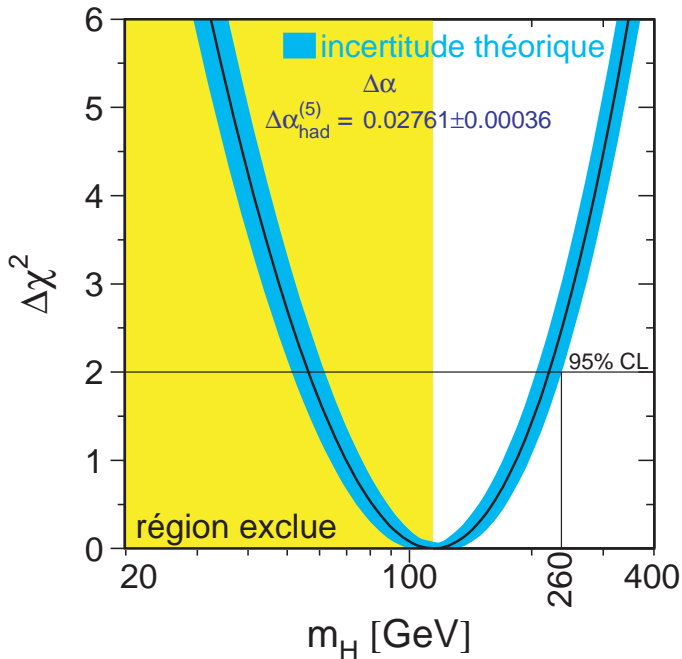


Remarks on the value of the Higgs mass ~ 125 GeV

- consistent with expectation from precision tests of the SM
- favors perturbative physics quartic coupling $\lambda = m_H^2/v^2 \simeq 1/8$

Window to new physics

- compatible with supersymmetry
but appears fine-tuned in its minimal version [8]
early to draw a general conclusion before LHC13/14
e.g. an extra singlet or split families can alleviate the fine tuning [9]
- very important to measure its properties and couplings [13]
any deviation of its couplings to top, bottom and EW gauge bosons
implies new light states involved in the EWSB altering the fine-tuning



Fine-tuning in MSSM

Upper bound on the lightest scalar mass:

$$m_h^2 \lesssim m_Z^2 \cos^2 2\beta + \frac{3}{(4\pi)^2} \frac{m_t^4}{v^2} \left[\ln \frac{m_{\tilde{t}}^2}{m_t^2} + \frac{A_t^2}{m_{\tilde{t}}^2} \left(1 - \frac{A_t^2}{12m_{\tilde{t}}^2} \right) \right] \lesssim (130 \text{ GeV})^2$$

$$m_h \simeq 126 \text{ GeV} \Rightarrow m_{\tilde{t}} \simeq 3 \text{ TeV or } A_t \simeq 3m_{\tilde{t}} \simeq 1.5 \text{ TeV}$$

\Rightarrow % to a few ‰ fine-tuning

$$\text{minimum of the potential: } m_Z^2 = 2 \frac{m_1^1 - m_2^2 \tan^2 \beta}{\tan^2 \beta - 1} \sim -2m_2^2 + \dots$$

$$\text{RG evolution: } m_2^2 = m_2^2(M_{\text{GUT}}) - \frac{3\lambda_t^2}{4\pi^2} m_{\tilde{t}}^2 \ln \frac{M_{\text{GUT}}}{m_{\tilde{t}}} + \dots \quad [25]$$

$$\sim m_2^2(M_{\text{GUT}}) - \mathcal{O}(1)m_{\tilde{t}}^2 + \dots \quad [6]$$

MSSM with dim-5 and 6 operators

I.A.-Dudas-Ghilencea-Tziveloglou '08, '09, '10

parametrize new physics above MSSM by higher-dim effective operators

relevant super potential operators of dimension-5:

$$\mathcal{L}^{(5)} = \frac{1}{M} \int d^2\theta (\eta_1 + \eta_2 S) (H_1 H_2)^2$$

η_1 : generated for instance by a singlet

$$W = \lambda \sigma H_1 H_2 + M \sigma^2 \quad \rightarrow \quad W_{\text{eff}} = \frac{\lambda^2}{M} (H_1 H_2)^2$$

Strumia '99 ; Brignole-Casas-Espinosa-Navarro '03

Dine-Seiberg-Thomas '07

η_1 : corresponding soft breaking term spurion $S \equiv m_S \theta^2$

Physical consequences of MSSM₅: Scalar potential

$$\mathcal{V} = m_1^2 |h_1|^2 + m_2^2 |h_2|^2 + B\mu(h_1 h_2 + \text{h.c.}) + \frac{g_2^2 + g_Y^2}{8} (|h_1|^2 - |h_2|^2)^2 \\ + (|h_1|^2 + |h_2|^2) (\eta_1 h_1 h_2 + \text{h.c.}) + \frac{1}{2} [\eta_2 (h_1 h_2)^2 + \text{h.c.}] + \mathcal{O}(\eta_i^2)$$

- $\eta_{1,2} \Rightarrow$ quartic terms along the D-flat direction $|h_1| = |h_2|$
- potential stability $\Rightarrow \eta_2 \geq 4|\eta_1|$

requiring η -corrections to be smaller than MSSM mass matrix elements \Rightarrow

only η_2 can change the tree-level bound $m_h \leq m_Z$ but marginally

Relevance of dim-6 operators

Relaxing the condition on potential positivity: guaranteed by dim-6 ops

only one dim-6 along the D-flat direction induced by dim-5: $\propto \eta_1^2$

$$W = \eta_1 (H_1 H_2)^2 \longrightarrow V = \left| \frac{\partial W}{\partial H_i} \right|^2 \sim \eta_1^2 |H_1 H_2|^2 (|H_1|^2 + |H_2|^2)$$

- tree-level mass can increase significantly
- bigger parameter space for LSP being dark matter

Bernal-Blum-Nir-Losada '09

MSSM Higgs with dim-6 operators

dim-6 operators can have an independent scale from dim-5

Classification of all dim-6 contributing to the scalar potential

(without SUSY) \Rightarrow

large $\tan \beta$ expansion: $\delta_6 m_h^2 = f v^2 + \dots$

constant receiving contributions from several operators

$f \sim f_0 \times (\mu^2/M^2, m_S^2/M^2, \mu m_S/M^2, v^2/M^2)$

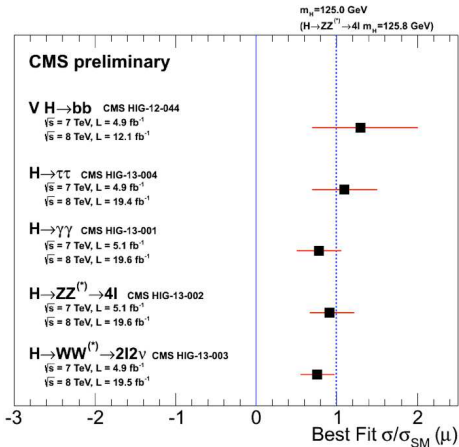
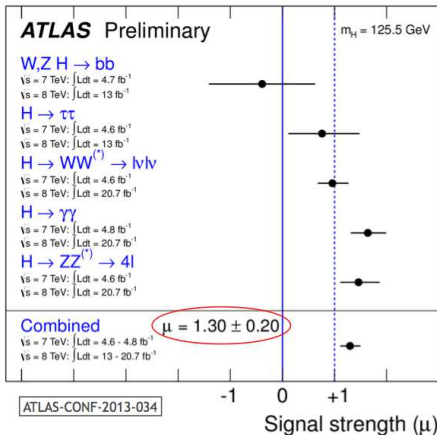
$m_S = 1 \text{ TeV}, M = 10 \text{ TeV}, f_0 \sim 1 - 2.5$ for each operator

$\Rightarrow m_h \simeq 103 - 119 \text{ GeV}$

\Rightarrow MSSM with dim-5 and dim-6 operators:

possible resolution of the MSSM fine-tuning problem [6]

Couplings of the new boson vs SM

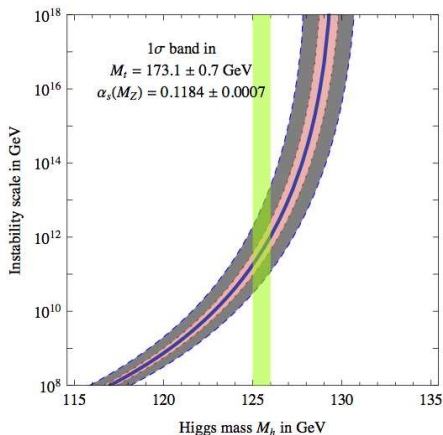
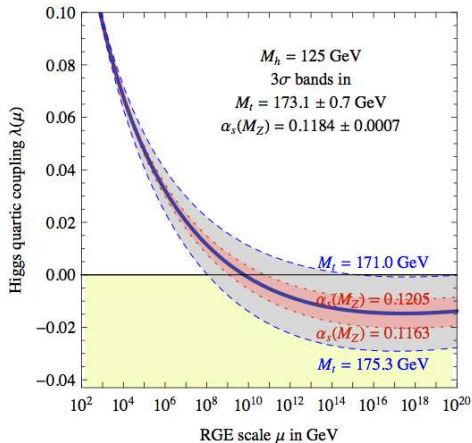


exclusion : spin 2 and pseudoscalar at $\gtrsim 95\%$ CL

Agreement with Standard Model expectation at $\sim 2\sigma$

Can the SM be valid at high energies?

Degrassi-Di Vita-Elias Miró-Espinosa-Giudice-Isidori-Strumia '12



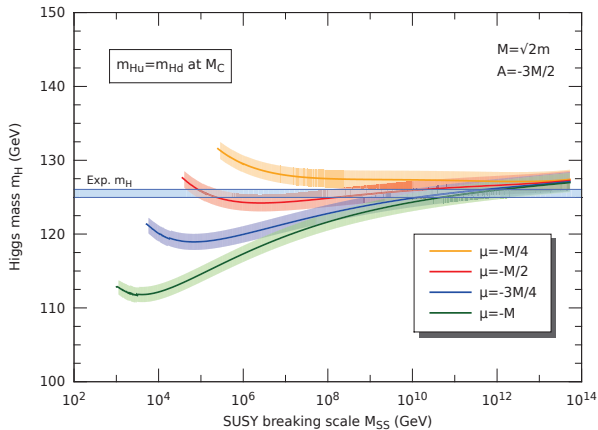
Instability of the SM Higgs potential \Rightarrow metastability of the EW vacuum

SUSY : $\lambda = 0 \Rightarrow \sin \beta = 1$

$$H_{SM} = \sin \beta H_u - \cos \beta H_d^* \quad \lambda = \frac{1}{8}(g_2^2 + g'^2) \cos^2 2\beta$$

$\lambda = 0$ at a scale $\geq 10^{10}$ GeV $\Rightarrow m_H = 126 \pm 3$ GeV

Ibanez-Valenzuela '13



e.g. for universal $\sqrt{2}m = M = M_{SS}$, $A = -3/2M$

If the weak scale is tuned \Rightarrow split supersymmetry is a possibility

Arkani Hamed-Dimopoulos '04, Giudice-Romanino '04

- natural splitting: gauginos, higgsinos carry R-symmetry, scalars do not
- main good properties of SUSY are maintained
 - gauge coupling unification and dark matter candidate
- also no dangerous FCNC, CP violation, ...
- experimentally allowed Higgs mass \Rightarrow 'mini' split

$m_S \sim \text{few} - \text{thousands TeV}$

gauginos: a loop factor lighter than scalars ($\sim m_{3/2}$)

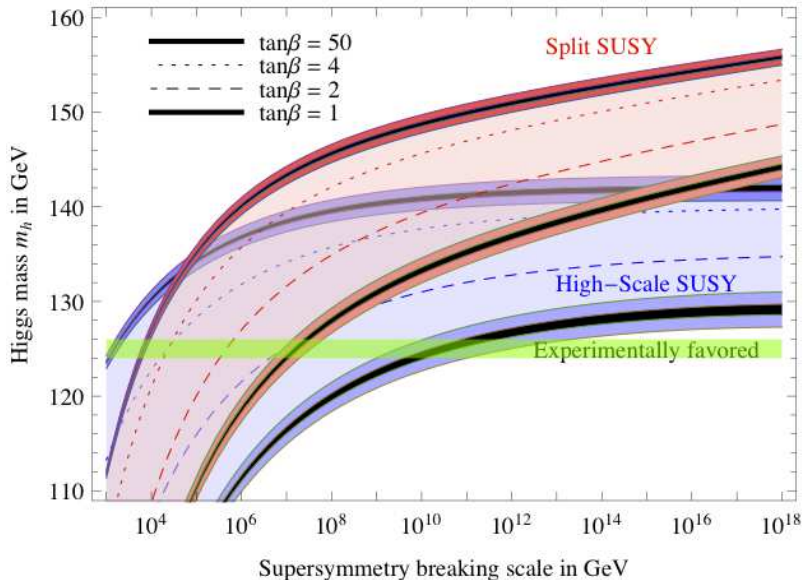
- natural string framework: intersecting (or magnetized) branes

IA-Dimopoulos '04

D-brane stacks are supersymmetric with massless gauginos

intersections have chiral fermions with broken SUSY & massive scalars

Predicted range for the Higgs mass



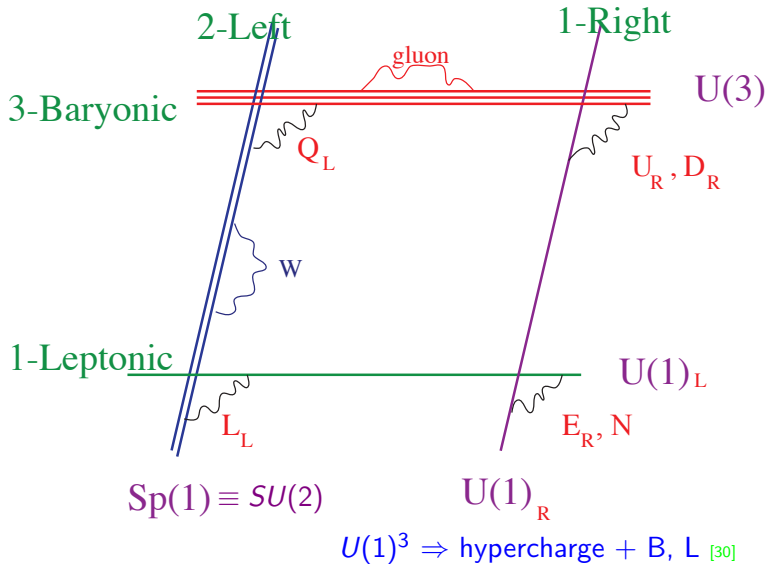
An extra $U(1)$ can also cure the instability problem

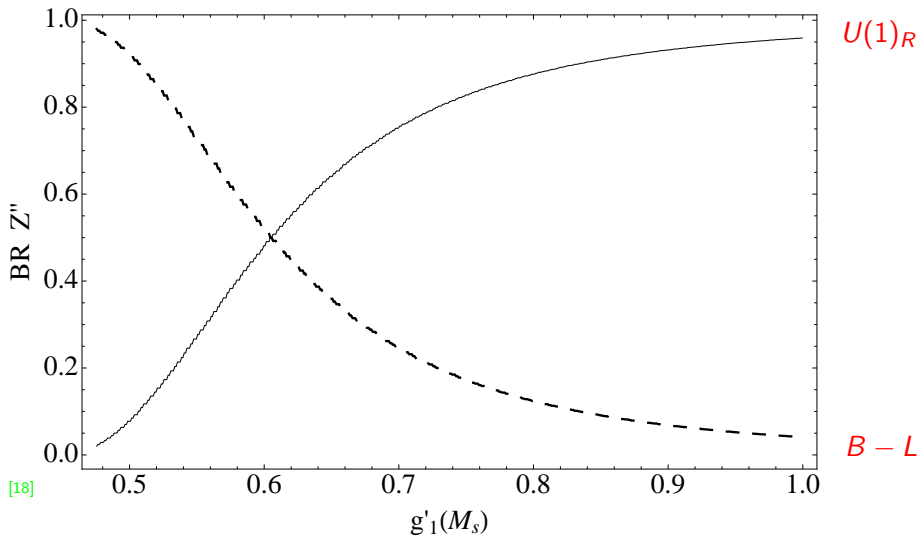
Anchordoqui-IA-Goldberg-Huang-Lüst-Taylor-Vlcek '12

usually associated to known global symmetries of the SM: B, L, \dots

- B anomalous and superheavy
- $B - L$ massless at the string scale (no associated 6d anomaly)
but broken at TeV by a scalar VEV with the quantum numbers of N_R
- L -violation from higher-dim operators suppressed by the string scale
- $U(3)$ unification, Y combination \Rightarrow 2 parameters: 1 coupling + m_Z
- perturbativity $\Rightarrow 0.5 \lesssim g_{U(1)_R} \lesssim 1$ [20]
- interesting LHC phenomenology and cosmology [21]

Standard Model on D-branes : SM⁺⁺





[18]

- Rotation of $U(1)$'s from the string to low energy basis Z, Z', Z'' :
completely fixed in terms of the couplings
 - Decoupling of anomalous $Z' \simeq B$
 - Z'' linear combination of $B - L$ and $U(1)_R$
- Recent cosmological observations indicate extra relativistic component
dark radiation parametrized by an effective ν -number close to 4 *
→ use the 3 ν_R 's interacting with SM fermions via Z''
data: their decoupling during the quark-hadron transition
⇒ $3.5 \lesssim M_{Z''} \lesssim 7 \text{ TeV}$ (within LHC14 discovery potential)

* before Planck results

Planck XVI

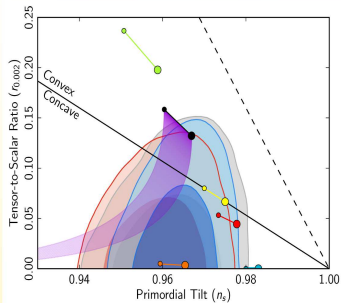
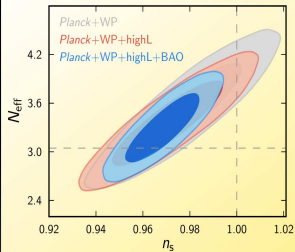
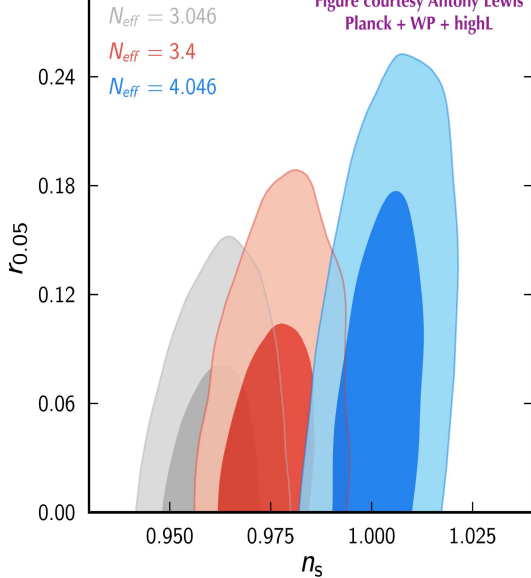


Fig. 1. Marginalized joint 68% and 95% CL regions for n_s and $r_{0.002}$ from *Planck* in combination with other data sets compared to the theoretical predictions of selected inflationary models.

Figure courtesy Antony Lewis Planck + WP + highL



Scalar potential:

$$V(H, H'') = \mu^2 |H|^2 + \mu'^2 |H''|^2 + \lambda_1 |H|^4 + \lambda_2 |H''|^4 + \lambda_3 |H|^2 |H''|^2$$

5 parameters $\Rightarrow v, m_h, v'', m_{h''}$ + a scalar mixing angle α

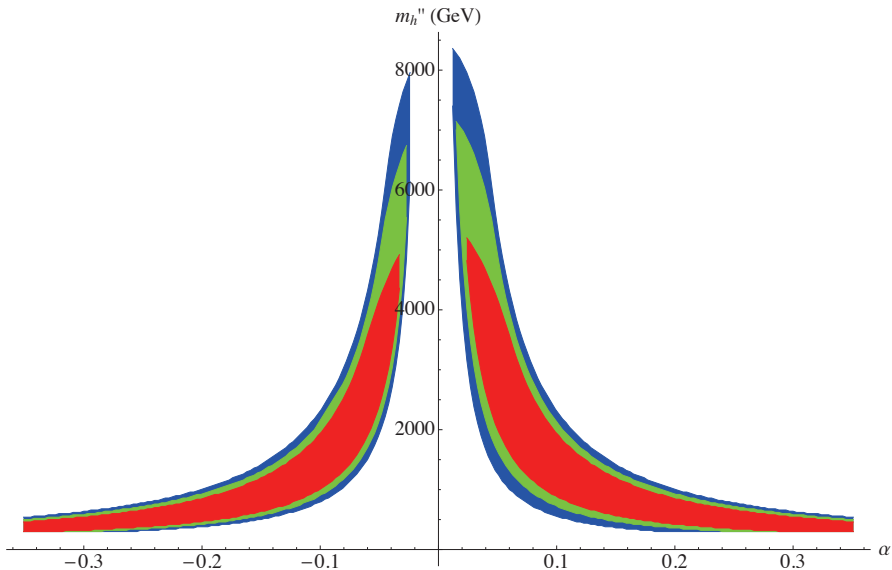
\Rightarrow 3 free parameters : $m_{h''}, \alpha, v'' \leftrightarrow M_{Z''}$

Stability conditions: $\lambda_1 > 0, \quad \lambda_2 > 0, \quad \lambda_1 \lambda_2 > \frac{1}{4} \lambda_3^2$

RGE analysis up to $M_s \Rightarrow$ stability is possible in SM⁺⁺

for $0.02 \lesssim |\alpha| \lesssim 0.35$ and $500 \text{ GeV} \lesssim m_{h''} \lesssim 5 \text{ TeV}$

$$M_{Z''} = 4.5 \text{ TeV}; \quad M_S = 10^{14}, 10^{16}, 10^{19} \text{ GeV}$$



Alternative answer: Low UV cutoff $\Lambda \sim \text{TeV}$

- low scale gravity \Rightarrow extra dimensions: large flat or warped
- low string scale \Rightarrow low scale gravity, ultra weak string coupling

$M_s \sim 1 \text{ TeV} \Rightarrow$ volume $R_{\perp}^n = 10^{32} l_s^n$ ($R_{\perp} \sim .1 - 10^{-13} \text{ mm}$ for $n = 2 - 6$)

- spectacular model independent predictions
- radical change of high energy physics at the TeV scale

Moreover no little hierarchy problem:

radiative electroweak symmetry breaking with no logs [8]

$\Lambda \sim$ a few TeV and $m_H^2 =$ a loop factor $\times \Lambda^2$

But unification has to be probably dropped

New Dark Matter candidates e.g. in the extra dims

Origin of EW symmetry breaking?

possible answer: radiative breaking

I.A.-Benakli-Quiros '00

$$V = \mu^2 H^\dagger H + \lambda (H^\dagger H)^2$$

$\mu^2 = 0$ at tree but becomes < 0 at one loop

non-susy vacuum

simplest case: one scalar doublet from the same brane

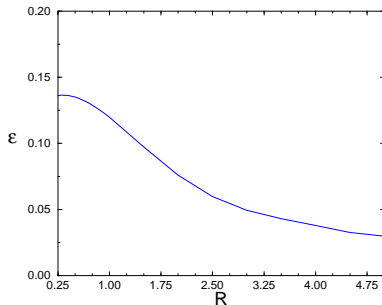
\Rightarrow tree-level V same as susy: $\lambda = \frac{1}{8}(g_2^2 + g'^2)$

D-terms

$\mu^2 = -g^2 \epsilon^2 M_5^2 \leftarrow$ effective UV cutoff

$$\epsilon^2(R) = \frac{R^3}{2\pi^2} \int_0^\infty dl l^{3/2} \frac{\theta_2^4}{16l^4 \eta^{12}} \left(il + \frac{1}{2} \right) \sum_n n^2 e^{-2\pi n^2 R^2 l}$$

Diagrammatic annotations for the equation above:
- A blue arrow labeled "UV" points to the upper limit of the integral (∞).
- A red arrow labeled "IR" points to the lower limit of the integral (0).
- A blue arrow labeled $e^{-\pi l}$ points to the exponential term in the summand.
- A red arrow labeled "1" points to the constant term $\frac{1}{2}$ in the summand's parentheses.



$R \rightarrow 0 : \varepsilon(R) \simeq 0.14$ large transverse dim $R_{\perp} = l_s^2/R \rightarrow \infty$

$R \rightarrow \infty : \varepsilon(R)M_s \sim \varepsilon_{\infty}/R$ $\varepsilon_{\infty} \simeq 0.008$ UV cutoff: $M_s \rightarrow 1/R$

Higgs scalar = component of a higher dimensional gauge field

$\Rightarrow \varepsilon_{\infty}$ calculable in the effective field theory

$\lambda = g^2/4 \sim 1/8 \quad \Rightarrow \quad M_H \simeq v/2 = 125 \text{ GeV}$

M_s or $1/R \sim$ a few or several TeV

Accelerator signatures: 4 different scales

- Gravitational radiation in the bulk \Rightarrow missing energy

present LHC bounds: $M_* \gtrsim 3 - 5$ TeV

- Massive string vibrations \Rightarrow e.g. resonances in dijet distribution

$$M_j^2 = M_0^2 + M_s^2 j \quad ; \quad \text{maximal spin : } j + 1$$

higher spin excitations of quarks and gluons with strong interactions

present LHC limits: $M_s \gtrsim 5$ TeV

- Large TeV dimensions \Rightarrow KK resonances of SM gauge bosons I.A. '90

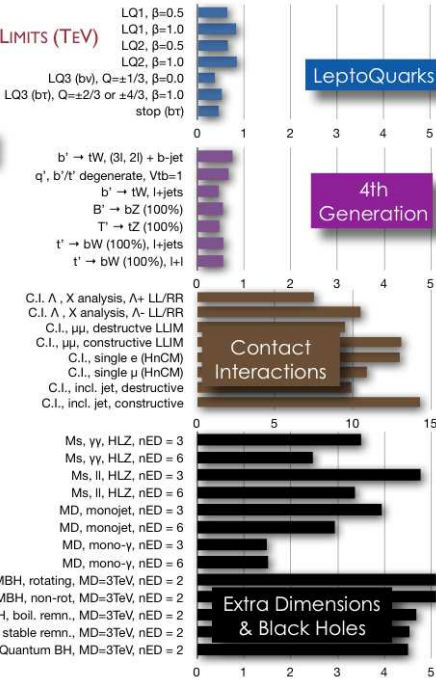
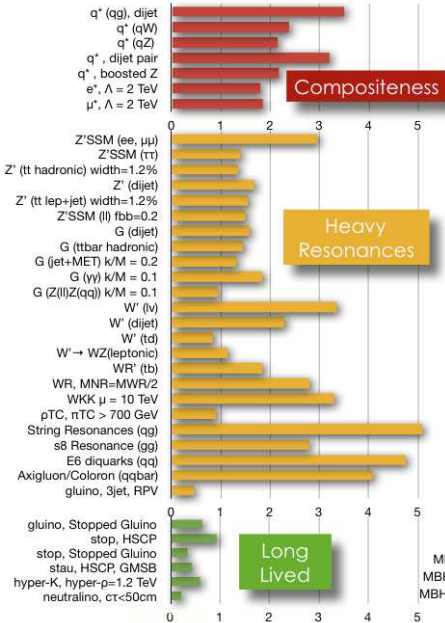
$$M_k^2 = M_0^2 + k^2/R^2 \quad ; \quad k = \pm 1, \pm 2, \dots$$

experimental limits: $R^{-1} \gtrsim 0.5 - 4$ TeV (UED - localized fermions)

- extra $U(1)$'s and anomaly induced terms

masses suppressed by a loop factor from M_s [30]

CMS EXOTICA 95% CL EXCLUSION LIMITS (TeV)



Extra $U(1)$'s and anomaly induced terms

masses suppressed by a loop factor

usually associated to known global symmetries of the SM

(anomalous or not) such as (combinations of)

Baryon and Lepton number, or PQ symmetry

Two kinds of massive $U(1)$'s:

I.A.-Kiritsis-Rizos '02

- 4d anomalous $U(1)$'s: $M_A \simeq g_A M_s$

- 4d non-anomalous $U(1)$'s: (but masses related to 6d anomalies)

$$M_{NA} \simeq g_A M_s V_2 \leftarrow (6d \rightarrow 4d) \text{ internal space} \Rightarrow M_{NA} \geq M_A$$

or massless in the absence of such anomalies [19]

- B and L become massive due to anomalies

Green-Schwarz terms

- the global symmetries remain in perturbation

- Baryon number \Rightarrow proton stability

- Lepton number \Rightarrow protect small neutrino masses

no Lepton number $\Rightarrow \frac{1}{M_s} LLHH \rightarrow$ Majorana mass: $\frac{\langle H \rangle^2}{M_s} LL$

\sim GeV

- $B, L \Rightarrow$ extra Z' 's

with possible leptophobic couplings leading to CDF-type Wjj events

$Z' \simeq B$ lighter than 4d anomaly free $Z'' \simeq B - L$

Conclusions

- Confirmation of the EWSB scalar at the LHC:
important milestone of the LHC research program
 - Precise measurement of its couplings is of primary importance
 - Hint on the origin of mass hierarchy and of BSM physics
 - natural or unnatural SUSY?
 - low string scale in some realization?
 - something new and unexpected?
- all options are still open
- LHC enters a new era with possible new discoveries