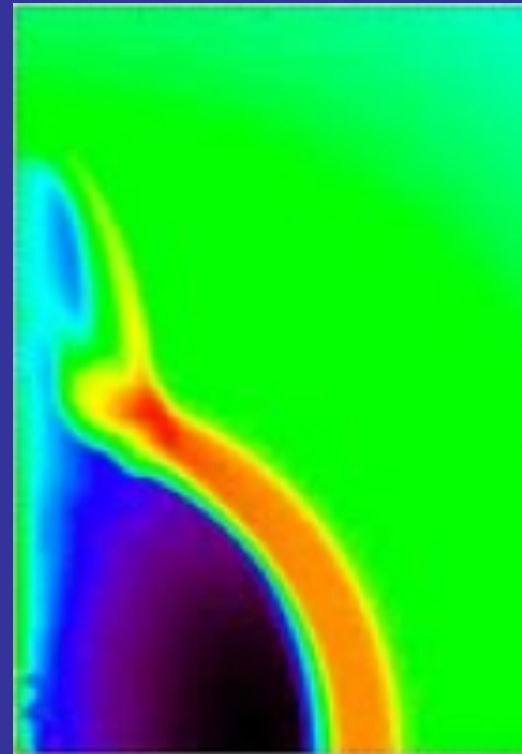
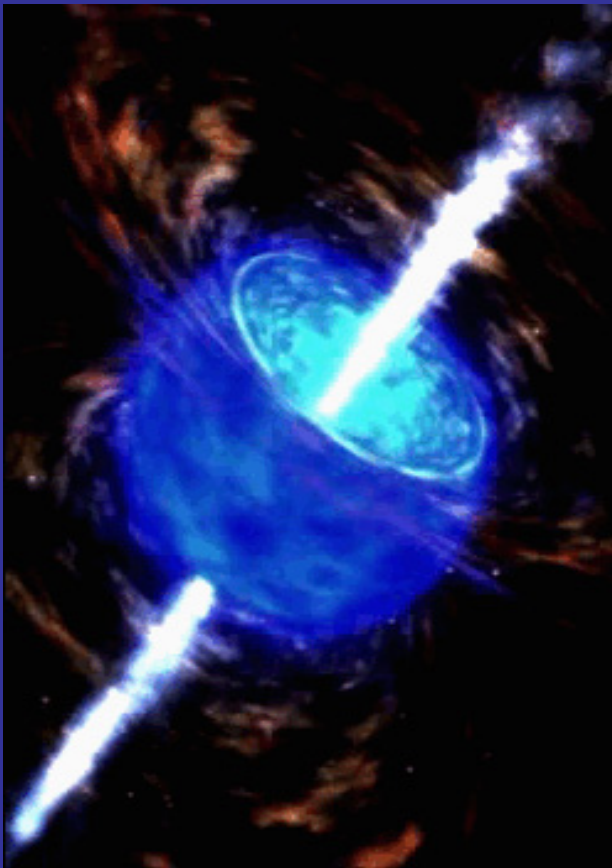


Gamma-Ray Bursts:

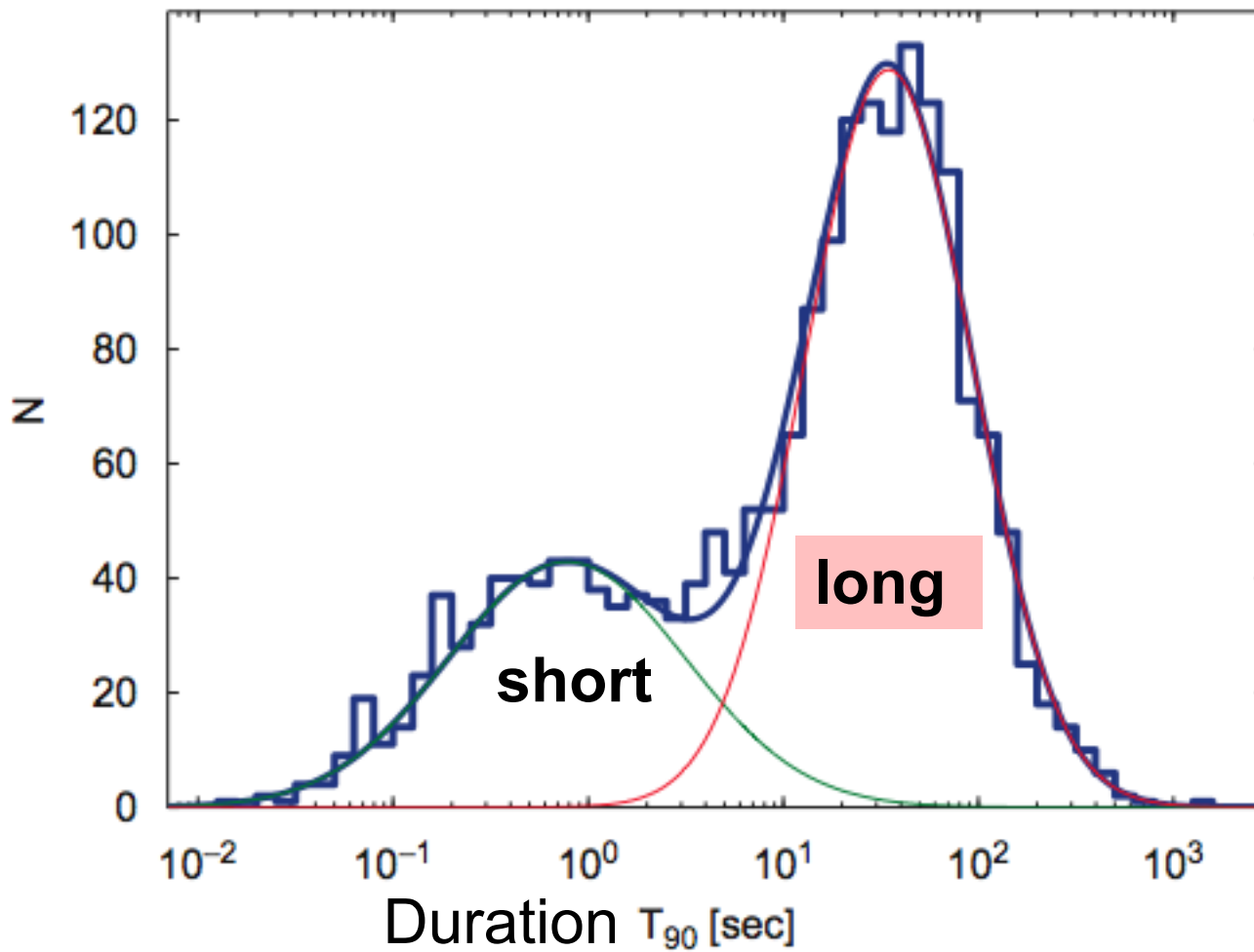
2. Long GRBs

Brian Metzger, Columbia University



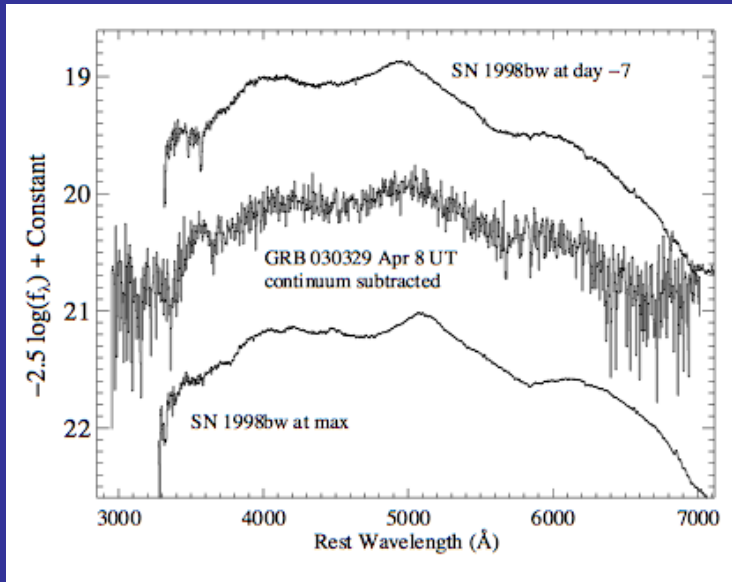
Gamma-Ray Burst Durations

E. Nakar / Physics Reports 442 (2007) 166–236

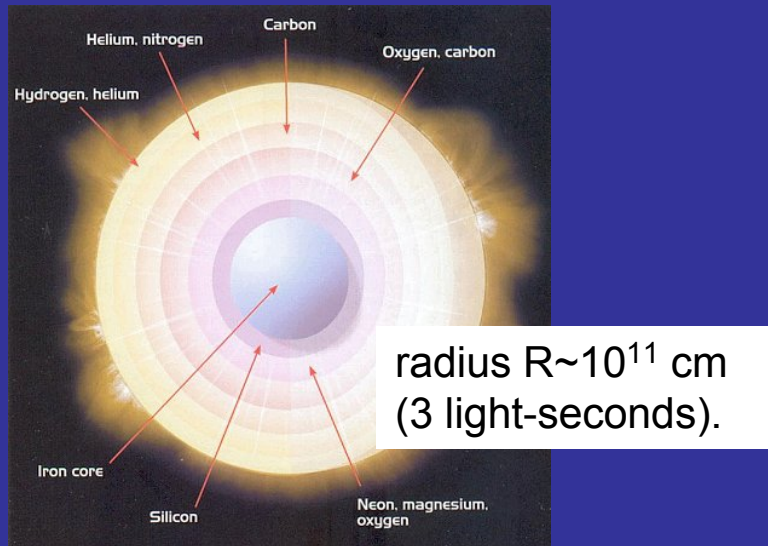


BATSE Bursts (from Nakar 2007)

GRB 030329 and the Supernova Connection

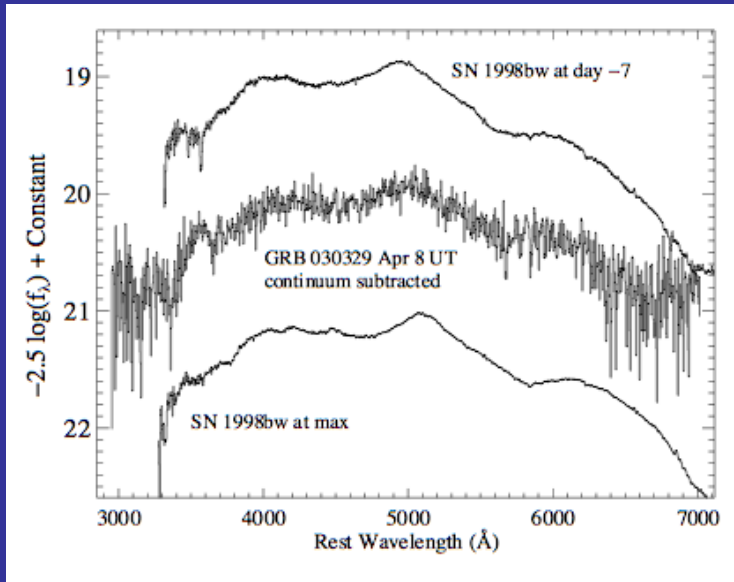


Exploding “Wolf-Rayet” Star

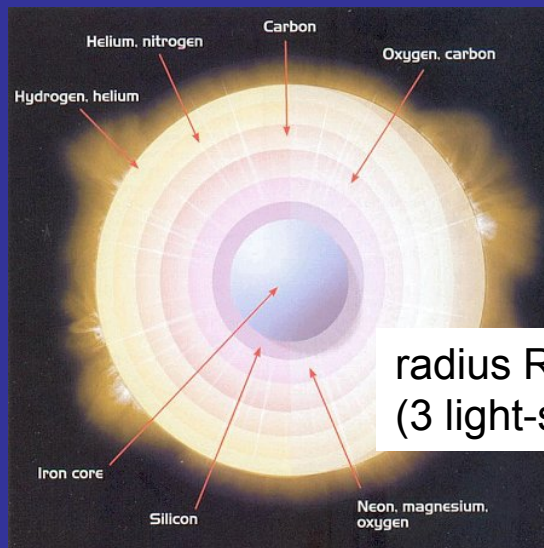


GRB 030329 and the Supernova Connection

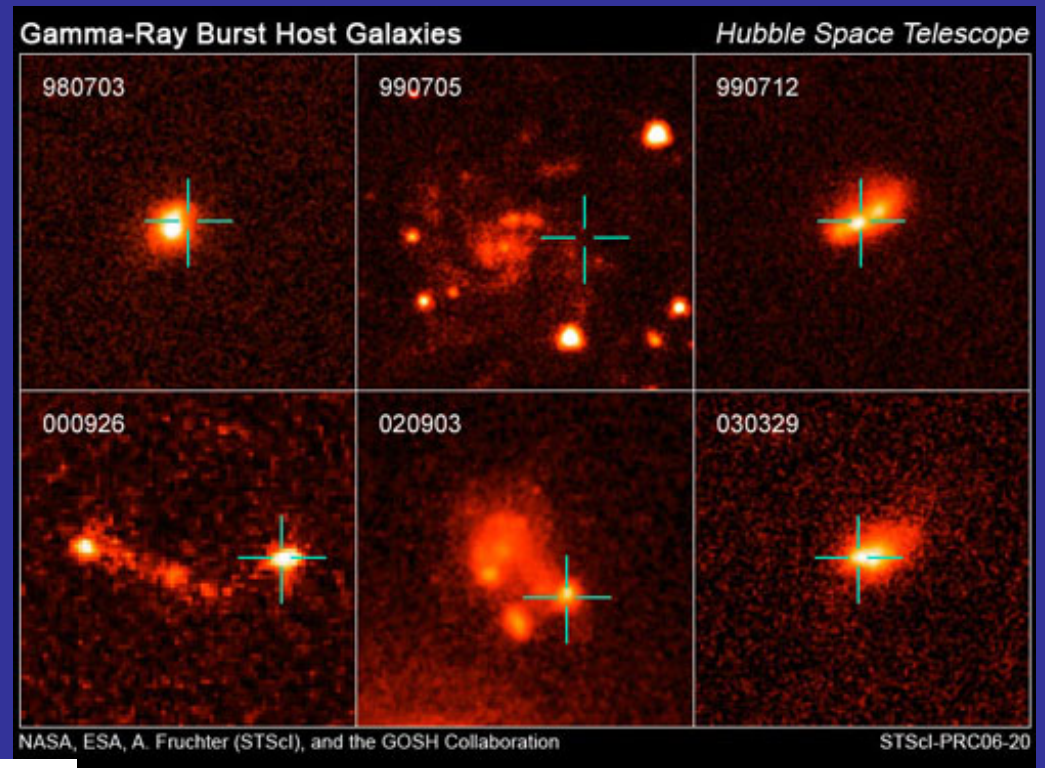
⇒ Long GRBs come from the deaths of massive Stars



Exploding “Wolf-Rayet” Star



radius $R \sim 10^{11}$ cm
(3 light-seconds).



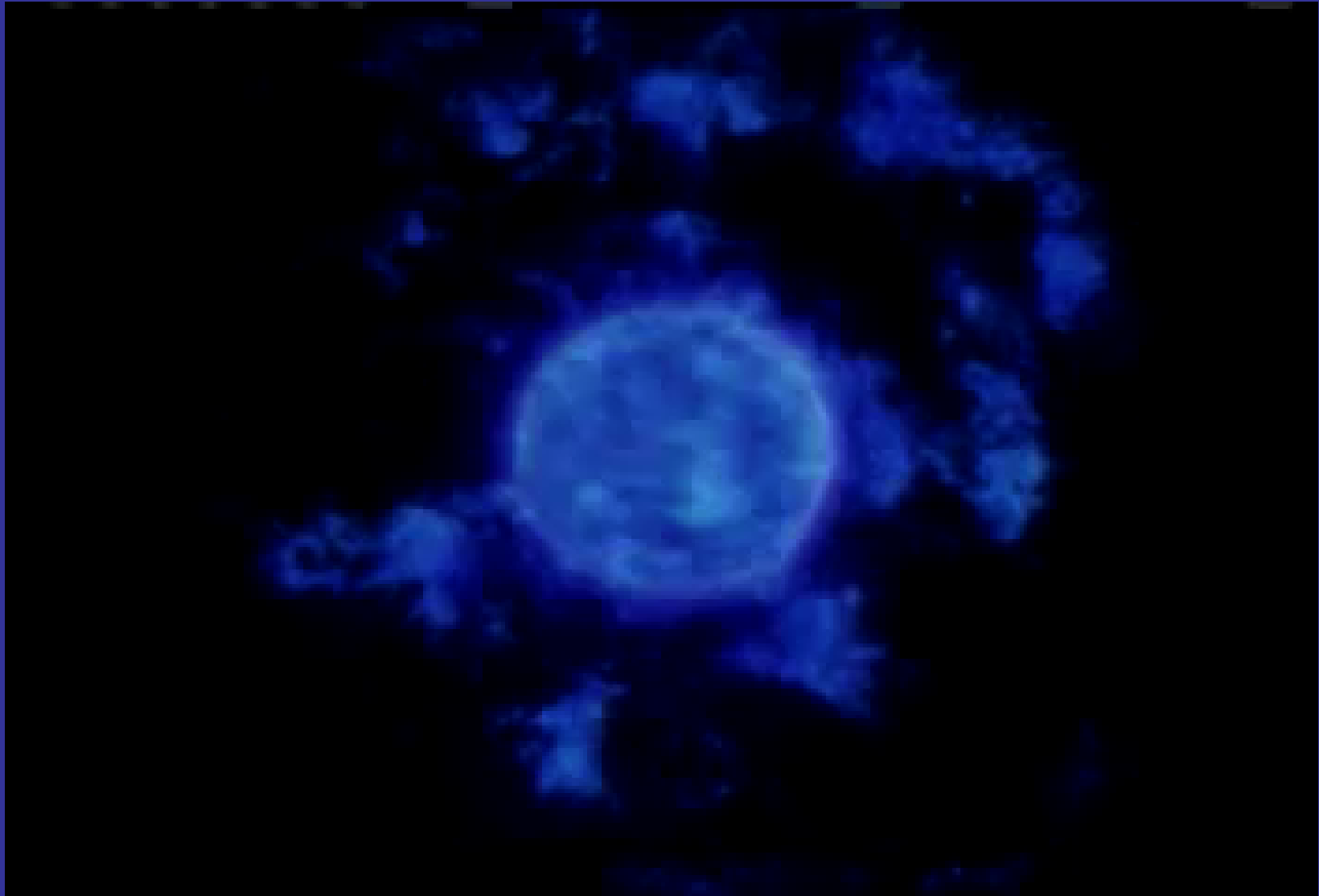
Gamma-Ray Burst Galaxies
(courtesy A. Fruchter)

GRB/Supernovae Rates and Energetics

	GRBs	SUPERNOVAE
<i>UNIVERSE-WIDE RATE</i>	100 - >1000/day	100000/day (all types) 1000-10000/day (Ic)
<i>RATE PER GALAXY</i>	1/10 ⁵ years	1/50-100 years
<i>ENERGY</i>	10 ⁵¹⁻⁵² erg	10 ⁵¹⁻⁵² erg

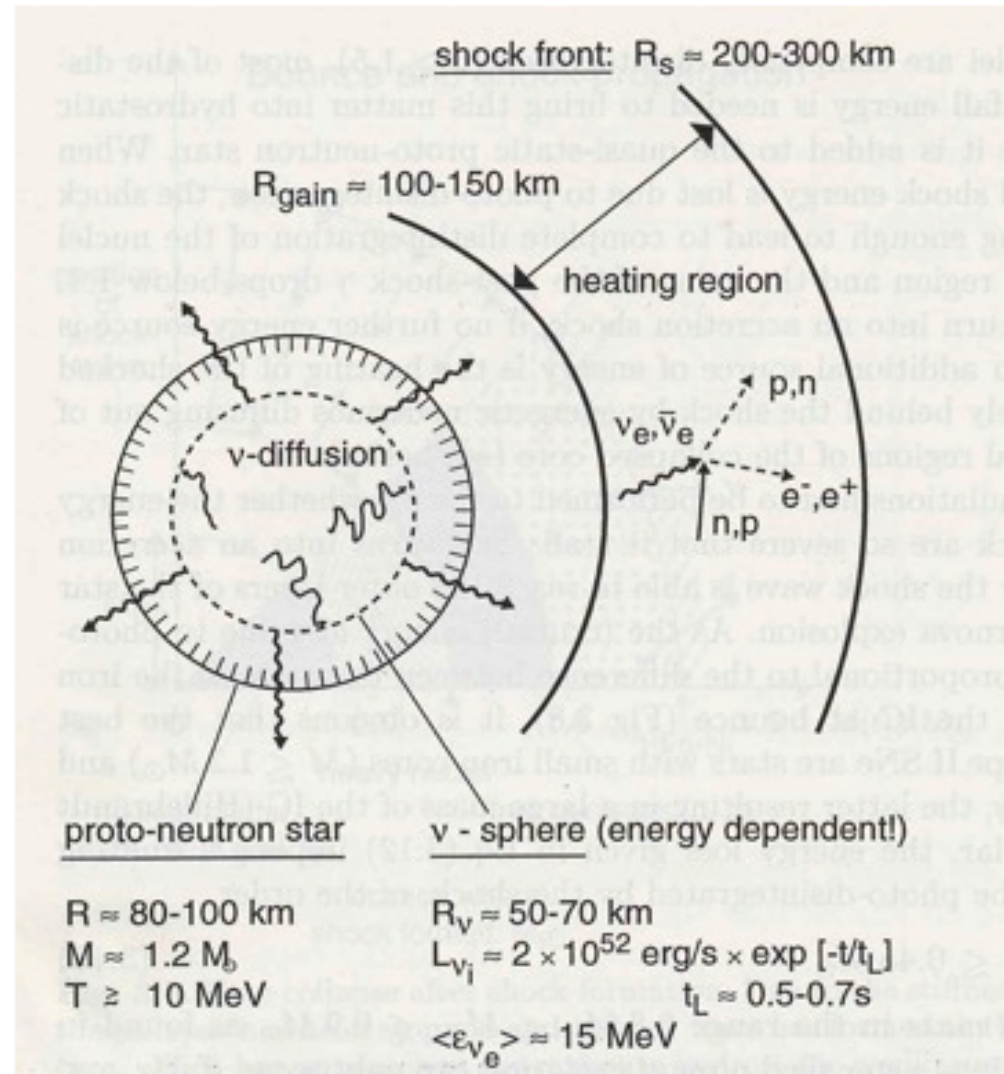
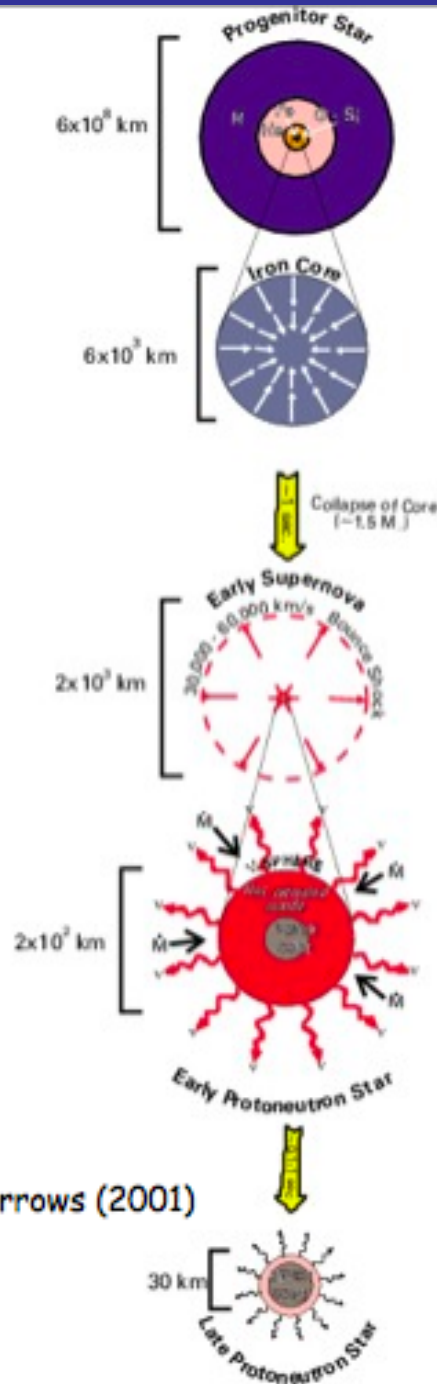
The 'Collapsar' Model (Woosley 1993)

“GRBs are powered by accretion onto a new formed black hole”



"Delayed" SN Explosion

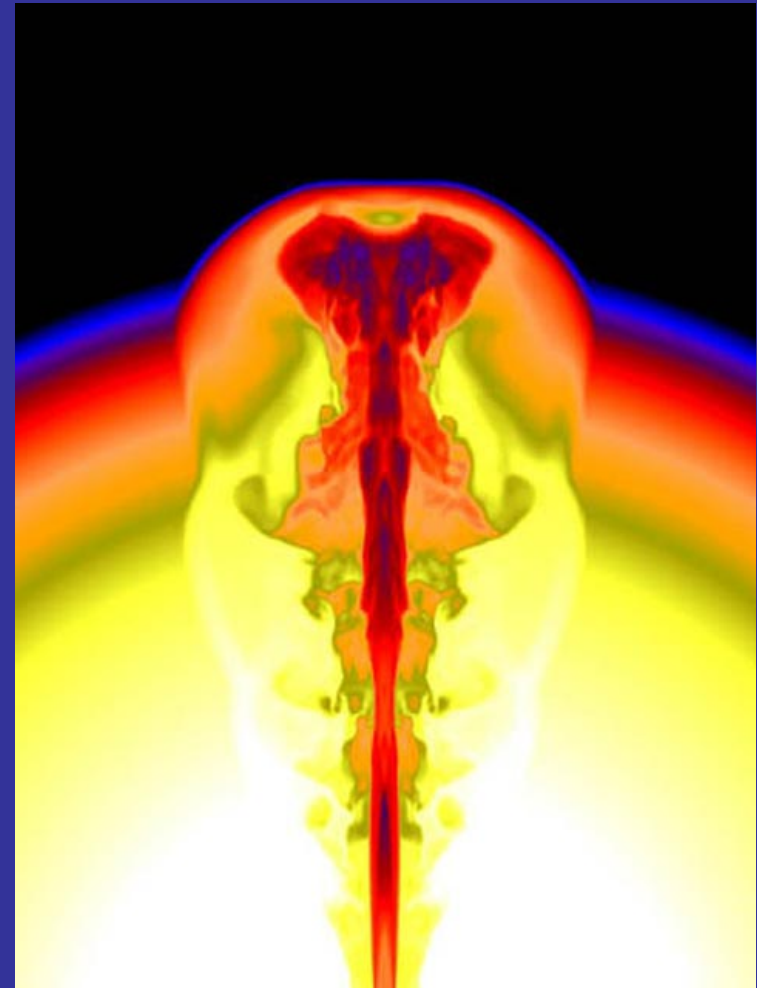
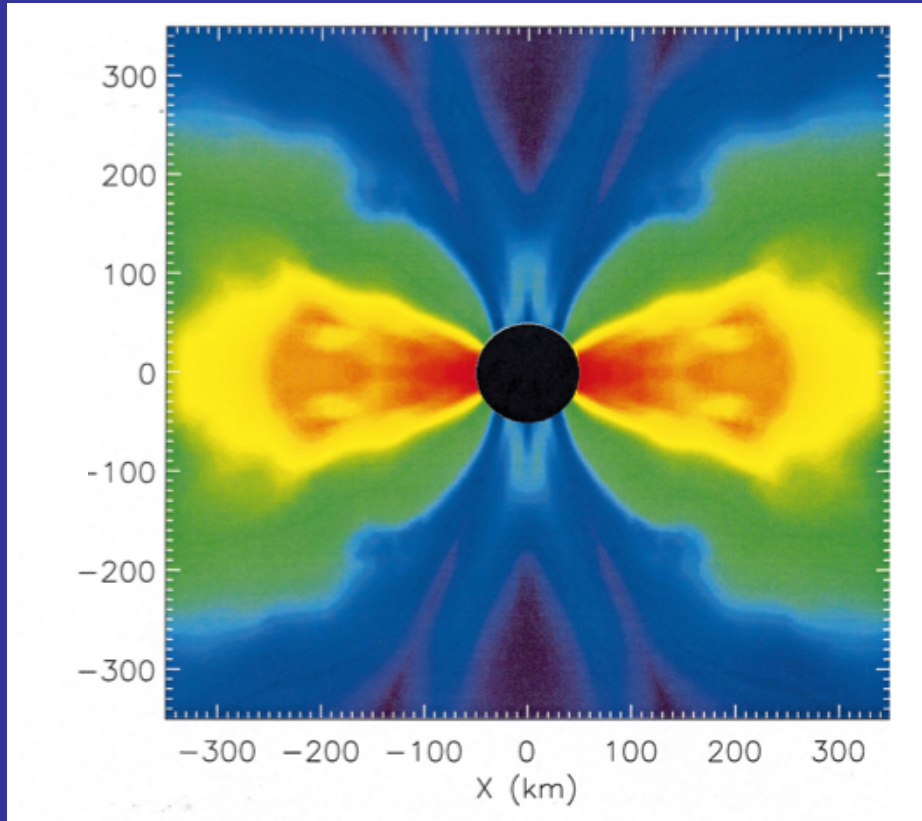
Accretion vs. Neutrino heating



Black Hole Model

(Woosley 93; MacFadyen & Woosley 1999)

MacFadyen & Woosley 1999



Zhang, Woosley & Heger 2004

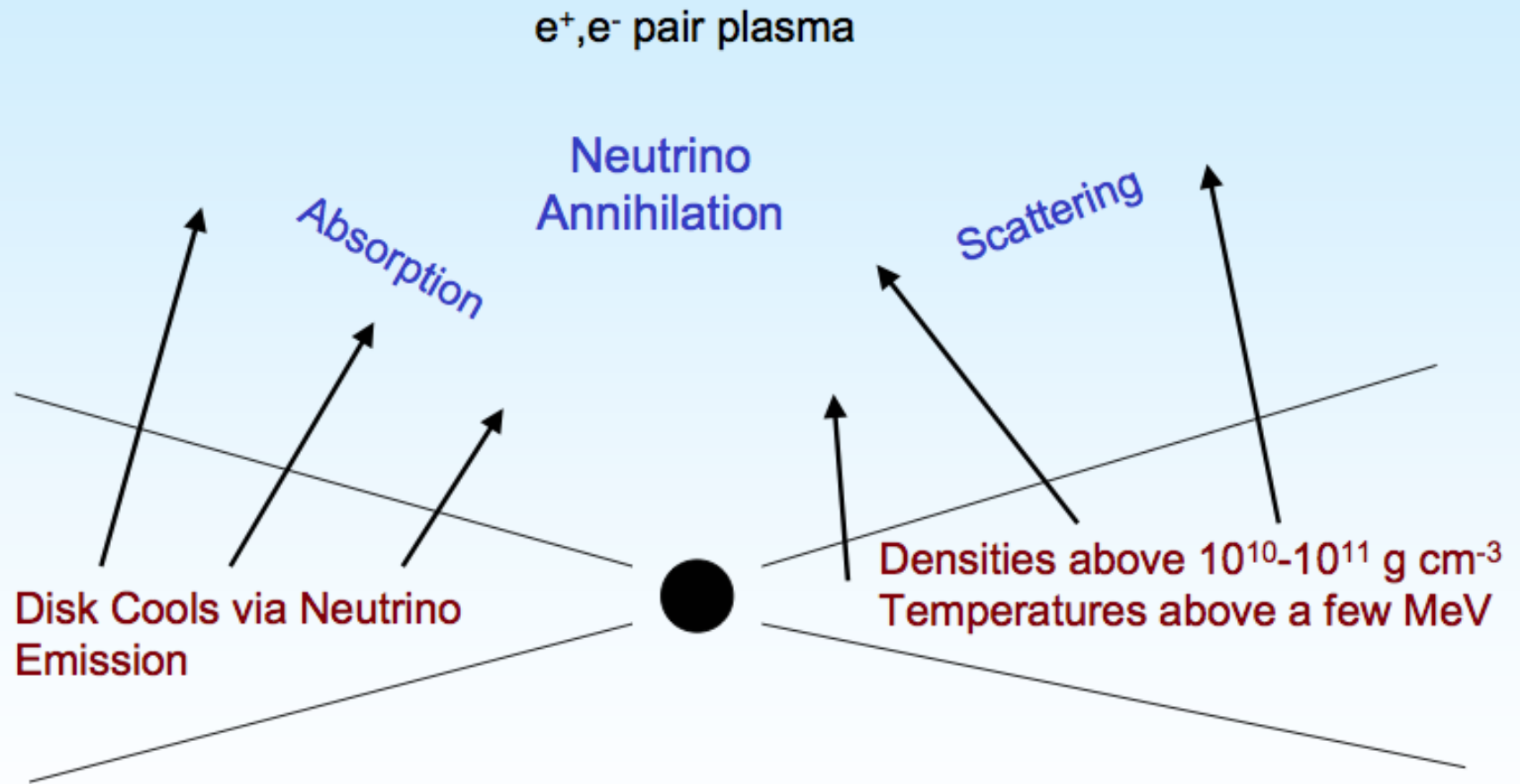
• Energy - **Accretion Power**

$$E \sim \epsilon_j M_* c^2 \sim 10^{51} \text{ ergs} \left(\frac{\epsilon_j}{10^{-3}} \right) \left(\frac{M_*}{M_\odot} \right)$$

• Duration - **Collapse Time of Star**

$$t_{ff} \sim \left(\frac{3\pi}{32G\rho_*} \right)^{1/2} \sim 100 \text{ s}$$

Neutrino-Driven Jets

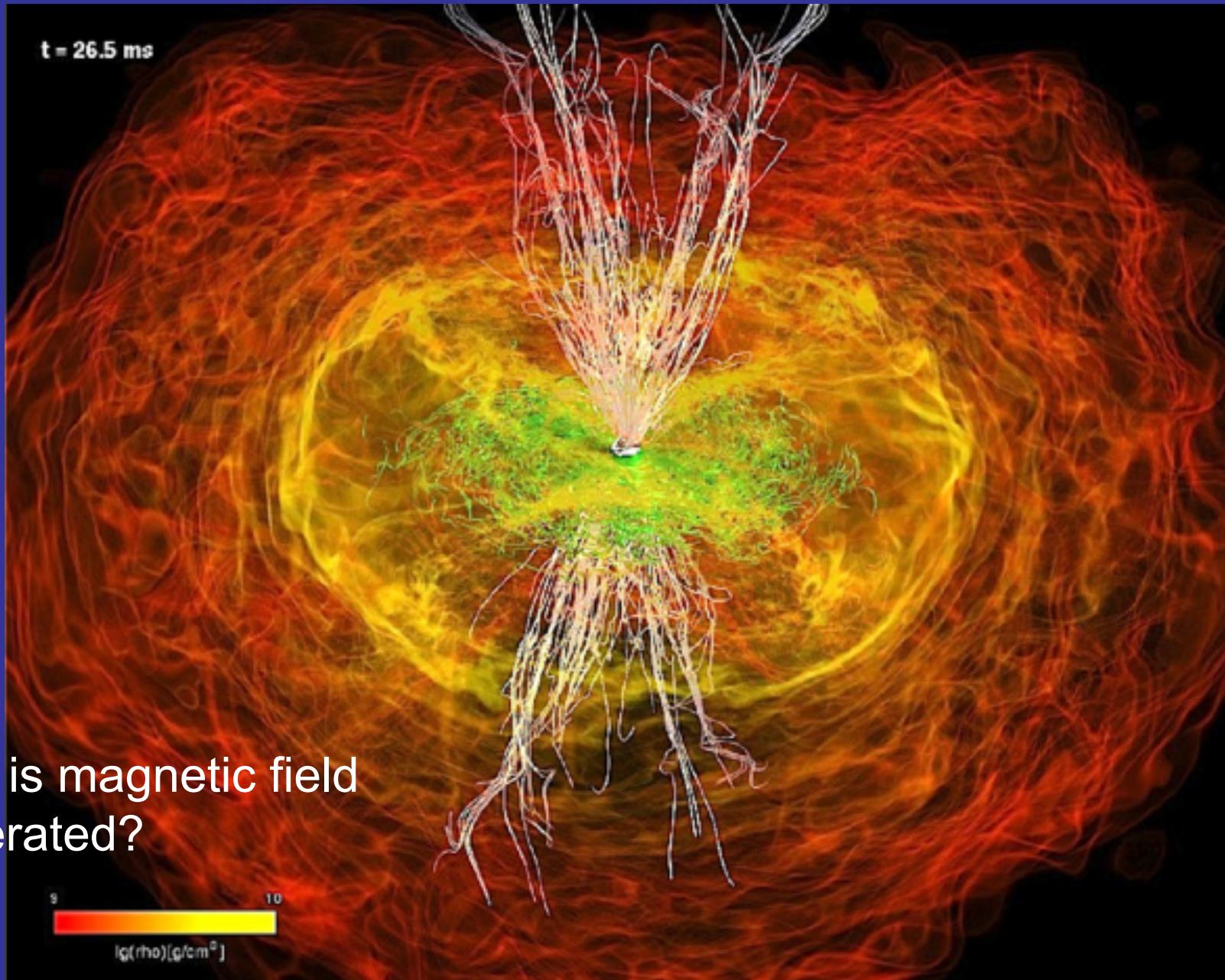


Neutrinos from accretion disk deposit their energy above the disk.
This deposition can drive an explosion.

Generally inefficient: $L_j < 10^{-3} \dot{M} c^2$

MHD Powered Jets (e.g. Blandford & Znajek 1978)

Rezzolla et al. 2010



How is magnetic field generated?

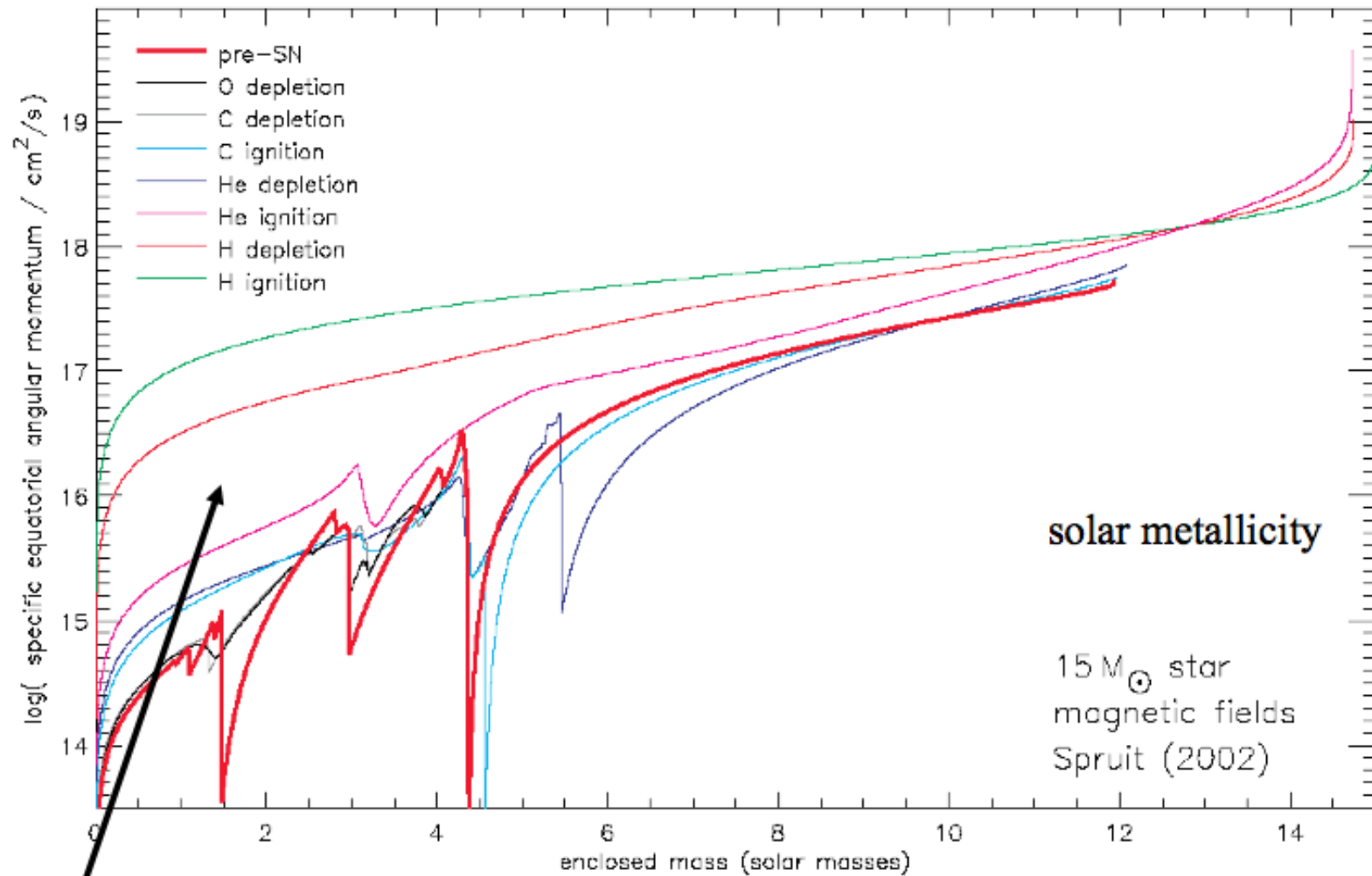
For the last stable orbit around a black hole in the collapsar model (i.e., the minimum j to make a disk)

$$j_{LSO} = 2\sqrt{3} GM / c = 4.6 \times 10^{16} M_{BH} / 3 M_{\odot} \text{ cm}^2 \text{ s}^{-1} \quad \text{non-rotating}$$

$$j_{LSO} = 2 / \sqrt{3} GM / c = 1.5 \times 10^{16} M_{BH} / 3 M_{\odot} \text{ cm}^2 \text{ s}^{-1} \quad \text{Kerr } a = 1$$

credit: Stan Woosley

It is somewhat easier to produce a magnetar model!



Much of the spin down occurs as the star evolves from H depletion to He ignition, i.e. forming a red supergiant.

Heger, Woosley, & Spruit (2004)

Need iron core rotation at death to correspond to a pulsar of < 5 ms period if rotation and B-fields are to matter to the explosion. Need a period of ~ 1 ms to make GRBs. This is much faster than observed in common pulsars.

Total rotational kinetic energy for a neutron star

$$E_{\text{rot}} \sim 2 \times 10^{52} (1 \text{ ms}/P)^2 (R/10 \text{ km})^2 \text{ erg}$$

$$j = R^2 \Omega \sim 6.3 \times 10^{15} (1 \text{ ms}/P) (R/10 \text{ km})^2 \text{ cm}^2 \text{ s}^{-1} \text{ at } M \approx 1.4 M_{\odot}$$

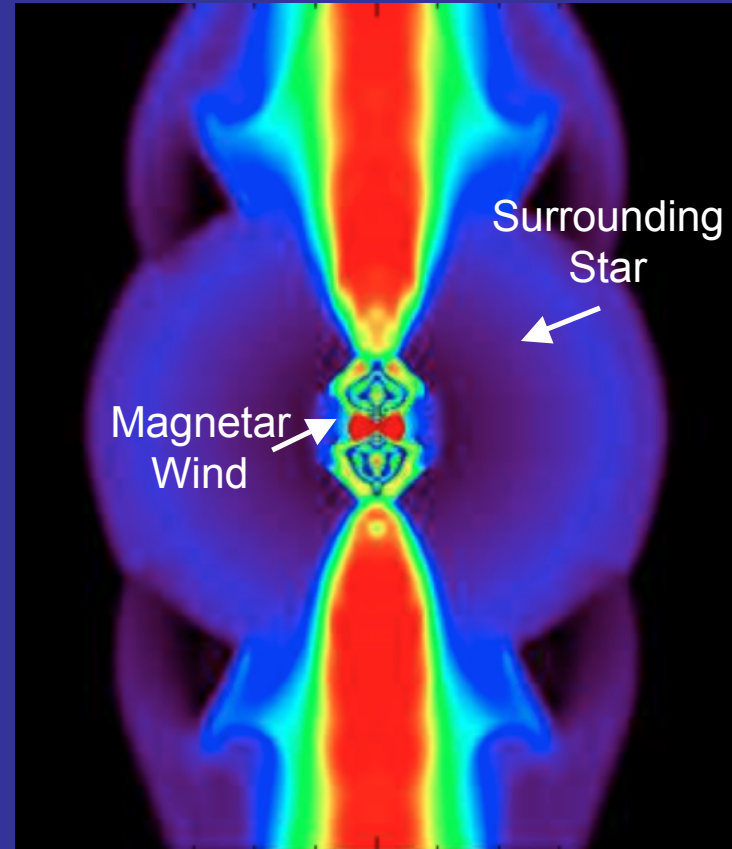
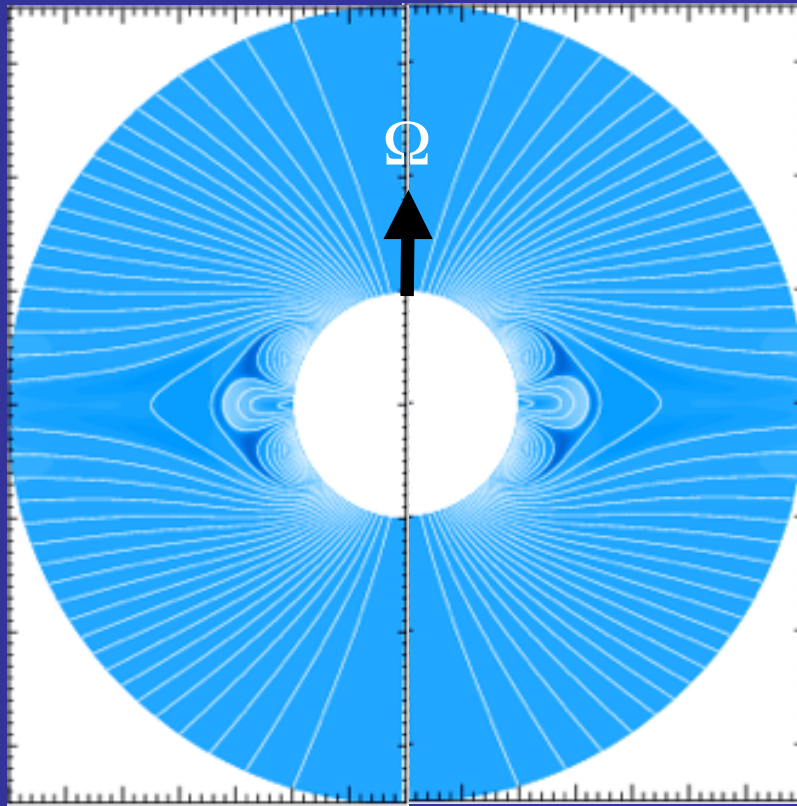
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It is somewhat easier to produce a magnetar model!

Millisecond Magnetar Model (e.g. Usov 1992; Metzger et al. 2011)



Bucciantini, Metzger et al. 2011

• Energy - **Rotation**

$$E_{\text{rot}} \sim \frac{1}{2} I \Omega^2 \sim 3 \times 10^{52} \text{ ergs} \left(\frac{P}{1 \text{ ms}} \right)^{-2}$$

• Luminosity - **Dipole Radiation**

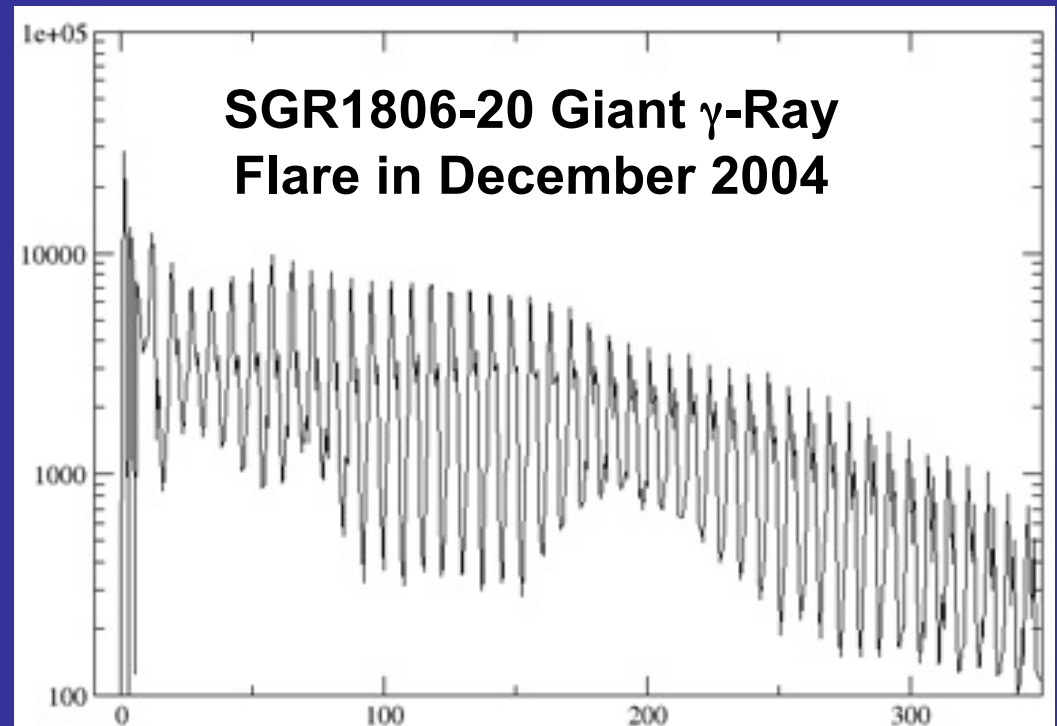
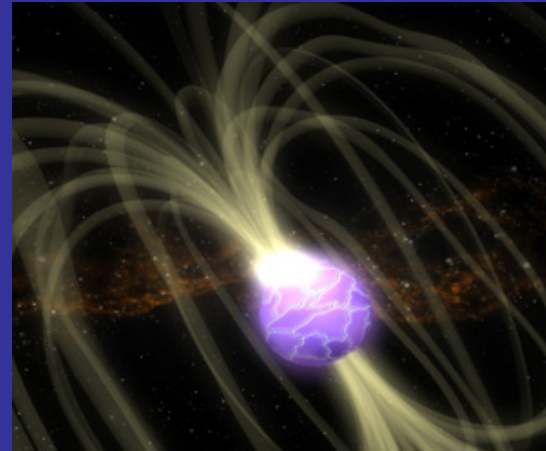
$$L_{\text{sd}} = \frac{\mu^2 \Omega^4}{c^3} \approx 6 \times 10^{49} \left(\frac{P}{1 \text{ ms}} \right)^{-4} \left(\frac{B_{\text{dip}}}{10^{15} \text{ G}} \right)^2 \text{ erg s}^{-1}$$

• Duration - **Spin-Down Time**

$$\tau_{\text{sd}} = \frac{E_{\text{rot}}}{L_{\text{sd}}} \approx 10 \left(\frac{P_0}{1 \text{ ms}} \right)^2 \left(\frac{B_{\text{dip}}}{10^{15} \text{ G}} \right)^{-2} \text{ min}$$

Magnetars: Super-Magnetized Neutron Stars

- **Surface Magnetic Field**
 10^{14} - 10^{15} G (would erase your credit card at distance of Sun).
- Observationally classified as “soft gamma-ray repeaters” and “anomalous X-ray pulsars”
- **“Giant Flares”** every ~10-100 years.
- 12 in Milky Way
- **Age:** 10^3 - 10^4 yrs
- rotation period $P \sim$ secs



What Produces Magnetar Fields?

All neutron stars form as hot, differentially-rotating 'proto-neutron stars'

rotational energy:

$$E_{\text{rot}} = \frac{1}{2} I \Omega^2 \sim 3 \times 10^{52} \left(\frac{P}{1 \text{ ms}} \right)^{-2} \text{ ergs}$$

$$\Delta E_{\text{rot}} = \frac{B^2}{8\pi} \times \frac{4\pi}{3} R_{\text{ns}}^3 \Rightarrow B_{\text{eq}} \sim 10^{17} \left(\frac{\Delta\Omega}{\Omega/2} \right)^2 \left(\frac{P}{1 \text{ ms}} \right)^{-2} \text{ G}$$

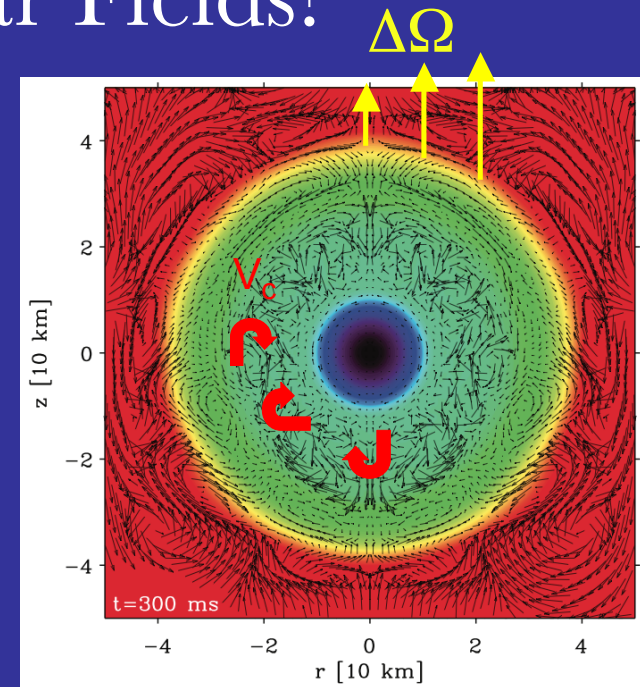
Field amplification:

- Shear instabilities (talk by Zrake)
- Magneto-rotational instability
- α - Ω dynamo (Thompson & Duncan 1993)

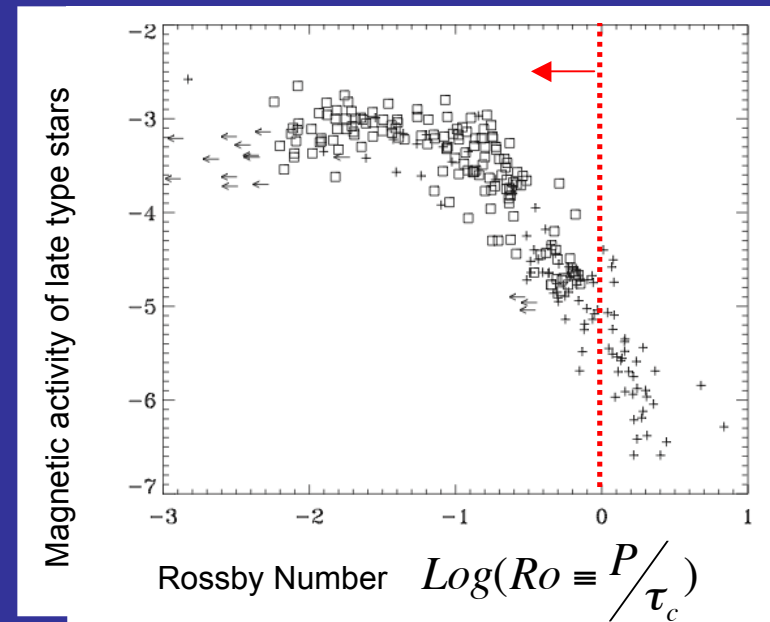
$$L_v \sim 4\pi R^2 \rho V_c^3, \quad l_p \sim 0.1 R_{\text{NS}}$$

$$\tau_c \sim \frac{l_p}{V_c} \sim 1 \text{ ms} \left(\frac{l_p}{0.1 R_{\text{NS}}} \right) \left(\frac{R_{\text{NS}}}{12 \text{ km}} \right)^{5/3} \left(\frac{\rho}{10^{14} \text{ g cm}^{-3}} \right)^{1/3} \left(\frac{L_v}{10^{52} \text{ erg s}^{-1}} \right)^{-1/3}$$

$Ro \sim 1$ for $P \sim 1 \text{ ms}$



Dessart et al. 2006



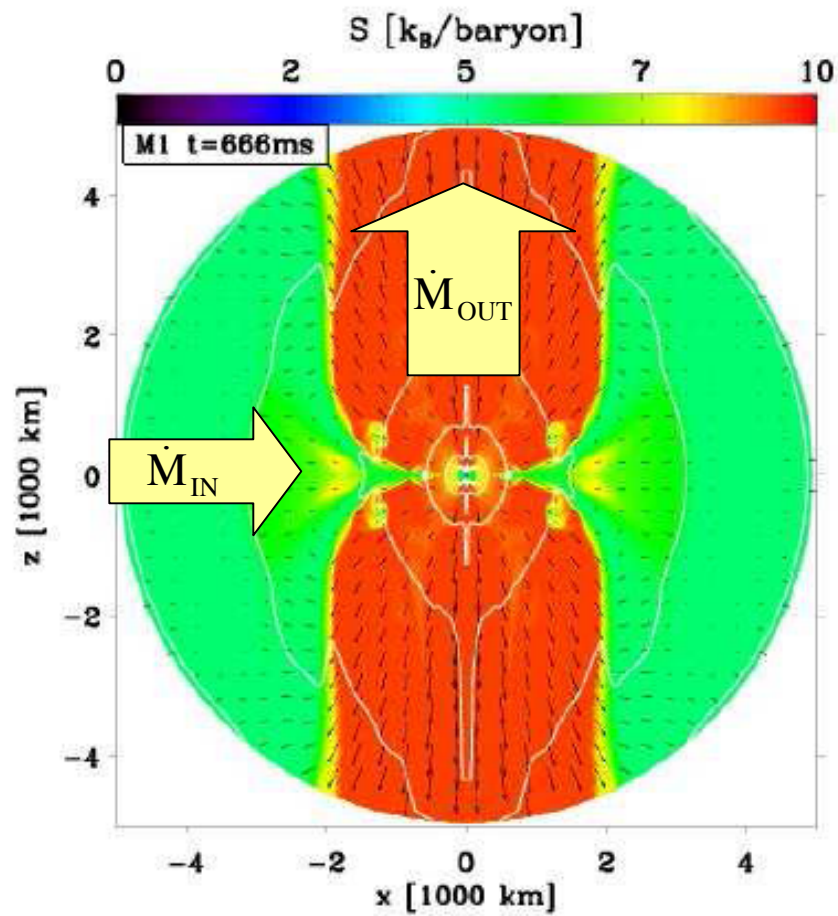
Pizzolato et al. 2003

Core Collapse with Magnetic Fields & Rotation

(e.g. LeBlanc & Wilson 1970)

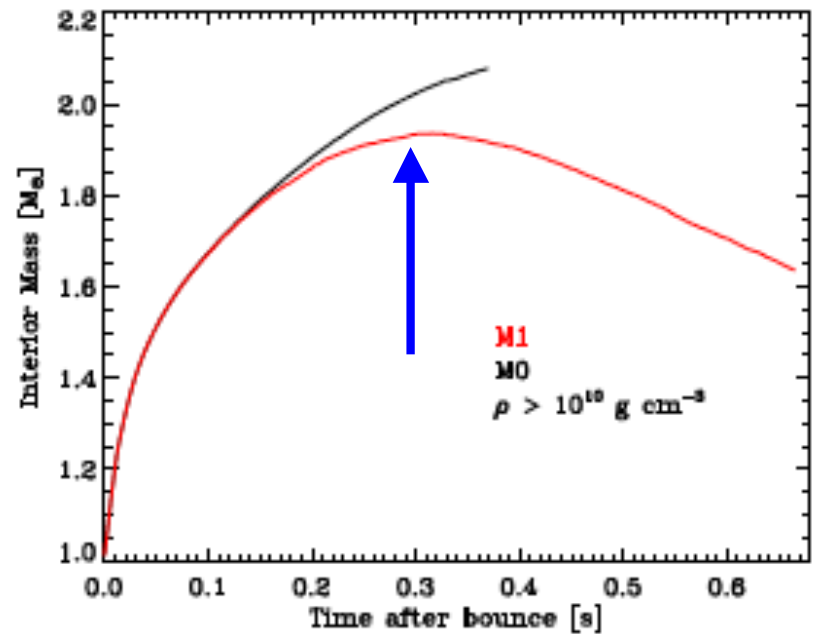
THE PROTO-NEUTRON STAR PHASE OF THE COLLAPSAR MODEL AND THE ROUTE TO LONG-SOFT GAMMA-RAY BURSTS AND HYPERNOVAE

L. DESSART¹, A. BURROWS¹, E. LIVNE², AND C.D. OTT¹



“Failed Collapsar”

Neutron Star Mass



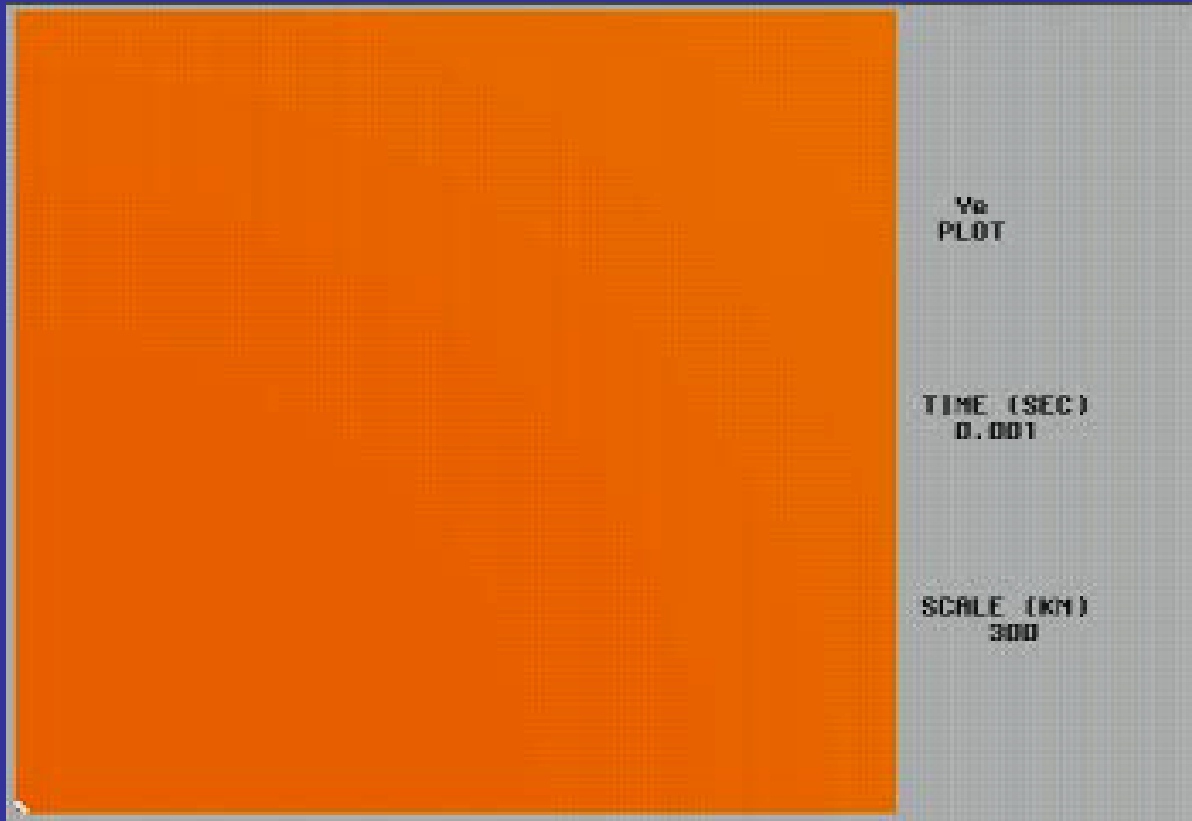
Time

Neutrino Driven Wind

Neutrinos heat proto-NS atmosphere (e.g. $\nu_e + n \Rightarrow p + e^-$)

\Rightarrow drives wind behind outgoing supernova shock (e.g. Qian & Woosley 96)

Burrows, Hayes, & Fryxell 1995



$$\dot{M} \sim 10^{-4} \left(\frac{L_\nu}{10^{52} \text{ erg s}^{-1}} \right)^{5/3} \left(\frac{\epsilon_\nu}{10 \text{ MeV}} \right)^{10/3} M_\odot \text{ s}^{-1} \Rightarrow \text{crucial to baryon loading}$$

Neutrino Driven Wind

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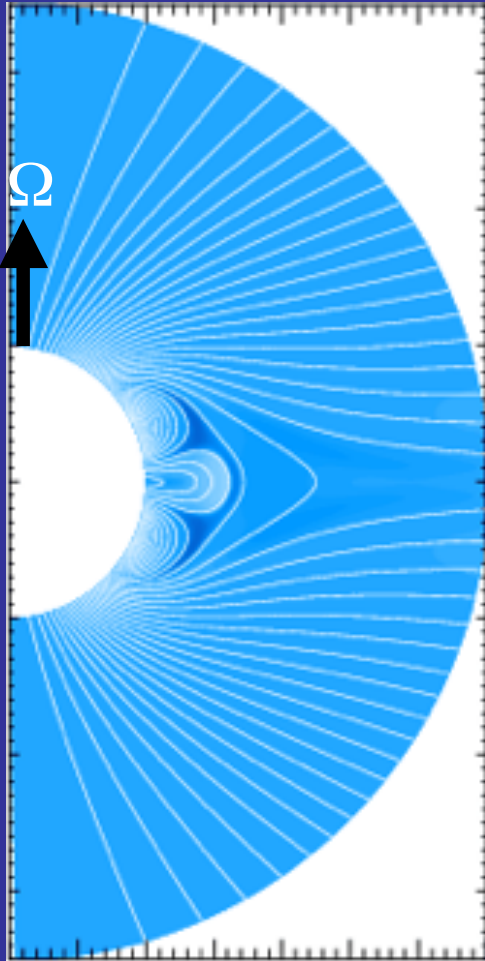
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$$\dot{M} \sim 10^{-4} \left(\frac{L_\nu}{10^{52} \text{ erg s}^{-1}} \right)^{5/3} \left(\frac{\epsilon_\nu}{10 \text{ MeV}} \right)^{10/3} M_\odot \text{ s}^{-1} \Rightarrow \text{crucial to baryon loading}$$

Effects of Strong Magnetic Fields

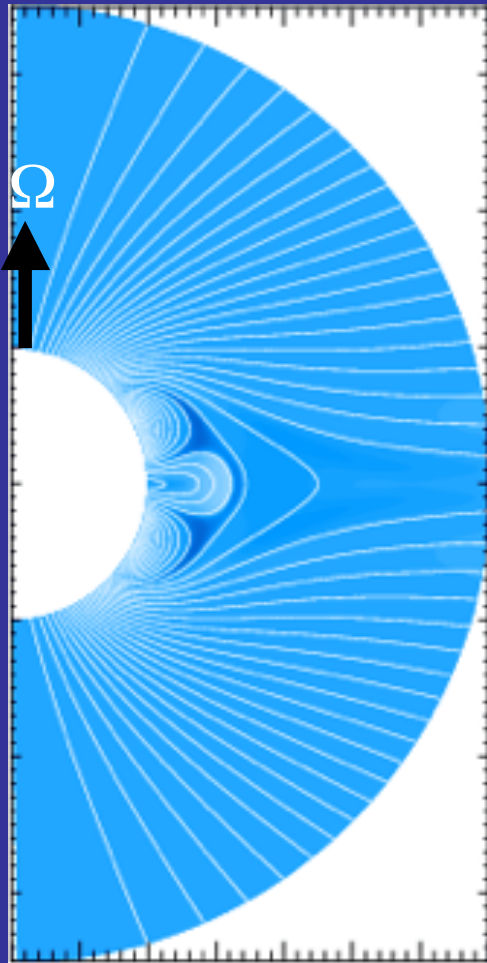
“Helmet - Streamer”



- Microphysics (EOS, ν Heating & Cooling)
 - Important for $B \geq 10^{16}$ G (Duan & Qian 2005)

Effects of Strong Magnetic Fields

“Helmet - Streamer”



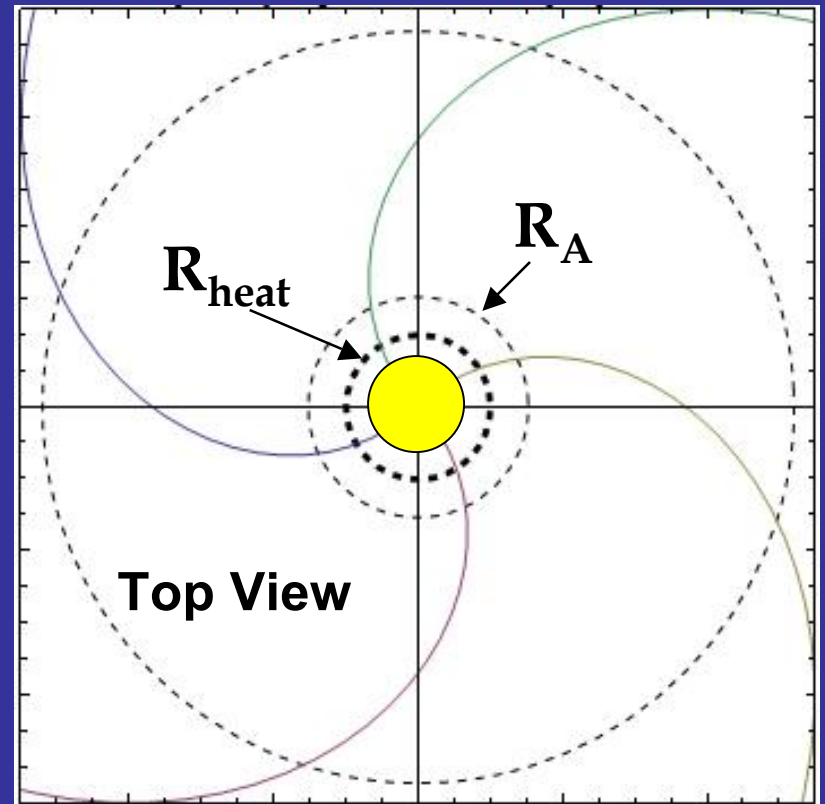
- **Microphysics (EOS, ν Heating & Cooling)**
 - Important for $B \geq 10^{16}$ G (Duan & Qian 2005)
- **Magneto-Centrifugal Slingshotting**
(Weber & Davis 1967; Thompson, Chang & Quataert 2004)

Outflow Co-Rotates
with Neutron Star when

$$\frac{B^2}{8\pi} > \frac{1}{2}\rho v_r^2$$

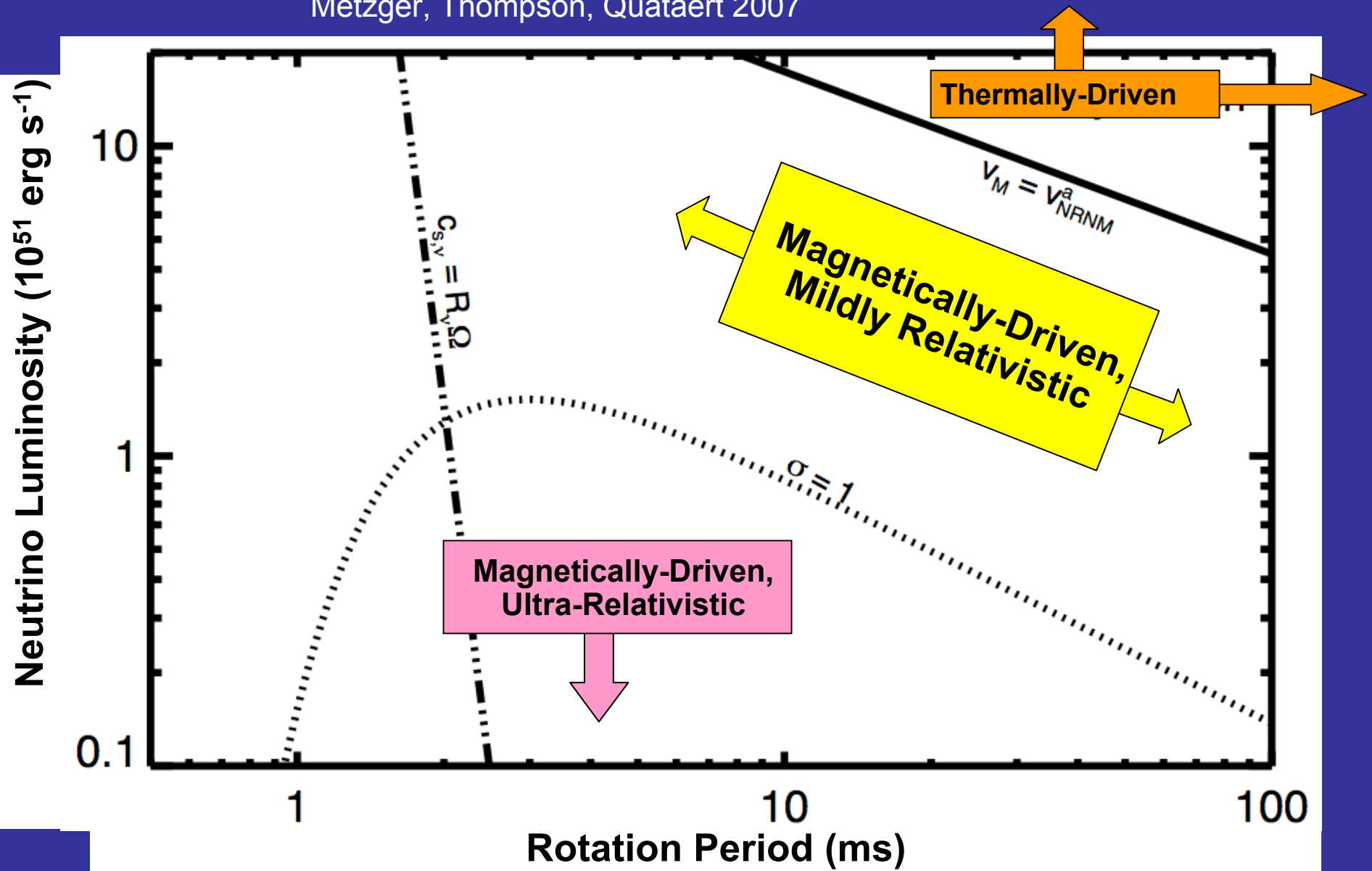
\Rightarrow

**Magneto-Centrifugal
Acceleration
 (“Beads on a Wire”)**



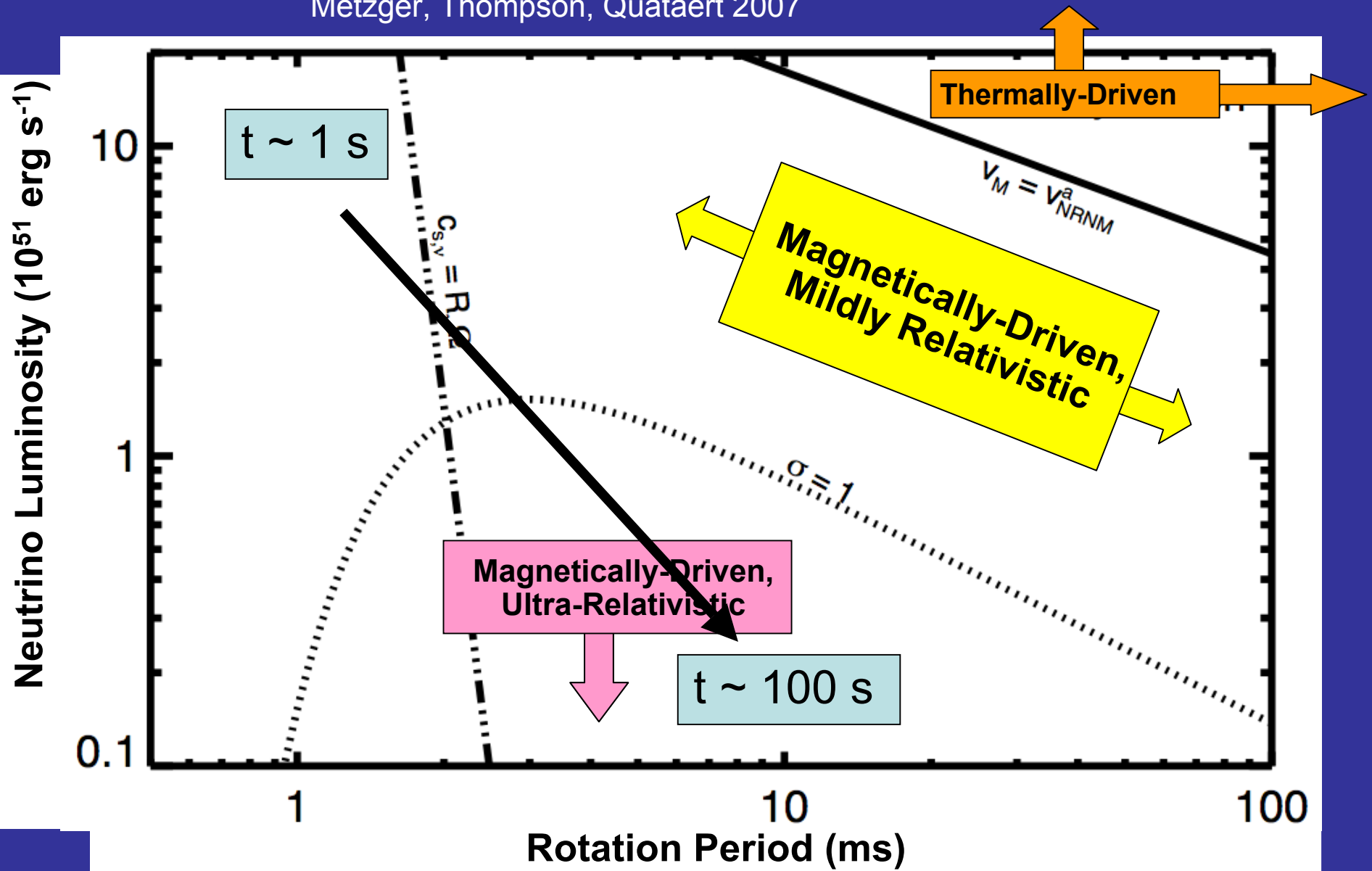
Regimes of Magnetized PNS Winds ($B = 3 \times 10^{14}$ G)

Metzger, Thompson, Quataert 2007



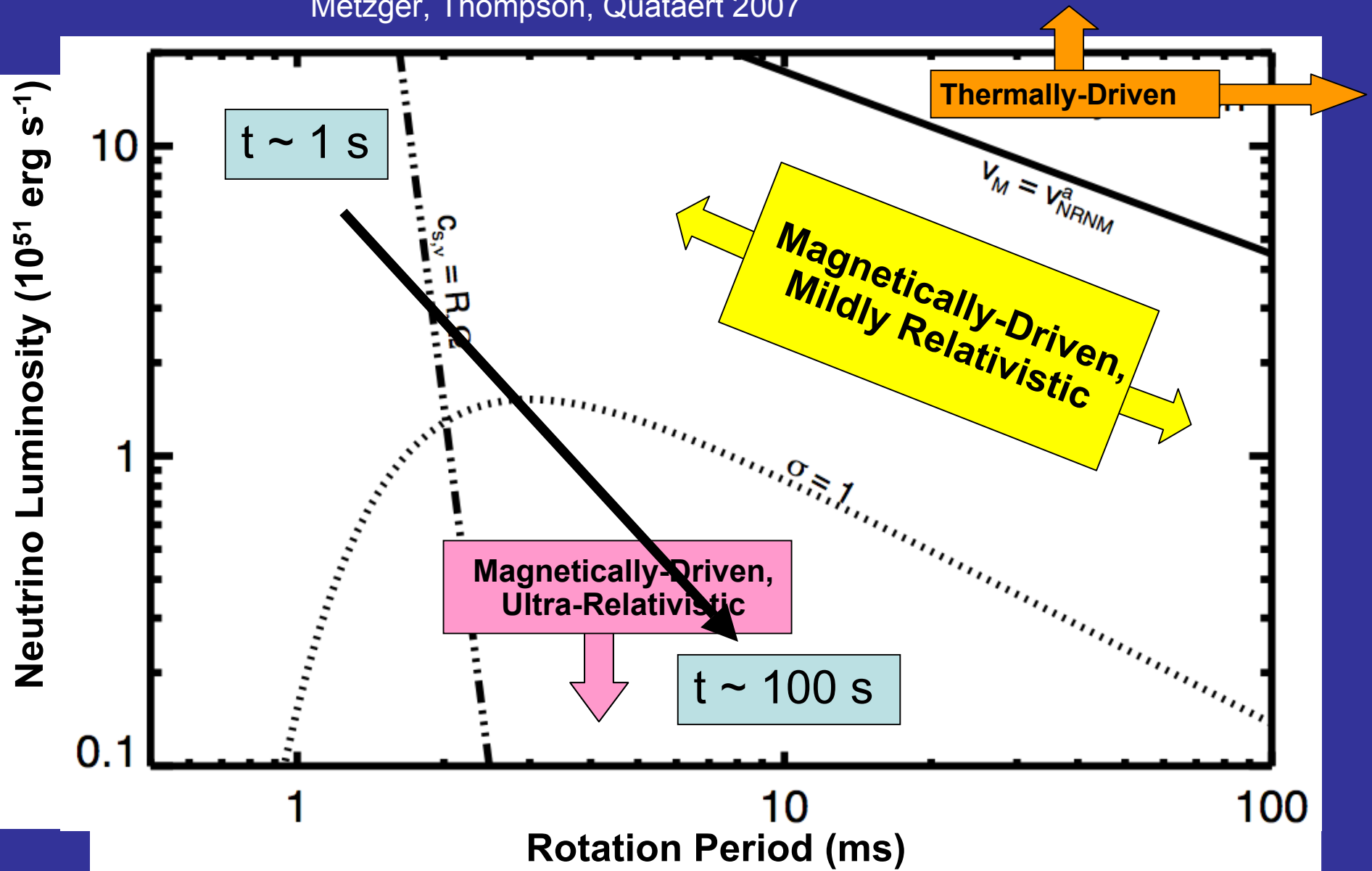
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Metzger, Thompson, Quataert 2007



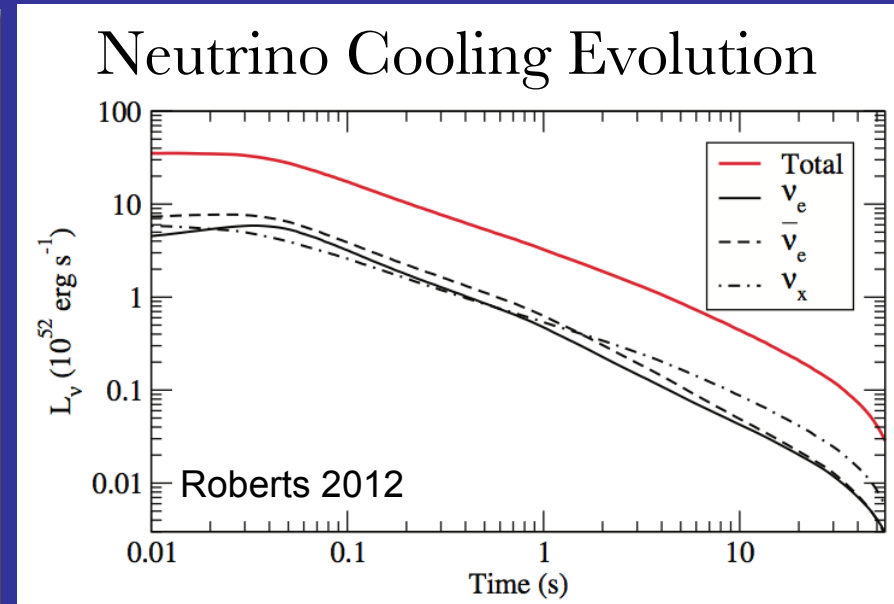
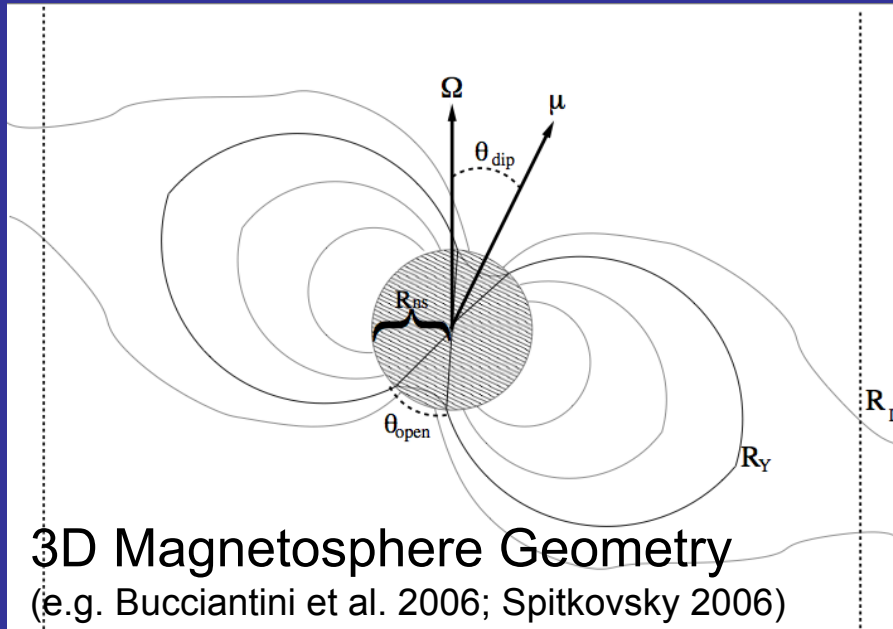
Regimes of Magnetized PNS Winds ($B = 3 \times 10^{14}$ G)

Metzger, Thompson, Quataert 2007



Evolution of Proto-Magnetar Outflows

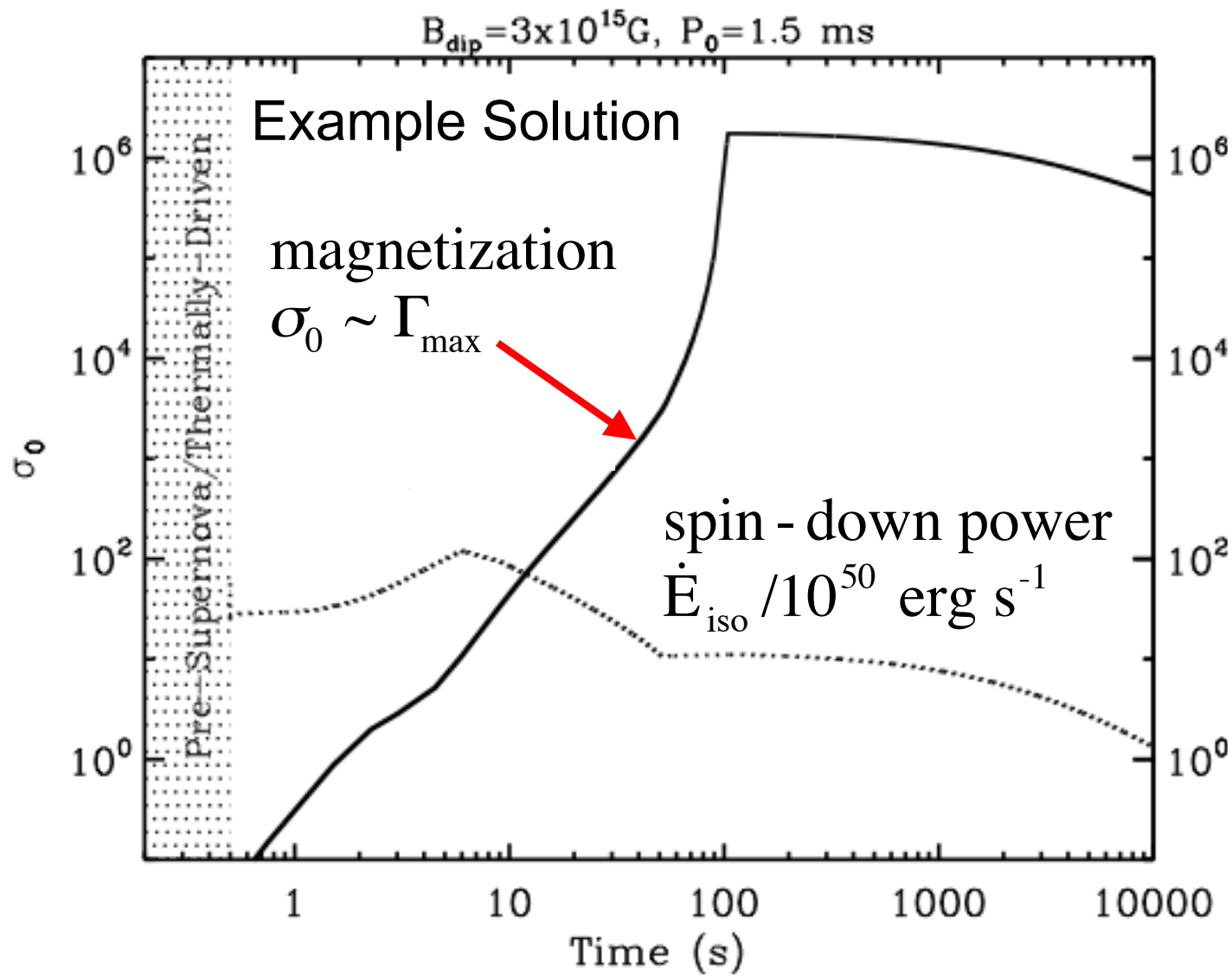
(BDM et al. 2007, 2011)



Calculate: Wind Power $\dot{E}(t)$, Mass Loss Rate $\dot{M}(t)$,
 \Rightarrow 'Magnetization' $\sigma(t) \sim \frac{\dot{E}}{\dot{M}c^2} = \Gamma_{\max}(t)$

In terms of

Initial rotation period P_0 , dipole field B_{dip} & obliquity θ_{dip}

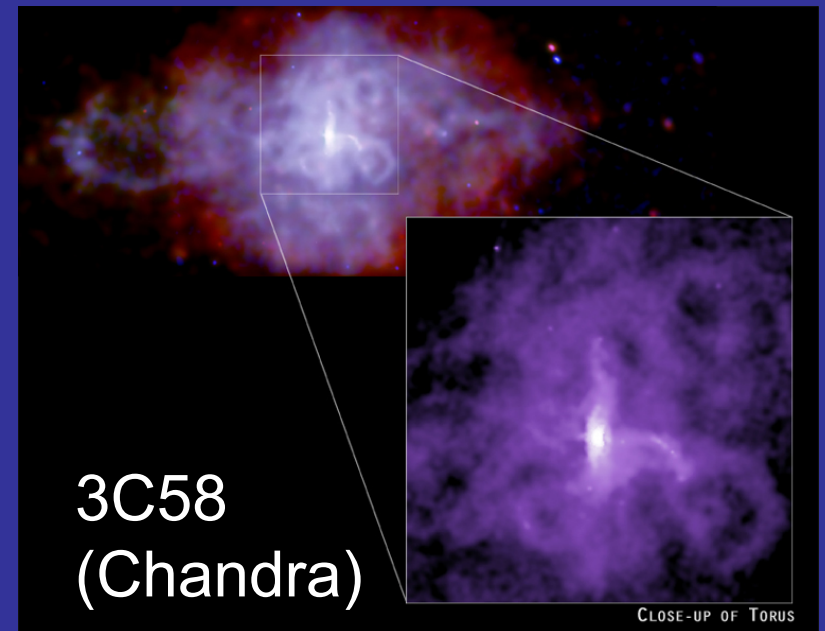
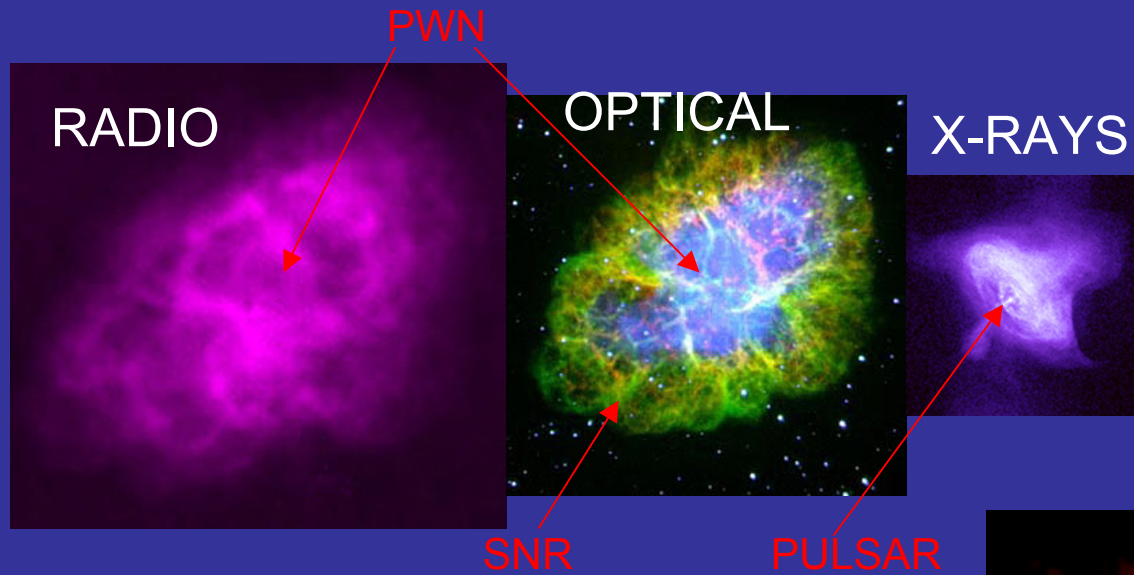


$$\sigma_0 \sim \Gamma_{\text{max}} = \frac{\dot{E}}{\dot{M}c^2} \propto \frac{B^2 \Omega^4}{L_v^{5/3} T^{10/3}}$$

increases as magnetar cools

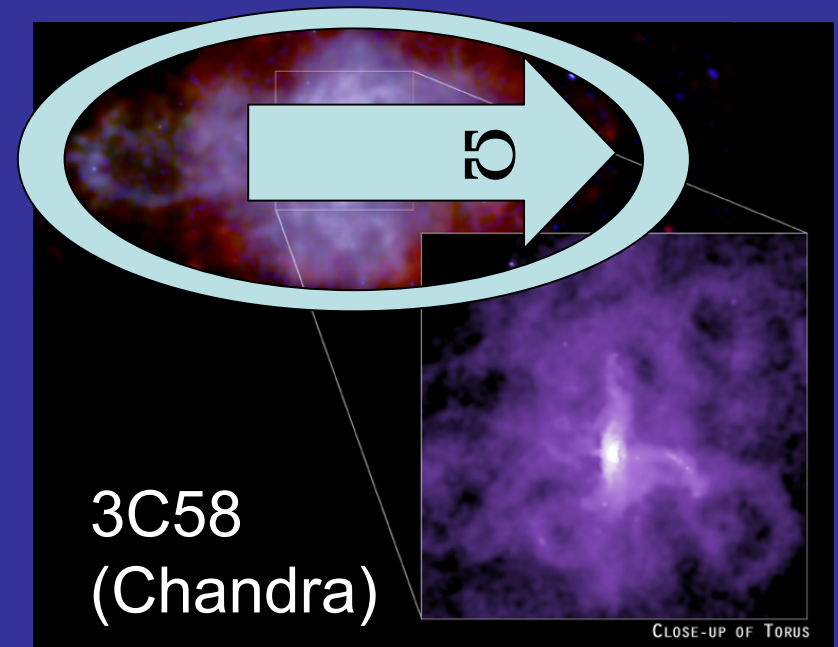
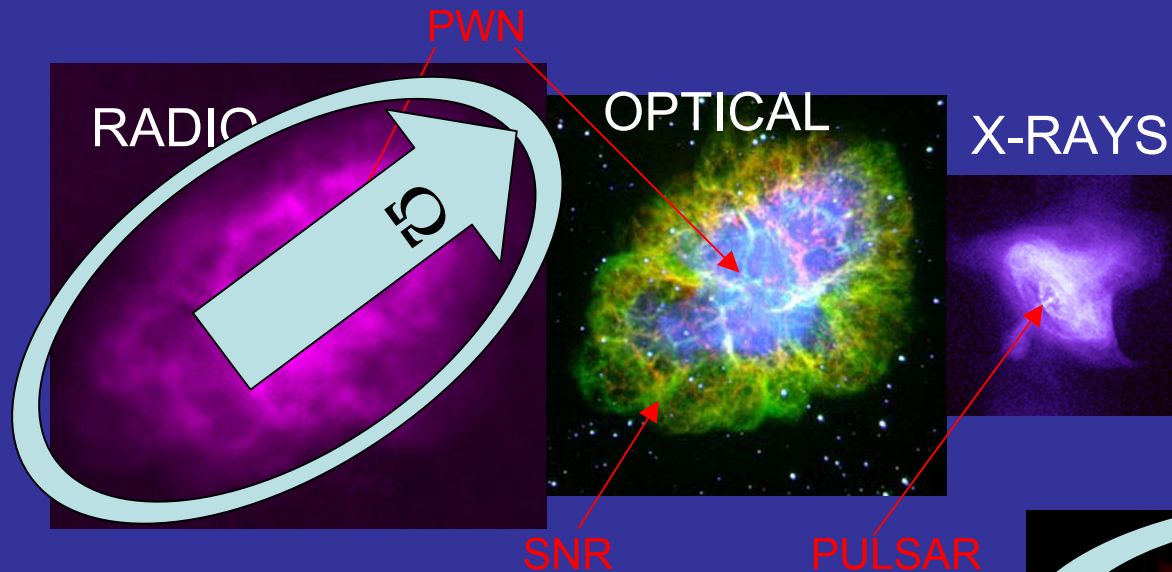
Collimation via Stellar Confinement

Multi-Wavelength Crab Nebula



Collimation via Stellar Confinement

Multi-Wavelength Crab Nebula

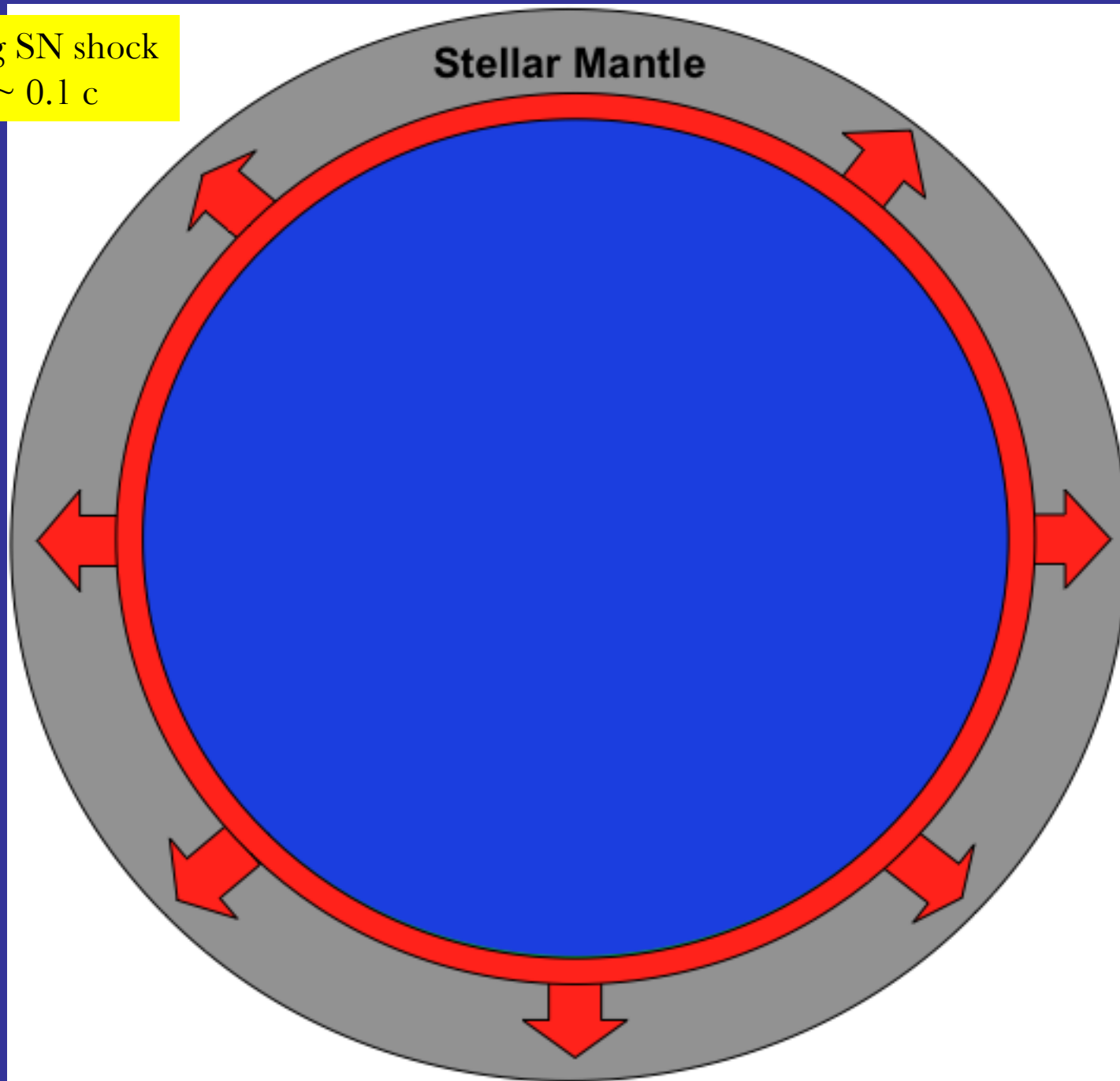


Supernova remnant elongated by **anisotropic magnetic stresses** in pulsar nebula? (Begelman & Li 1992)



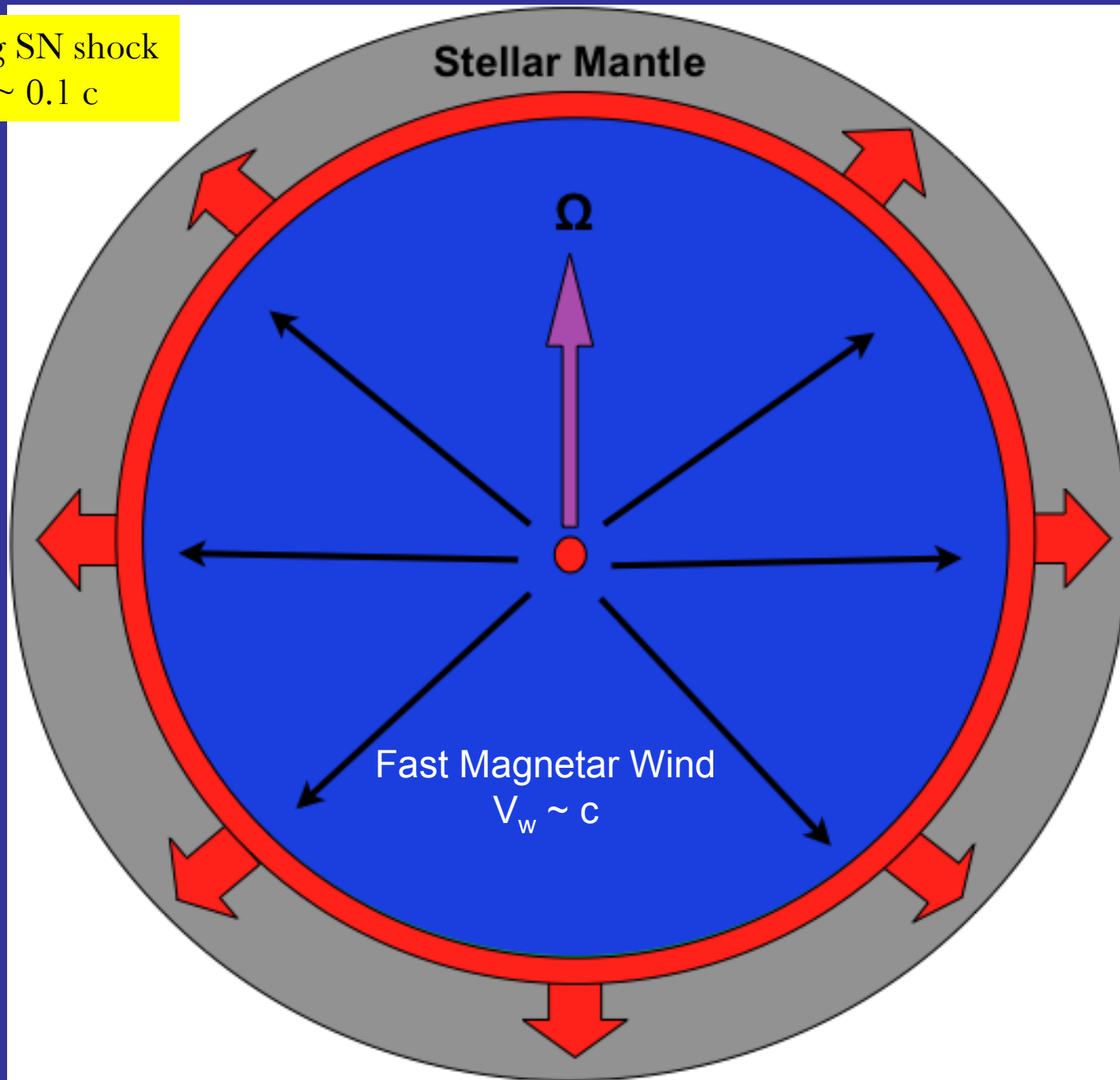
Stellar Mantle

Outgoing SN shock
 $V_{\text{SN}} \sim 0.1 c$



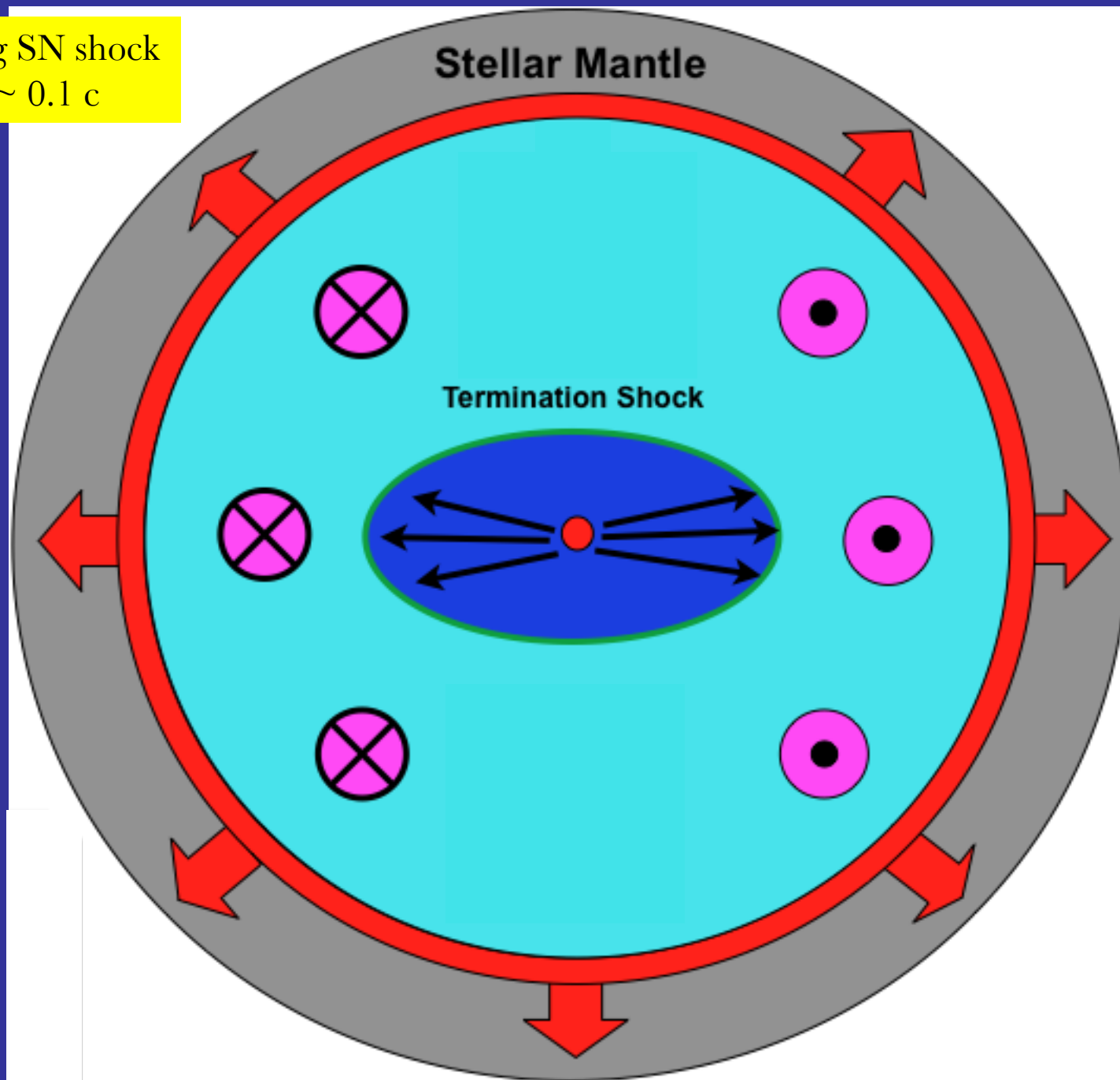
Outgoing SN shock

$$V_{\text{SN}} \sim 0.1 c$$



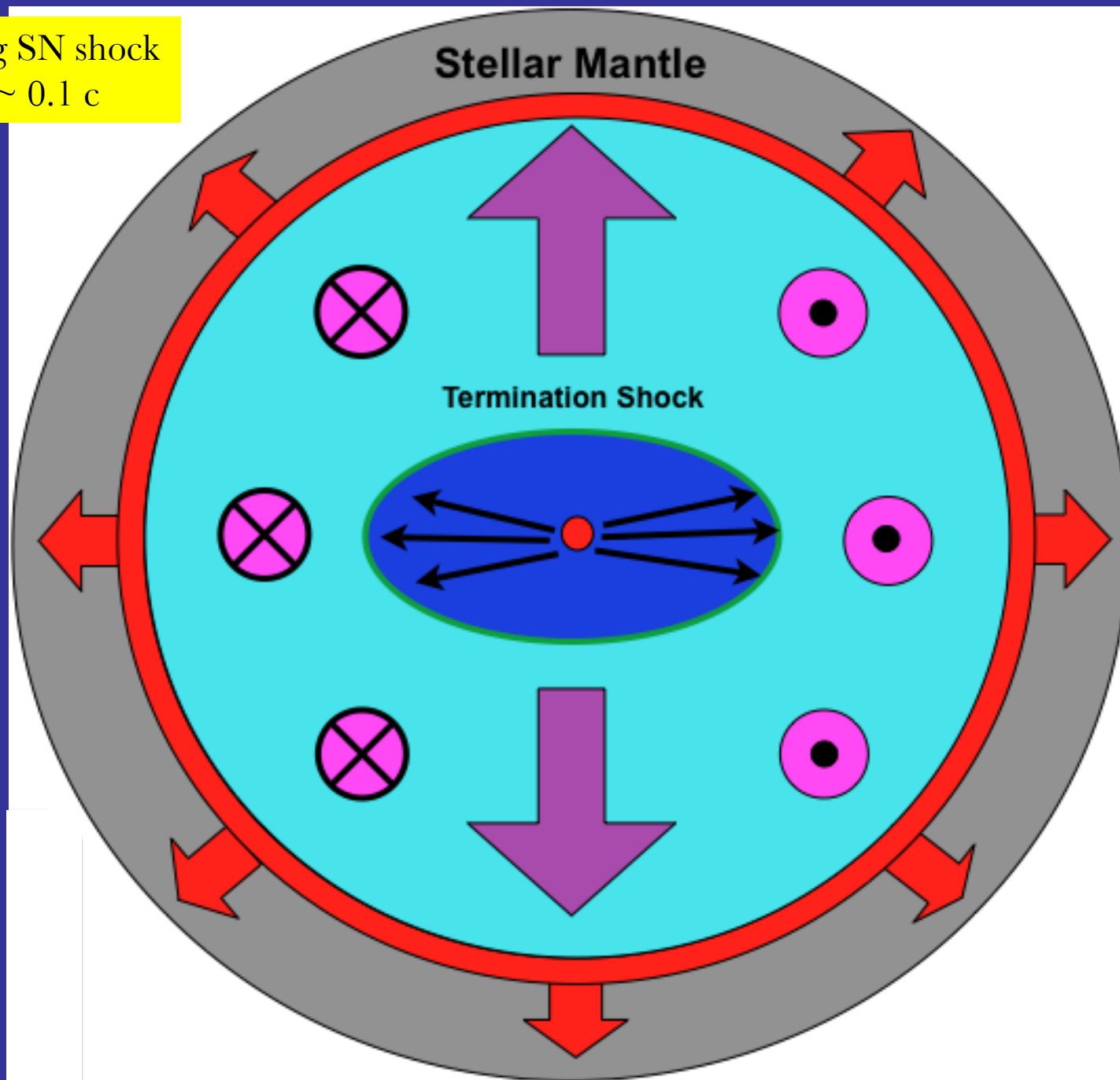
Outgoing SN shock

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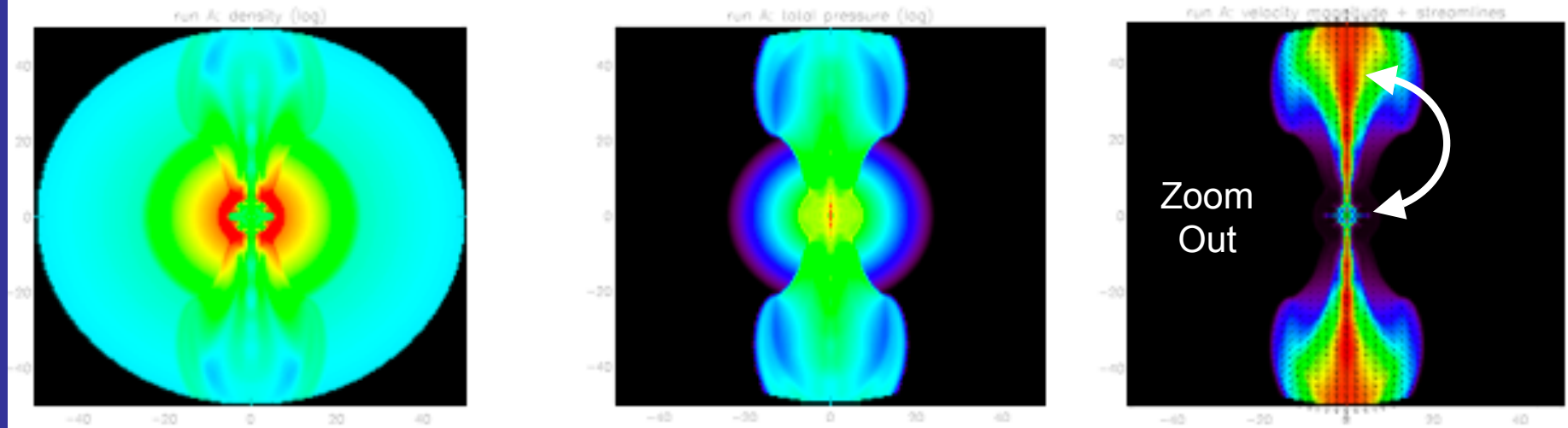
Outgoing SN shock

$$V_{\text{SN}} \sim 0.1 c$$

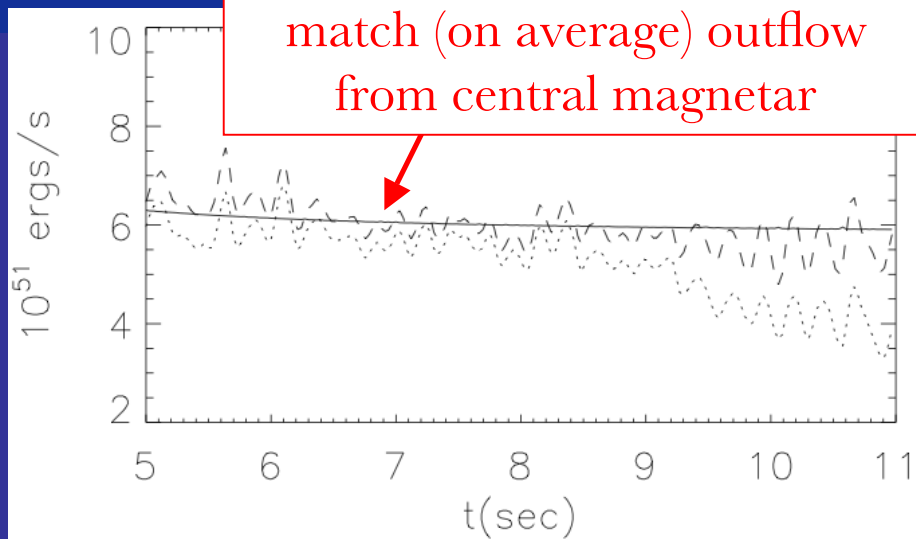


Jet Formation via Stellar Confinement

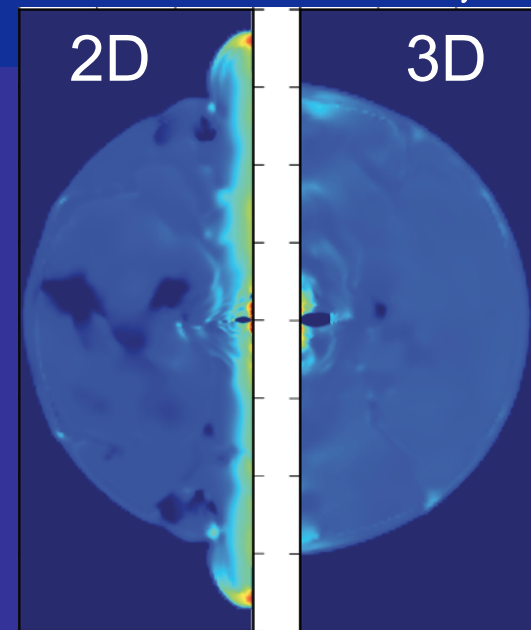
(Bucciantini et al. 2007, 08, 09; cf. Uzdensky & MacFadyen 07; Komissarov & Barkov 08)



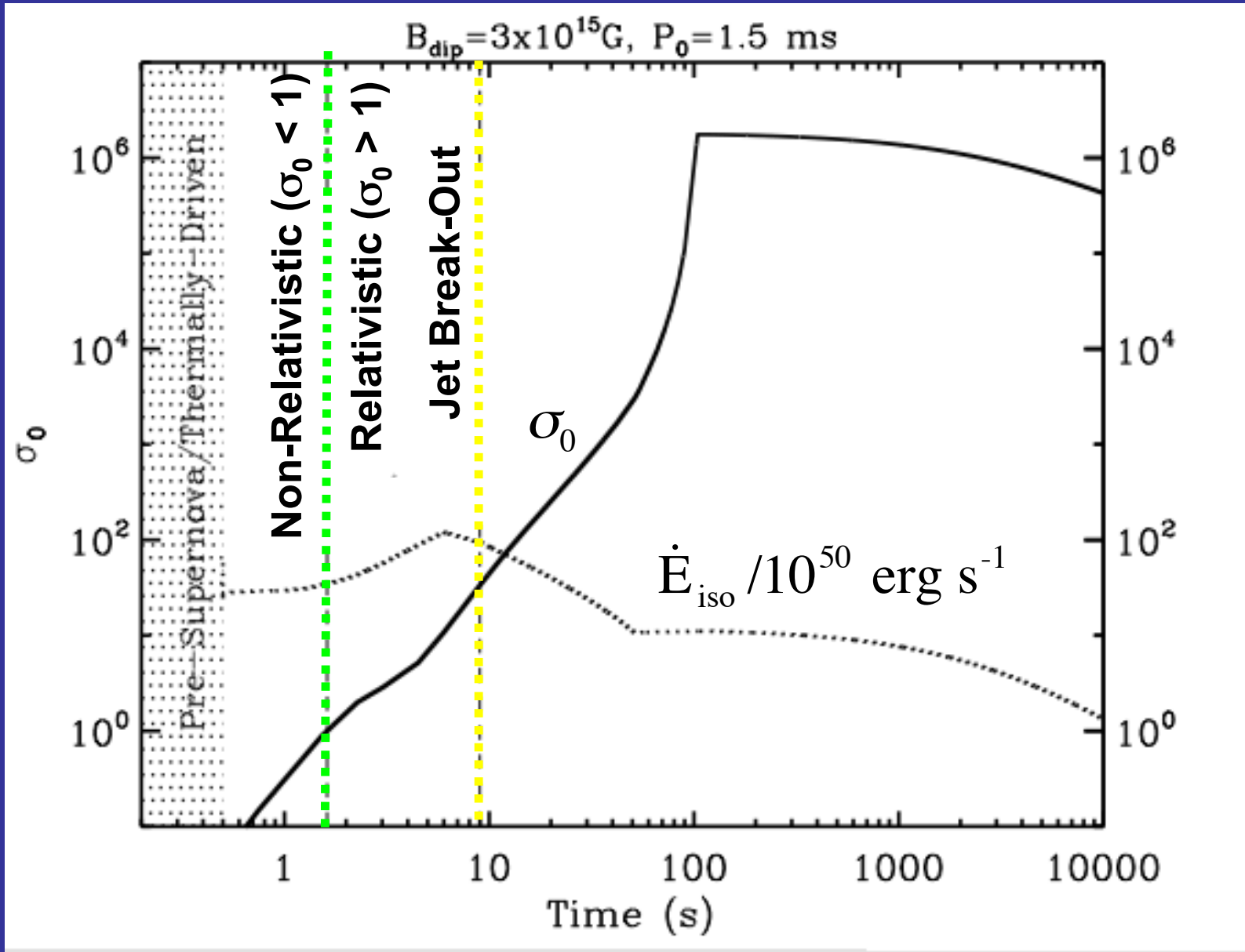
Jet power & mass-loading
match (on average) outflow
from central magnetar



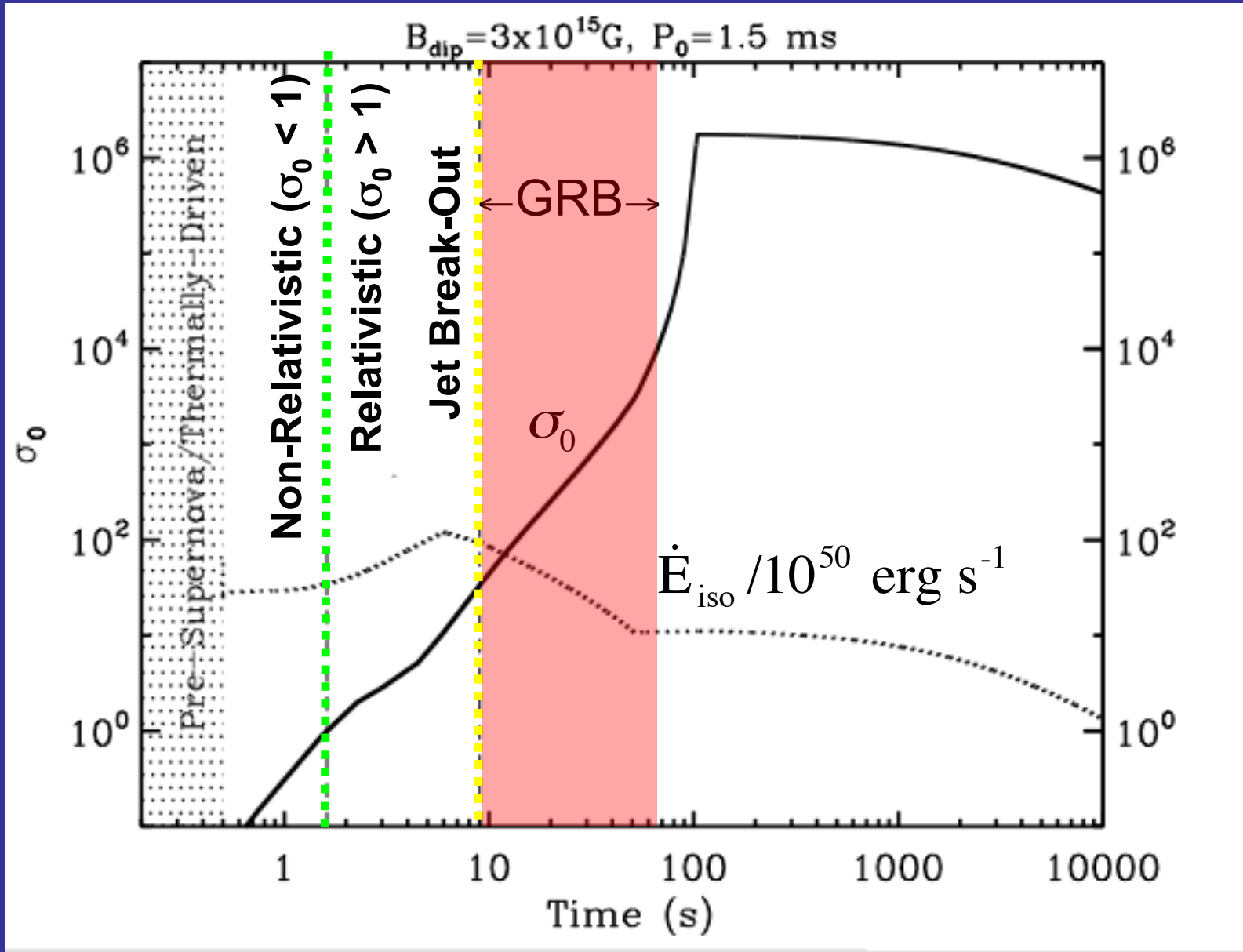
Kink Instability



Porth, Komissarov, & Keppens 13

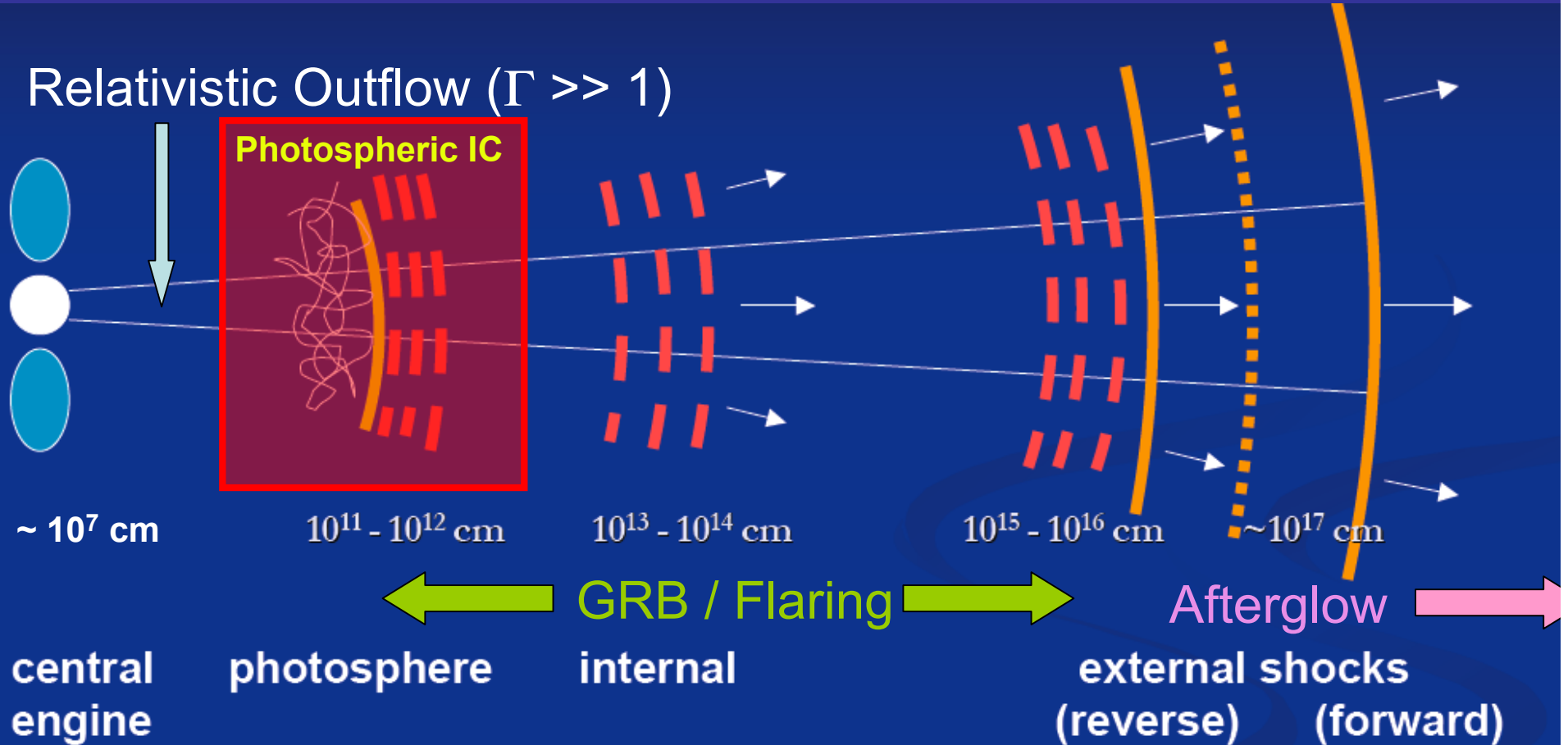


Outflow becomes relativistic at $t \sim 2$ seconds;
 Jet breaks out of star at $t_{\text{bo}} \sim R_{\star} / \beta c \sim 10$ seconds

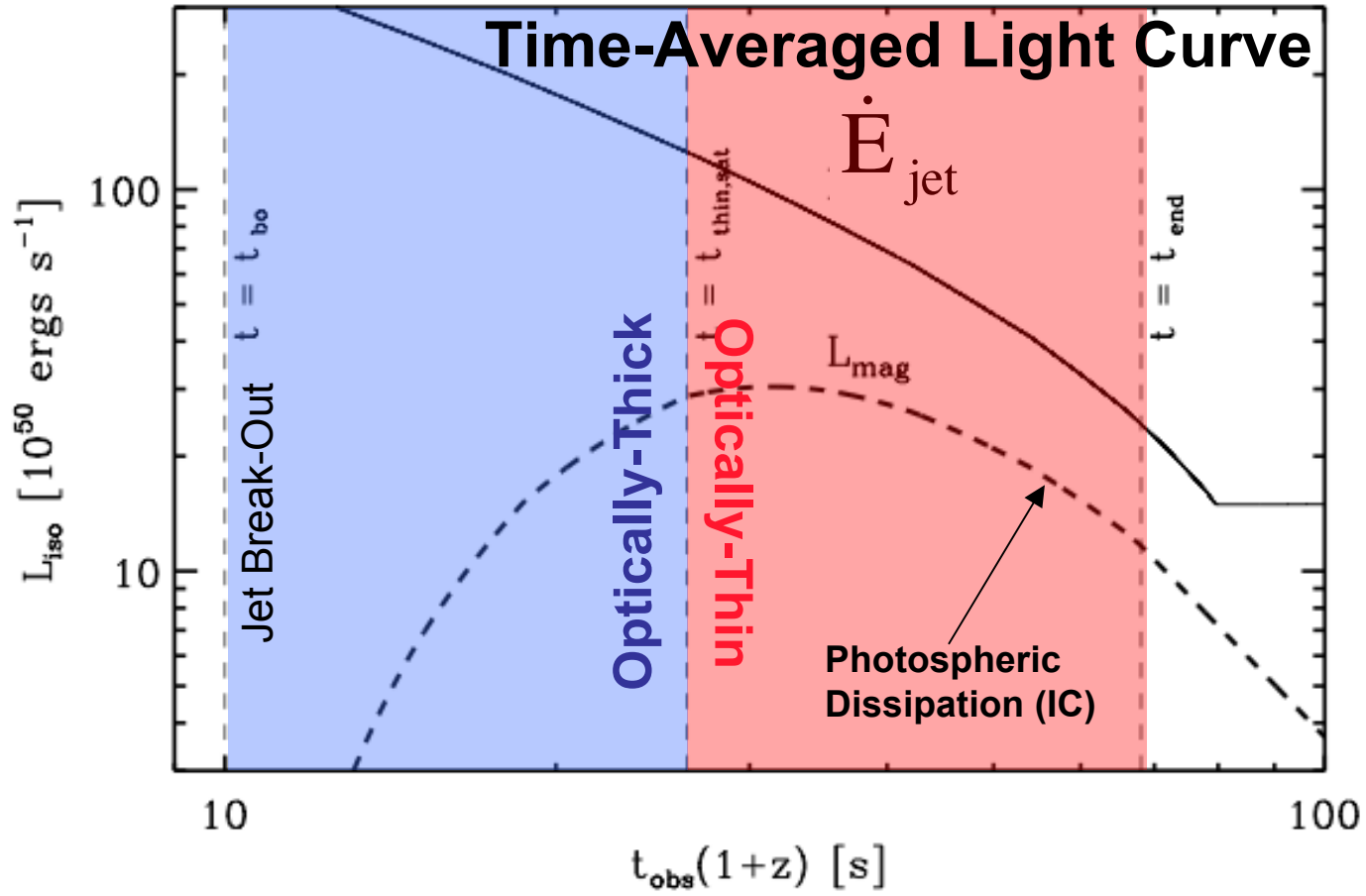


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 Jet breaks out of star at $t_{\text{bo}} \sim R_{\star} / \beta c \sim 10$ seconds

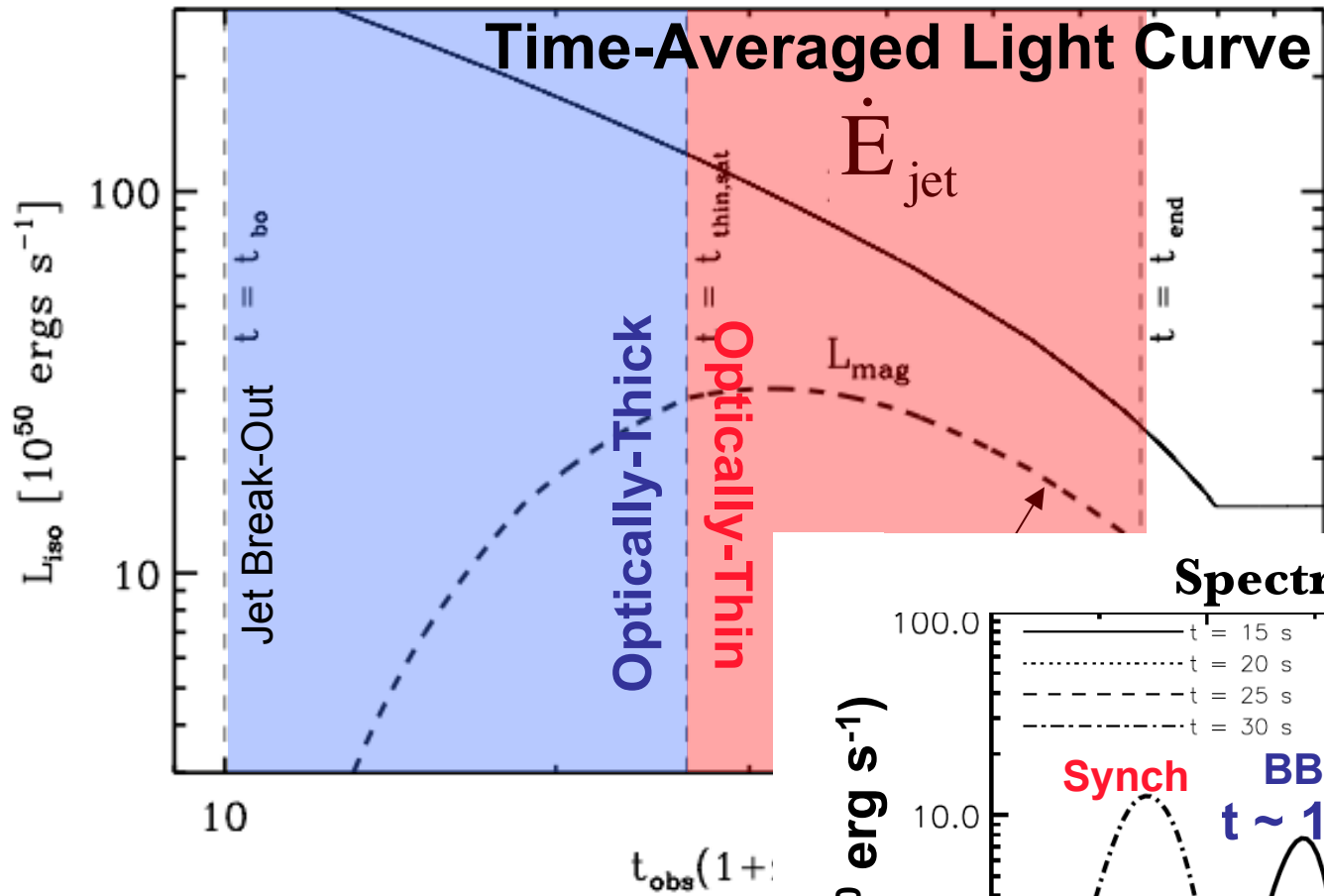
GRB Emission - What, Where, How?



1. **What** is jet's composition? (kinetic or magnetic?)
2. **Where** is dissipation occurring? (photosphere? deceleration radius?)
3. **How** is radiation generated? (synchrotron, IC, hadronic?)

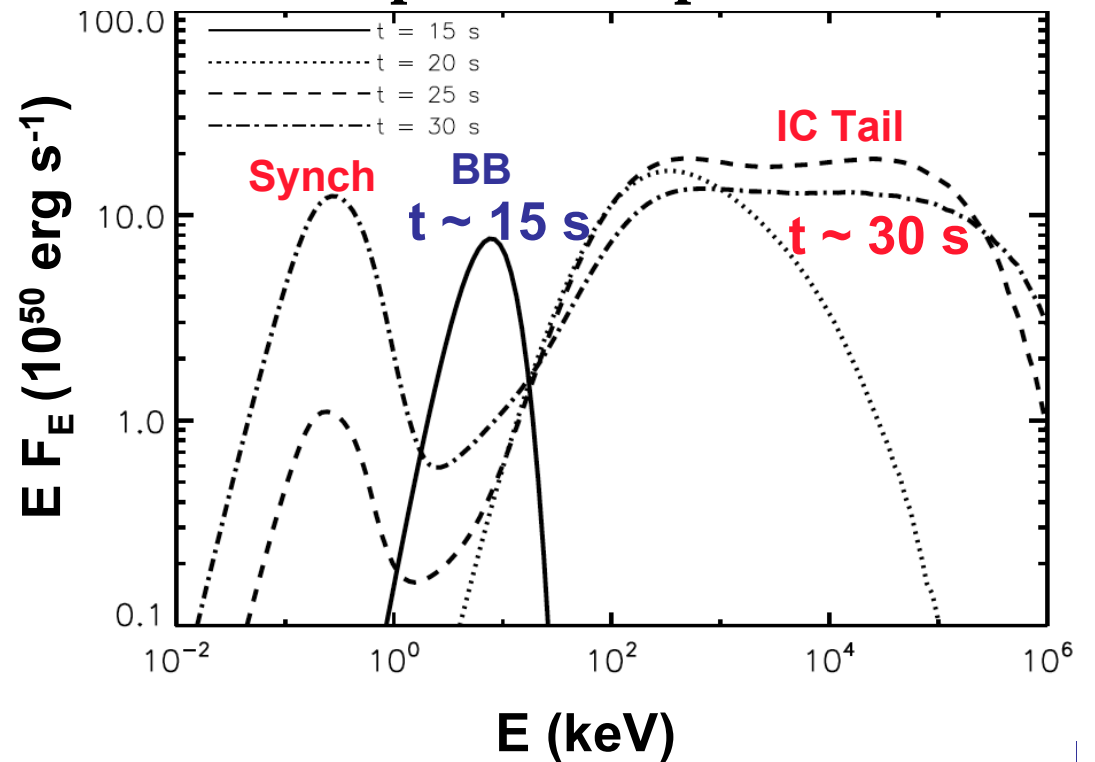


Metzger et al. 2011



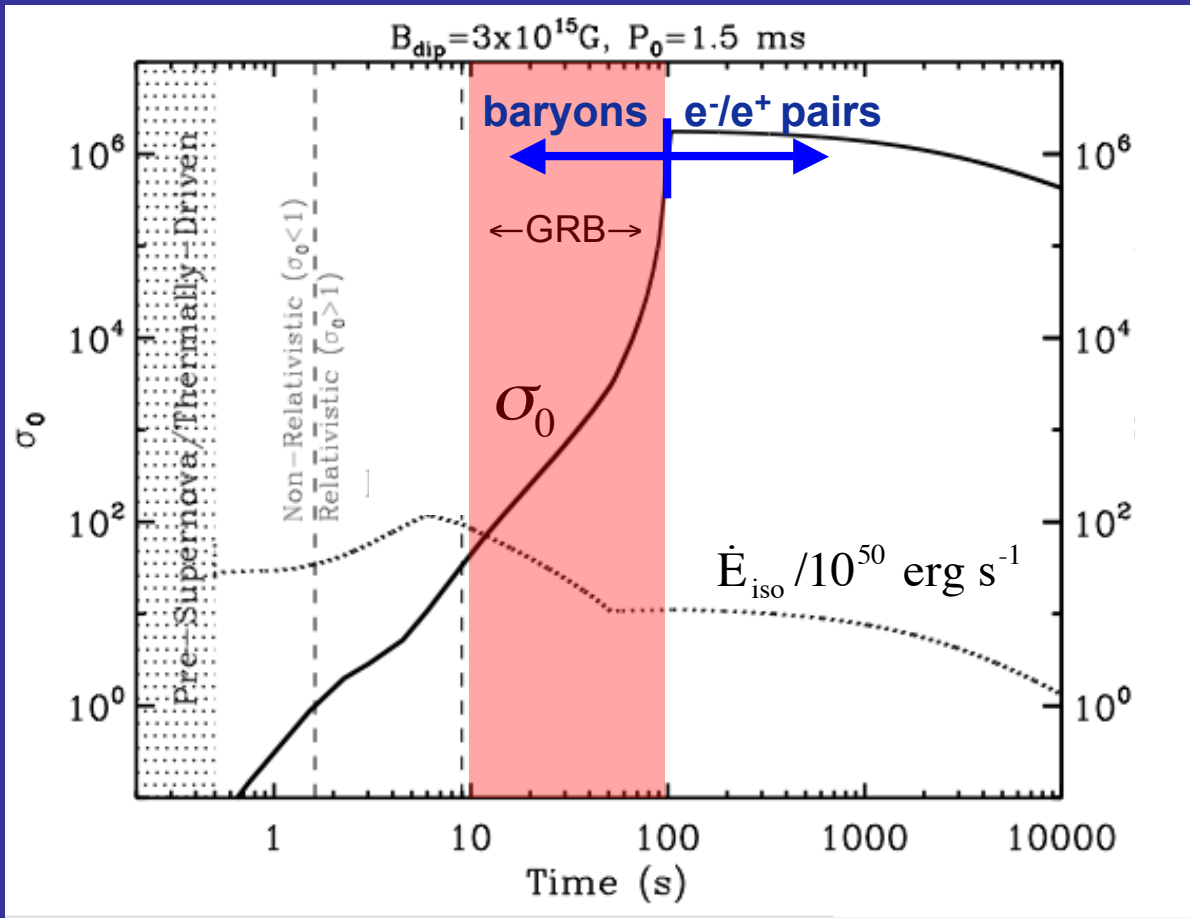
Metzger et al. 2011

Spectral Snapshots



Hot Electrons \Rightarrow
 IC Scattering (γ -rays)
 and Synchrotron (optical)

End of the GRB = Neutrino Transparency?



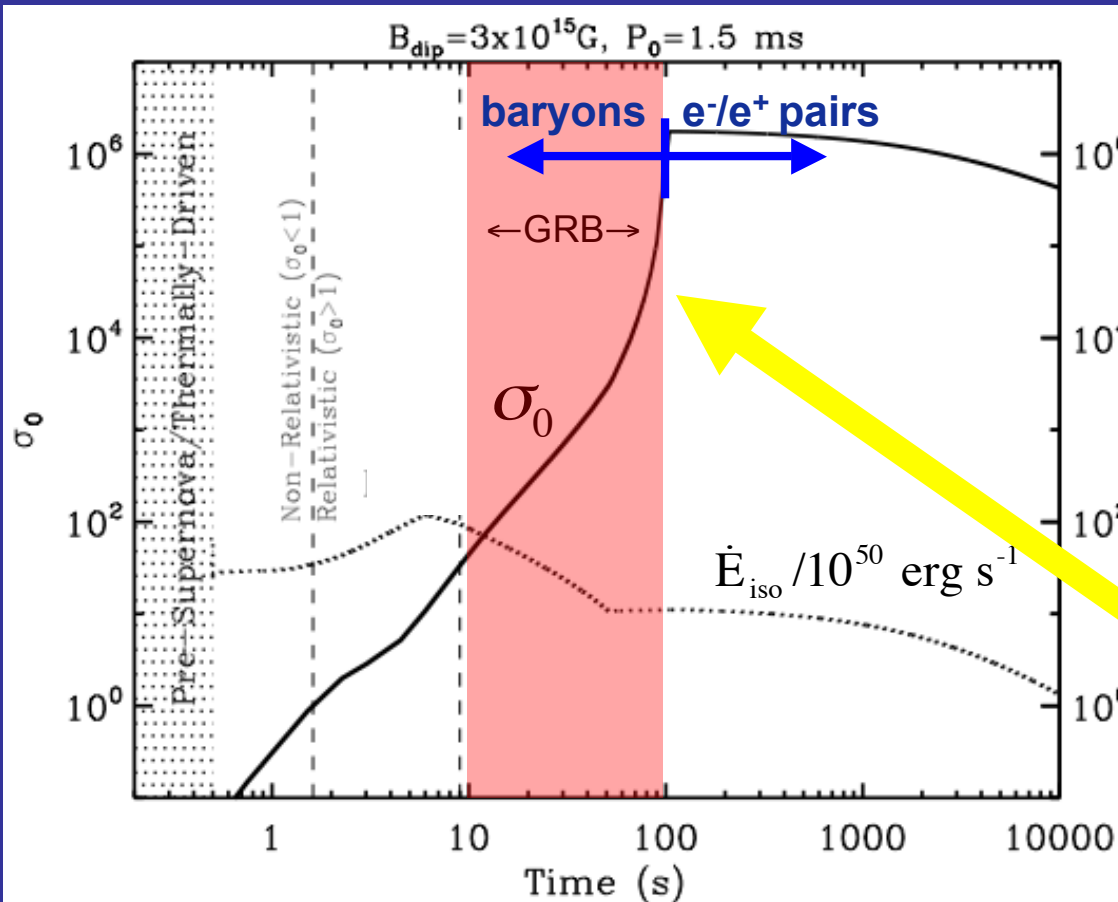
Ultra High- σ Outflows

⇒

- Acceleration is Inefficient (e.g. Tchekhovskoy et al. 2009)
- Internal Shocks are Weak (e.g. Kennel & Coroniti 1984)
- Reconnection is Slow (e.g. Drenkahn & Spruit 2002)

$$T_{\text{GRB}} \sim T_{\text{v thin}} \sim 20 - 100 \text{ s}$$

End of the GRB = Neutrino Transparency?

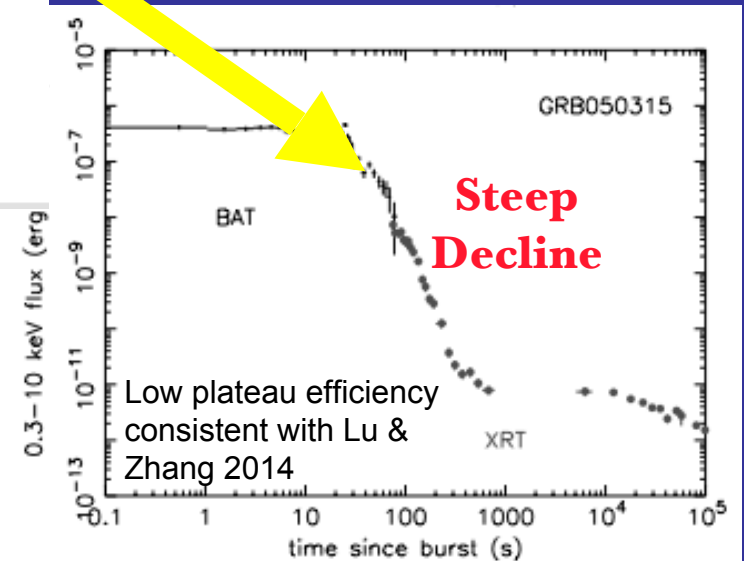


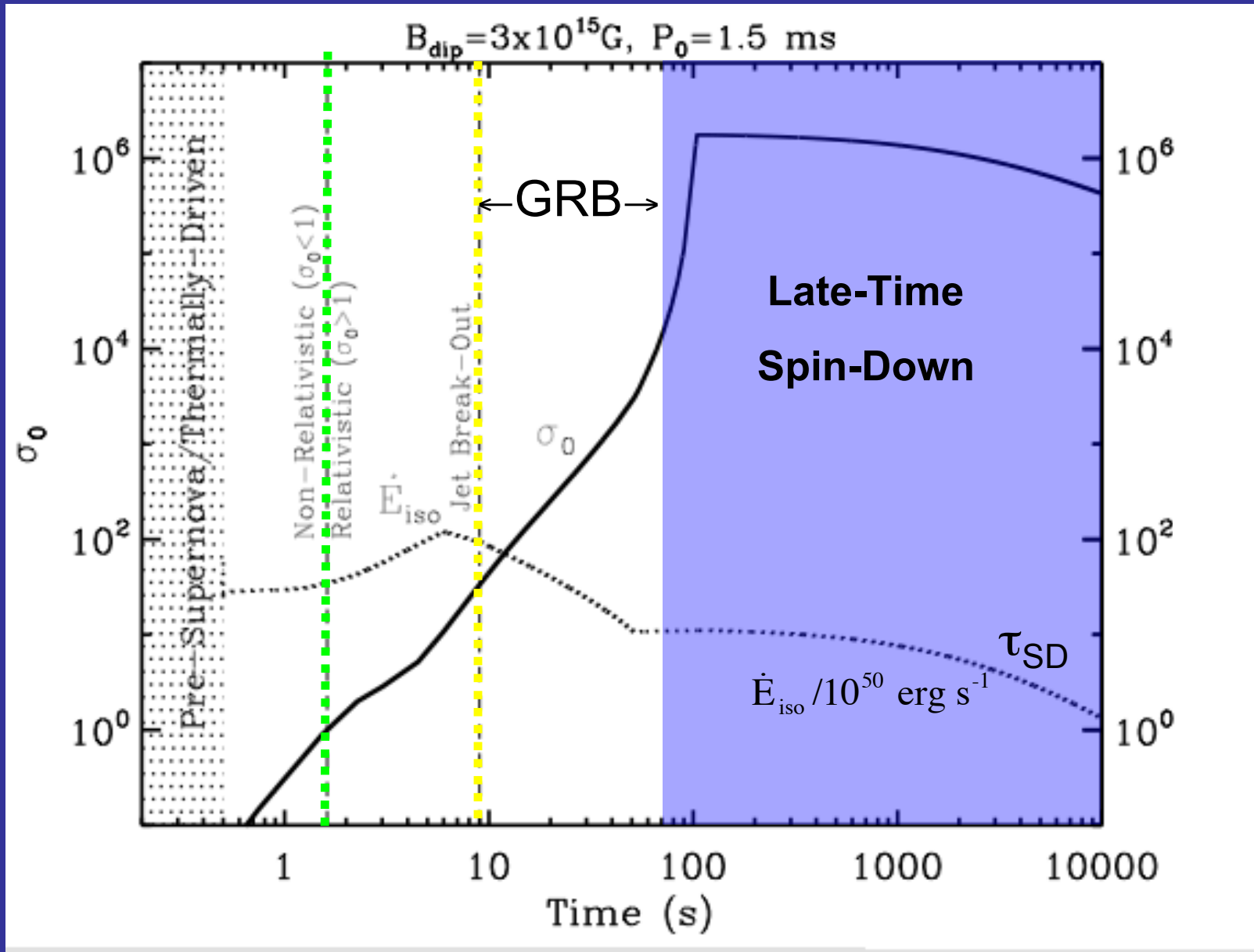
Ultra High- σ Outflows



- Acceleration is Inefficient (e.g. Tchekhovskoy et al. 2009)
- Internal Shocks are Weak (e.g. Kennel & Coroniti 1984)
- Reconnection is Slow (e.g. Drenkahn & Spruit 2002)

$$T_{\text{GRB}} \sim T_{\text{v thin}} \sim 20 - 100 \text{ s}$$





e.g. Zhang & Meszaros 2001; Troja et al. 2007; Yu et al. 2009; Lyons et al. 2010

$B_{\text{dip}} = 3 \times 10^{15} \text{G}$, $P_0 = 1.5 \text{ ms}$

10^6

Given

← GRB →

X-ray Afterglow

GRB060729

z 0.540

E_{peak} 116.

E_{iso} 0.4

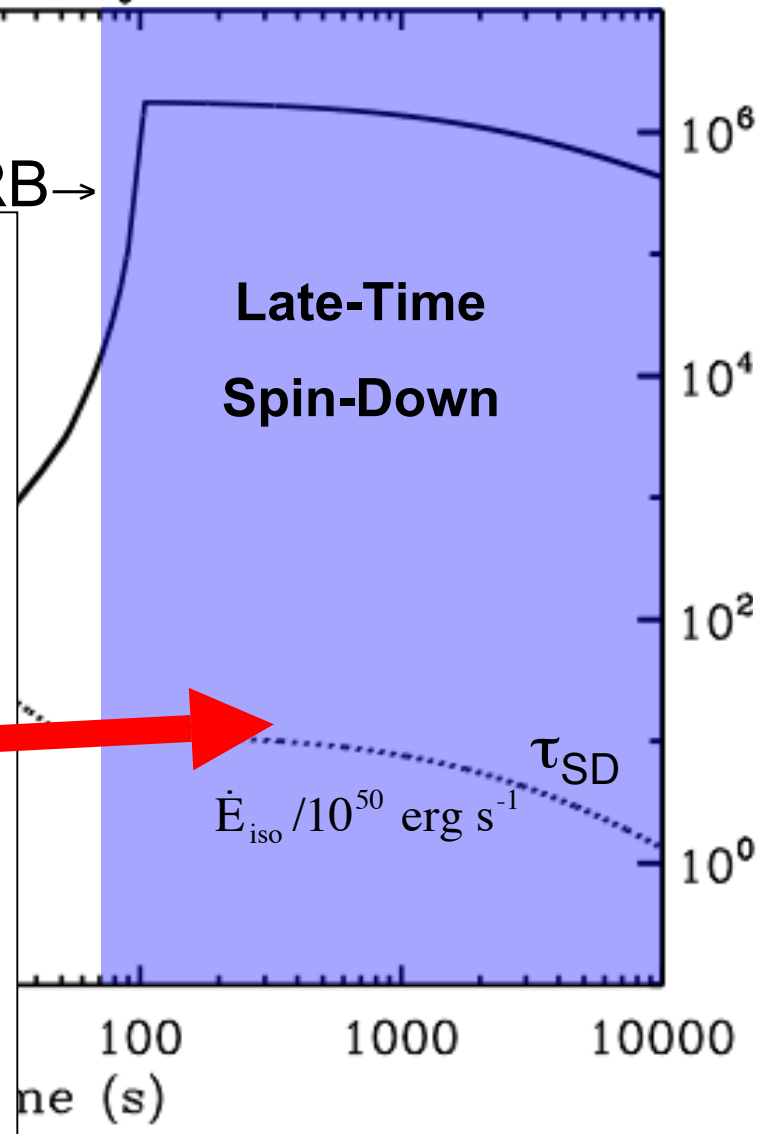
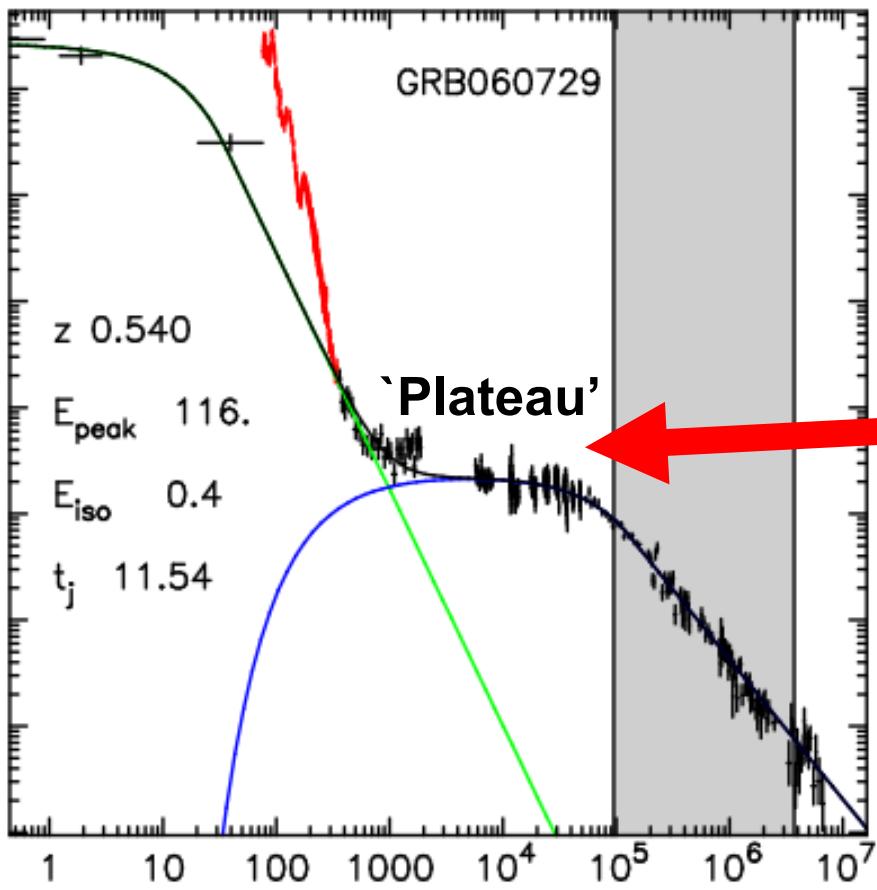
t_j 11.54

'Plateau'

Late-Time Spin-Down

$\dot{E}_{\text{iso}} / 10^{50} \text{ erg s}^{-1}$ τ_{SD}

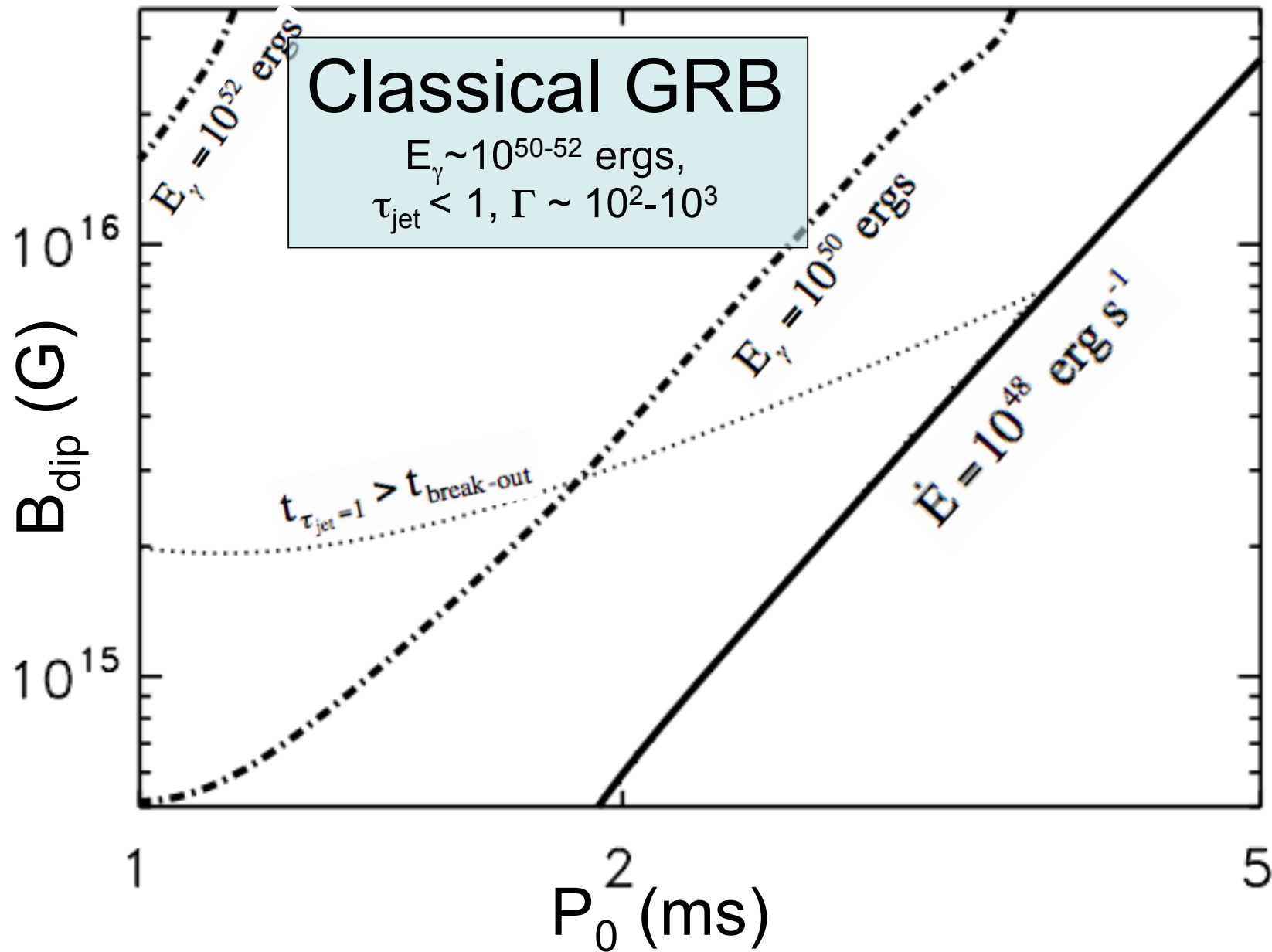
Willingale et al. 2007



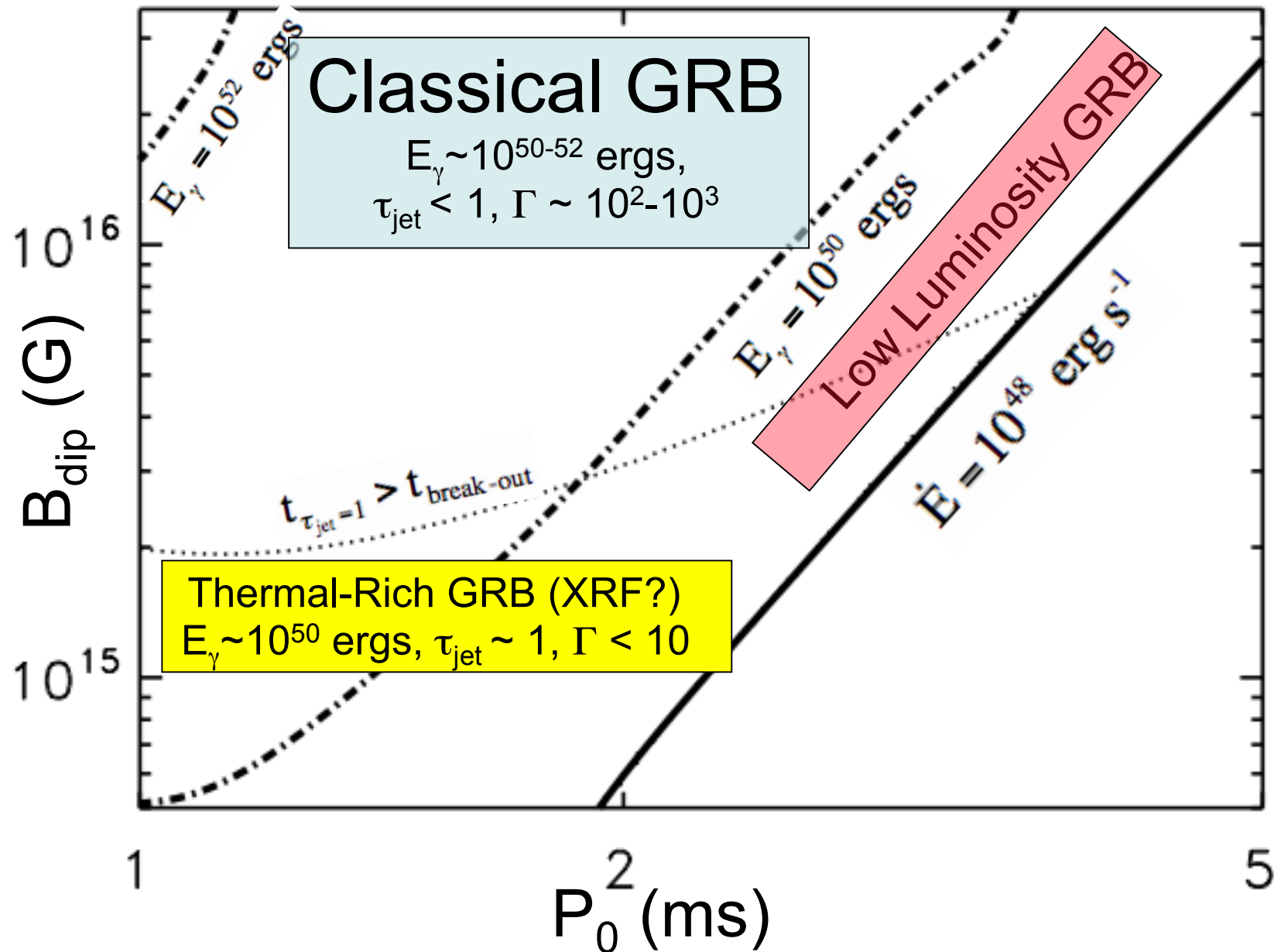
Time after trigger (s)

e.g. Zhang & Meszaros 2001; Troja et al. 2007; Yu et al. 2009; Lyons et al. 2010; Rowlinson et al. 2010, 2013; Gompertz et al. 2013

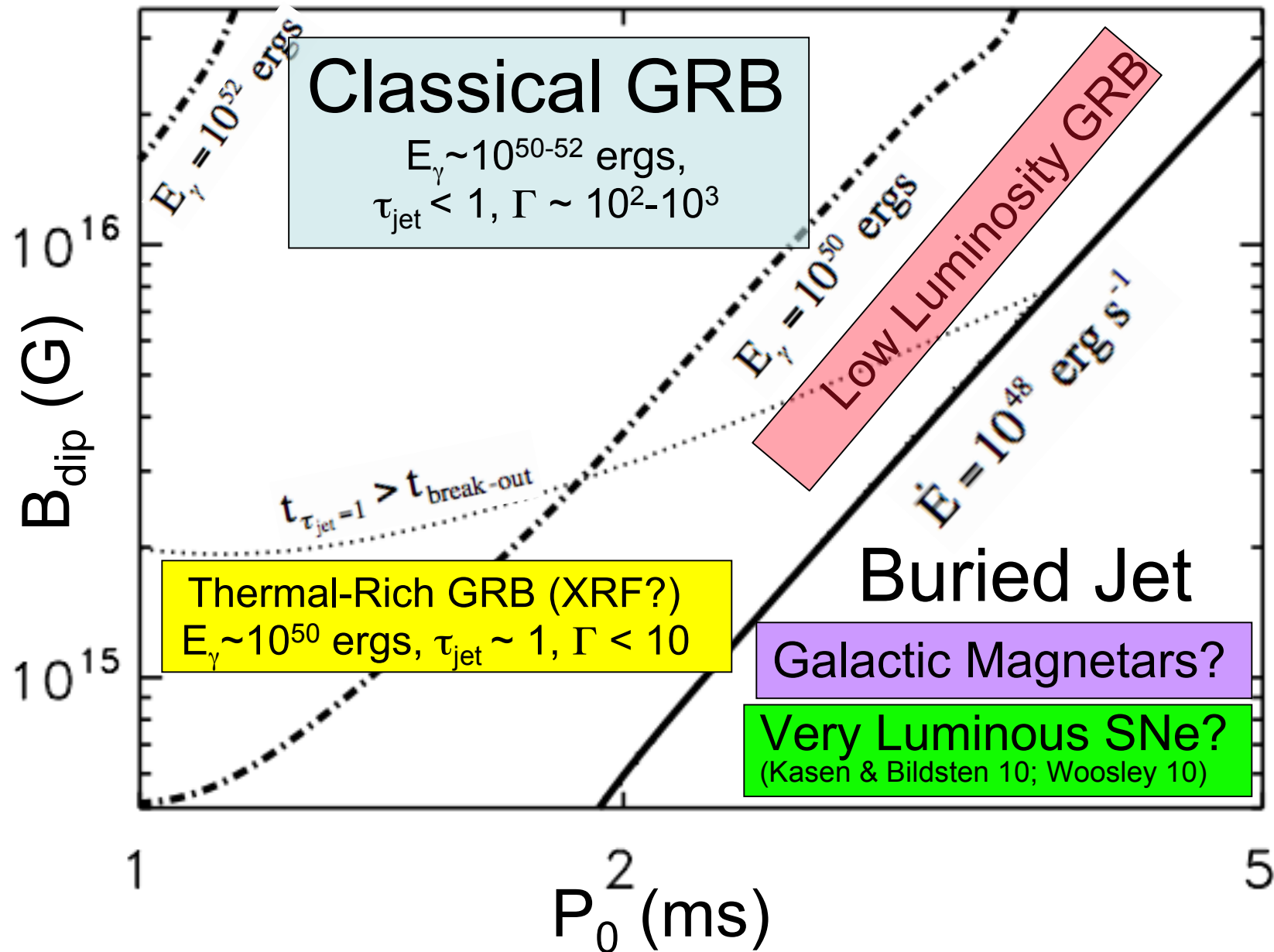
A Diversity of Magnetar Birth



A Diversity of Magnetar Birth

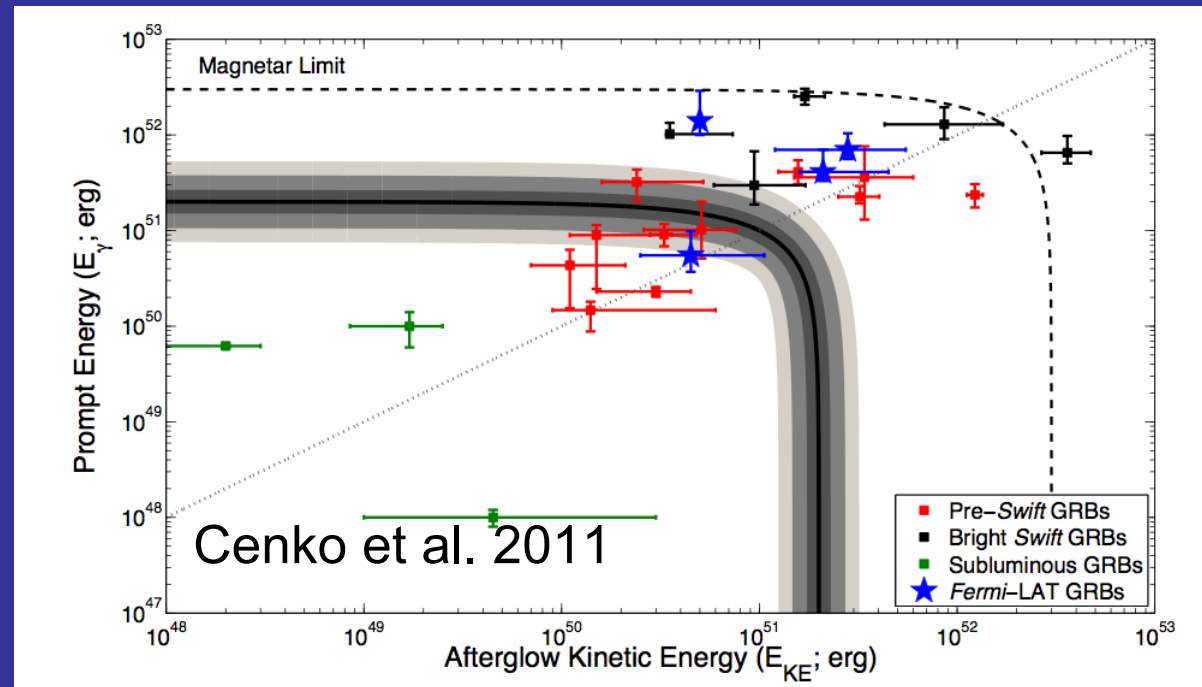


A Diversity of Magnetar Birth



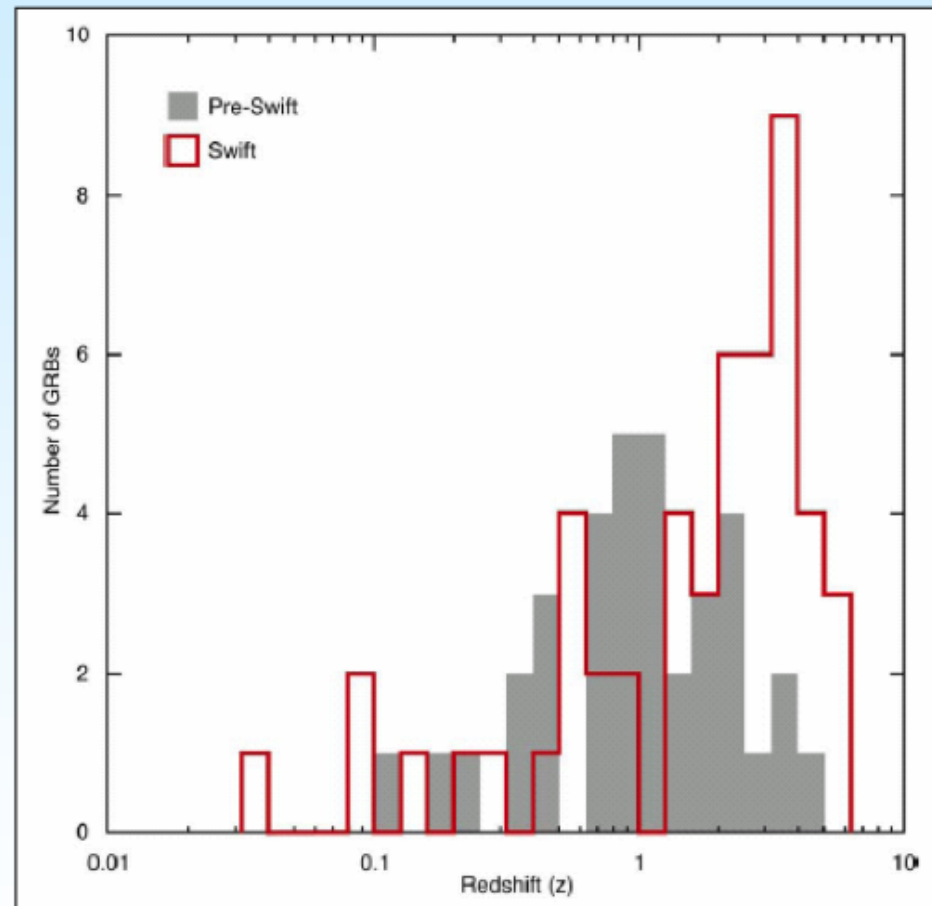
Observational Tests & Constraints

- **Max Energy*** -
 $E_{KE} + E_{\gamma} < 3 \times 10^{52}$ ergs
*subject to uncertainties in
afterglow modeling.
(e.g. Zhang & MacFadyen 09).



- **Long GRB** *always* accompanied by bright, energetic
 - Consistent with observations thus far (Woosley & Bloom 2006).
- **Γ increases during GRB and correlates with E_{γ}**
 - translate jet luminosity/magnetization into unique prediction for gamma-ray light curves and spectra.

Redshift Distribution of Swift Long-Duration GRBs



Mean redshift $\langle z \rangle \sim 2$ vs. $\langle z \rangle \sim 1$ for pre-Swift GRBs

Most Distant GRBs

$z = 8.26$

$t = 630$ Myr

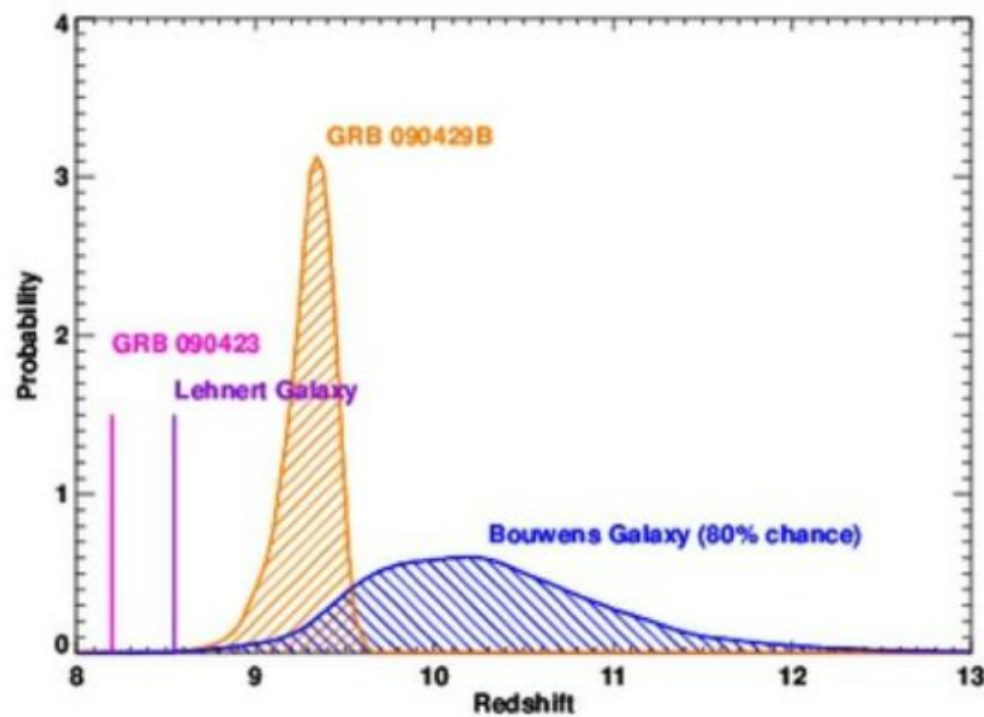
$z = 10$

$t = 480$ Myr



GRB 090423

ESO



GRBs As Probes of Chemical Evolution

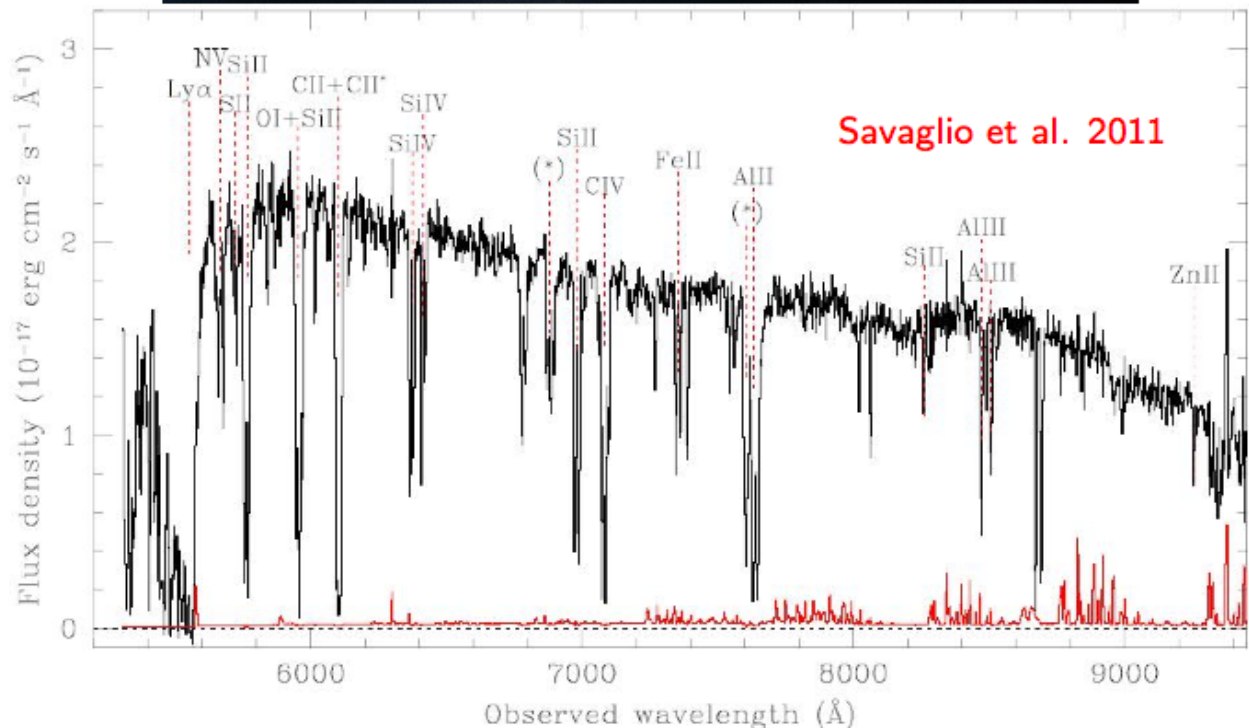
GRB light is absorbed by intervening galaxies.

Two systems, $z = 3.5673$ and $z = 3.5774$, probably merging galaxies, are illuminated.

GRB could have a progenitor formed in star formation triggered by merger.

$[Zn/H] = 0.29$ and $[S/H] = 0.67$ are highest metallicities recorded for $z > 3$ objects.

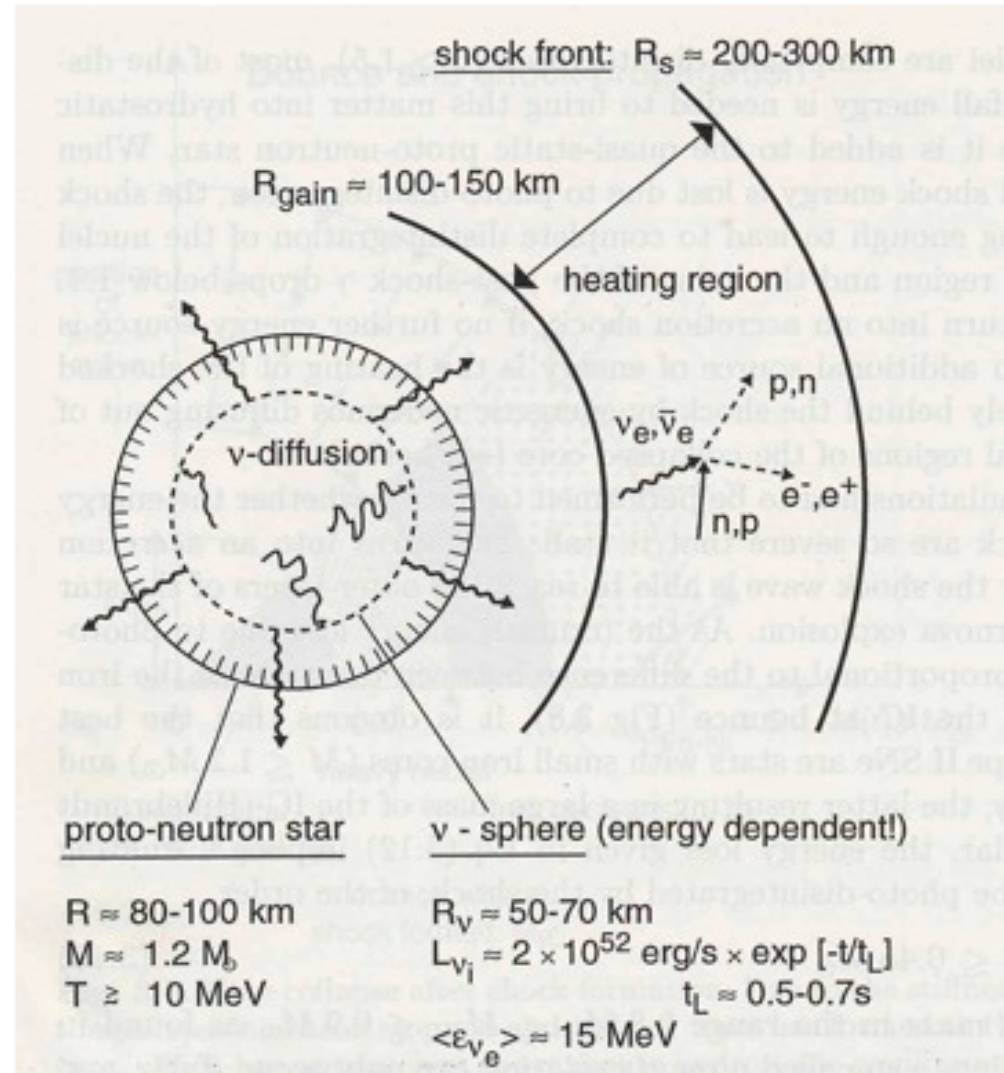
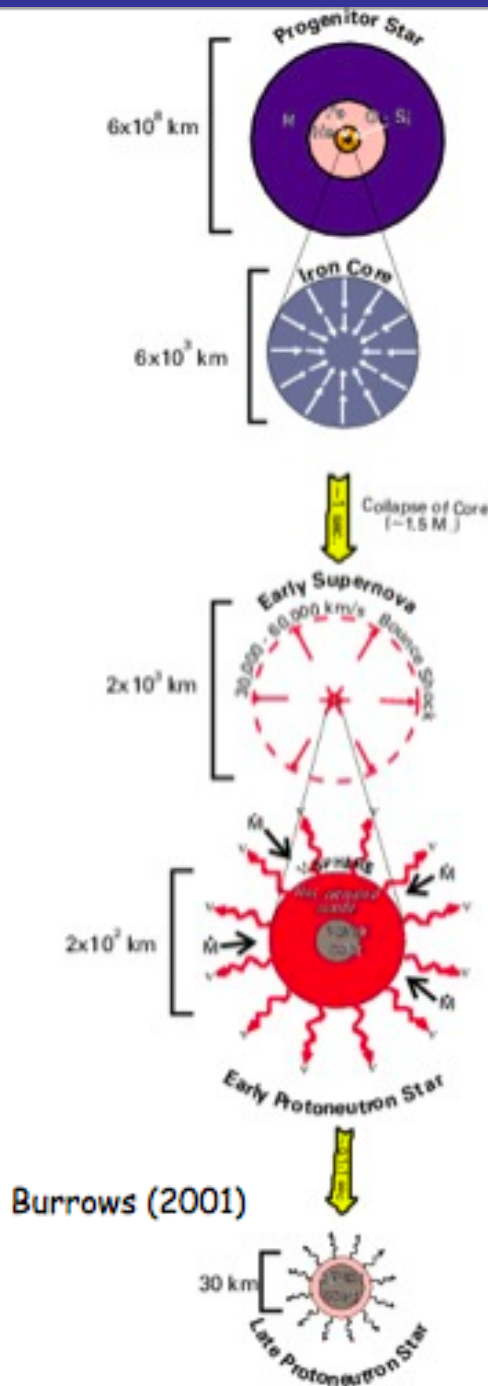
Shows star formation and metallicities heightened by interaction of galaxies.



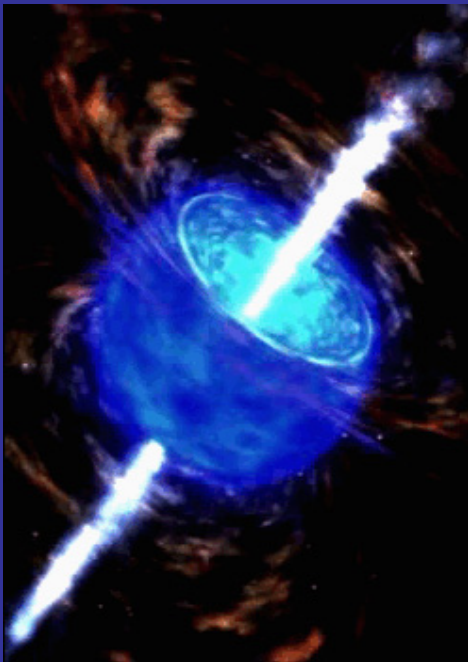
"Delayed" SN Explosion

Accretion vs. Neutrino heating

Courtesy A. MacFadyen



The Proto-Magnetar Model for Gamma-Ray Bursts



Brian Metzger
NASA Einstein Fellow
Princeton University

In collaboration with

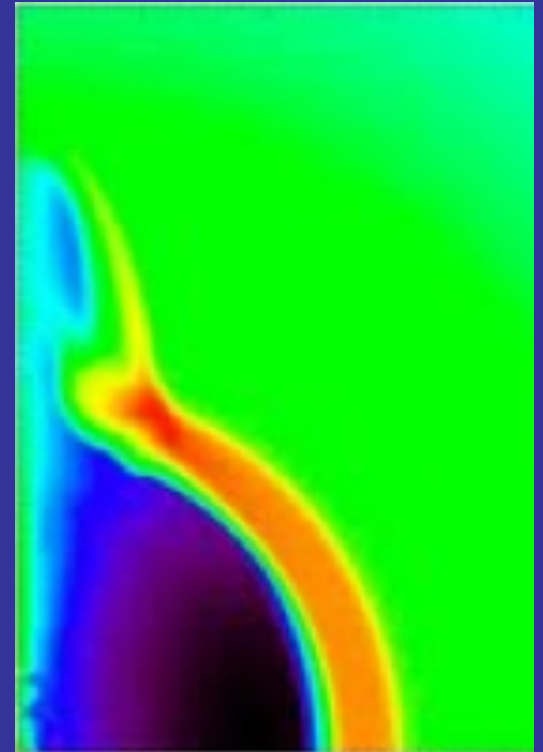
Eliot Quataert (UC Berkeley)

Todd Thompson (Ohio State)

Dimitrios Giannios (Princeton)

Niccolo Bucciantini (Nordita)

Jon Arons (UC Berkeley)

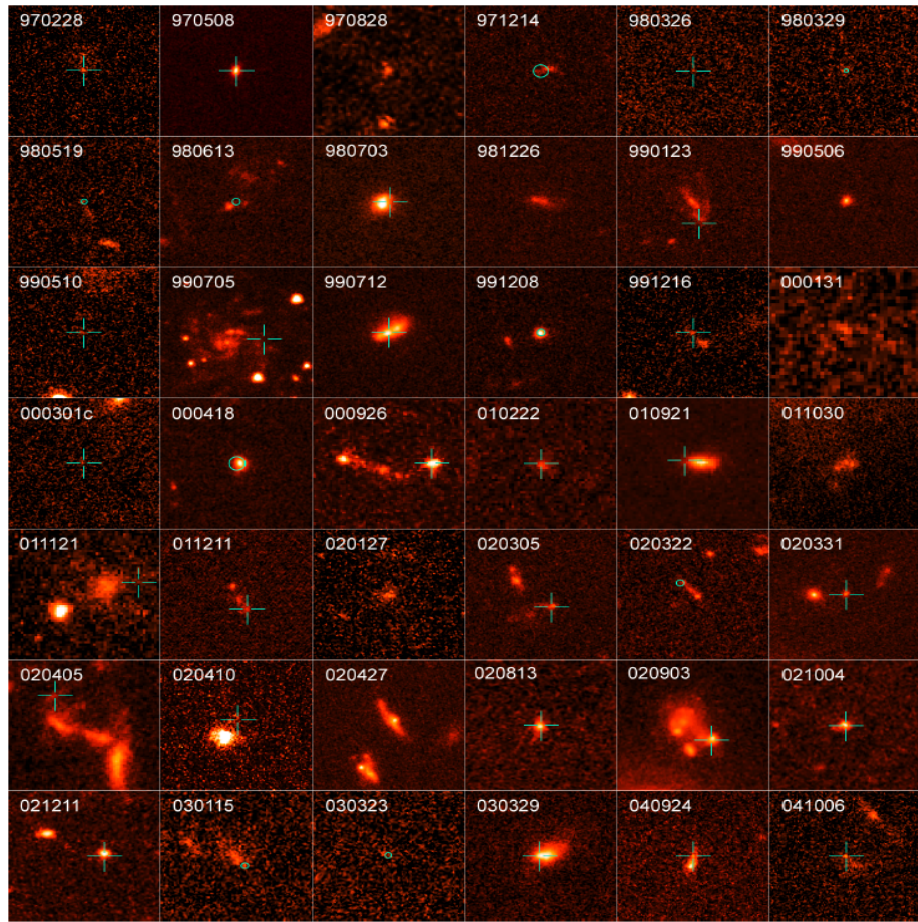


University of Minnesota Astronomy Colloquium,
October 22, 2010

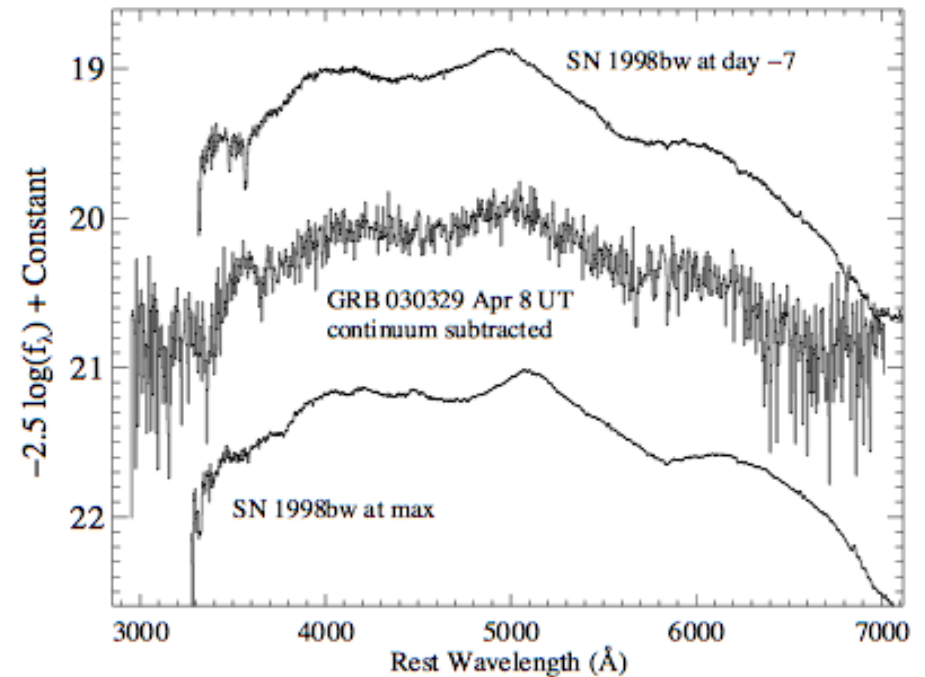
Long GRBs = Massive Stellar Death

(Paczynski 98, Galama et al. 98, Bloom et al. 99, Pian et al. 06, Modjaz et al. 06, Woosley & Bloom 06)

GRB 030329 \leftrightarrow SN 2003dh (BL Type Ic)



HST, Fruchter+ 2006

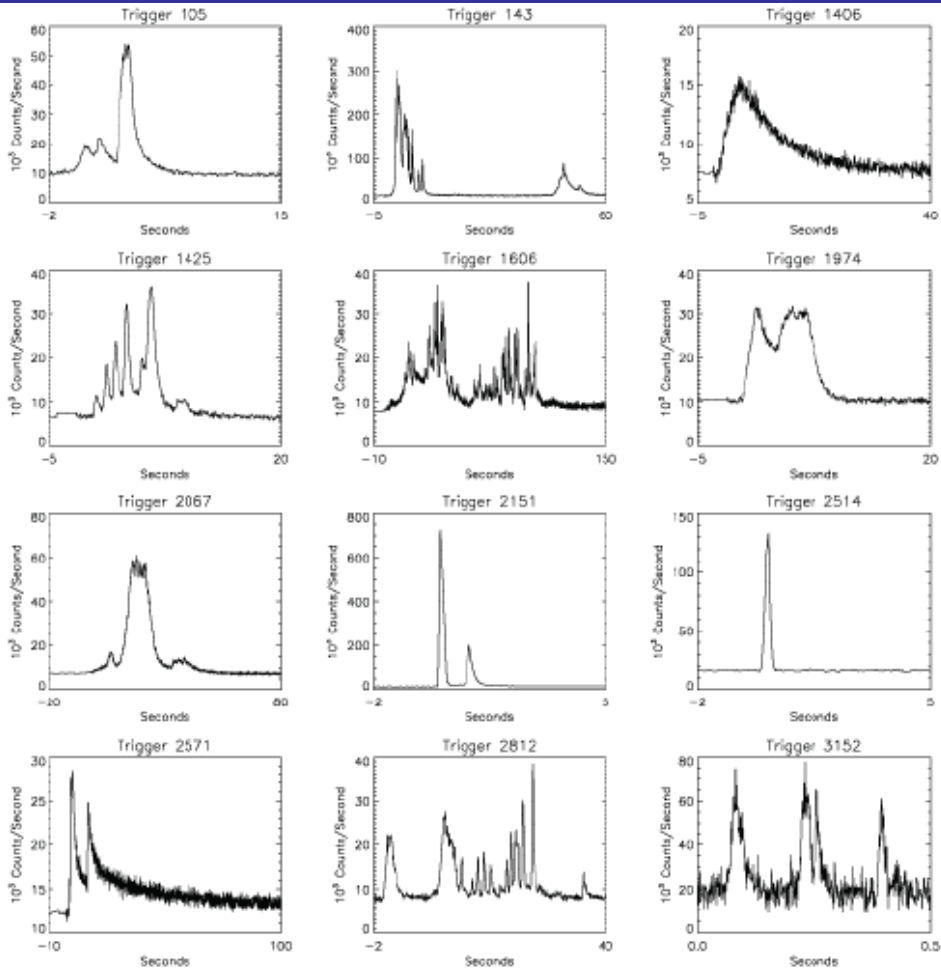


Stanek+ 2003

Broad Spectral Lines $\Rightarrow E_{\text{SN}} \sim 10^{52}$ ergs (“Hyper-Novae”)

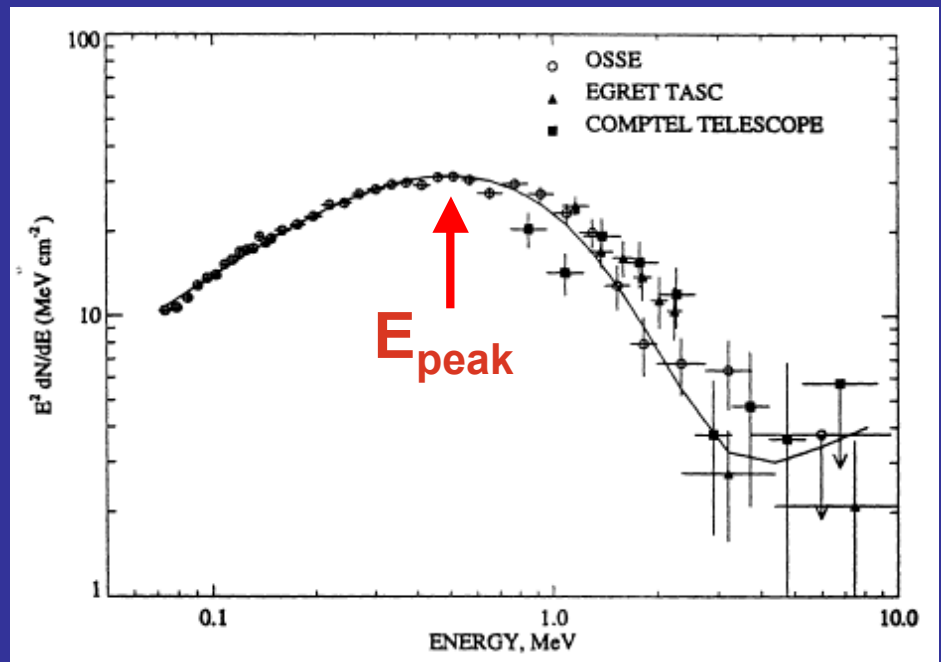
Long Duration Gamma-Ray Bursts

Light Curves



Average Duration $\sim 10 - 30$ s
 Typical Variability ~ 1 s

Spectra



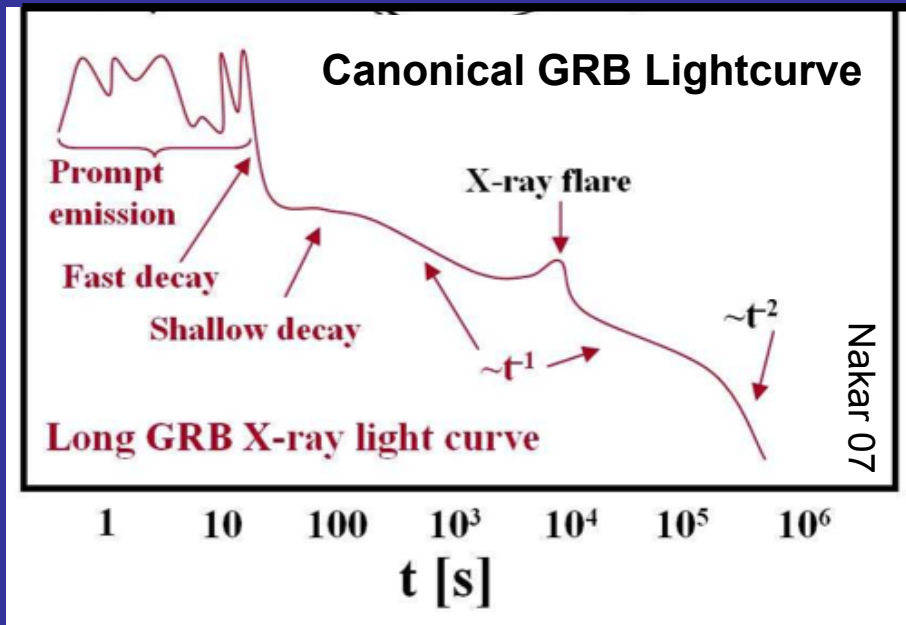
Broken power-law ('Band')
 spectrum w $E_{\text{peak}} \sim 300$ keV

Amati-Yonetoku Relations:

$$E_{\text{peak}} \propto E_{\text{iso}}^{0.4}, L_{\text{iso}}^{0.5}$$

Constraints on the Central Engine

- Energies - $E_\gamma \sim 10^{49-52}$ ergs
- Rapid Variability (down to ms)
 - But Most 'Power' on Timescales ~ 1 s
- Duration - $T_\gamma \sim 10-100$ seconds
- Steep Decay after GRB



- Ultra-Relativistic, Collimated Outflow $\Gamma \sim 100-1000$
- Association w Energetic Core Collapse Supernovae
- Late-Time Central Engine Activity (Plateau & Flaring)

BH

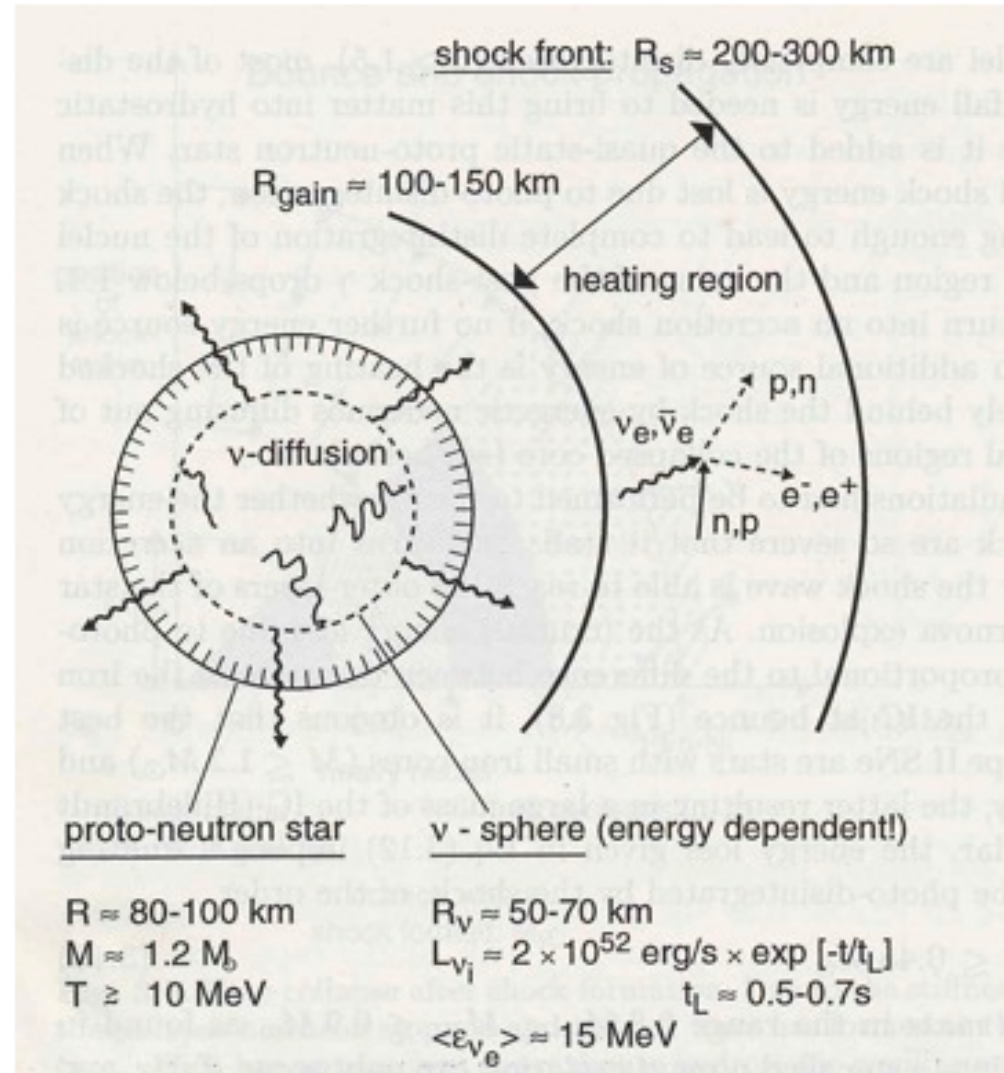
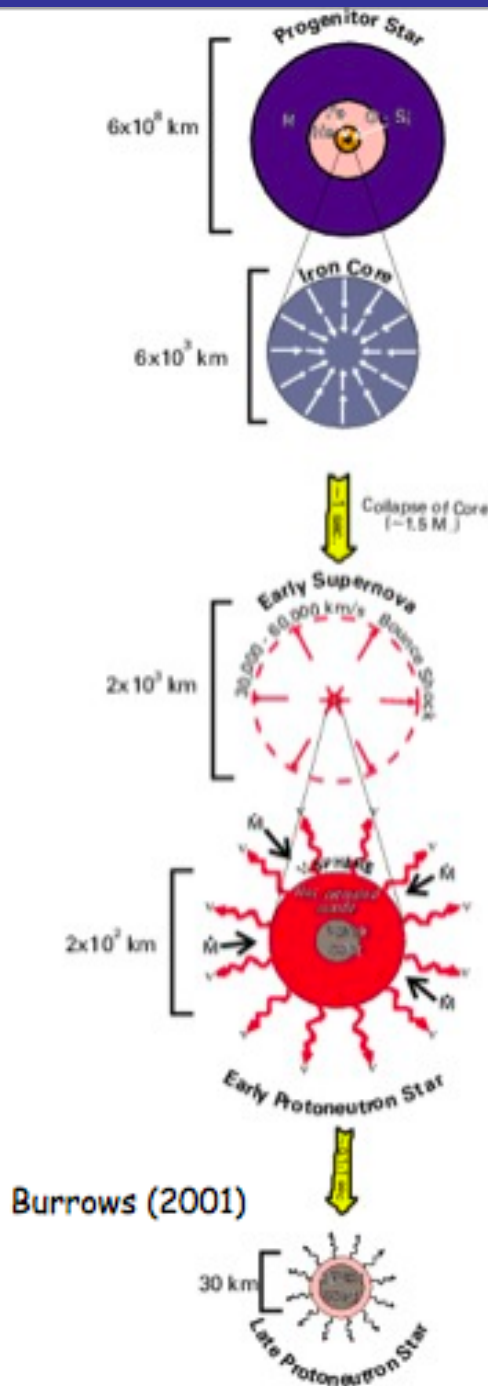
versus

NS

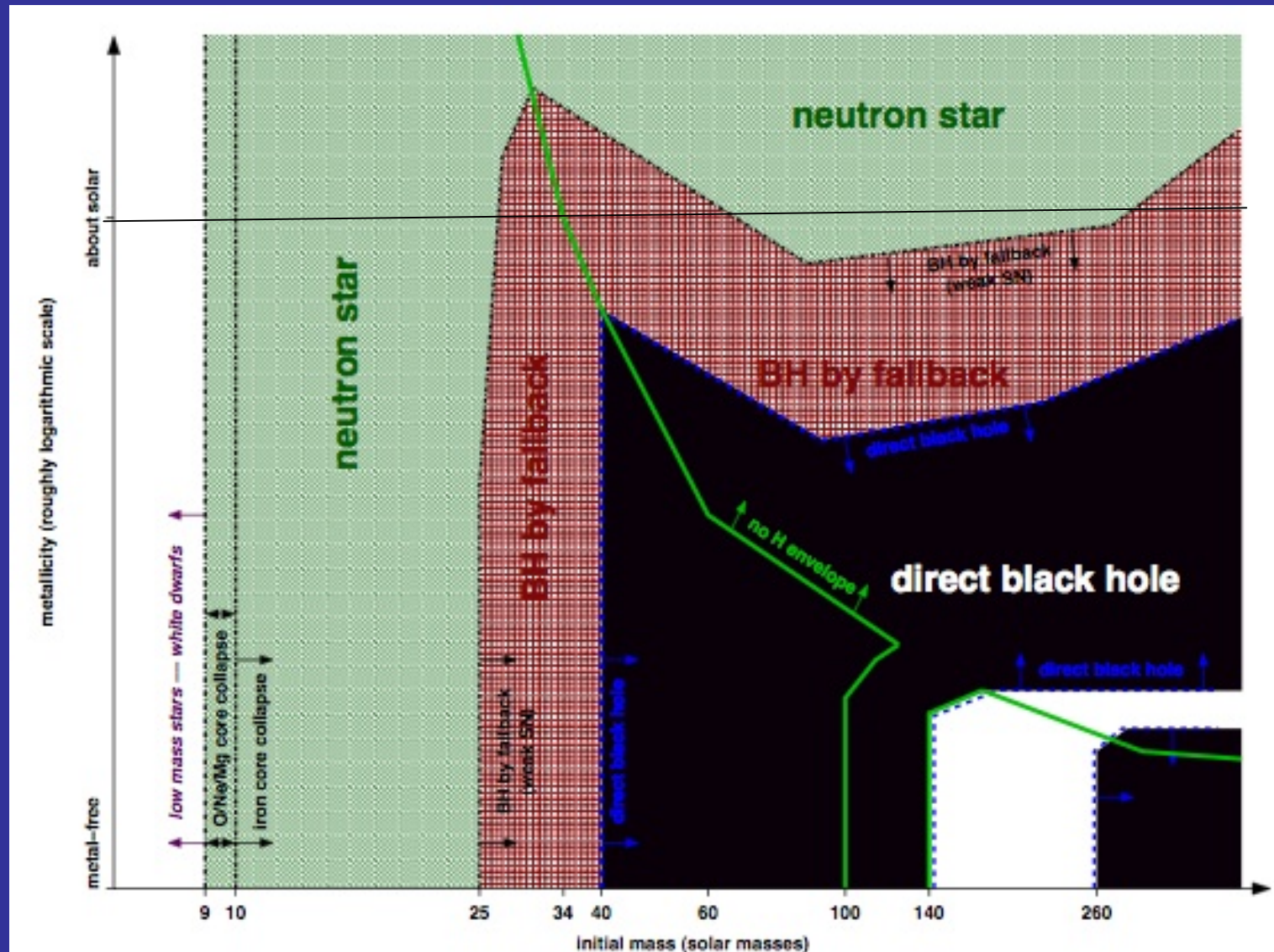
"Delayed" SN Explosion

Accretion vs. Neutrino heating

Courtesy A. MacFadyen



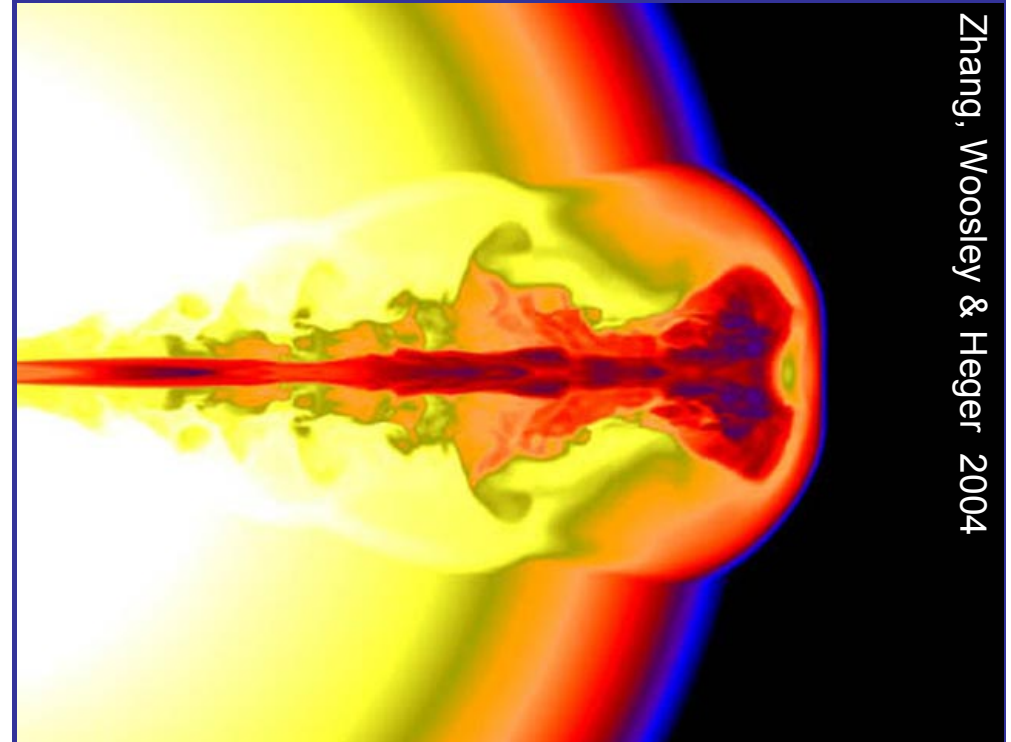
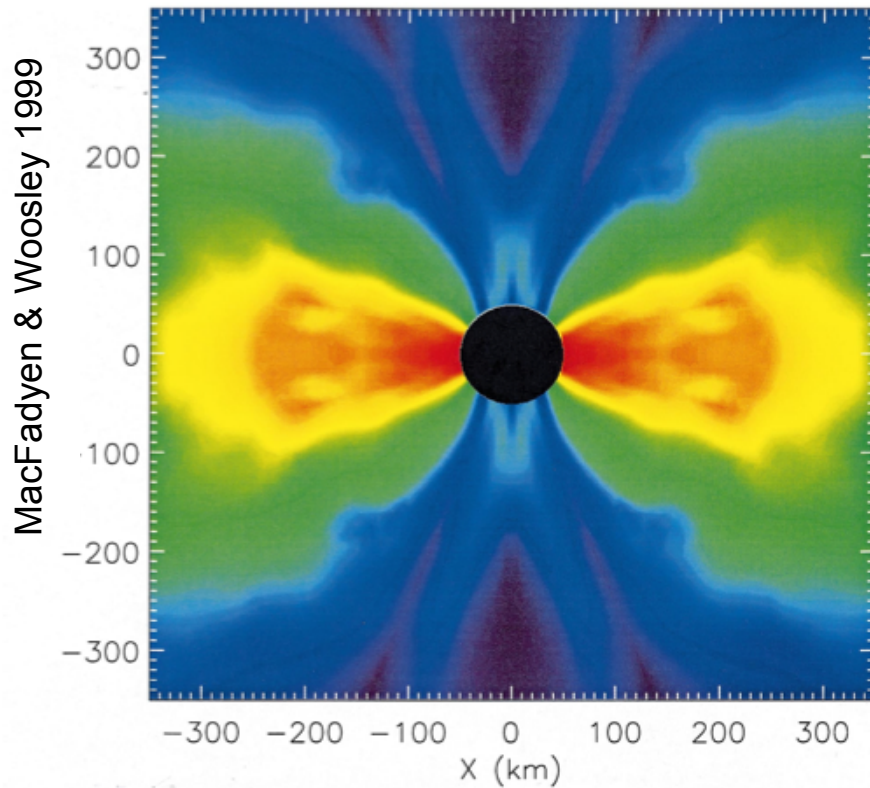
The Fates of Massive Stars (Heger et al. 2003)



Assumes neutrino-powered supernova with energy $\sim 10^{51}$ ergs

The Collapsar (Failed Supernova) Model (Woosley 93)

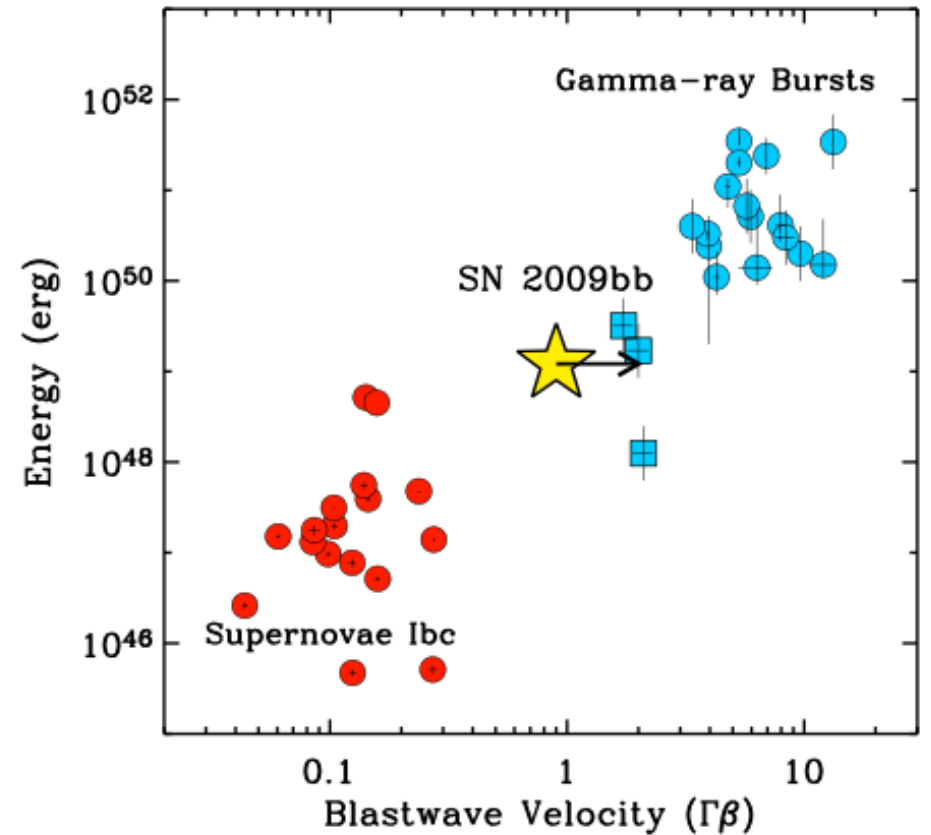
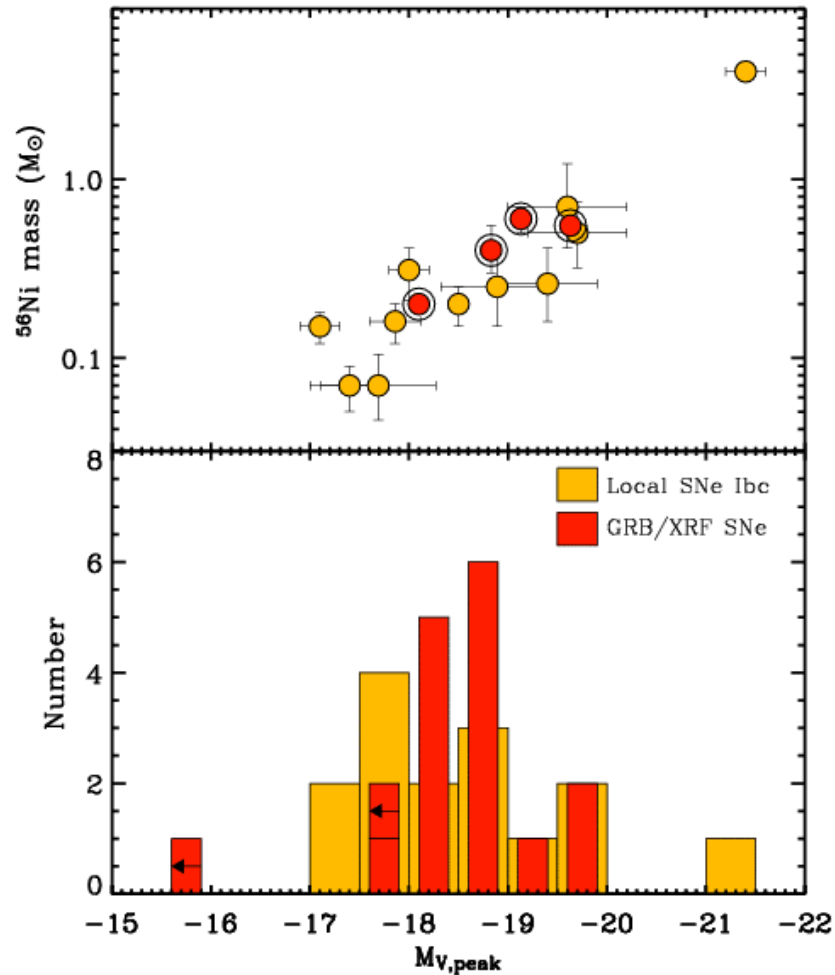
(e.g. Aloy et al. 2000; MacFadyen et al. 2001; Proga & Begelman 2003; Takiwaki et al. 2008; Barkov & Komissarov 2008; Nagataki et al. 2007; Lindler et al. 2010)



- Energy -
- Duration -
- Hyper-Energetic SNe -
- Late-Time Activity -

Accretion / Black Hole Spin
Stellar Envelope In-Fall
Delayed Black Hole Formation
or Accretion Disk Winds
Fall-Back Accretion

What Distinguishes GRB Supernovae?



Soderberg et al. 2006, 2007, 2009

“GRB-SNe are **not** clearly distinguished from ordinary SNe Ibc either by optical luminosity or photospheric velocities.”

Core Collapse with Magnetic Fields & Rotation

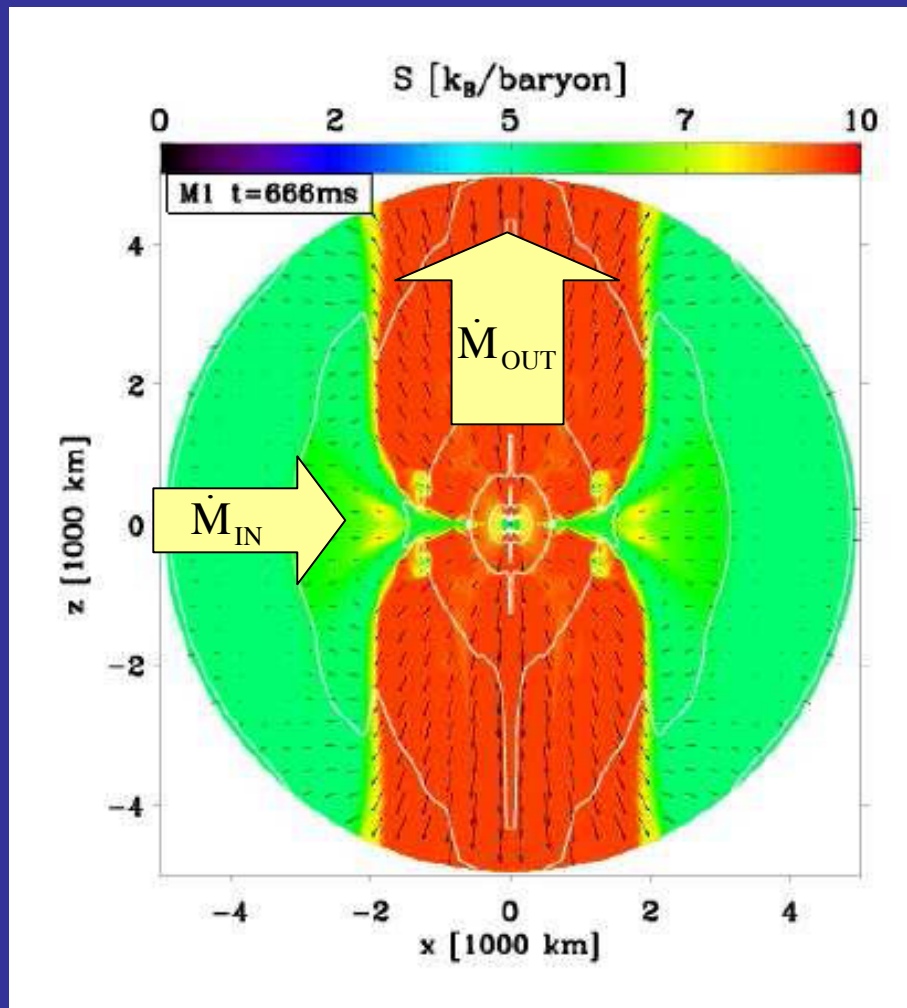
(e.g. LeBlanc & Wilson 1970; Bisnovaty-Kogan 1971; Akiyama et al. 2003)

THE PROTO-NEUTRON STAR PHASE OF THE COLLAPSAR MODEL AND THE ROUTE TO LONG-SOFT GAMMA-RAY BURSTS AND HYPERNOVAE

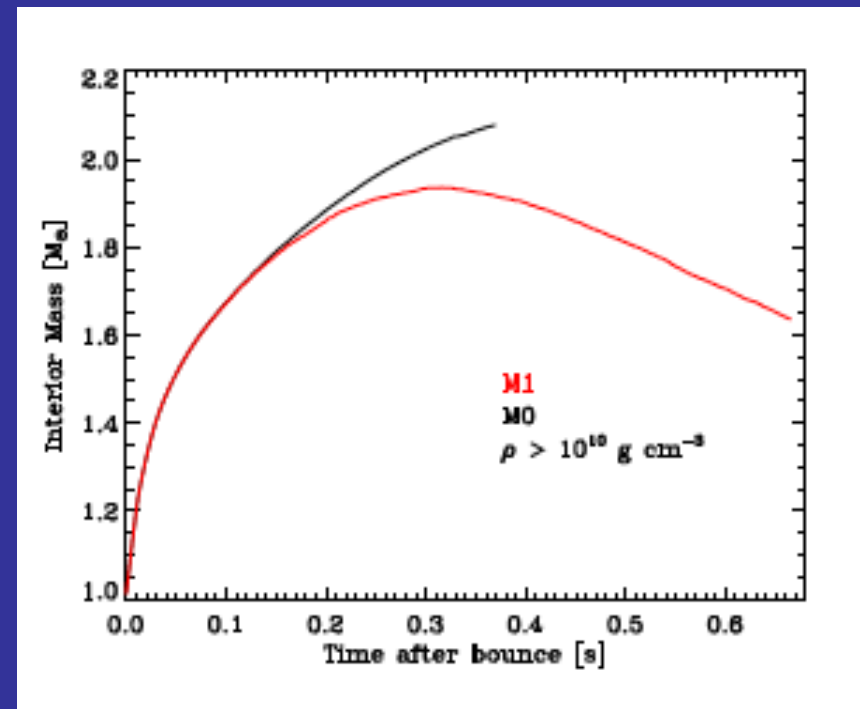
L. DESSART¹, A. BURROWS¹, E. LIVNE², AND C.D. OTT¹

Collapsar Requirements:

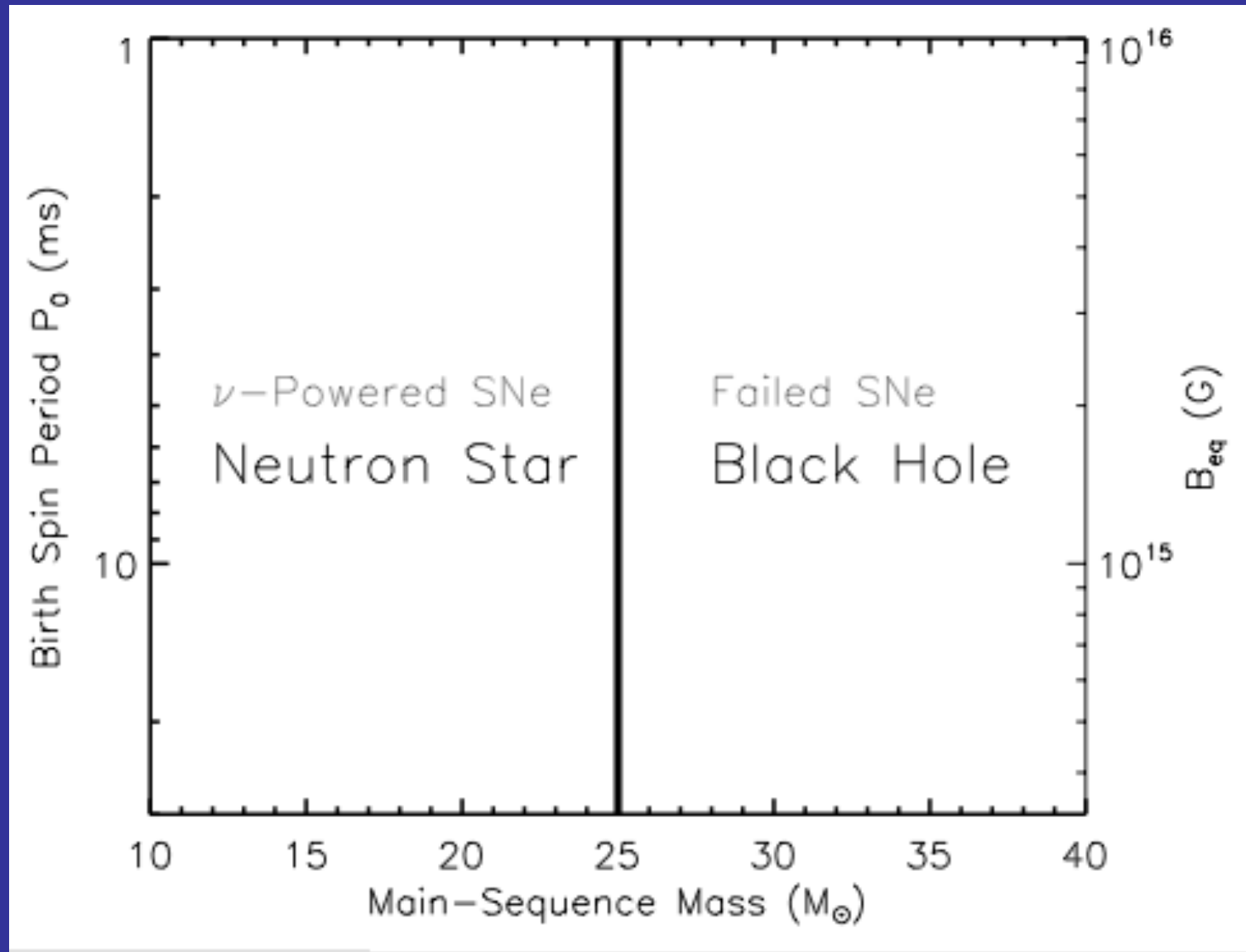
- Angular Momentum
- Strong, Ordered Magnetic Field
(e.g. Proga & Begelman 2003; McKinney 2006)



Neutron Star Mass



Heger et al. Black Hole-Neutron Star Dichotomy (at Sub-Solar Metallicity)



Millisecond Magnetar Model (Usov 92; Thompson 94)

$$E_{\text{Rot}} \approx 2 \times 10^{52} \left(\frac{P}{1 \text{ ms}} \right)^{-2} \text{ ergs}$$

$$\dot{E} \approx 10^{49} \left(\frac{P}{1 \text{ ms}} \right)^{-4} \left(\frac{B_{\text{Dip}}}{10^{15} \text{ G}} \right)^2 \text{ ergs s}^{-1}$$

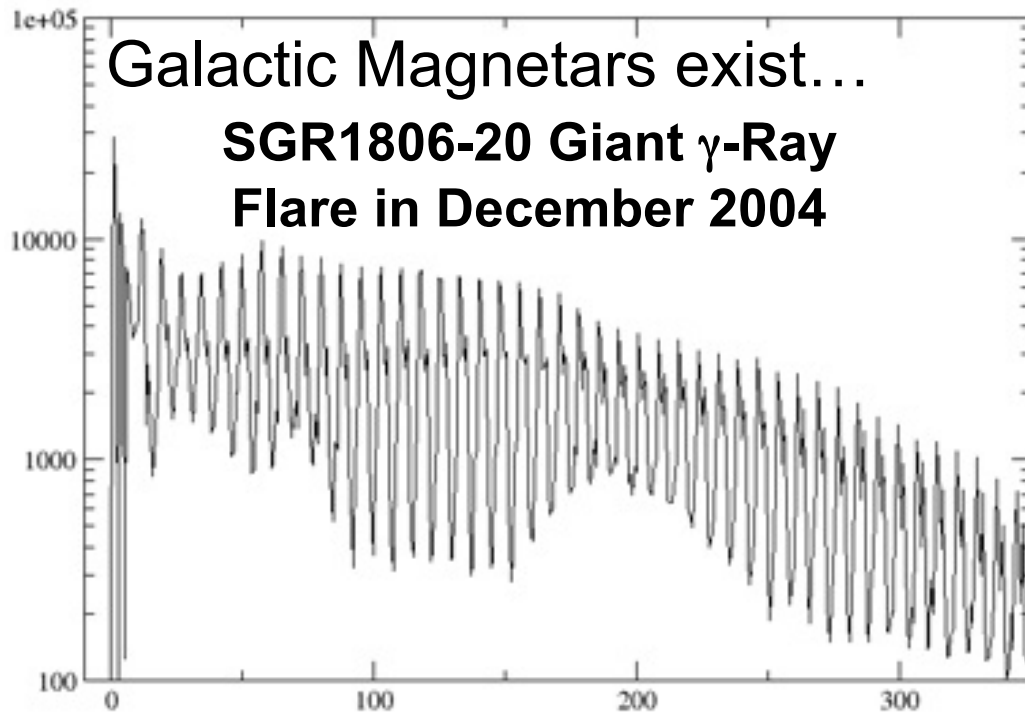
- Rapid Rotation \Leftrightarrow Efficient α - Ω Dynamo \Leftrightarrow Strong B-Field at P \sim 1 ms
(Duncan & Thompson 1992; Thompson & Duncan 1993)

Millisecond Magnetar Model (Usov 92; Thompson 94)

$$E_{\text{Rot}} \approx 2 \times 10^{52} \left(\frac{P}{1 \text{ ms}} \right)^{-2} \text{ ergs}$$

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- Rapid Rotation \leftrightarrow Efficient α - Ω Dynamo \leftrightarrow Strong B-Field at $P \sim 1 \text{ ms}$
(Duncan & Thompson 1992; Thompson & Duncan 1993)



...and can have massive progenitors

Magnetar

Westerlund I: O7 Stars still present!

Muno +06

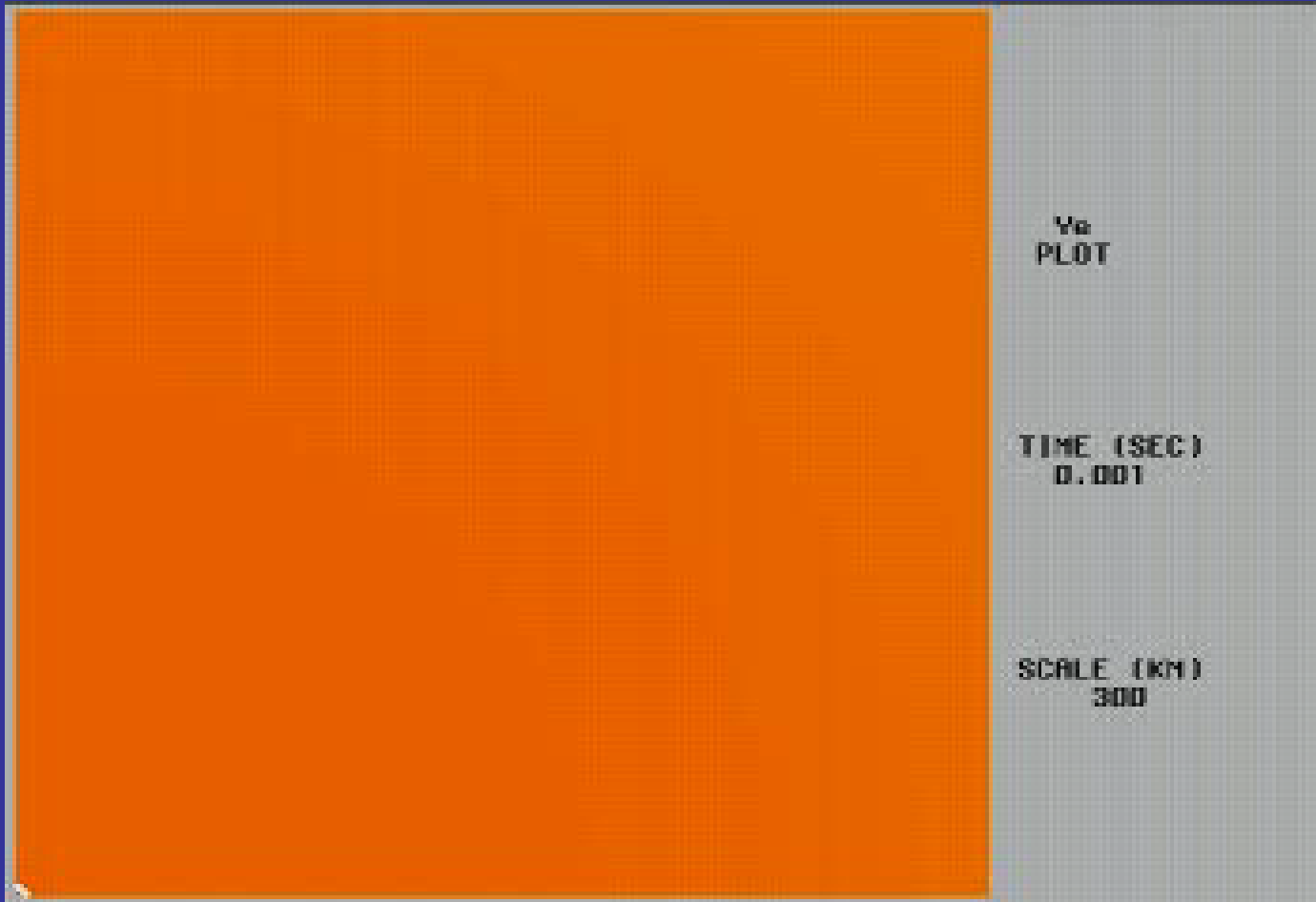
Key Insight :

(Thompson, Chang & Quataert 04)

Neutron Stars are Born **Hot**,
Cool via ν -Emission:
 $\sim 10^{53}$ ergs in $\tau_{\text{KH}} \sim 10\text{-}100$ s

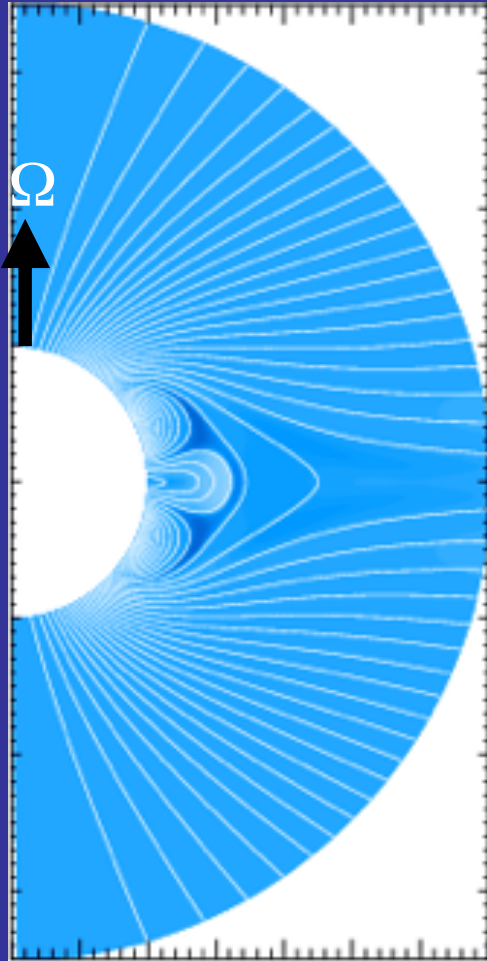
- Neutrinos Heat Proto-NS Atmosphere (e.g. $\nu_e + n \Rightarrow p + e^-$)
⇒ **Drives Thermal Wind Behind SN Shock** (e.g. Qian & Woosley 96)

Burrows, Hayes, & Fryxell 1995



Effects of Strong Magnetic Fields

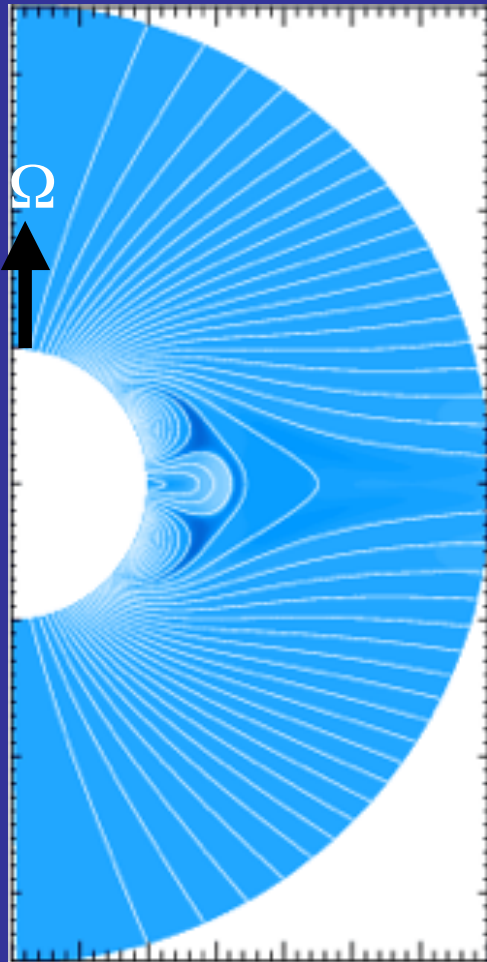
“Helmet - Streamer”



- Microphysics (EOS, ν Heating & Cooling)
 - Important for $B \geq 10^{16}$ G (Duan & Qian 2005)

Effects of Strong Magnetic Fields

“Helmet - Streamer”



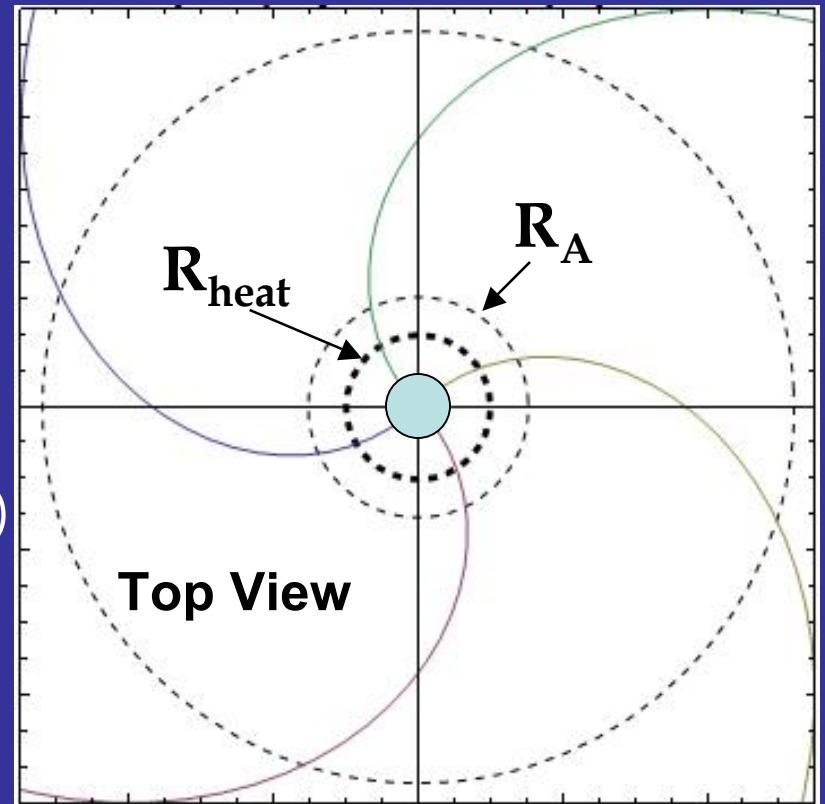
- **Microphysics (EOS, ν Heating & Cooling)**
 - Important for $B \geq 10^{16}$ G (Duan & Qian 2005)
- **Magneto-Centrifugal Outflows**
(Weber & Davis 1967)

Outflow Co-Rotates
with Neutron Star when

$$\frac{B^2}{8\pi} > \frac{1}{2}\rho v_r^2$$

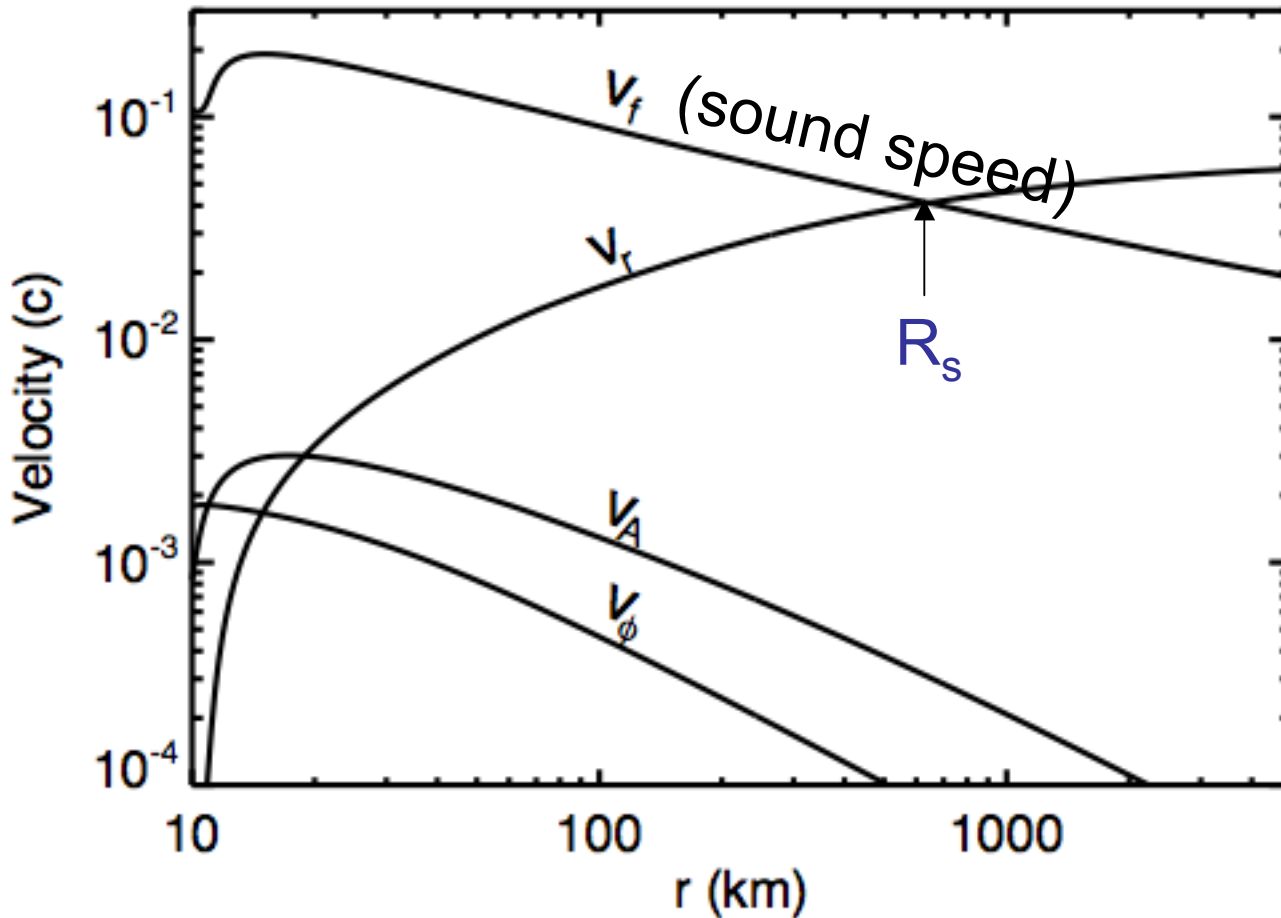
⇒

- **Magneto-Centrifugal Acceleration**
 (“Bead on a Wire”)
- **Enhanced Mass Loss Rate**



Example Thermally-Driven Wind

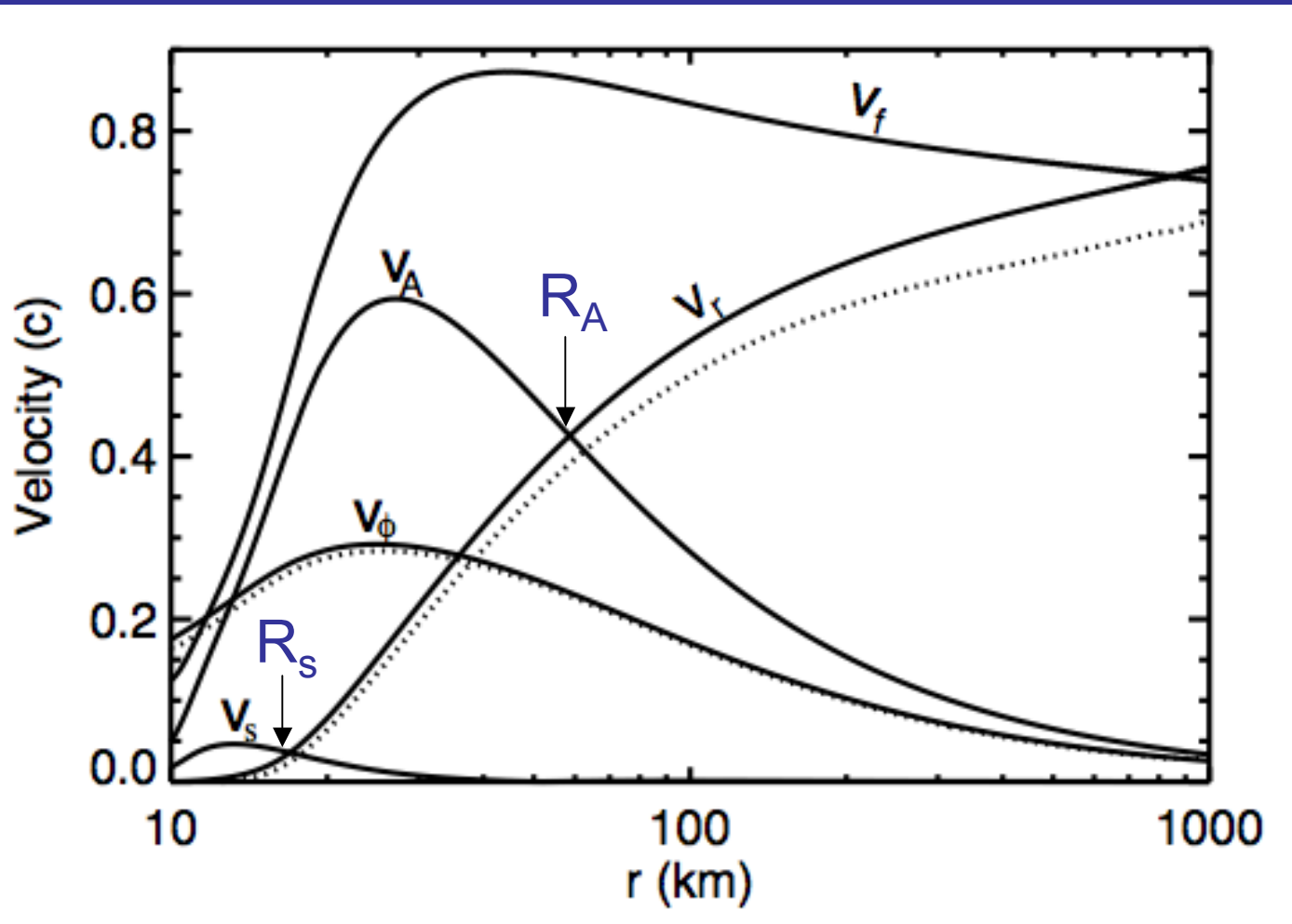
$$L_{\bar{\nu}_e} \sim 8 \times 10^{51} \text{ ergs s}^{-1}; \quad B_0 = 10^{13} \text{ G}; \quad P = 100 \text{ ms}$$



$$\dot{M} \sim 10^{-4} M_\odot \text{ s}^{-1}$$

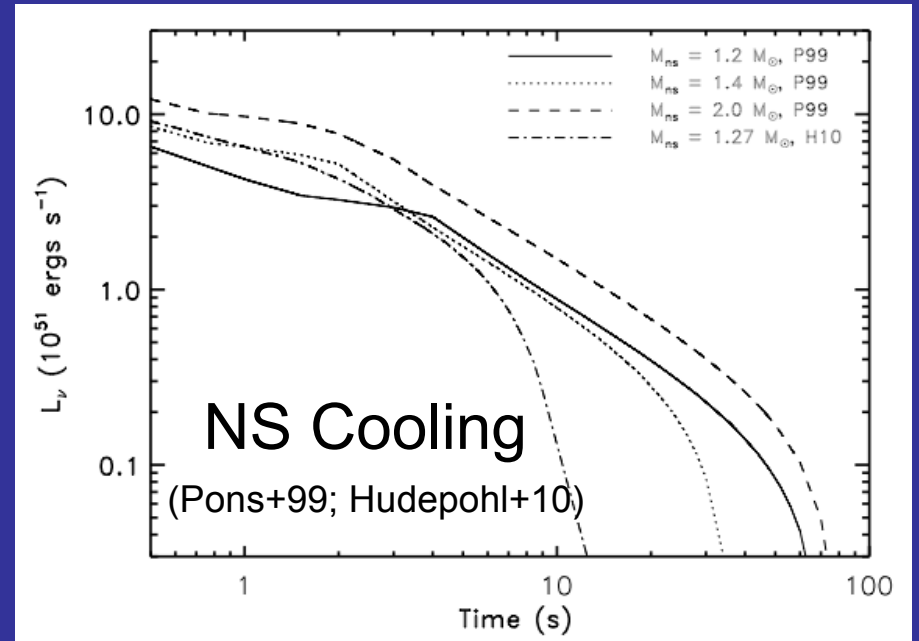
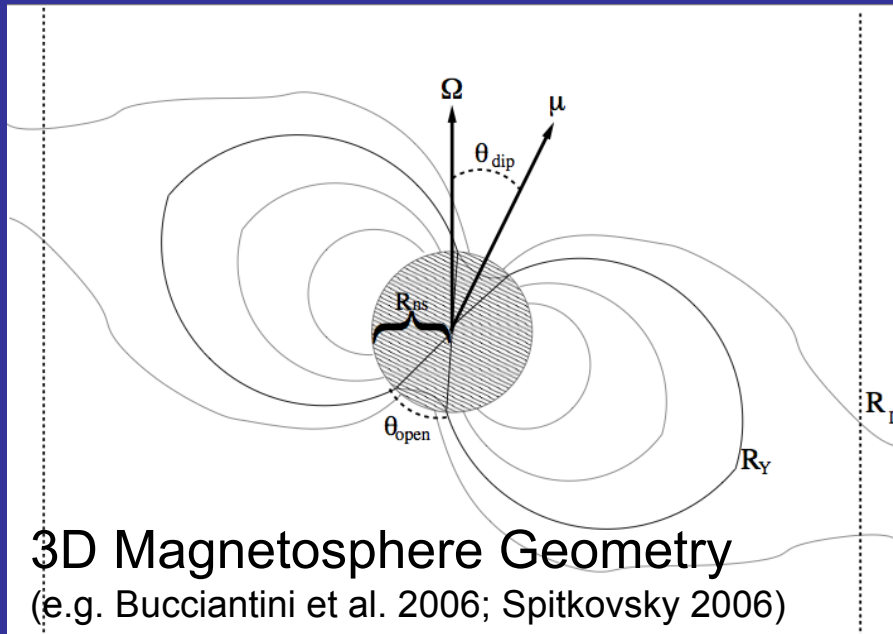
Magnetically-Driven Wind

$$L_{\bar{v}_e} \sim 8 \times 10^{51} \text{ ergs s}^{-1}; B_0 = 10^{15} \text{ G}; P = 1.2 \text{ ms}$$



$$\dot{M} \sim 3 \times 10^{-3} M_\odot \text{ s}^{-1}$$

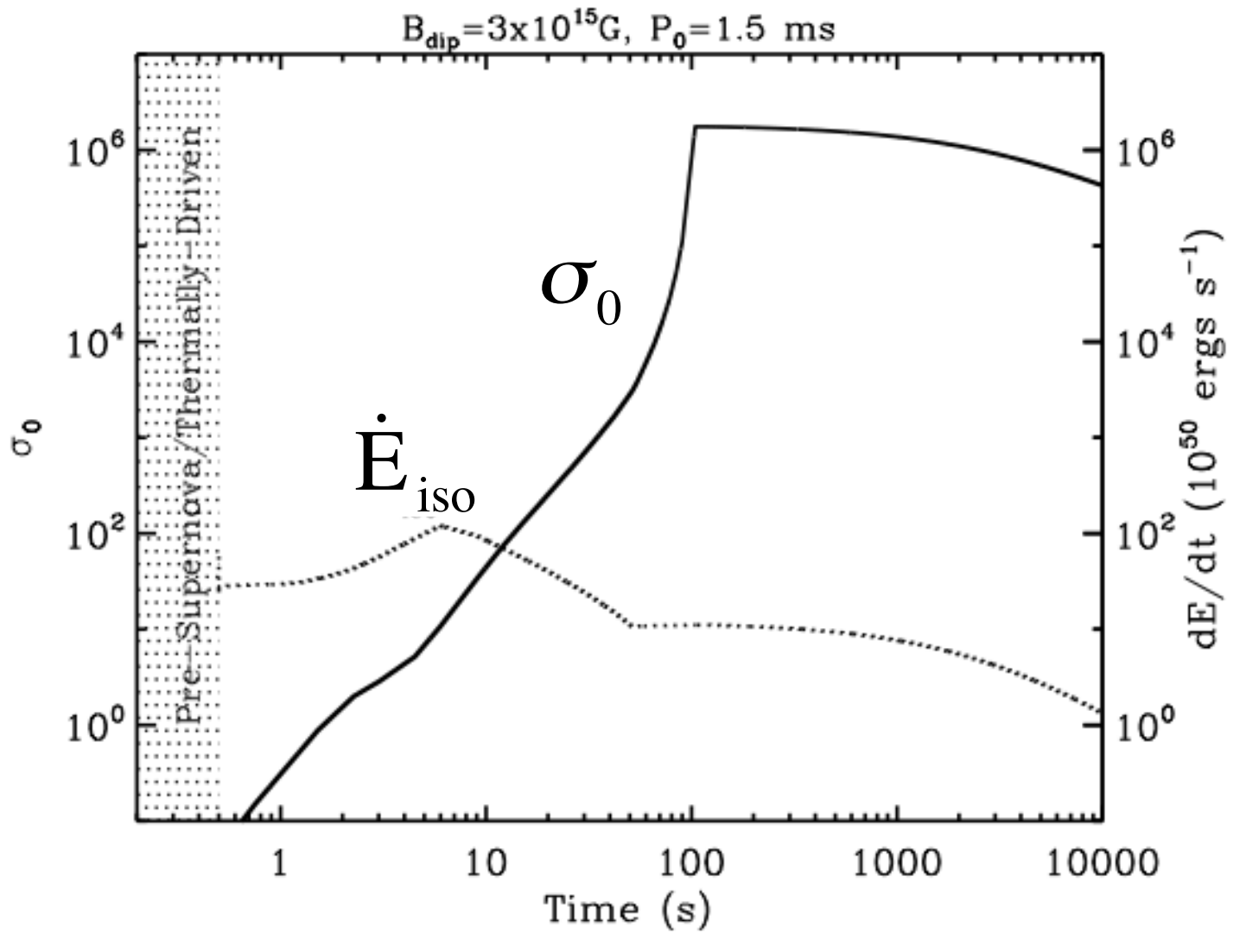
Evolutionary Wind Models (BDM et al. 2007, 2008, 2010)



Calculate: Wind Power $\dot{E}(t)$, Mass Loss Rate $\dot{M}(t)$,
 \Rightarrow 'Magnetization' $\sigma(t) \sim \frac{\dot{E}}{\dot{M}c^2} = \Gamma_{\max}(t)$

In terms of

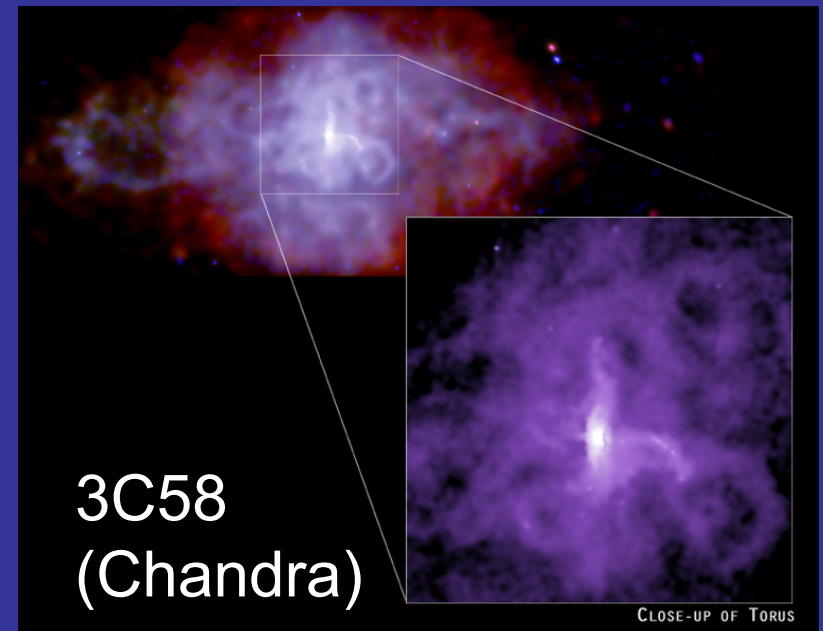
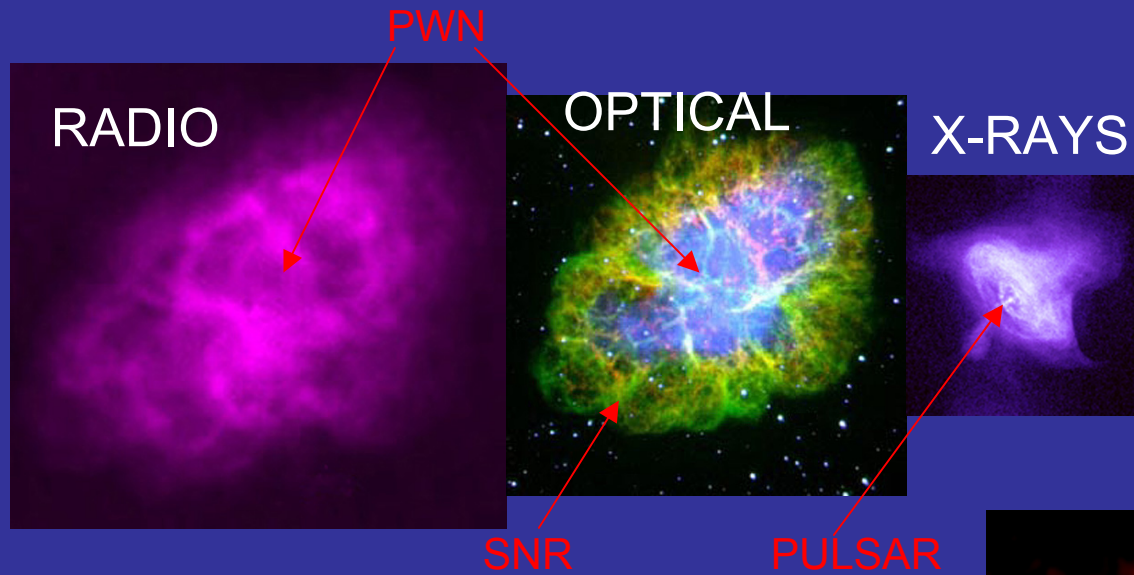
Initial Rotation Period P_0 , Dipole Field Strength B_{dip} & Obliquity θ_{dip}



$$\sigma \sim \Gamma_{\text{max}} = \frac{\dot{E}}{\dot{M}c^2} \propto \frac{B^2 \Omega^4}{L_v^{5/3} T^{10/3}}$$

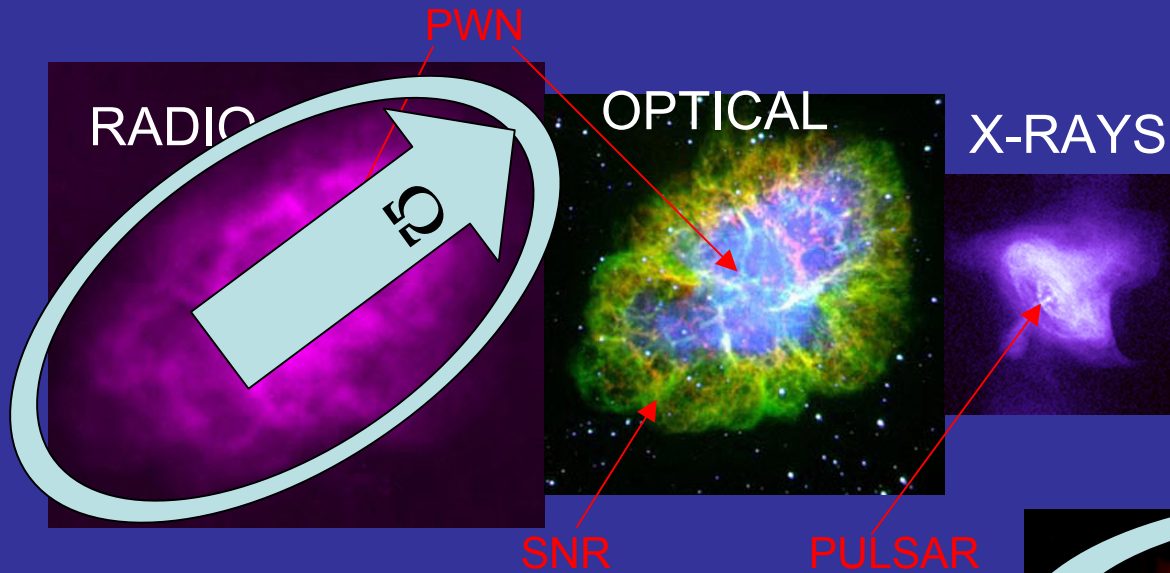
Collimation via Stellar Confinement

Multi-Wavelength Crab Nebula

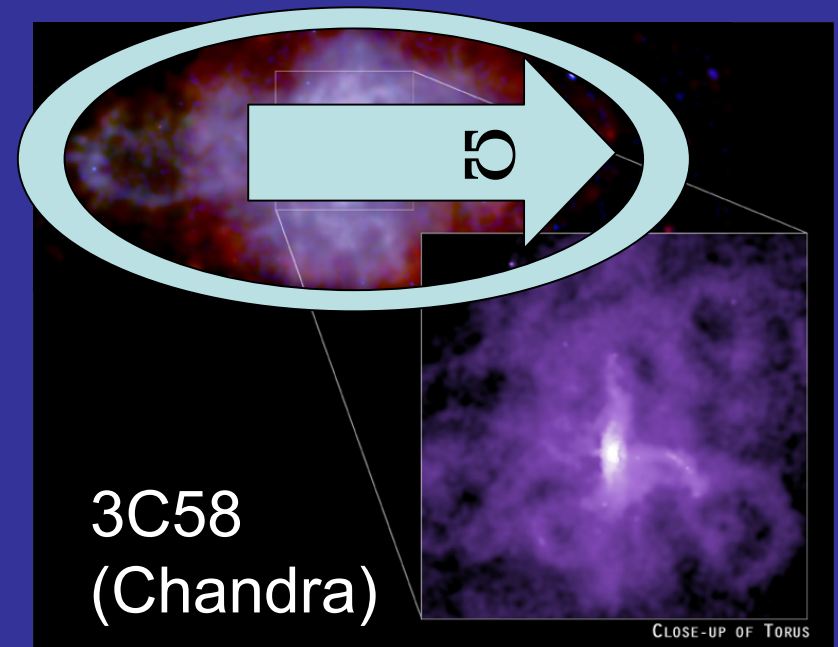


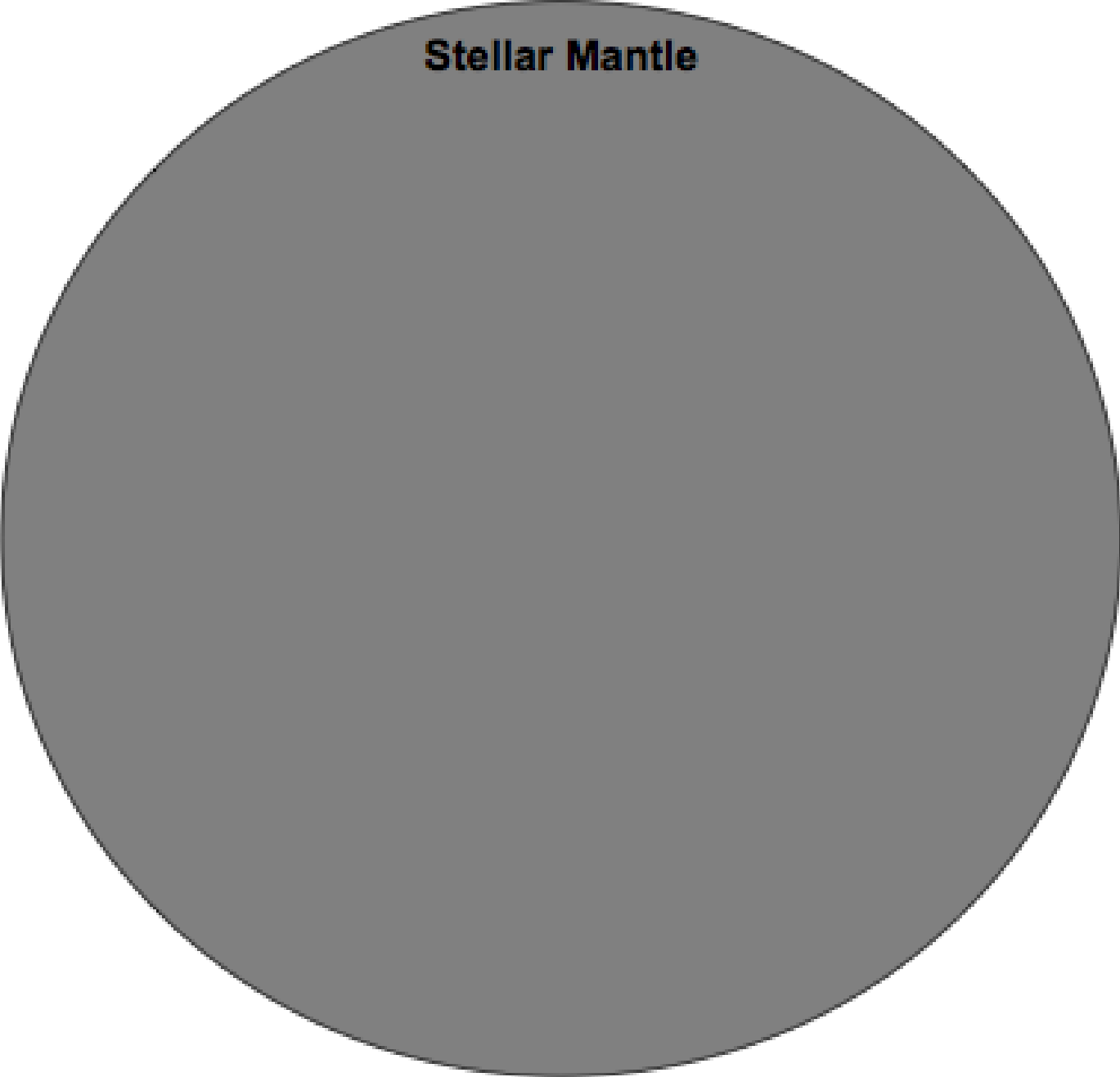
Collimation via Stellar Confinement

Multi-Wavelength Crab Nebula



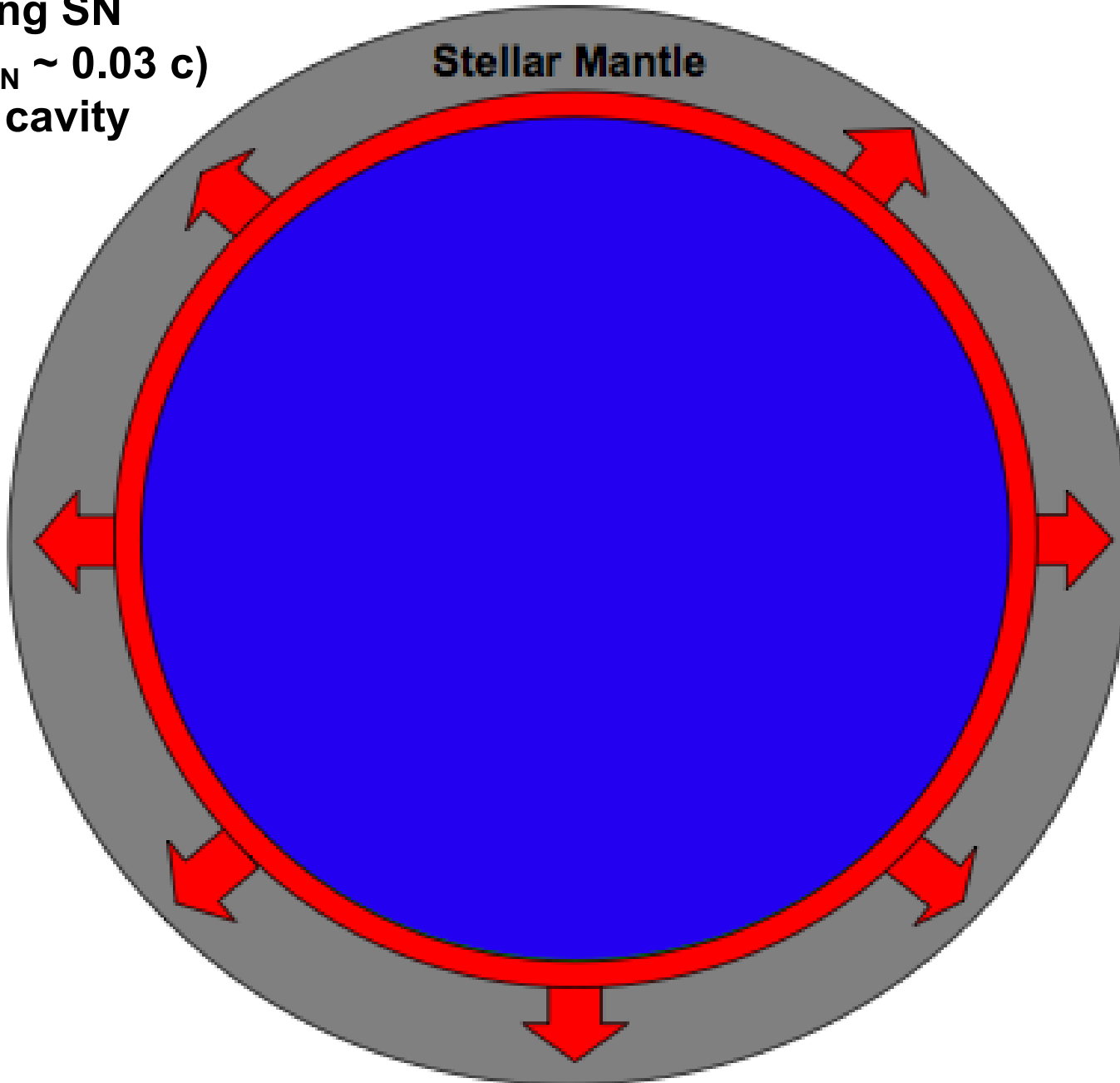
Magnetic Stresses in
Pulsar Nebula
Elongates SN Remnant!
(Begelman & Li 1992)





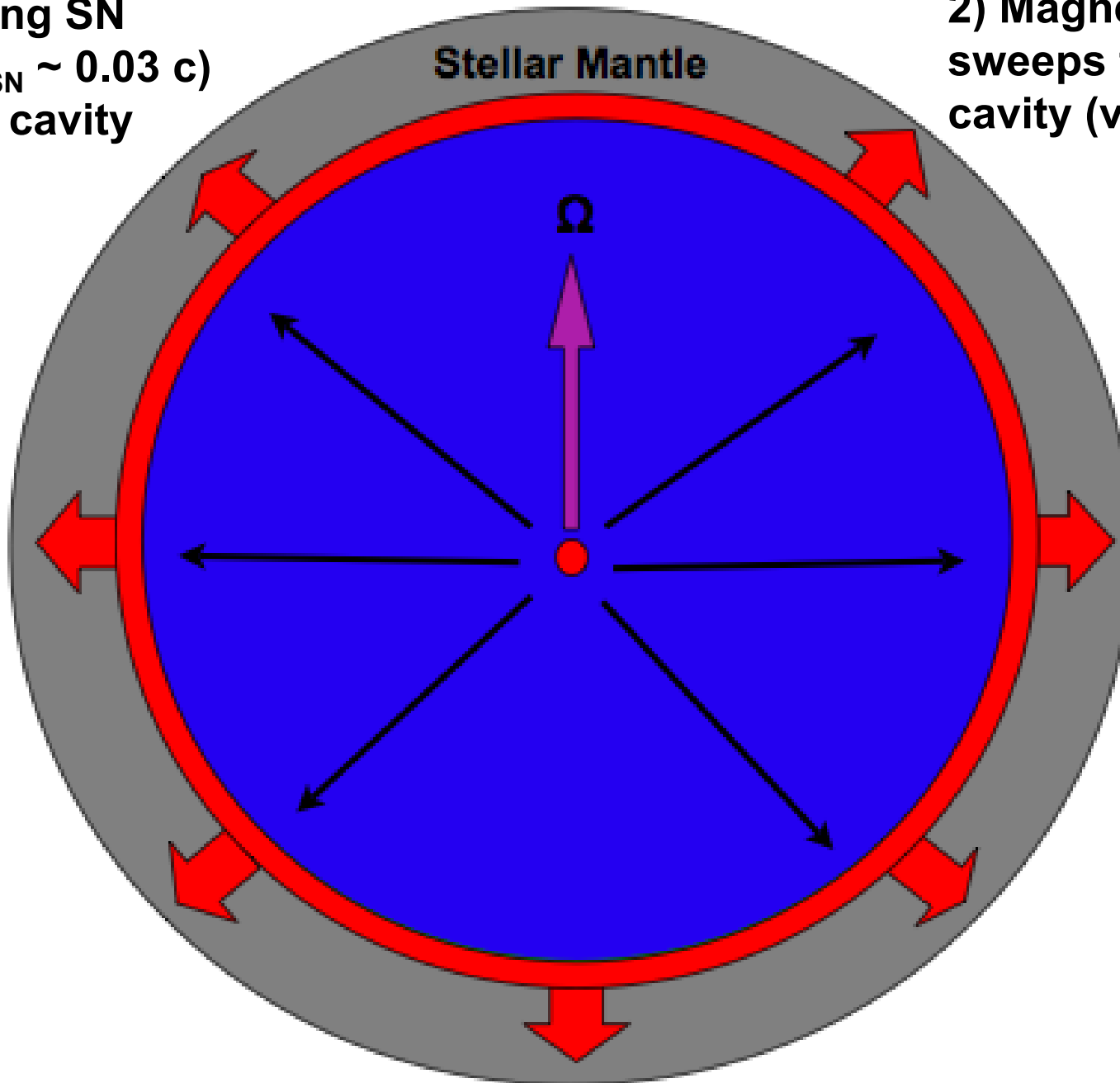
Stellar Mantle

1) Outgoing SN
shock ($v_{\text{SN}} \sim 0.03 c$)
creates a cavity



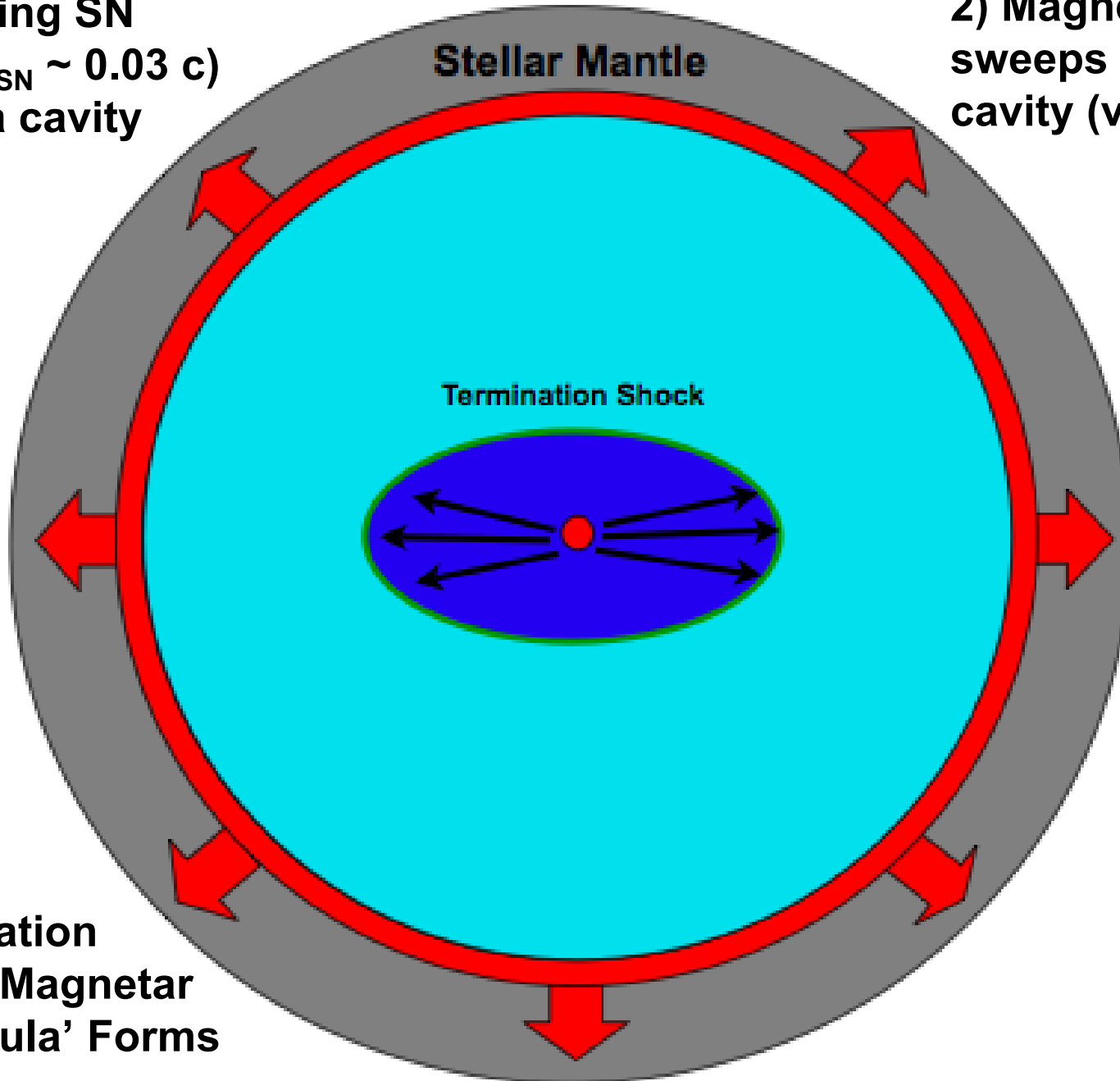
1) Outgoing SN shock ($v_{\text{SN}} \sim 0.03 c$) creates a cavity

2) Magnetar wind sweeps through cavity ($v_w \sim c$)



1) Outgoing SN shock ($v_{\text{SN}} \sim 0.03 c$) creates a cavity

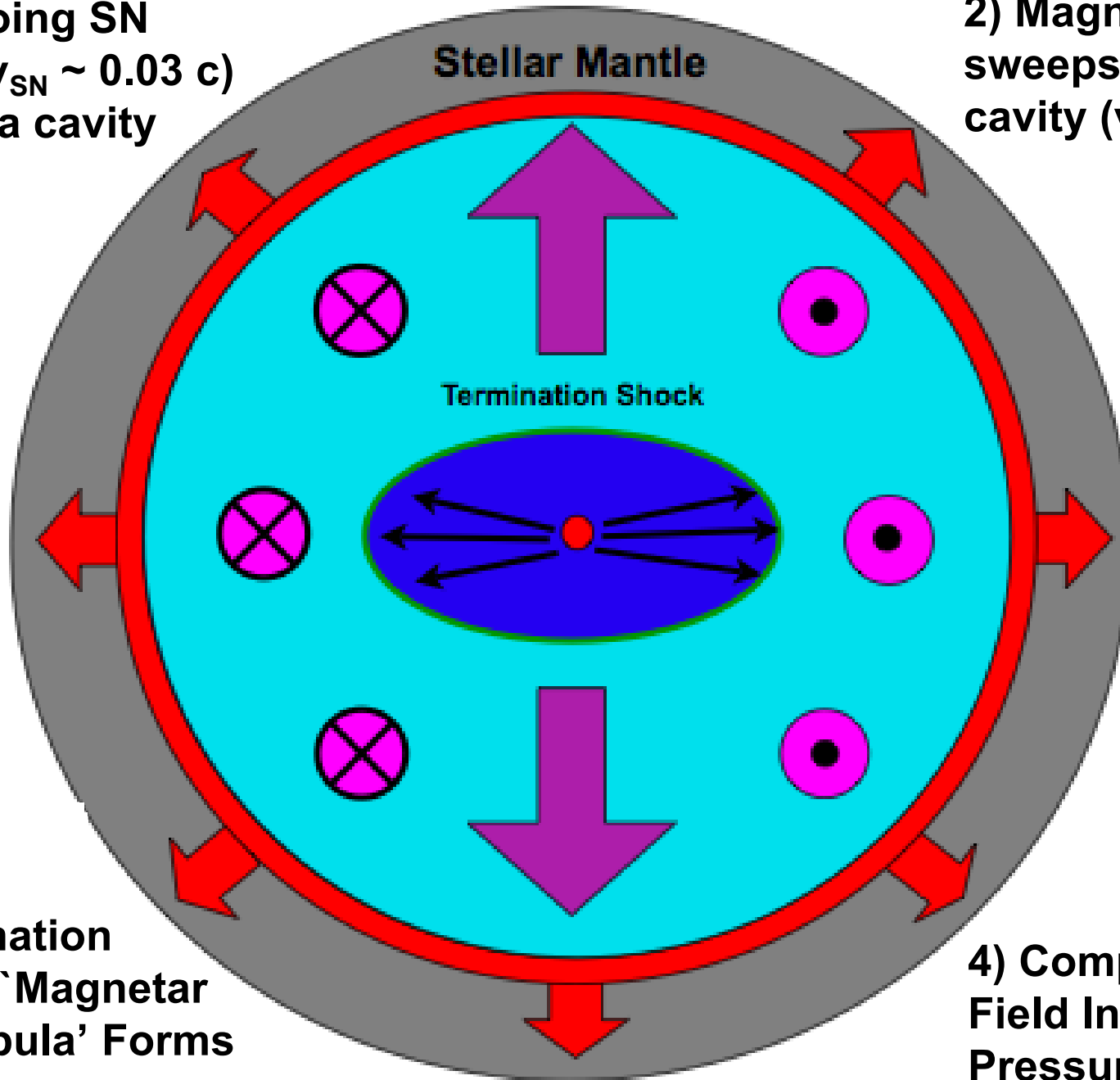
2) Magnetar wind sweeps through cavity ($v_w \sim c$)



3) Termination Shock & 'Magnetar Wind Nebula' Forms

1) Outgoing SN shock ($v_{SN} \sim 0.03 c$) creates a cavity

2) Magnetar wind sweeps through cavity ($v_W \sim c$)

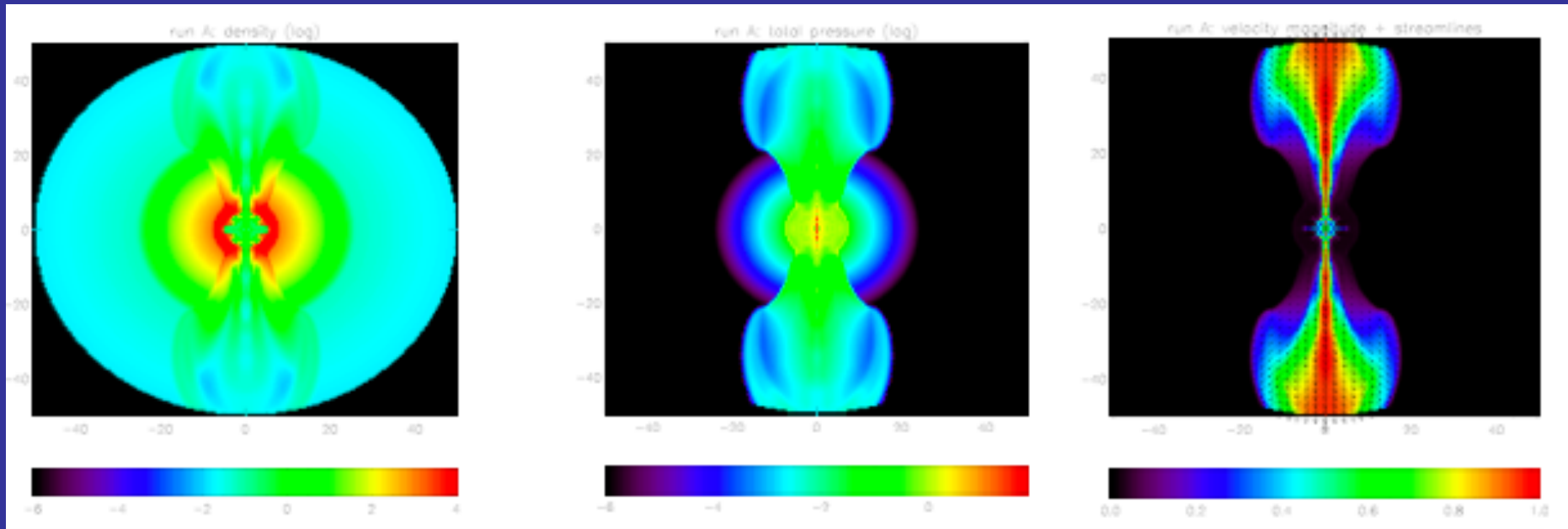


3) Termination Shock & 'Magnetar Wind Nebula' Forms

4) Compressed Field Increases Pressure on Axis

Proto-Magnetar Jet Formation

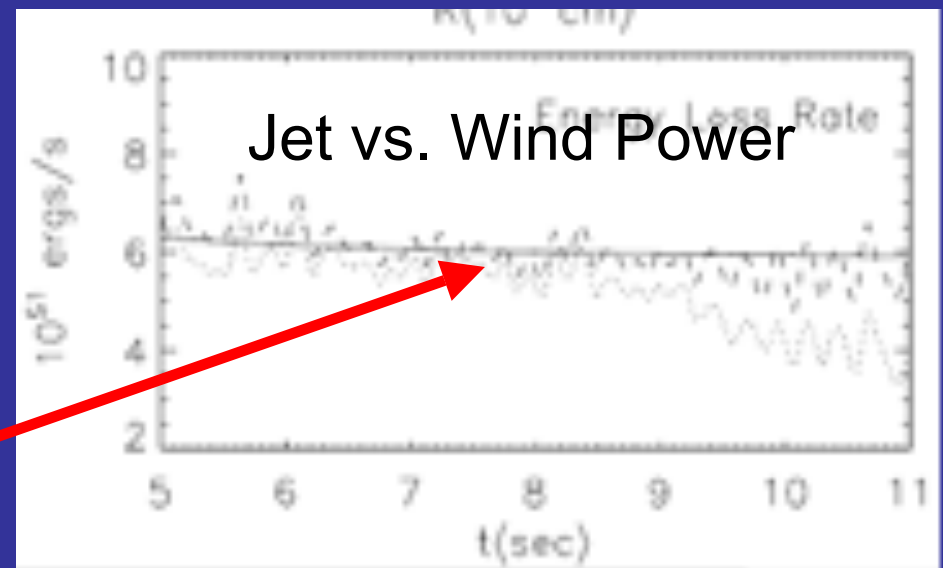
(Bucciantini et al. 2007, 2008, 2009)

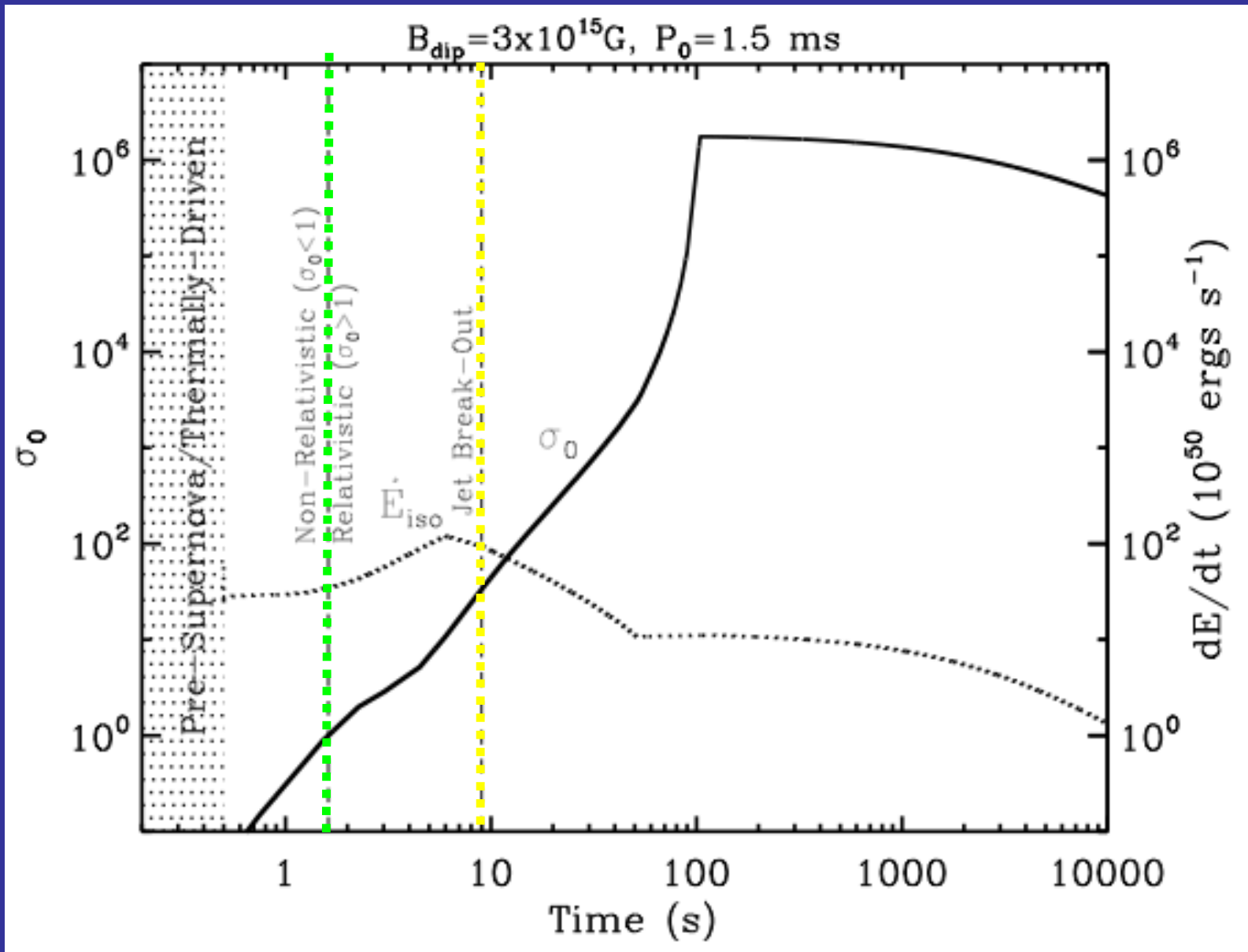


➤ Assume Successful Supernova
(35 M_{\odot} ZAMS Progenitor; Woosley & Heger 06)

➤ Inner BC from Proto-Magnetar Wind
Calcs for $B_{\text{dip}} = 3 \times 10^{15}$ G and $P_0 = 1$ ms

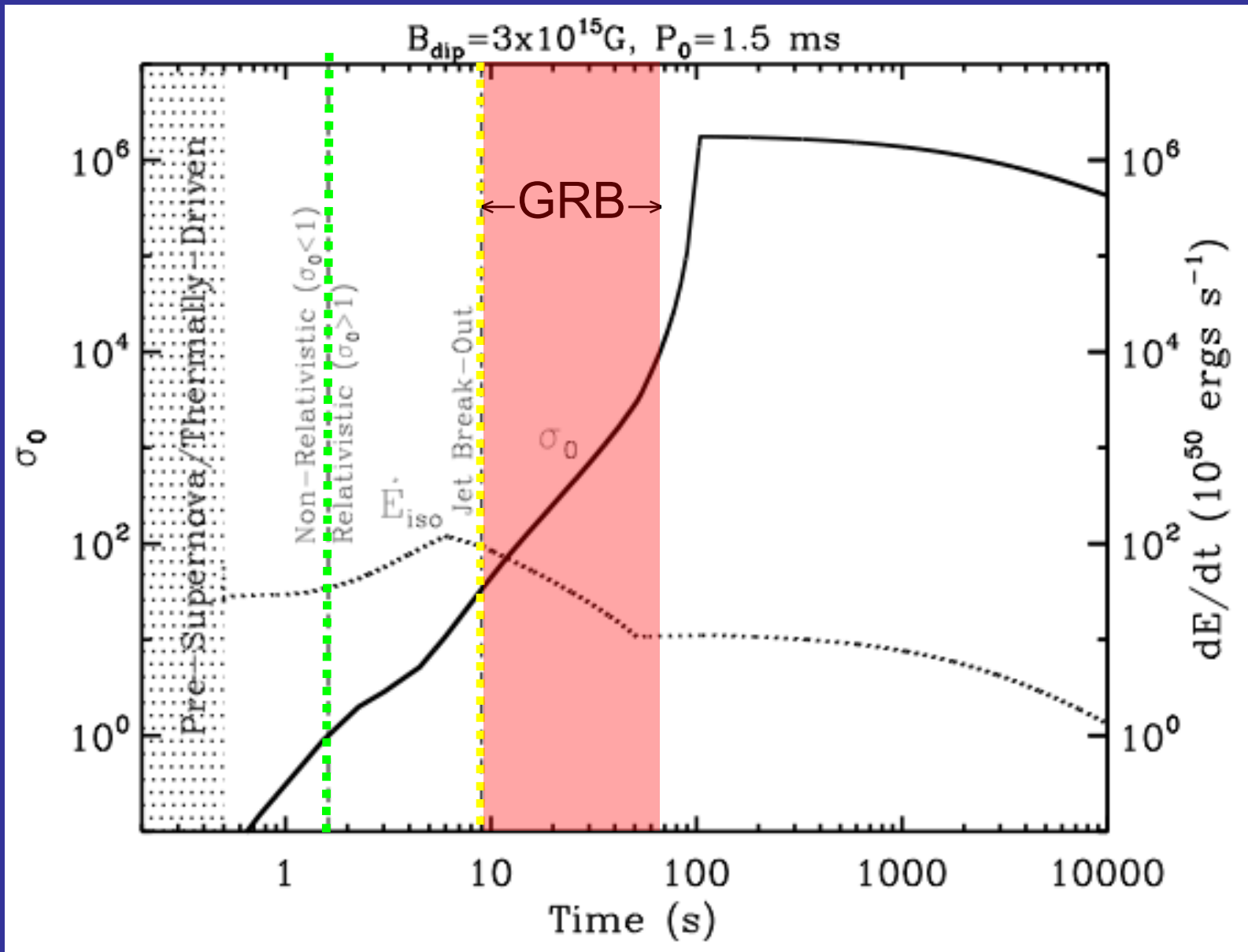
**Average Power and Mass Loss
Rate of Jet Leaving Star Match
Those Set by Magnetar Wind**





Wind becomes relativistic at $t \sim 2$ seconds;

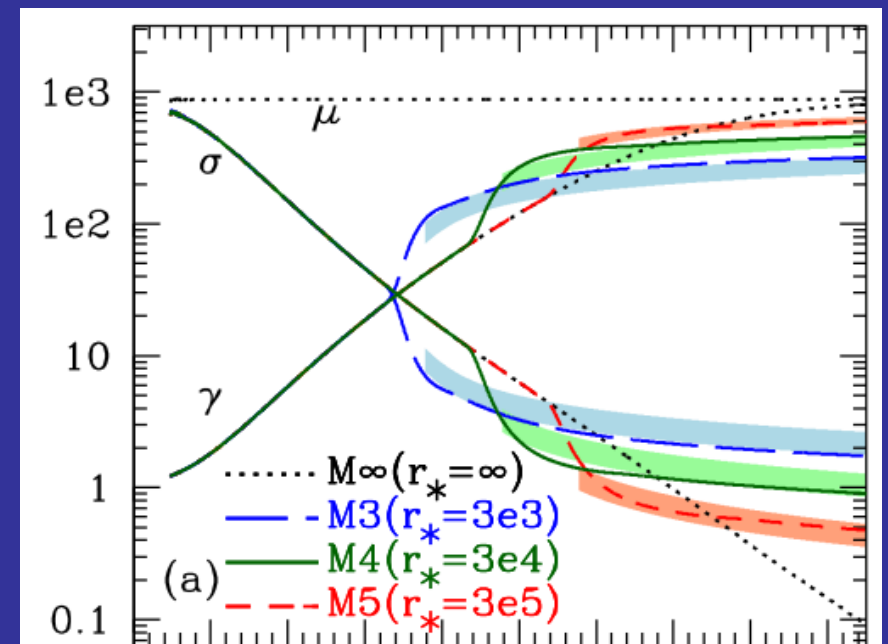
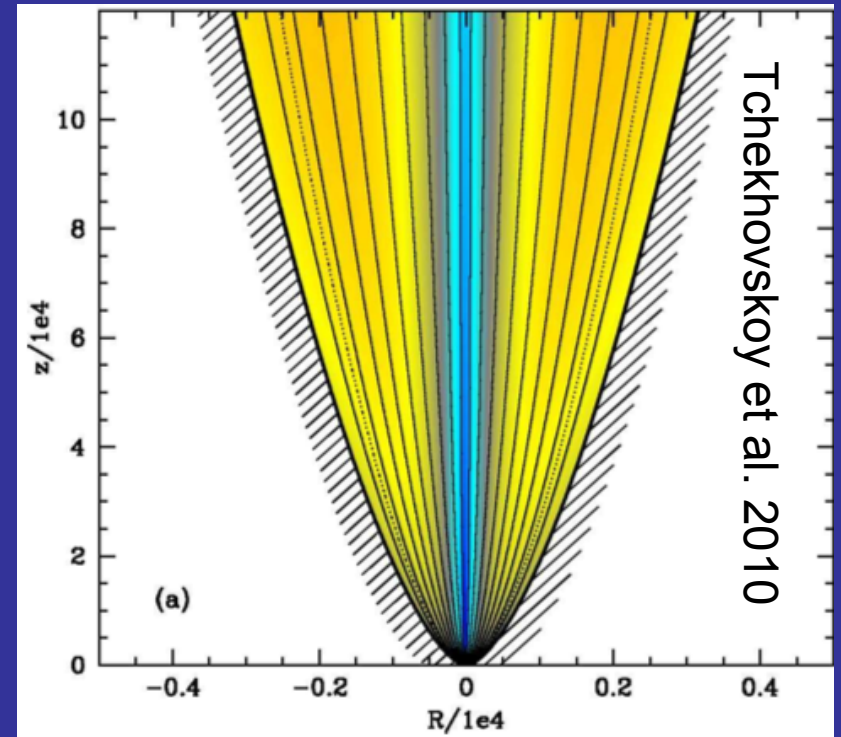
Jet breaks out of star at $t_{\text{bo}} \sim R_{\star} / \beta c \sim 10$ seconds



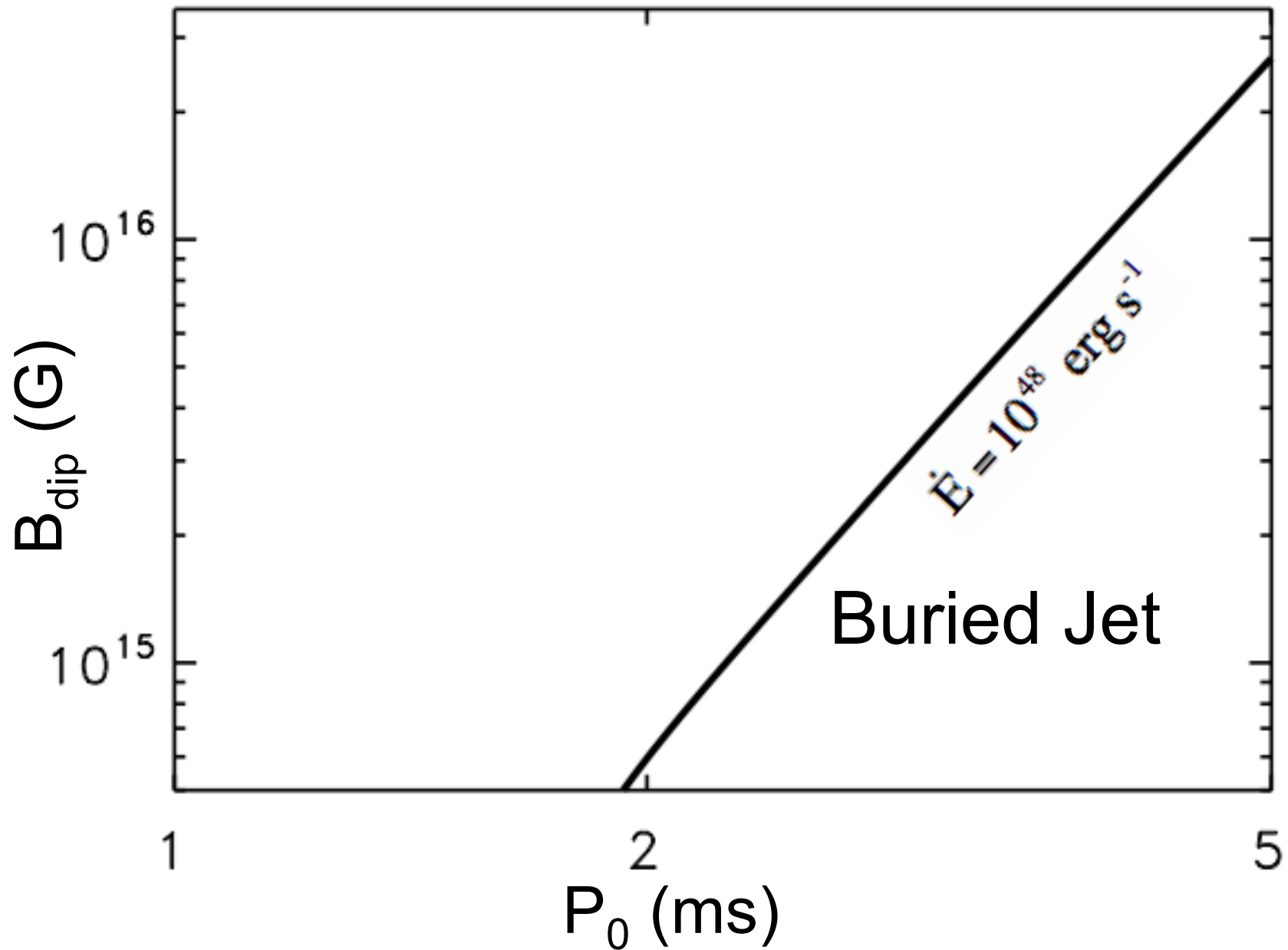
High Energy Emission (GRB) from $t \sim 10$ to ~ 100 s as Magnetization Increases from $\sigma_0 \sim \Gamma \sim 30$ to $\sim 10^3$

Acceleration in GRB Jets

- Jet's energy at small radii is mostly carried by magnetic field (Poynting flux)
- Magnetic energy must be transferred to kinetic energy to reach $\Gamma_\infty \sim \sigma_0$
- Time stationary, unconfined outflows in ideal MHD attain $\Gamma_\infty \sim \sigma_0^{1/3} \ll \sigma_0$ (e.g. Goldreich & Julian 1970)
- Proposed Solutions -
 - 1) converging (parabolic) geometry
 - 2) time-variable outflow
 - 3) non-ideal MHD (e.g. reconnection)
- All predict power law acceleration $\Gamma \propto R^\alpha$ ($\alpha < 1$) with max Lorentz factor $\Gamma_{\max} \sim 10^2 - 10^3$



Magnetar Birth - A Variety of Phenomena



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