Wγ production at the LHC at NLOPS accuracy*

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* arXiv:1408.5766

Outline

★ Basic motivations

- Precision test of the Standard Model
- New physics searches

★ Method

- Simulation of the photon fragmentation contribution using POWHEG+PYTHIA
- Practical implementation

★ Results

- Comparisons with MCFM results
- Comparison with experimental data at 7 TeV (ATLAS 2011)

Electroweak boson pairs production at the LHC

Anomalies in triple gauge couplings

Search for ATGC: high pT(γ) region in exclusive Njet=0 selection



No deviations found from SM predictions



- New Physics phenomena
 - Background for SuperSymmetry processes
 - Search for narrow resonances: techni-mesons ($\omega_T \rightarrow Z\gamma$, $a_T \rightarrow W\gamma$)

LO Feynman diagrams for Wy production



QCD NLO corrections to Wy

NLO corrections are very sizeable K-factor ~ 2 for some event selections*



- Gluon-induced processes strongly enhanced at the energies of the LHC (large gluon PDFs)
- Divergent for collinear photon emission (no virtual counterpart)

* J. M. Campbell et al., Vector boson pair production at the LHC, JHEP 1107 (2011)

Isolated photons in hadronic collisions (I)



- Higher orders in QCD: a series of consecutive collinear splittings from a high- p_T parton, ending up with a quark-photon splitting.
 - Collinear singularities factorized to all orders in as and absorbed into photon fragmentation functions*

 $D_{q \to \gamma}(z, M_F)$ $D_{g \to \gamma}(z, M_F)$ Fragmentation scale Longitudinal momentum fraction

* S. Catani, M. Fontannaz, J. Ph. Guillet and E. Pilon, Cross section of isolated prompt photons in hadron-hadron collisions, JHEP 0205 (2002)

Isolated photons in hadronic collisions (II)

Direct and fragmentation contributions (depending on *M_F*)

$$\sigma(p_{\gamma}) = \sigma_{\gamma}^{(D)}(p_{\gamma}; M_F) + \sum_{a} \int_0^1 \frac{dz}{z} \sigma_a^{(F)}(p_{\gamma}/z; M_F) D_{a \to \gamma}(z, M_F)$$

For $M_F > 1$ GeV, fragmentation contribution $O(a_{em}/a_s(M_F))$

Same order of the Born term

- Non perturbative contributions arising in parton fragmentation
 - Fragmentation functions extracted from experimental data
 - Large uncertainties
 - Pure collinear approximation

Isolated photons in hadronic collisions (III)

Smooth isolation prescription*

$$\sum_{j} E_{Tj} \Theta(R - R_{j\gamma}) \le \epsilon_h E_{T\gamma} \left(\frac{1 - \cos R}{1 - \cos R_0} \right)$$
$$\forall R < R_0$$



- Fragmentation contribution fully removed
- Soft radiation can be integrated over the whole phase space. Cancellation of QCD infrared singularities in inclusive R+V not spoiled
- Difficult experimental application

 $\sum_{j} E_{Tj} \Theta(R_0 - R_{j\gamma}) \le \epsilon_h E_{T\gamma}$

* S. Frixione, Isolated photons in perturbative QCD, Phys.Lett. B429 (1998)

The POWHEG method*

Cross section for the hardest emission

$$\begin{split} d\sigma &= \sum_{f_b} \bar{B}^{f_b}(\Phi_n) d\Phi_n \bigg\{ \Delta^{f_b}(\Phi_n, p_T^{min}) + \sum_{\alpha_r \in \{\alpha_r \mid f_b\}} \frac{\left[d\Phi_{rad} \Theta(k_T - p_T^{min}) \Delta^{f_b}(\Phi_n, k_T) R(\Phi_{n+1}) \right]_{\alpha_r}^{\Phi_n^{\alpha_r} = \Phi_n}}{B^{f_b}(\Phi_n)} \bigg\} \\ \bar{B}^{f_b}(\Phi_n) &= \left[B(\Phi_n) + V(\Phi_n) \right]_{f_b} + \sum_{\alpha_r \in \{\alpha_r \mid f_b\}} \int \left[d\Phi_{rad} \left(R(\Phi_{n+1}) - C(\Phi_{n+1}) \right) \right]_{\alpha_r}^{\Phi_n^{\alpha_r} = \Phi_n}}{NLO} \\ &+ \sum_{\alpha_{\oplus} \in \{\alpha_{\oplus} \mid f_b\}} \int \left[\frac{dz}{z} G_{\oplus}^{\alpha_{\oplus}}(\Phi_{n, \oplus}) \right]^{\Phi_n = \Phi_n} + \sum_{\alpha_{\ominus} \in \{\alpha_{\ominus} \mid f_b\}} \int \left[\frac{dz}{z} G_{\ominus}^{\alpha_{\ominus}}(\Phi_{n, \ominus}) \right]^{\Phi_n^{\alpha_r} = \Phi_n}}{\text{inclusive}} \\ \Delta^{f_b}(\Phi_n, p_T) &= exp \bigg\{ - \sum_{\alpha_r \in \{\alpha_r \mid f_b\}} \int \frac{\left[d\Phi_{rad} R(\Phi_{n+1}) \Theta(k_T(\Phi_{n+1}) - p_T) \right]_{\alpha_r}^{\Phi_n^{\alpha_r} = \Phi_n}}{B^{f_b}(\Phi_n)} \bigg\} \begin{array}{c} \text{Sudakov} \\ \text{form factors} \end{array} \end{split}$$

Then add softer radiation in the shower approximation

 * P. Nason, A new method for combining NLO QCD with SMC algorithms, JHEP 11 (2004)
 S. Frixione, P. Nason, and C. Oleari, Matching NLO QCD computations with Parton Shower simulations: the POWHEG method, JHEP 11 (2007)

Fragmentation contribution in POWHEG+PYTHIA (I)

Apply the POWHEG method to both QCD and QED radiation



- Remove photon collinear singularity through the subtraction method*
- Perform hardest radiation within the POWHEG framework also for photon emission from a Wj underlying Born

QED corrections to Wj

▶ ISR and lepton-FSR are also included

Simulation of the fragmentation contribution through QED emission in POWHEG + QCD+QED shower (PYTHIA)

Reproduce perturbative and non-perturbative effects of the fragmentation mechanism

* S. Frixione, Z. Kunszt and A. Signer, Nucl. Phys. B 467 (1996)

Fragmentation contribution in POWHEG+PYTHIA (II)

Separate QCD and QED singular regions:



 $R = \sum_{\alpha_r \in QCD} R^{\alpha_r}_{W\gamma;j} + \sum_{\alpha_r \in QED} R^{\alpha_r}_{Wj;\gamma}$



Wγ underlying Born + parton emission Wj underlying Born + photon emission

$$d\sigma = \bar{B}_{W\gamma} \left(\Delta^{QCD} + \sum_{\alpha_r \in QCD} \left[\frac{\Delta^{QCD}(k_T) R_{W\gamma;j}}{B_{W\gamma}} \right]_{\alpha_r} \right) + B'_{Wj} \left(\Delta^{QED} + \sum_{\alpha_r \in QED} \left[\frac{\Delta^{QED}(k_T) R_{Wj;\gamma}}{B_{Wj}} \right]_{\alpha_r} \right)$$

Separation of phase space domains

QCD or QED hardest emission

Wy in POWHEG+MiNLO

- ► Wj Born: divergent for collinear parton emission
 - → need a cut at the generation level
- But: must be inclusive over colored radiation

The MiNLO method*

- Reweight with Sudakov form factors and as evaluated at dynamical scales (based on CKKW)
- NLO accuracy preserved
 - Integrate over the whole phase space: $p_T(W)$ reliable and smooth behavior down to very small p_T
 - \rightarrow Dynamically motivated choice of μ_R and μ_F scales

* K. Hamilton, P. Nason, G. Zanderighi, MINLO: Multi-scale improved NLO, arXiv:1206.3572

Wy implementations in POWHEG+MiNLO

- **POWHEG-NC** Wγ NLO cross section + parton emission
 - ▶ Wj cross section + photon emission
- Need to account for colored radiation in competition with photons
 - → Can be done in an effective way

- **POWHEG-C-LO**

 Photon/parton emission in competition
 - ► Wj cross section at QCD LO accuracy

POWHEG-C-NLO > Wj cross section at QCD NLO accuracy

Wy in POWHEG+MiNLO: NC method

Wγ underlying Born (photon harder than parton)

$$\bar{B}_{W\gamma} = B_{W\gamma} + V_{W\gamma(j)} + \sum_{\alpha_r \in QCD} \int \left[d\Phi_{rad} (R_{W\gamma;j} - C) \right]^{\alpha_r}$$

$$\Delta_{W\gamma}^{QCD} = exp \left\{ -\sum_{\alpha_r \in QCD} \int \left[d\Phi_{rad} \frac{R_{W\gamma;j}}{B_{W\gamma}} \right]^{\alpha_r} \right\}$$

$$\sum_{\alpha_r \in QCD} \int \left[d\Phi_{rad} \frac{R_{W\gamma;j}}{B_{W\gamma}} \right]^{\alpha_r}$$

$$\sum_{\alpha_r \in QCD} \int \left[d\Phi_{rad} \frac{R_{W\gamma;j}}{B_{W\gamma}} \right]^{\alpha_r}$$

Wj underlying Born (parton harder than photon)

$$B'_{Wj} = B_{Wj} + \sum_{\alpha_r \in QED} \int \left[d\Phi_{rad} (R_{Wj;\gamma} - C) \right]^{\alpha_r} \quad \text{Effe}$$
$$\Delta_{Wj}^{QED} = exp \left\{ -\sum_{\alpha_r \in QED} \int \left[d\Phi_{rad} \frac{R_{Wj;\gamma}}{B_{Wj}} \right]^{\alpha_r} \right\}$$

Effective choice for *scalup**

- ▶ *pτ*(j_{Born}) QCD shower
- *p*^{rel}(γ_{rad}) QED shower

* L. D'Errico, P. Richardson, Next-to-Leading-Order Monte Carlo Simulation of Diphoton Production in Hadronic Collisions, JHEP 1202 (2012)

Wy in POWHEG+MiNLO: C-LO method

For Wj underlying Born include also colored radiation dynamics (in competition with photon radiation)



$$B_{Wj}^{\prime} \Big(\Delta_{Wj}^{QED} + \sum_{\alpha_r \in QED} \Big[\frac{\Delta_{Wj}^{QED}(k_T) R_{Wj;\gamma}}{B_{Wj}} \Big]_{\alpha_r} + \Delta_{Wj}^{QCD} + \sum_{\alpha_r \in QCD} \Big[\frac{\Delta_{Wj}^{QCD}(k_T) R_{Wj;j}}{B_{Wj}} \Big]_{\alpha_r} \Big)$$

- Generated events:
 - $W\gamma(+j)$ photon harder than partons
 - $W_{j+\gamma}$ photon softer than one parton
 - Wj+j+γPs photon softer than at least two partons
- Scalup choice: POWHEG standard (dynamically motivated)
- ▶ Wj normalization at QCD LO (same as for NC)

Wy in POWHEG+MiNLO: C-NLO method

→ Wj component normalized at QCD NLO

$$\bar{B}_{Wj} = B_{Wj} + V_{Wj(j)} + \sum_{\alpha_r \in QED} \int \left[d\Phi_{rad} (R_{Wj;\gamma} - C) \right]^{\alpha_r} + \sum_{\alpha_r \in QCD} \int \left[d\Phi_{rad} (R_{Wj;j} - C) \right]^{\alpha_r} d\Phi_{rad} (R_{Wj;j} - C) d\Phi_{rad} (R_{Wj;j} - C) \right]^{\alpha_r} + \sum_{\alpha_r \in QCD} \int \left[d\Phi_{rad} (R_{Wj;j} - C) \right]^{\alpha_r} d\Phi_{rad} (R_{Wj;j} - C) d\Phi_{rad} (R_{Wj;j} - C) \right]^{\alpha_r} d\Phi_{rad} (R_{Wj;j} - C) d\Phi_{rad}$$

- Radiation dynamics same as for C-LO
- NLO corrections to the fragmentation contribution
 - Very important, as fragmentation $\sim O(a_{em}/a_s)$
 - gg channel opened (high gluon PDFs)
- Goes in the direction of Wγ at NNLO+PS, but...
 - Exact real Wγjj and 2-loop virtual still missing
 → NNLO calculations available* (+20% for leptonic cuts + smooth isolation)

*M. Grazzini, Vector-boson pair production at NNLO, arXiv:1407.1618

Results: selection cuts

Selection cuts: ATLAS setup*

 $p_T(\gamma) > 15 \text{GeV}, \quad |\eta_\gamma| < 2.37, \quad \Delta R_{l\gamma} > 0.7$

 $p_T(l) > 25 \text{GeV}, \quad |\eta_l| < 2.47, \quad E_T^{miss} > 35 \text{GeV}$

$$R_0 = 0.4, \quad \epsilon_h = 0.5$$

$$\sum_j E_{Tj} \Theta(R_0 - R_{j\gamma}) \le \epsilon_h E_{T\gamma}$$

Jet cuts for exclusive N_{jet} selections (anti-kt jet-finding algorithm, r=0.4)

 $E_T(jet) > 30 \text{GeV}, \quad |\eta_{jet}| < 4.4, \quad \Delta R(e/\mu/\gamma, jet) > 0.3$

Results obtained with POWHEG+PYTHIA8

* The ATLAS Collaboration, Phys. Rev. D 87 (2013)

Results: comparisons with MCFM



Results: comparisons with ATLAS data* (I)

[pb]	Exclusive $N_{jet} = 0$	Inclusive $N_{jet} \ge 0$
Exp.	$1.77^{\pm 0.04 \text{stat}}_{\pm 0.08 \text{lumi}} \pm 0.24 \text{ syst}$	$2.74^{\pm 0.05 \text{stat}}_{\pm 0.14 \text{lumi}} \pm 0.32 \text{ syst}$
MCFM	1.39 ± 0.17	1.96 ± 0.17
PWG-C-LO	$1.42^{+0.15}_{-0.15}$	$2.25_{-0.24}^{+0.24}$
PWG-C-NLO	$1.69^{+0.11}_{-0.22}$	$2.95_{-0.38}^{+0.20}$



$$pp \to l\nu_l \gamma + X \qquad \sqrt{s} = 7 \text{TeV}$$

Electron and muon channels combined

$$M_{\perp}^{2}(l\nu_{l}\gamma) = \left(\sqrt{m_{l\gamma}^{2} + |\mathbf{p}_{T}^{\gamma} + \mathbf{p}_{T}^{l}|^{2}} + p_{T}^{\nu_{l}}\right)^{2} - |\mathbf{p}_{T}^{\gamma} + \mathbf{p}_{T}^{l} + \mathbf{p}_{T}^{\nu_{l}}|^{2}$$

$$p_T(\gamma) > 40 \text{ GeV}$$

* The ATLAS Collaboration, Phys. Rev. D 87 (2013)

Results: comparisons with ATLAS data (II)

$$pp \to l\nu_l\gamma + X$$
 $\sqrt{s} = 7 \text{TeV}$ Data 2011 $\int \mathcal{L}dt = 4.6 \text{ fb}^{-1}$

Photon transverse momentum



Exclusive selection N_{jet}=0

n



Results: comparisons with ATLAS data (III)

$$pp \rightarrow l\nu_l\gamma + X$$
 $\sqrt{s} = 7 \text{TeV}$ Data 2011 $\int \mathcal{L}dt = 4.6 \text{ fb}^{-1}$

Jet multiplicity

 $p_T(\gamma) > 15 \text{ GeV}$

 $p_T(\gamma) > 60 \text{ GeV}$

p



Conclusions and outlook

\star NLO QCD corrections to Wy production in pp collisions:

- NLO+PS description within POWHEG-MINLO+PYTHIA
- Photon fragmentation contribution included
- NLO corrections to the fragmentation contribution
- ★ Improvements in data/theory comparison
- General treatment of isolated photons in hadronic collisions
 - Future applications to Zγ/γγ production
- The code will be available on the POWHEG-BOX-V2 repository http://powhegbox.mib.infn.it

Backup slides

Results: comparisons with MCFM (II)



Results: uncertainties from scale variations



Upper and lower boundings from variations in opposite direction

Photon + lepton cuts

 $(K_R,K_F)\in \left\{\left(\frac{1}{2},2\right),\left(2,\frac{1}{2}\right)\right\}$

 $\sqrt{s} = 14 \text{TeV}$