

# Vector-like Fermions and the Electroweak Phase Transition

**Eduardo Pontón**

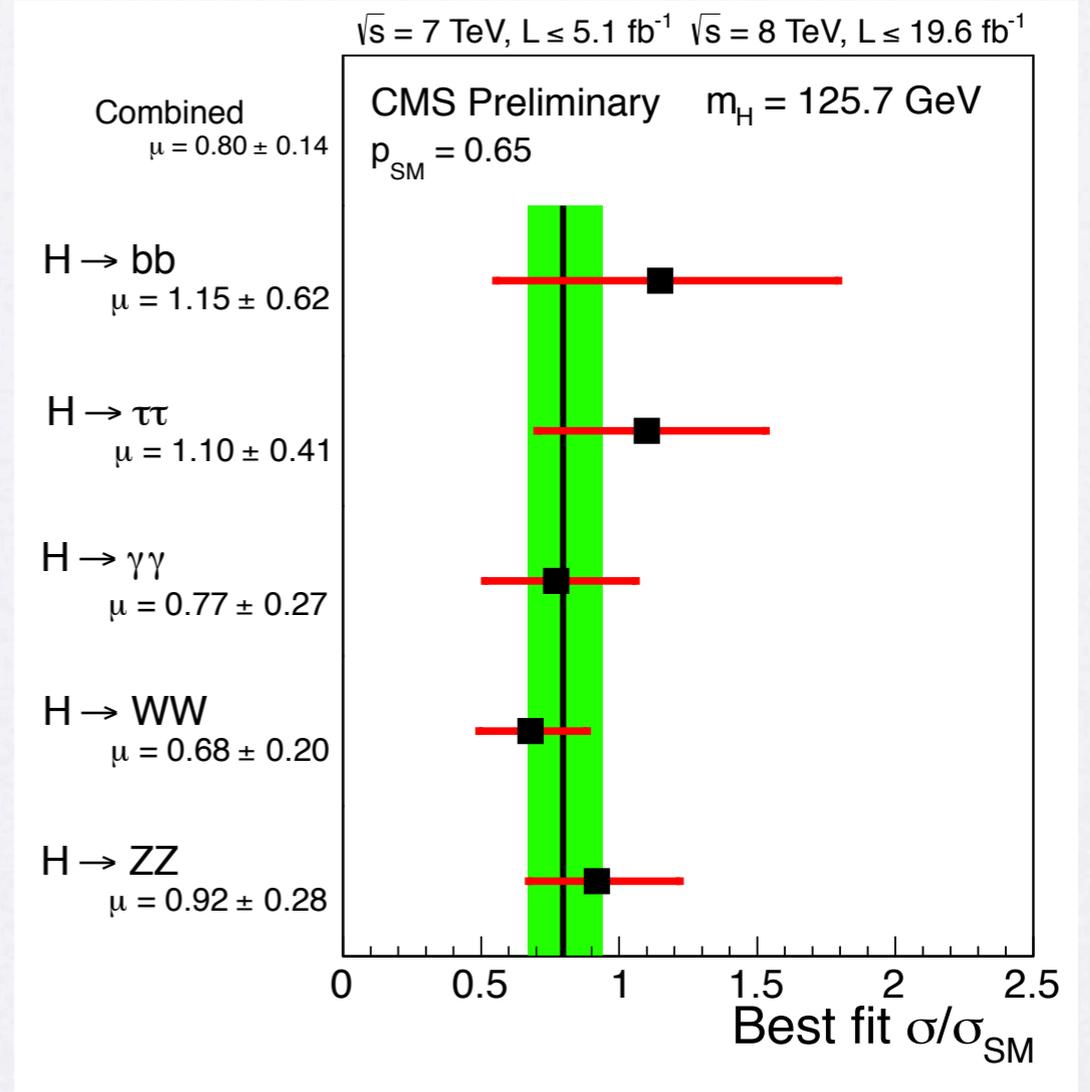
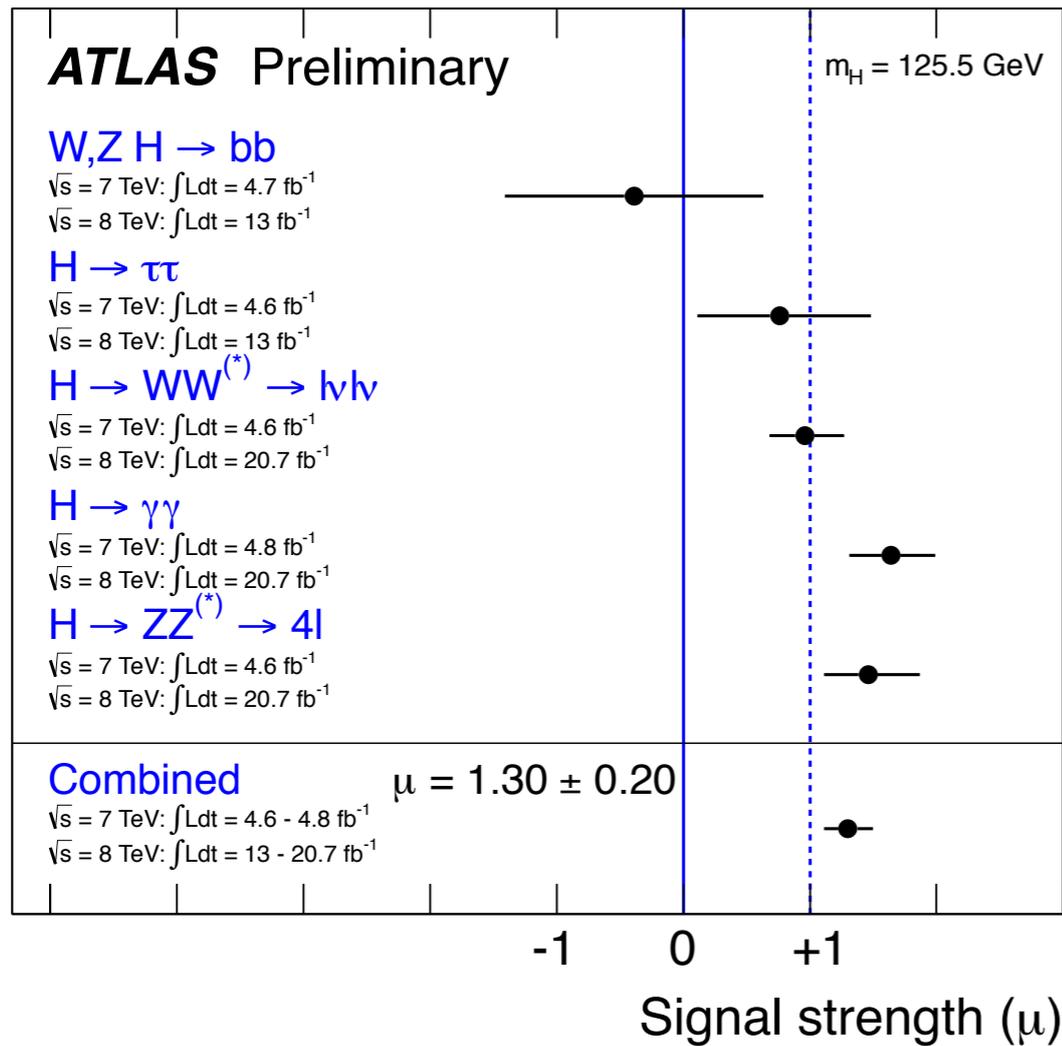
**Instituto de Física Teórica -UNESP  
& ICTP-SAIFR**

(Work with H. Davoudiasl & I. Lewis, 1211.3449)

**Beyond the SM after the first LHC run, GGI**

**July 11, 2013**

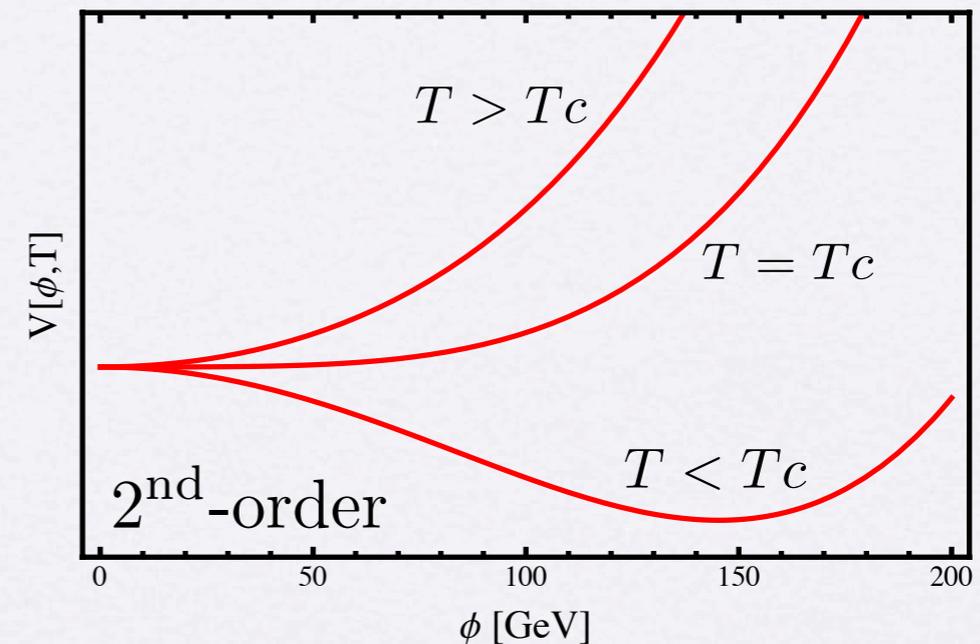
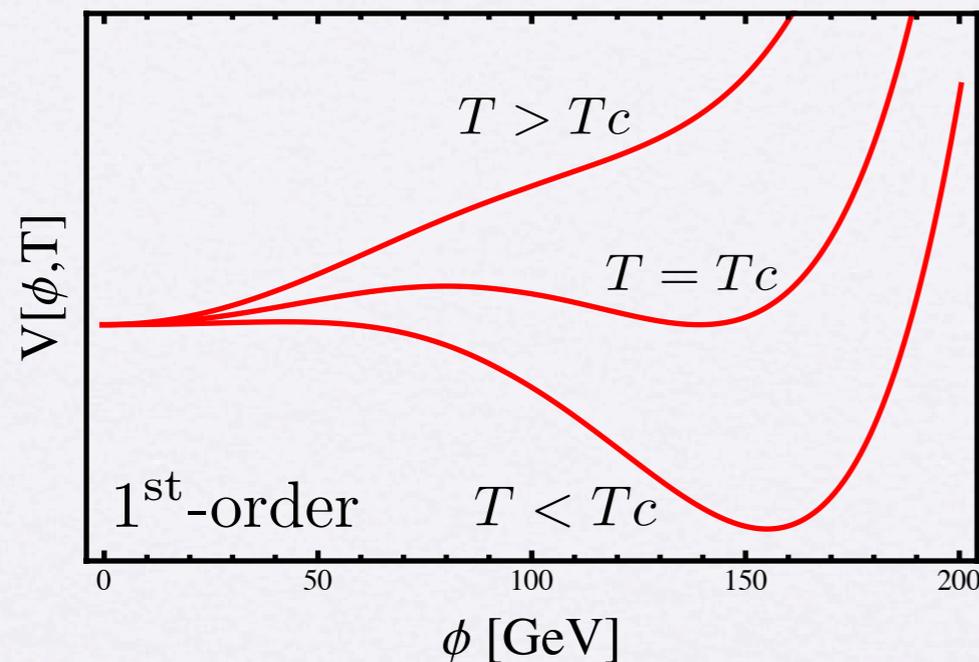
# A SM-like Higgs



Overall good agreement with SM expectations... but still room for interesting deviations.

# The Nature of the EWPhT

Precise measurement of Higgs properties can illuminate nature of EW phase transition



In the SM, with  $m_H = 125$  GeV : a smooth crossover (1<sup>st</sup>-order PT only for  $m_H \lesssim 80$  GeV)

(Rummukainen et. al., hep-lat/9805013)

But deviations (due to new physics) might have potentially important implications

- Baryon Asymmetry of the Universe
- Gravitational wave signals

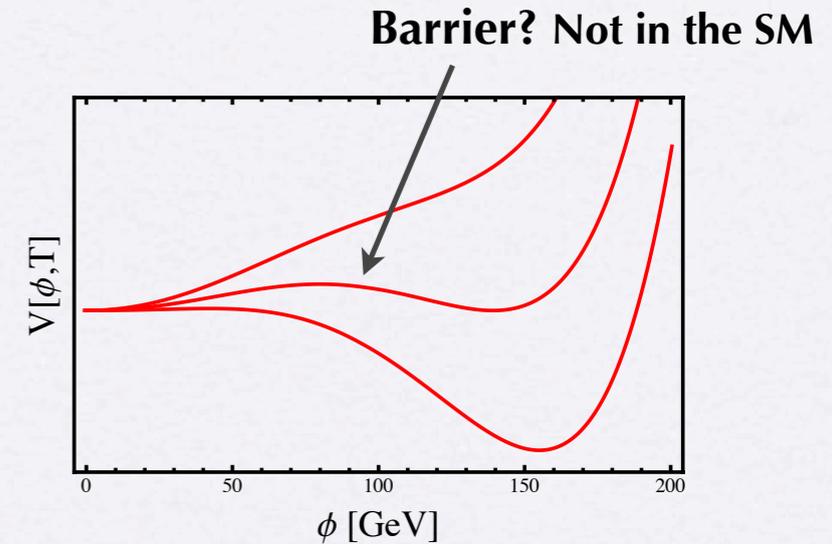
# The EWPhT in the SM

Finite temperature: thermal masses + new cubic term in  $\phi$  :

$$V(\phi, T) \sim \underbrace{m(T)^2 \phi^2}_{\text{when positive: min. at origin}} + \underbrace{ET\phi^3 + \lambda(T)\phi^4}_{\text{"far away min."}}$$

when positive:  
min. at origin

"far away min.":  $\phi \sim ET/\lambda$



$$\rightarrow \frac{\phi_c}{T_c} \sim \frac{E}{\lambda} \sim \frac{E v^2}{m_h^2}$$

**Small E:**  $\phi_c/T_c \ll 1$

(need non-perturbative methods to establish nature of phase transition)

Should be  $> 1$ , for EW baryogenesis

# BSM and the EWPhT

New physics required if the EWPhT plays a role in generation of the BAU

“Lore”: to strengthen the EWPhT, requires new *bosonic* degrees of freedom to either

- change the Higgs potential at tree-level (e.g. adding singlets)
- enhance the  $E$ -term at loop level

In light of Higgs discovery:

Cohen, Morrissey & Pierce, 1203.2924

Carena, Nardini, Quirós & Wagner, 1207.6330

Fok, Kribs, Martin & Tsai, 1208.2784

Chung, Long & Wang, 1209.1819

Huan, Shu & Zhang, 1210.0906

Laine, Nardini & Rummikainen, 1211.7344

Both cases can be thought as relying on effective cubic terms

Typically, new fermions not thought to be particularly useful for this purpose...

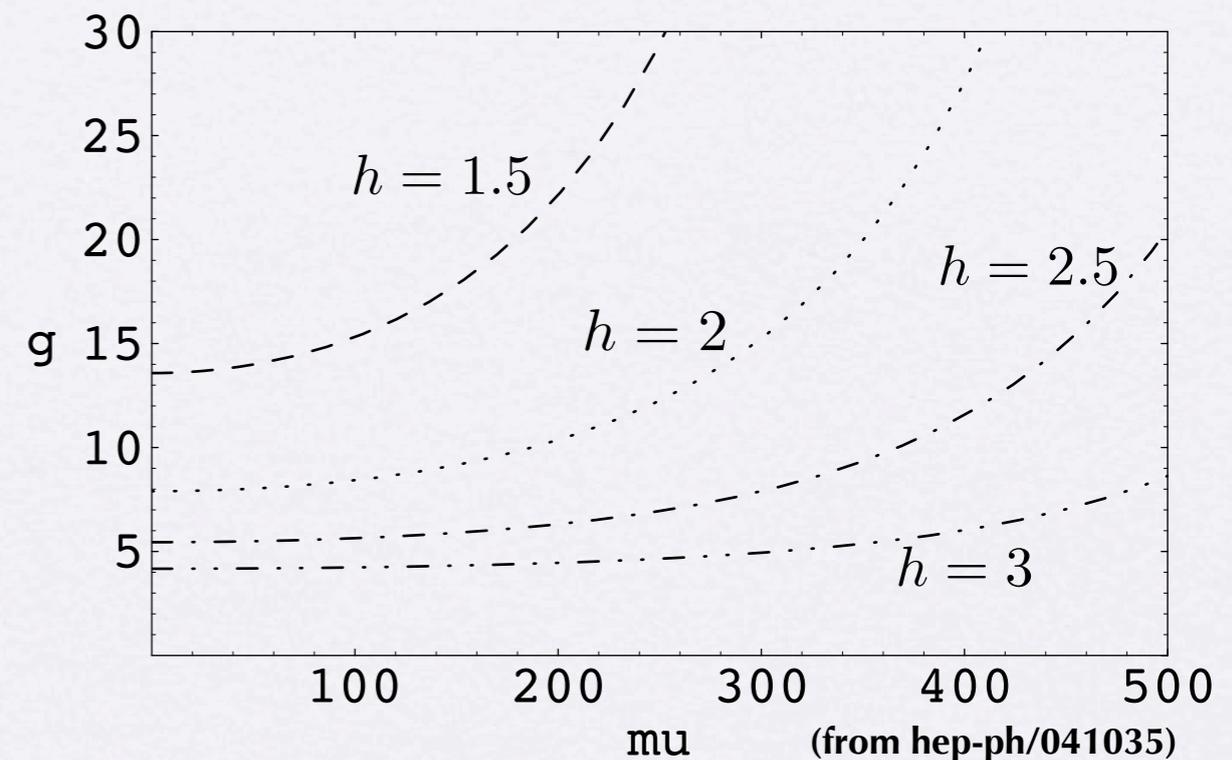
... they do not induce a cubic term

# Can Fermions Help?

However, I know of one previous example where fermions can help the EWPhT:

- Use the fact that when their mass depends on the Higgs vev, it is different in the broken and unbroken phases.
- **Decoupling** from thermal bath in **broken phase** leads to reheating, delaying the phase transition:  
→ increase in  $\phi_c/T_c$

Carena, Megevand, Quirós & Wagner,  
hep-ph/041035

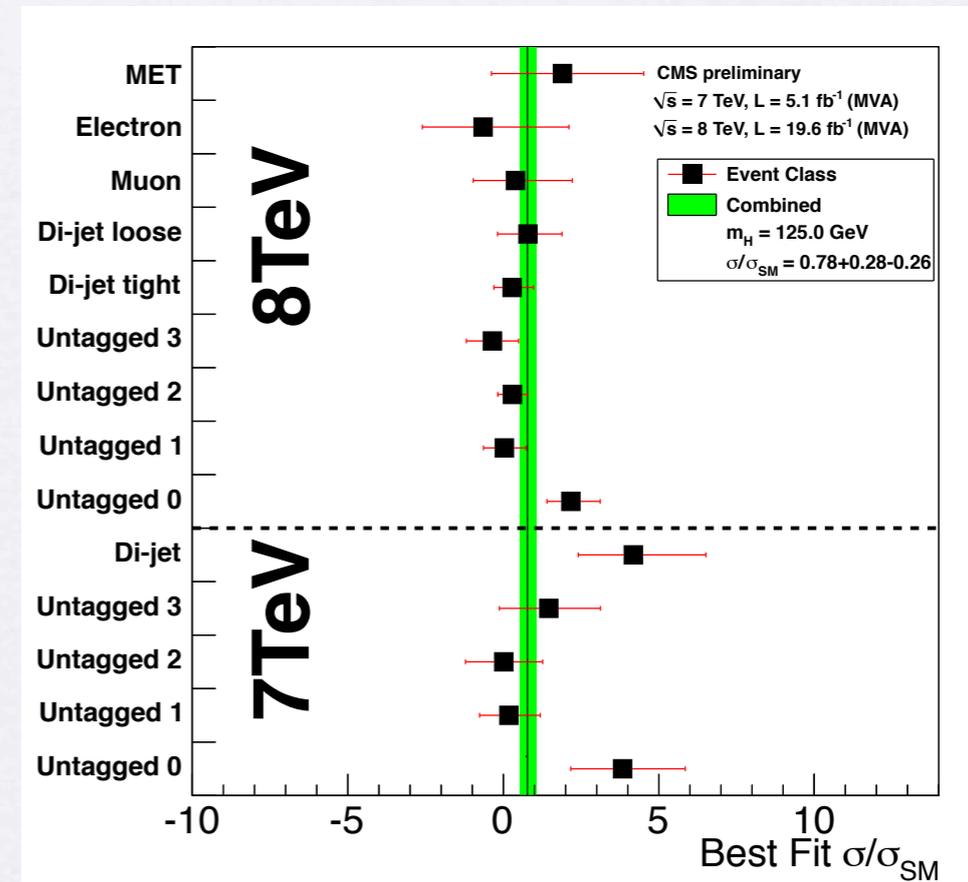
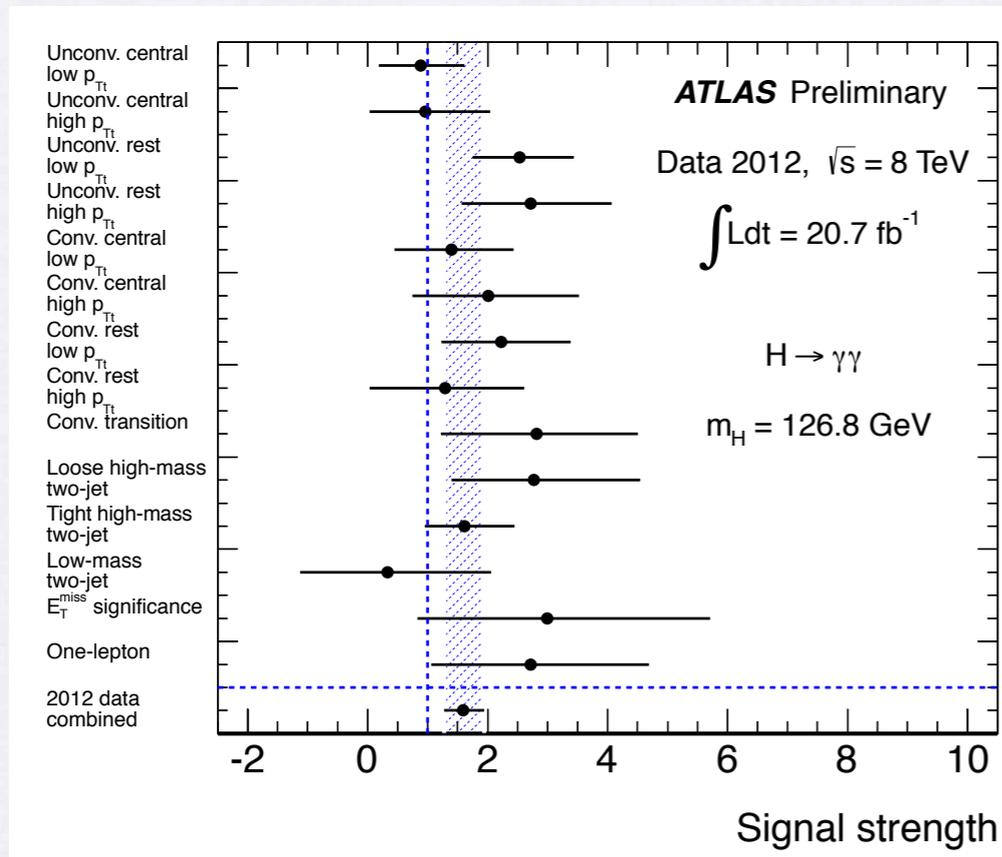


# Higgs Di-photon Rate

Loop-level processes prime suspects for deviations from SM

$$R_{\gamma\gamma}^{\text{ATLAS}} = 1.80^{+0.42}_{-0.36} \quad R_{\gamma\gamma}^{\text{CMS}} = 1.56 \pm 0.43 \quad (\text{pre-Moriond 2013})$$

$$R_{\gamma\gamma}^{\text{ATLAS}} = 1.55^{+0.33}_{-0.28} \quad R_{\gamma\gamma}^{\text{CMS}} = 0.78^{+0.28}_{-0.26} \left\{ \begin{array}{l} 1.69^{+0.65}_{-0.59} \quad (7 \text{ TeV}) \\ 0.55^{+0.29}_{-0.27} \quad (8 \text{ TeV}) \end{array} \right. \quad (\text{post-Moriond 2013})$$





# A Simple Model

Minimal extension (for illustration):

$$(\psi, \psi^c) \sim (1, 2)_{\pm 1/2} \quad (\chi, \chi^c) \sim (1, 1)_{\mp 1} \quad (\text{“vector-like leptons”})$$

Mass and Yukawa terms:

$$-\mathcal{L}_m = -m_\psi \psi \psi^c + m_\chi \chi \chi^c + y H \psi \chi + y_c H^\dagger \psi^c \chi^c + \text{h.c.}$$

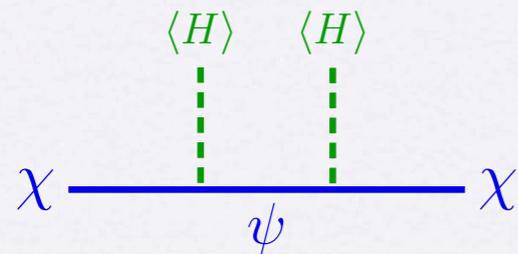
Interesting region for EWPhT when fermion masses hierarchical. Assume  $m_\psi \gg m_\chi$

EFT analysis with  $\psi$  integrated out: may be more general than simple model

# EFT Analysis

(Davoudiasl, Lewis & EP, 1211.3449)

For  $m_\psi \gg v, m_\chi$ , integrate out “ $\psi$ ” and obtain EFT for (SM + “ $\chi$ ”):



$$\Delta\mathcal{L} = 2G_m H^\dagger H \chi \chi^c$$

where  $G_m = \frac{y^2}{2(m_\psi - m_\chi)} > 0$

Light vector-like fermion mass is  $m_1 = m_\chi - \underbrace{G_m v^2}_{\text{“level repulsion”}}$

Interactions of light state with physical Higgs (after EWSB):

$$\mathcal{L}_{\text{Yuk}} = y_{\text{eff}} h \chi \chi^c + \text{h.c.}$$

$$y_{\text{eff}} = -2G_m v < 0$$

Interesting diphoton enhanc. when

$$y_{\text{eff}} \sim \mathcal{O}(1) \text{ and}$$

$$m_1 \sim 100 - 200 \text{ GeV}$$

# $T = 0$ Potential in EFT

(Davoudiasl, Lewis & EP, 1211.3449)

1-loop effective potential ( $T = 0$ )

$$\Delta V_{\text{Eff}} = \text{---} \bigcirc \text{---} + \text{---} \bigcirc \text{---} + \text{---} \bigcirc \text{---} + \dots$$

$$= -\frac{4}{64\pi^2} m_1(\phi)^4 \left[ \ln \frac{m_1(\phi)^2}{\mu^2} - \frac{3}{2} \right] + \text{counterterms} \quad (\text{in } \overline{MS} \text{ scheme})$$

Noting that  $\phi^6$  and  $\phi^8$  terms are divergent:

$$"V_{\text{Tree}}" = -\frac{1}{2} \bar{\mu}^2 \phi^2 + \frac{1}{4} \bar{\lambda} \phi^4 + \frac{1}{6} \bar{\gamma} \phi^6 + \frac{1}{8} \bar{\delta} \phi^8$$

$\bar{\gamma}$  and  $\bar{\delta}$   
arbitrary within EFT

However, within the UV model, they are determined...

# Matching and Running

(Davoudiasl, Lewis & EP, 1211.3449)

However, within the UV model, they are determined...

$$\bar{\gamma} = \bar{\gamma}_{\text{th}} + \bar{\gamma}_{\text{RG}} \qquad \bar{\delta} = \bar{\delta}_{\text{th}} + \bar{\delta}_{\text{RG}}$$

with threshold contributions

$$\bar{\gamma}_{\text{th}} = \frac{Z_\gamma y^6}{16\pi^2} \frac{m_\psi(m_\psi^2 + 7m_\chi m_\psi - 2m_\chi^2)}{(m_\psi - m_\chi)^5} \sim \frac{y^6}{16\pi^2} \frac{1}{m_\psi^2}$$
$$\bar{\delta}_{\text{th}} = -\frac{Z_\delta y^8}{48\pi^2} \frac{7m_\psi^3 + 27m_\chi m_\psi^2 - 4m_\chi^3}{(m_\psi - m_\chi)^7} \sim -\frac{7y^8}{48\pi^2} \frac{1}{m_\psi^4}$$

$\left( Z_\gamma, Z_\delta = 1 \right)$   
(at lowest order)

and running contributions

$$\bar{\gamma}_{\text{RG}} \approx -\frac{3G_m^3 m_\chi}{2\pi^2} \ln \left( \frac{m_\psi^2}{\mu^2} \right)$$
$$\bar{\delta}_{\text{RG}} \approx \frac{G_m^4}{2\pi^2} \ln \left( \frac{m_\psi^2}{\mu^2} \right)$$

- Determined by EFT only
- This is the only difference with naive CW potential in UV model (actually, just the sign)

# Quartic Instabilities?

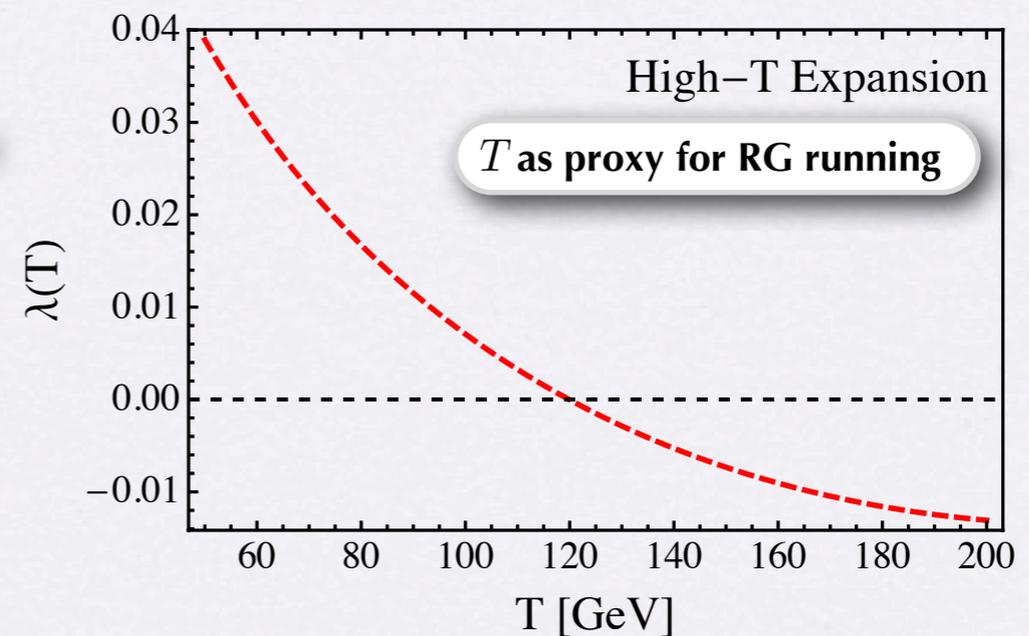
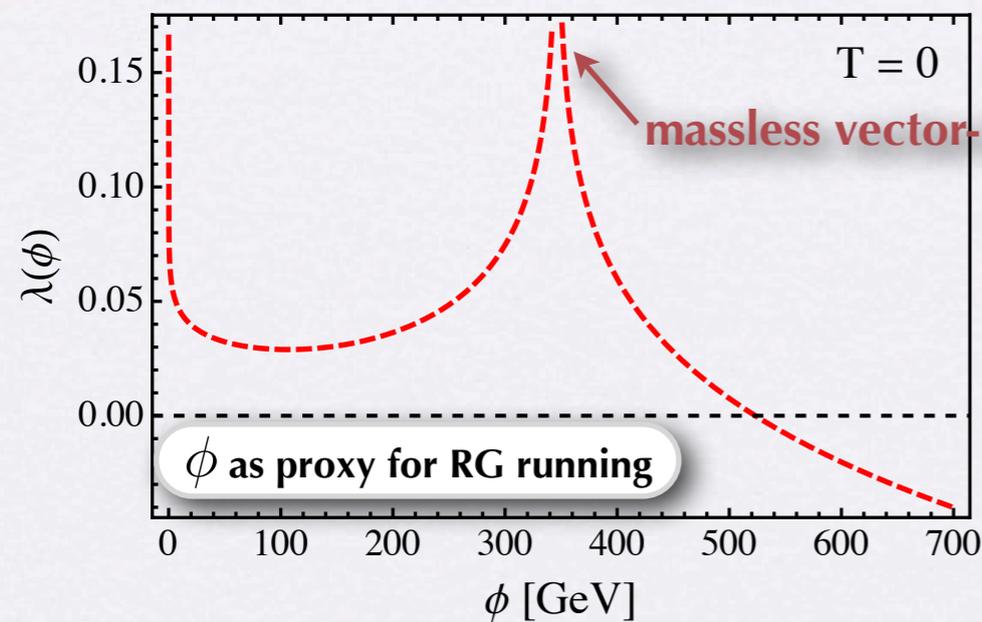
(Davoudiasl, Lewis & EP, 1211.3449)

RG running of Higgs quartic (below  $m_\psi$ ):

$$16\pi^2 \frac{d\lambda}{dt} = \lambda (6\lambda - 9g_2^2 - 3g_1^2 + 12y_t^2) \underbrace{-6y_t^4}_{\text{from new light fermion}} + \frac{3}{8} \left[ 2g_2^2 + (g_2^2 + g_1^2)^2 \right] \underbrace{-48G_m^2 m_\chi^2}_{\text{from new light fermion}}$$

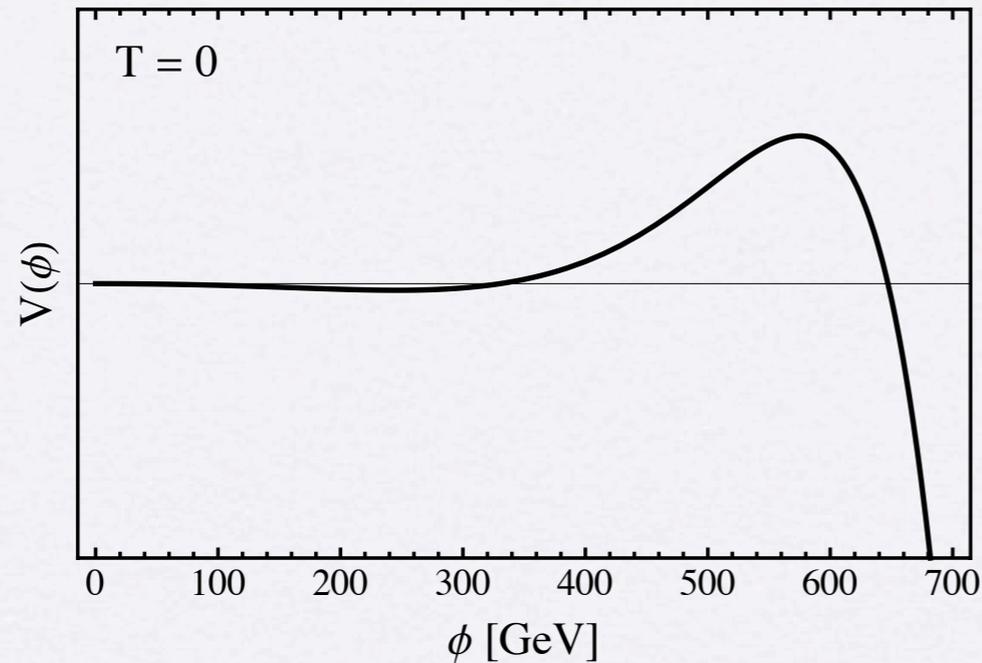
fermionic terms induce "instability"

Quartic coupling from effective potential at low and high temperatures:



# The Shape of the Eff. Pot.

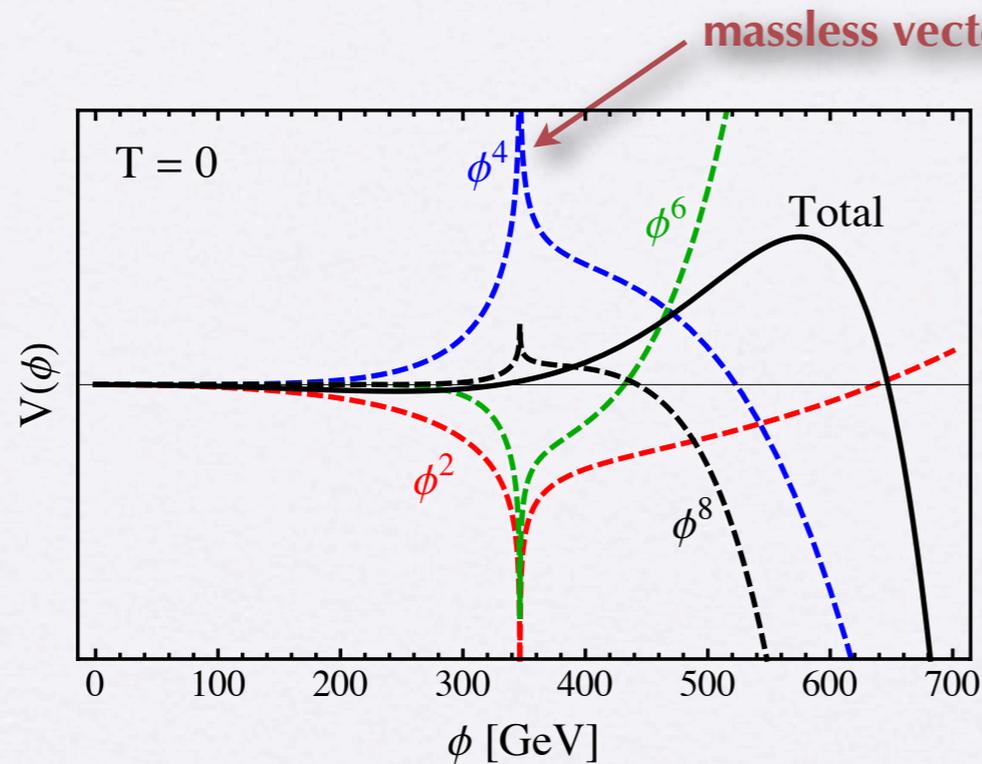
(Davoudiasl, Lewis & EP, 1211.3449)



- **Effective Potential in EFT shows instability at**  
 $\sim 600 \text{ GeV} \gg \text{EW scale}$

# The Shape of the Eff. Pot.

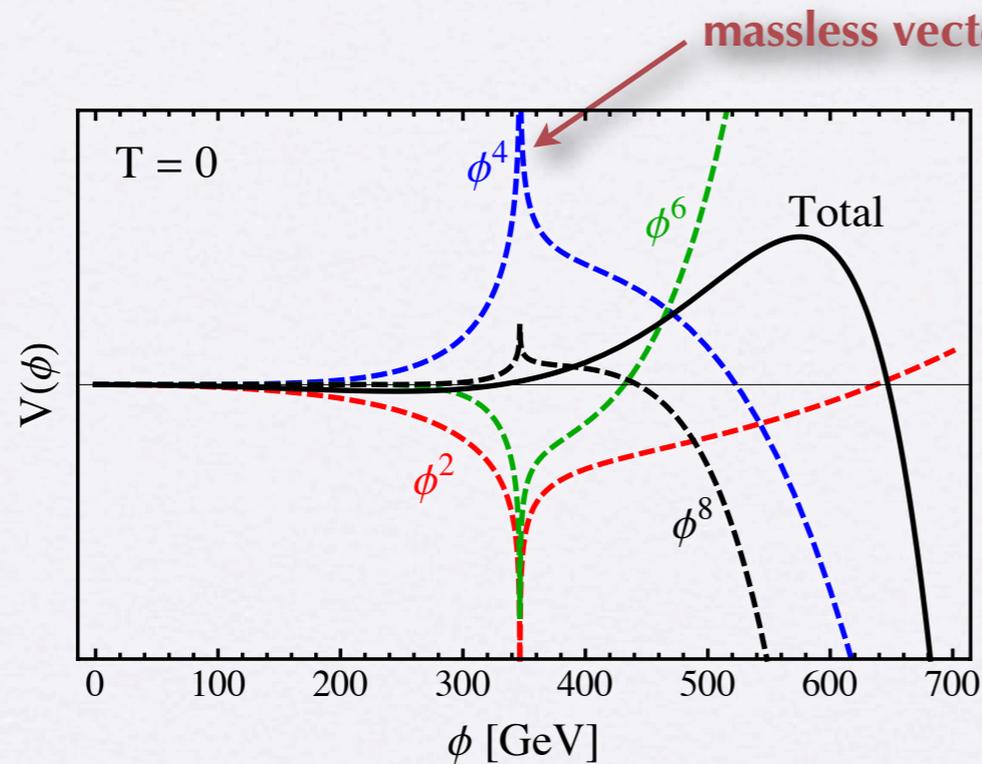
(Davoudiasl, Lewis & EP, 1211.3449)



- Effective Potential in EFT shows instability at  $\sim 600 \text{ GeV} \gg \text{EW scale}$

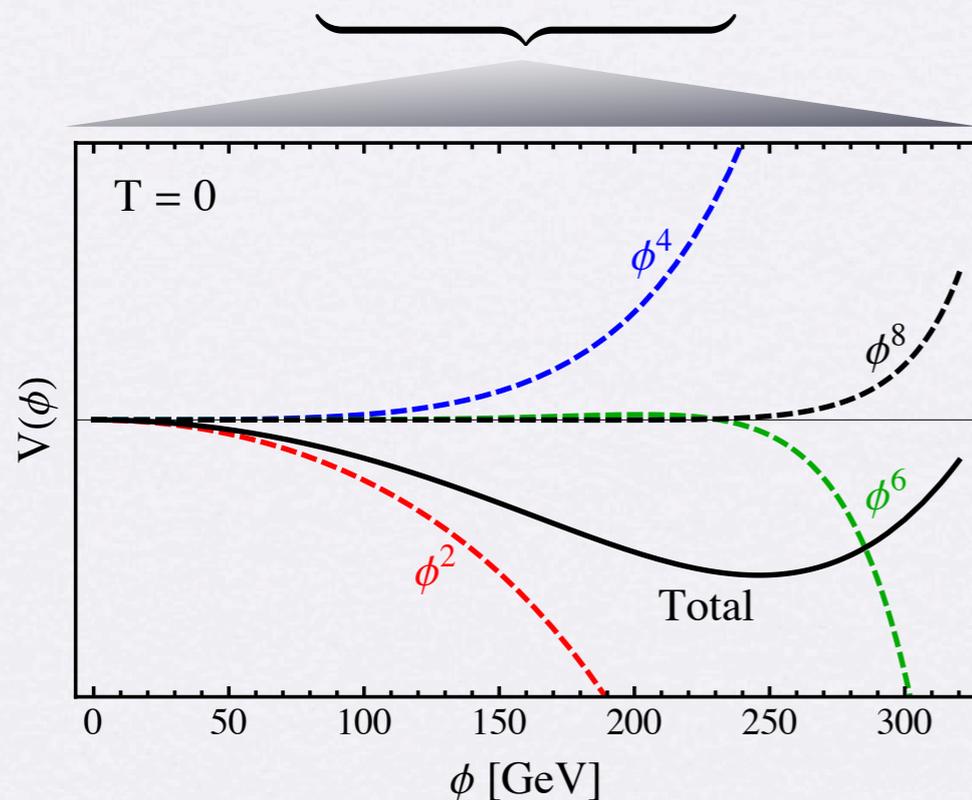
# The Shape of the Eff. Pot.

(Davoudiasl, Lewis & EP, 1211.3449)



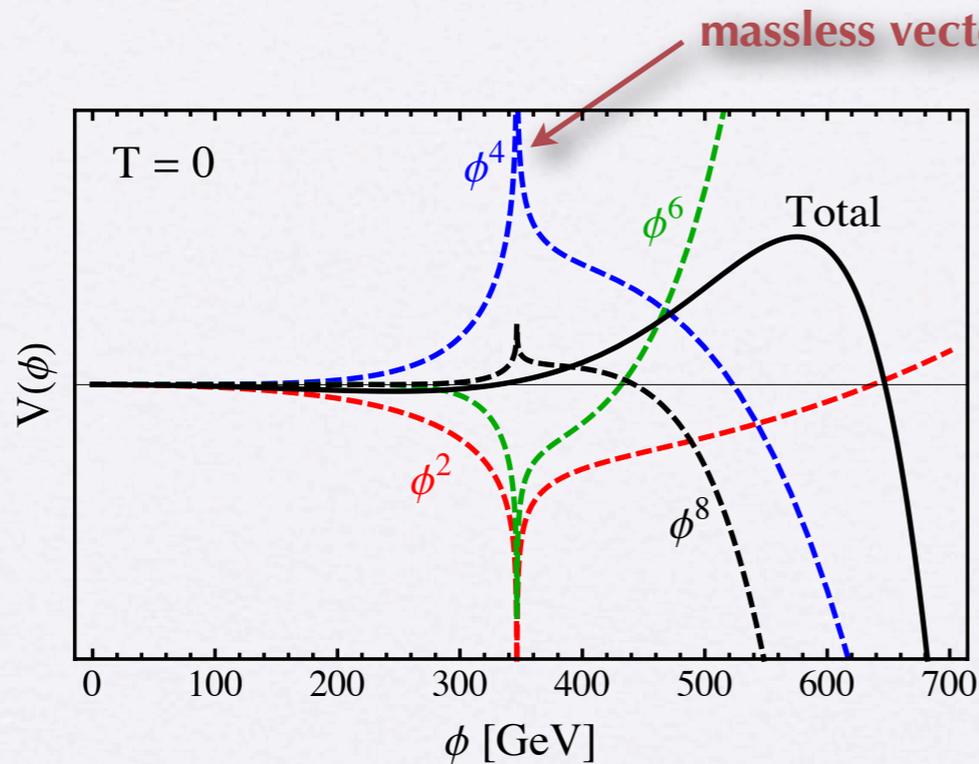
- Effective Potential in EFT shows instability at  $\sim 600 \text{ GeV} \gg \text{EW scale}$

- At  $T = 0$  and  $\phi \sim \text{EW}$ :  $m^2 < 0$ ,  $\lambda > 0$



# The Shape of the Eff. Pot.

(Davoudiasl, Lewis & EP, 1211.3449)



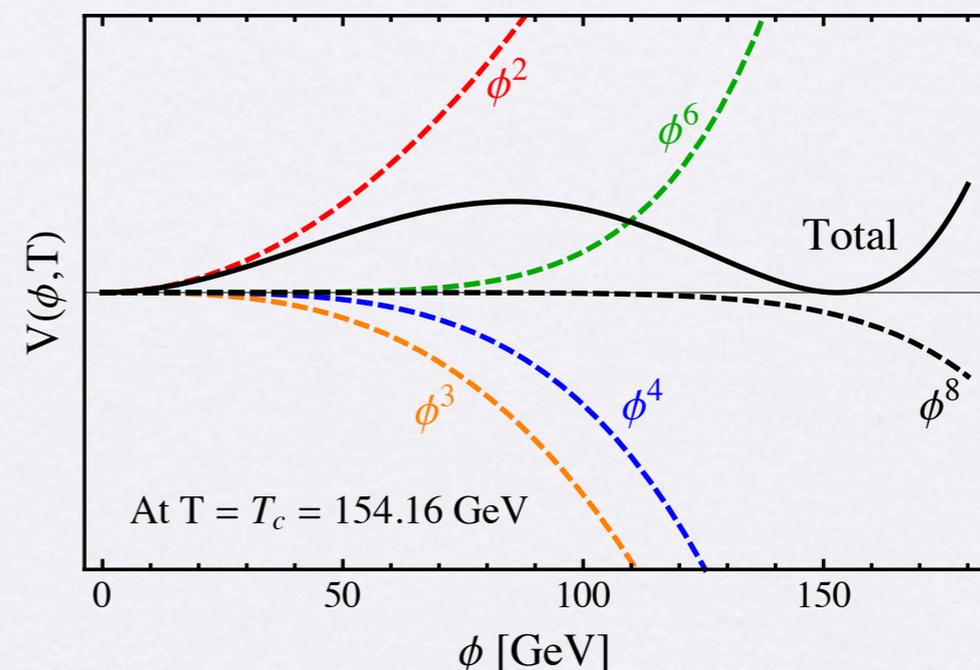
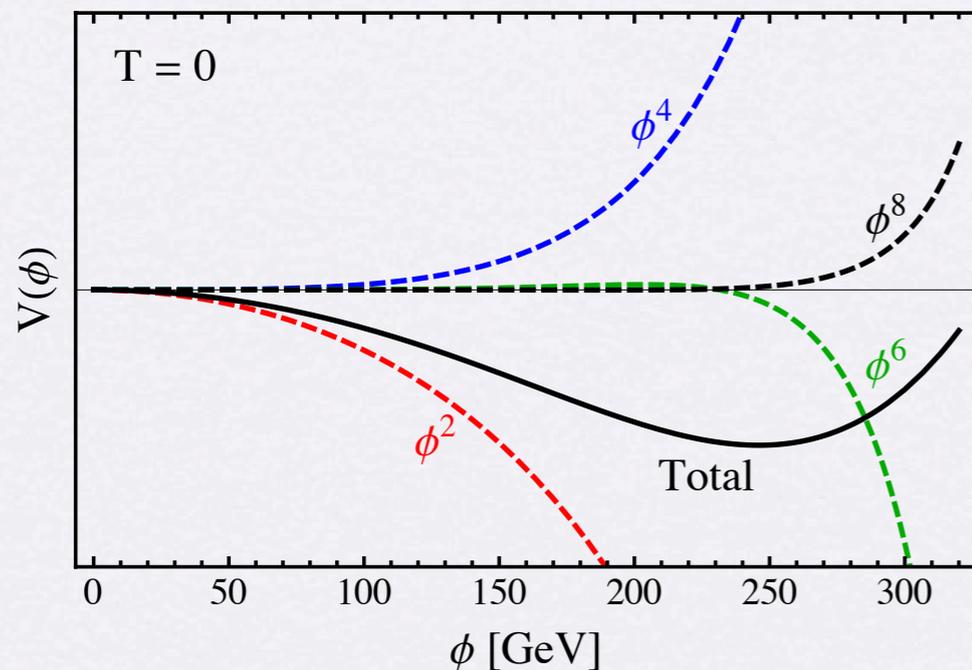
- Effective Potential in EFT shows instability at  $\sim 600 \text{ GeV} \gg \text{EW scale}$

- At  $T = 0$  and  $\phi \sim \text{EW}$ :  $m^2 < 0$ ,  $\lambda > 0$

- At  $T \sim \text{EW}$ :  $m^2 > 0$ ,  $\lambda < 0$

→ stabilization by higher-dim operators

(Realization of idea in hep-ph/0407019)



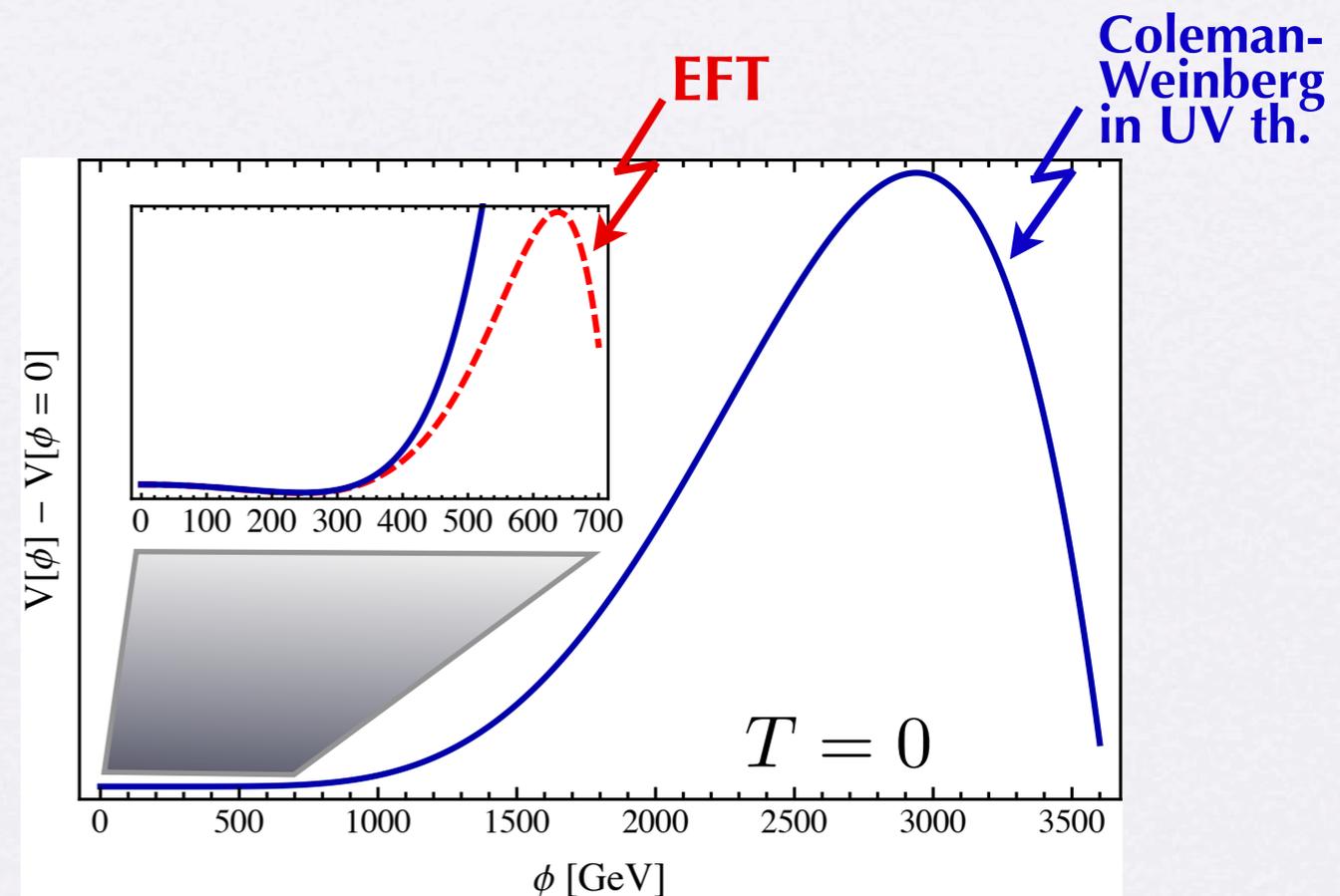
# Back to the Instability

(Davoudiasl, Lewis & EP, 1211.3449)

- EFT with  $\psi$  integrated out: match  $\phi$  correlators and run from  $m_\psi$  to  $m_\chi \sim v$ 
  - captures “small  $\phi$ ” behavior, but not large  
(finite radius of convergence of Taylor expansion of effective potential)

- Coleman-Weinberg potential in full UV model suggests instability delayed to multi-TeV scale ( $\sim m_\psi$ )

- Thus, in non-renormalizable theories one should be careful in interpreting the familiar quartic instability



# The Mechanism

Keep it simple by dropping non-crucial terms, e.g. cubic:

$$V(\phi, T) \sim \frac{1}{2} \bar{\mu}^2 \phi^2 + \underbrace{\frac{1}{4} \bar{\lambda} \phi^4 + \frac{1}{6} \bar{\gamma} \phi^6}_{\text{"far away min."}}$$

When  $\bar{\mu}^2 > 0, \bar{\lambda} < 0, \bar{\gamma} > 0$

"far away min.":  $\phi \sim \sqrt{-\bar{\lambda}/\bar{\gamma}}$

Degenerate with min. at origin when  $\bar{\lambda}^2 \sim 6\bar{\gamma}\mu^2 \ll 1$  (determines critical temp.)

Also estimate

$$\phi_c \sim \sqrt{\mu}/\gamma^{1/4} \rightarrow \text{may get sizeable } \phi_c/T_c !$$

Similar to proposal by Grojean, Servant & Wells. Here,  $\lambda < 0$  from fermions and finite Temp.  
(hep-ph/0407019)

# EFT Agnostic Analysis

(Davoudiasl, Lewis & EP, 1211.3449)

- Strength of the phase transition ( $\phi_c/T_c$ ) in

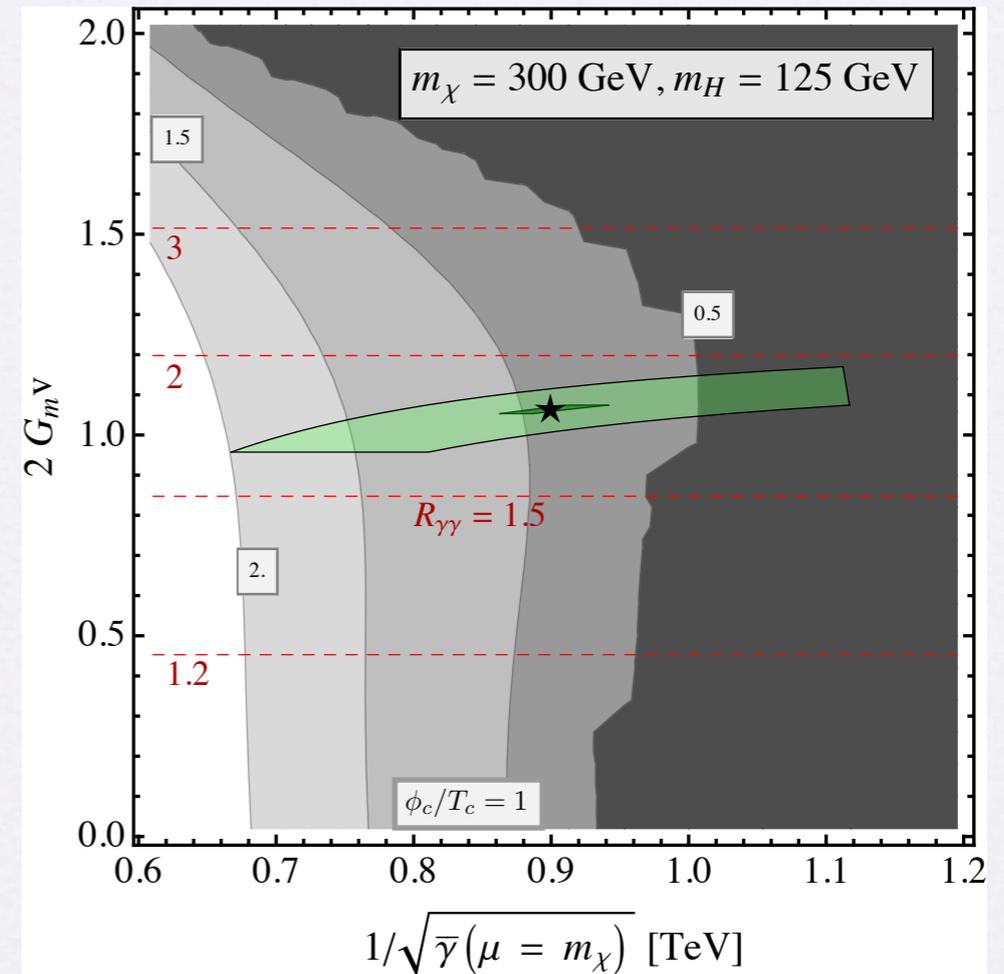
$$-y_{\text{eff}} = 2G_m v \quad (\text{coupling of light fermion to Higgs})$$

versus

$$\bar{\gamma}(\mu = m_\chi) \quad (\text{dim-6 stabilizing operator})$$

## Observations:

- Need sizeable underlying  $y, y_c \sim \text{few}$
- Sensitivity to UV completion through stabilizing higher-dim. operators
- Consistent with important Higgs diphoton rate *enhancement*



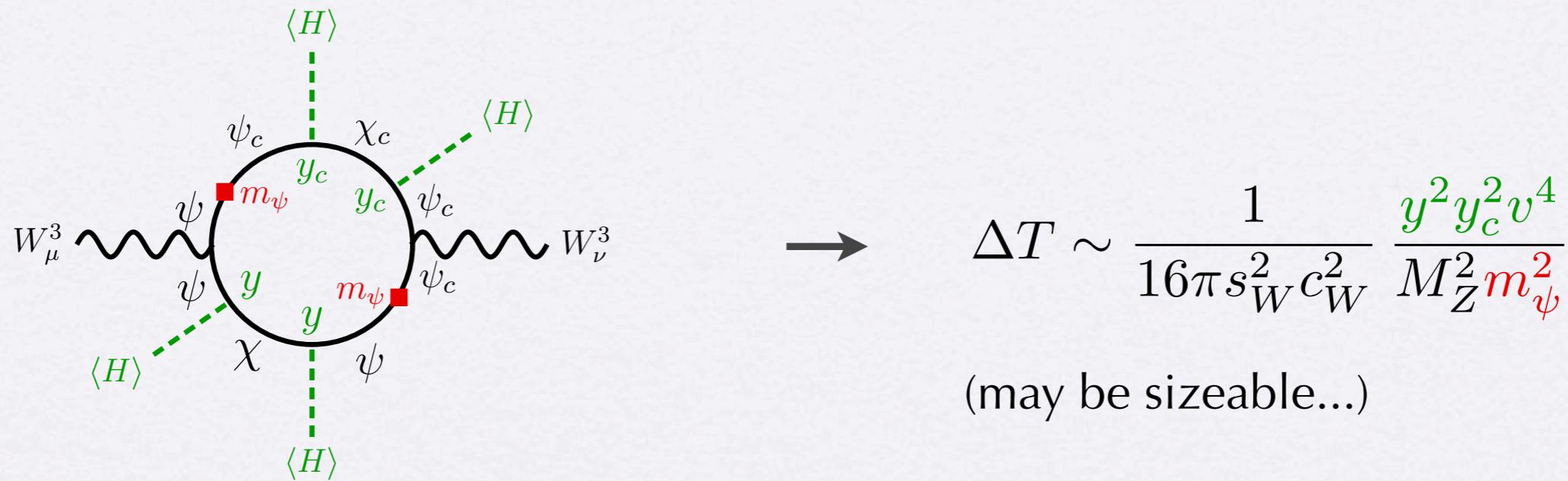
**Star** is a benchmark in UV model with:

$m_\psi = 4 \text{ TeV}$	$m_\chi = 300 \text{ GeV}$	$y = y_c = 4$
--------------------------	----------------------------	---------------

**Green regions:** effect of 10% (1%) higher-loop corrections at the matching scale

# EW Precision Tests

In the heavy doublet limit:  $m_\psi \gg m_\chi, v$ , leading contribution to T from



The S-parameter is less constraining in the same limit:

$$\Delta S \sim \frac{2y^2 v^2}{9\pi m_\psi^2} [6 \ln(m_\psi/m_\chi) - 7]$$

# A Custodial Extension

Two important shortcomings so far:

- In spite of large  $m_\psi$ , sizeable T parameter suggests imposing custodial symmetry
- Lightest charged state must decay (so far stable)

Both can be addressed by adding a vector-like “RH neutrino”,  $(n, n^c)$

$$-\Delta\mathcal{L}_m = m_n n n^c + \tilde{y} H^\dagger \psi n + \tilde{y}_c H \psi^c \chi^c + \text{h.c.}$$

When  $y = \tilde{y}$ ,  $y_c = \tilde{y}_c$  and  $m_n = -m_\chi$ , can rewrite as  $SU(2)_L \times SU(2)_R$  invariant:

$$-\mathcal{L}_{\text{Yuk}} = y \psi \Phi \xi + y_c \psi^c \Phi \xi^c + \text{h.c.} \quad \text{with} \quad \xi^{(c)} \equiv \begin{pmatrix} n^{(c)} \\ -\chi^{(c)} \end{pmatrix} \quad \Phi \equiv \begin{pmatrix} H^{0*} & H^+ \\ -H^- & H^0 \end{pmatrix}$$

$(0, 2) \qquad (2, 2)$

In this limit  $T = 0$ . Need to ensure neutral lightest state, which breaks custodial softly.

# A Detailed Example

## Input parameters:

$m_\psi = 4 \text{ TeV}$	$m_\chi = 300 \text{ GeV}$	$m_n = -250 \text{ GeV}$	$y = 4$	$\tilde{y} = 4$	$y_c = 3.5$	$\tilde{y}_c = 3.5$
--------------------------	----------------------------	--------------------------	---------	-----------------	-------------	---------------------

## Vector-like spectrum:

At  $T = 0$

Charged	Neutral
$m_2^\pm = 4.11 \text{ TeV}$	$m_2^0 = 4.11 \text{ TeV}$
$m_1^\pm = 189 \text{ GeV}$	$m_1^0 = 140 \text{ GeV}$

At  $T = T_c$

Charged	Neutral
$m_2^\pm = 4.06 \text{ TeV}$	$m_2^0 = 4.06 \text{ TeV}$
$m_1^\pm = 240 \text{ GeV}$	$m_1^0 = 191 \text{ GeV}$

## Phase transition:

$$\begin{aligned} \phi_c &= 179.3 \text{ GeV} \\ T_c &= 158.4 \text{ GeV} \\ \phi_c/T_c &= 1.13 \end{aligned}$$

Bubbles nucleate slightly below  $T_c$

## EWPT and diphoton enhancement:

$$\begin{aligned} \Delta T &\sim 10^{-4} \\ \Delta S &\approx 0.04 \\ \Delta U &\sim 10^{-6} \\ R_{\gamma\gamma} &\approx 1.5 \end{aligned}$$

} Consistent at 95% CL  
with current PDG ellipse

← Look forward to further  
developments!

# Conclusions & Outlook

- We illustrated in a simple model the potentially far-reaching consequences of deviations from the SM Higgs properties in answering long-standing questions:
  - The nature of the EWPhT itself
  - The relevance of EW scale physics in the generation of the BAU (details to be worked out)
- Within the model:
  - Triple Higgs coupling:  $V'''(v) = 3m_H^2/v + \overbrace{8\bar{\gamma}v^3}^{\mathcal{O}(1) \text{ correction}}$ 
    - expect 40-60% suppression in  $gg \rightarrow HH$
    - measurement of stabilizing effect (a crucial ingredient in underlying mechanism)
  - Measurement of lightest charged fermion mass + diphoton rate:  $(m_\chi, G_m v)$
- **Main ingredients:**
  - a fermion state in the few TeV scale
  - a parametrically lighter fermion state
  - Underlying Yukawa interactions with  $y \sim 3 - 4$

} Rather familiar from (warped) extra-dimensional constructions!