

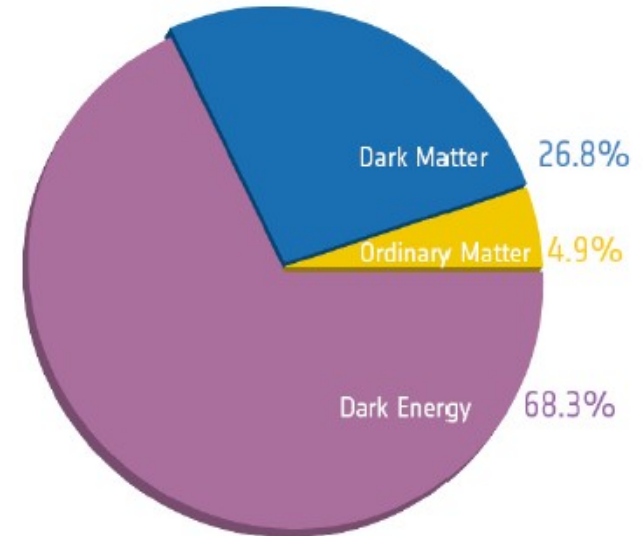
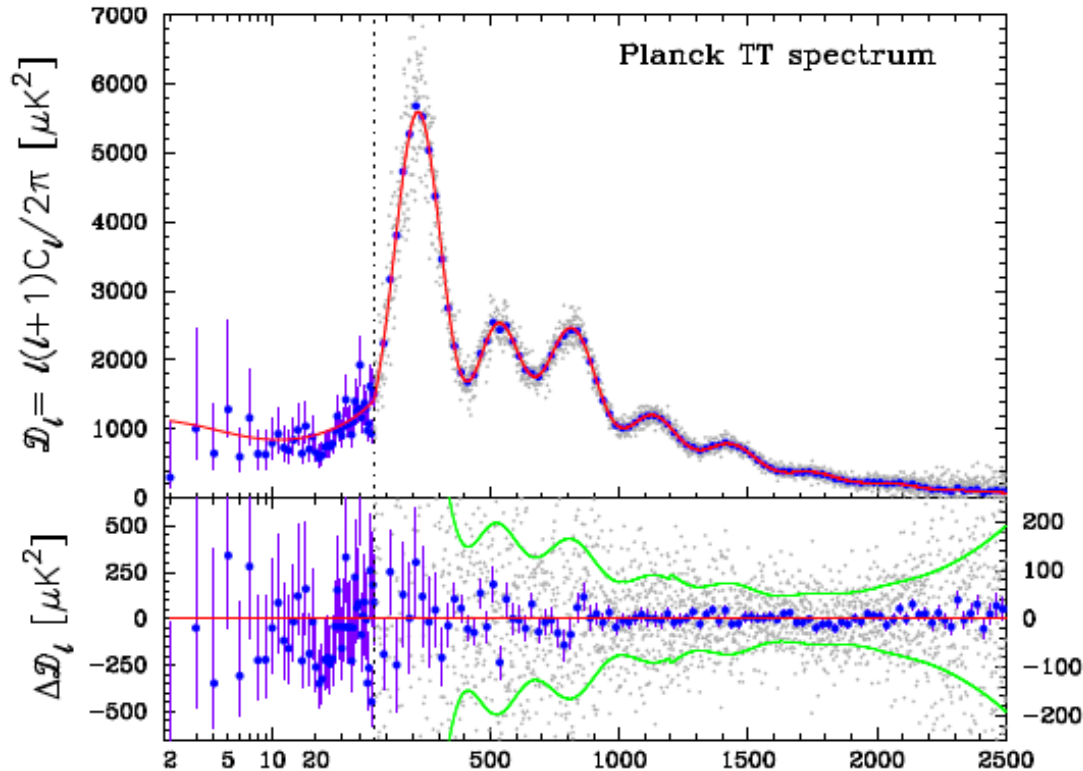
New physics at the LHC

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Motivations for going beyond Standard Model

- **Observations unexplained by SM**
 - Dark matter problem
 - Matter-antimatter asymmetry problem
- **Fine-tuning problems**
 - Hierarchy problem associated with Higgs
 - Flavour problem
 - Strong CP problem
- **“Why so” puzzles**
 - Charge quantisation
 - Gauge coupling unification
 - Proton stability
 - Fermion mass hierarchy
 - Why three generations

Amount of Dark matter in the universe



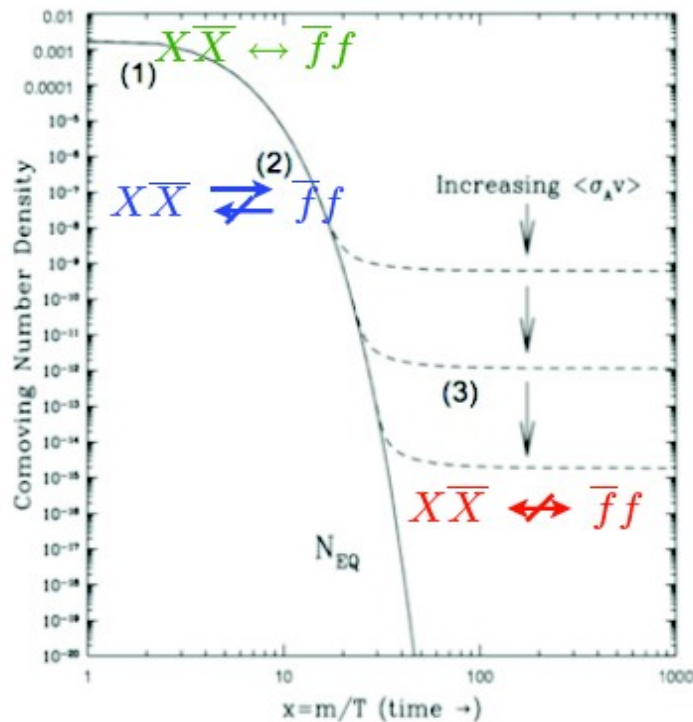
Extremely precise results on Dark Matter abundance from measurement of anisotropies in Cosmic Microwave Background (CMB)

If Dark Matter is made of Weakly Interacting Massive Particles (WIMP), what we observe is the relic abundance of these particles after the cooling of the universe

The “WIMP miracle”: DM may be relevant for LHC

The WIMP relic abundance follows from the generic thermal freeze-out mechanism in the expanding universe

$$\dot{n} + 3Hn = -\langle\sigma v\rangle(n^2 - n_T^2)$$



freeze-out :

$$H \sim \frac{\sqrt{g}T^2}{M_P} \sim \Gamma = n\sigma v$$

Thermal relic: $\Omega h^2 \propto 1/\langle\sigma_{\text{anni}} v\rangle$

$$\Rightarrow \langle\sigma_{\text{anni}} v\rangle \approx 1 \text{ pb}$$

$$\sigma \sim \alpha^2/m^2$$

$$\Rightarrow m \sim 100 \text{ GeV}$$

The “WIMP miracle”

$$\Omega_{\text{DM}} \approx \frac{O(1) \text{ pb}}{\sigma_{\text{anni}}}$$

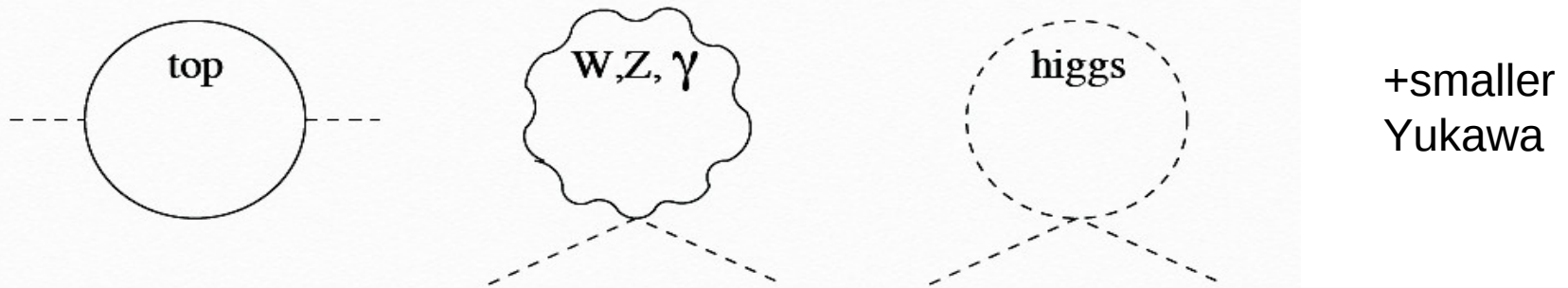
→ a particle with a typical EW-scale cross section

$\sigma_{\text{anni}} \approx 1 \text{ pb}$ leads to the correct dark matter abundance.

The naturalness problem

Key assumption: SM is Effective Field Theory valid up to scale $\Lambda \gg \text{TeV}$

Radiative corrections to Higgs mass:



$$\delta m_H^2 = \frac{3\Lambda^2}{8\pi^2 v^2} \left(2m_W^2 + m_Z^2 + m_H^2 - 4m_t^2 \right) \sim -(0.23 \Lambda)^2$$

If $\Lambda=5 \text{ TeV}$ already need cancellation between tree level and radiative corrections of 2 orders of magnitude

We have observed a 125 GeV scalar

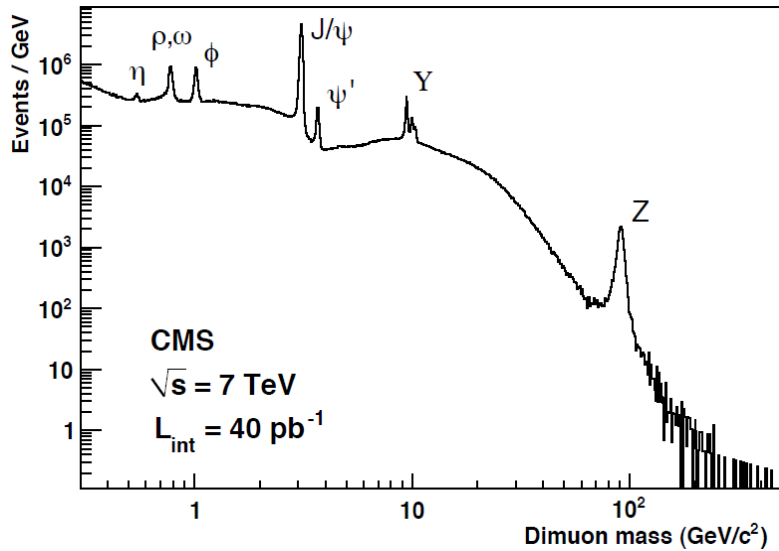
We need to understand why it is so light

All proposed solutions imply new physics at the TeV scale
Search for this physics high priority at the LHC

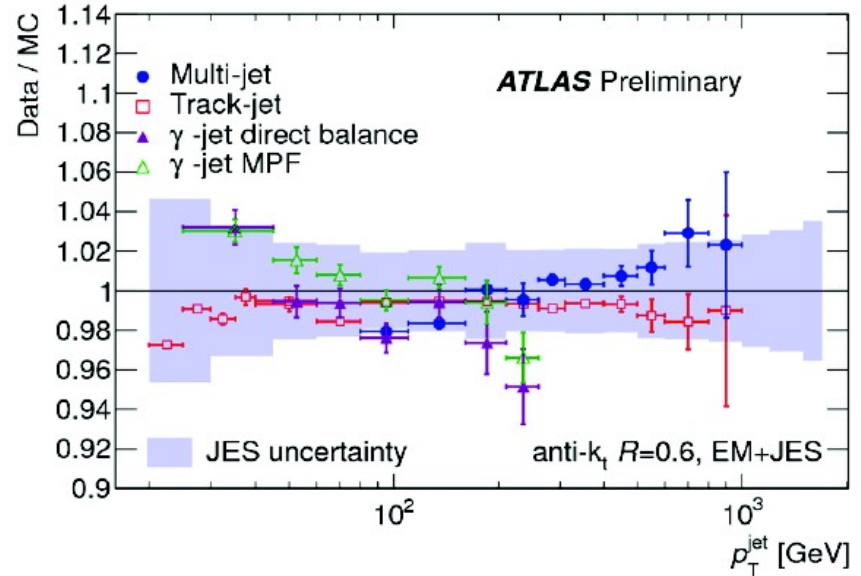
Discovering new physics: preliminaries

- Once good data on disk:
 - Calibration has to be determined and applied
 - Detector objects to be reconstructed
 - Reconstructed data to be made available on the grid
 - Complete calibration loop within 48 hours of data taking
- Starting from reconstructed data, two steps necessary before going for new physics searches:
 - Understanding of detector performance for main objects: leptons, jets, photons, b-jets, τ -jets, E_{miss}
 - Measurements of Standard Model processes to ensure that our detector understanding is adequate to look for deviations

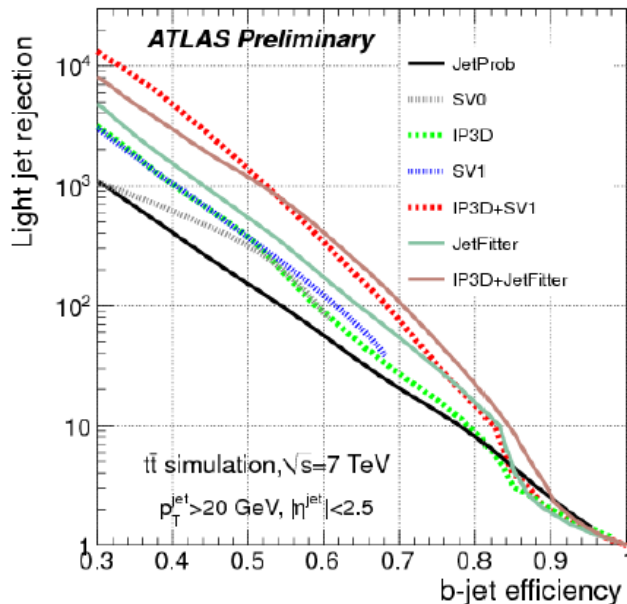
Performance examples



Leptons: need excellent id capabilities
And resolution

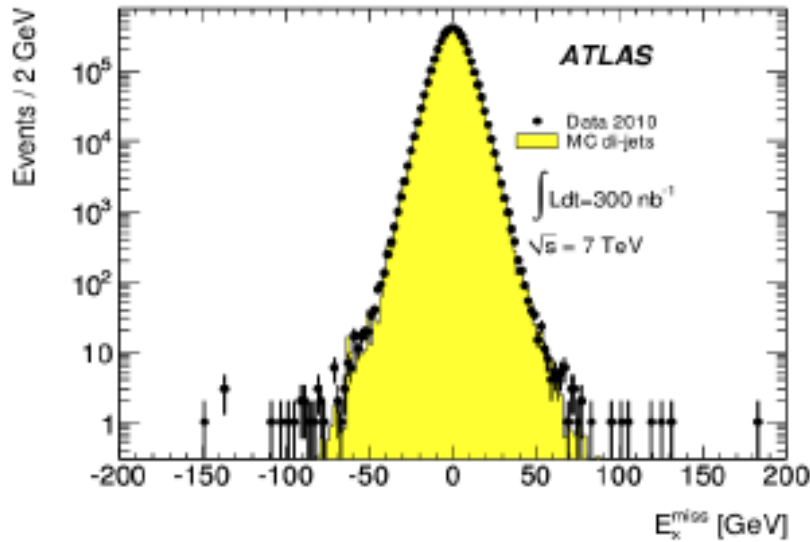


Jet energy scale to 2-4% for
Jet $p_T > 20 \text{ GeV}$



B-tagging: key to detailed searches
Advanced methods validated with 2011 data
For 60% efficiency rejection of several hundreds
On light jets

Etmiss measurement



Key ingredient in SUSY analysis

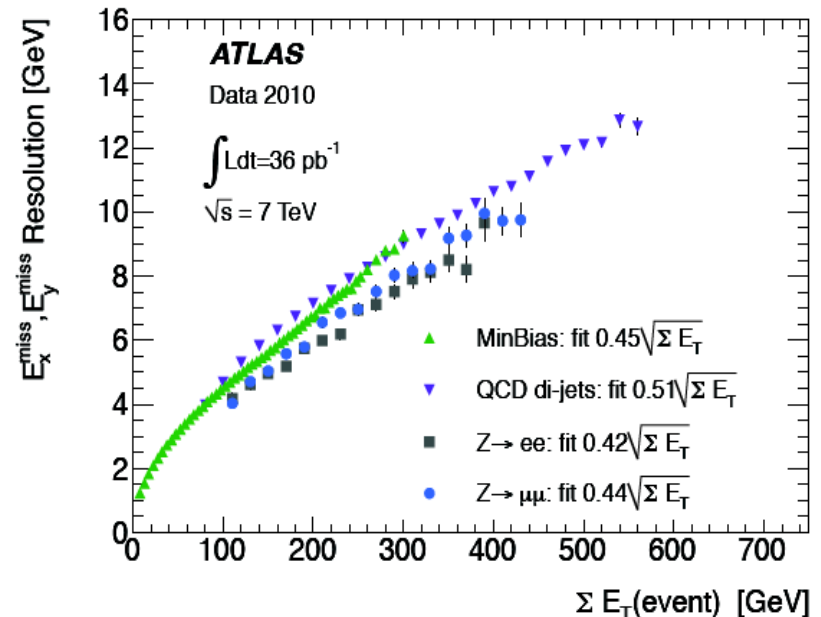
$$E_{x(y)}^{\text{miss}} = E_{x(y)}^{\text{miss,calo}} + E_{x(y)}^{\text{miss},\mu}$$

$$E_{x(y)}^{\text{miss,calo}} = E_{x(y)}^{\text{miss},e} + E_{x(y)}^{\text{miss},\gamma} + E_{x(y)}^{\text{miss},\tau} + E_{x(y)}^{\text{miss,jets}} \\ + E_{x(y)}^{\text{miss,softjets}} + (E_{x(y)}^{\text{miss,calo},\mu}) + E_{x(y)}^{\text{miss,CellOut}}$$

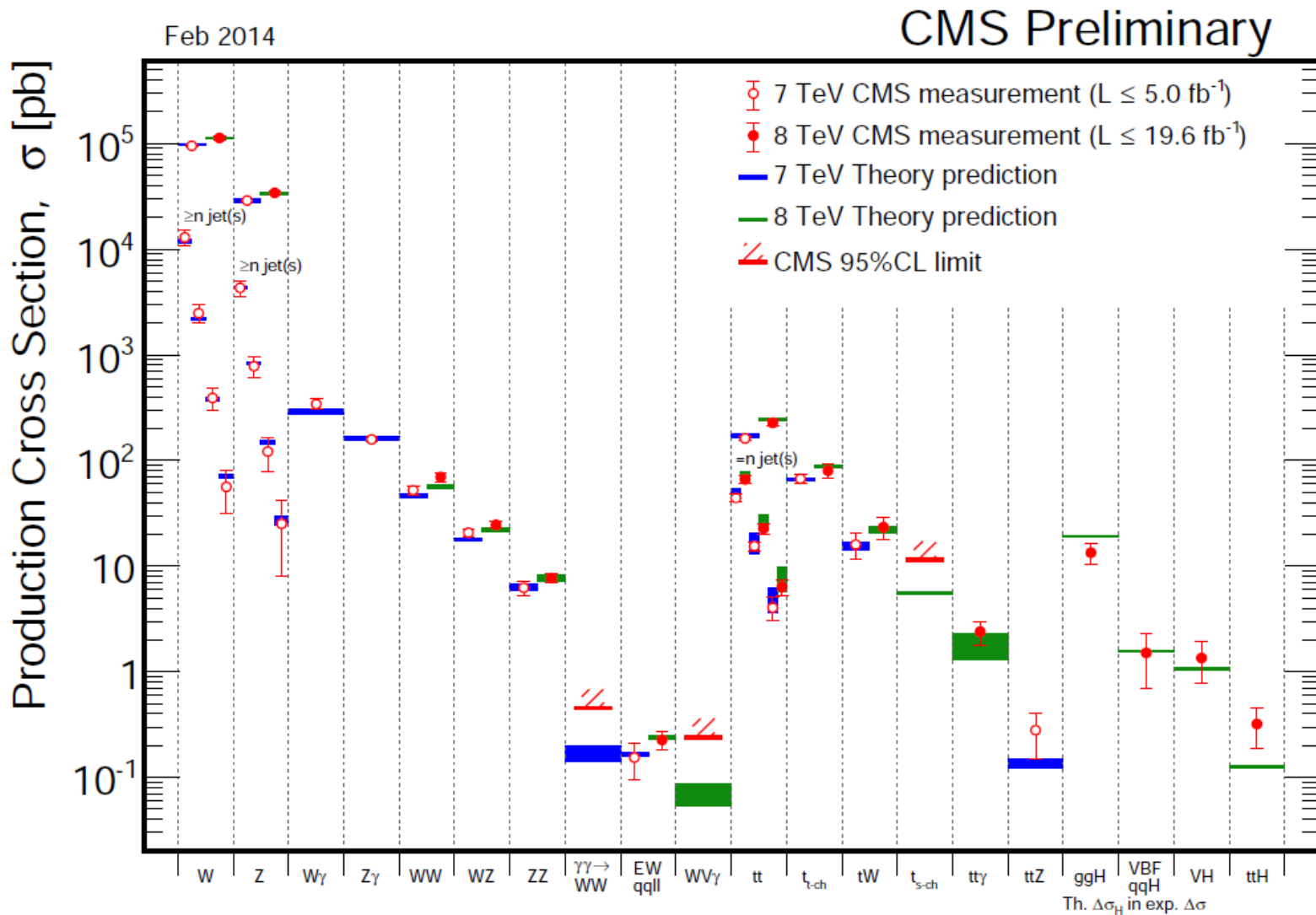
Vector sum of the measured energy deposit of all objects in the detector

Any local malfunction in the detector would be registered as a tail in Etmiss distribution

From early data taking tails under control and measurement resolution in agreement with expected value



Standard Model measurements



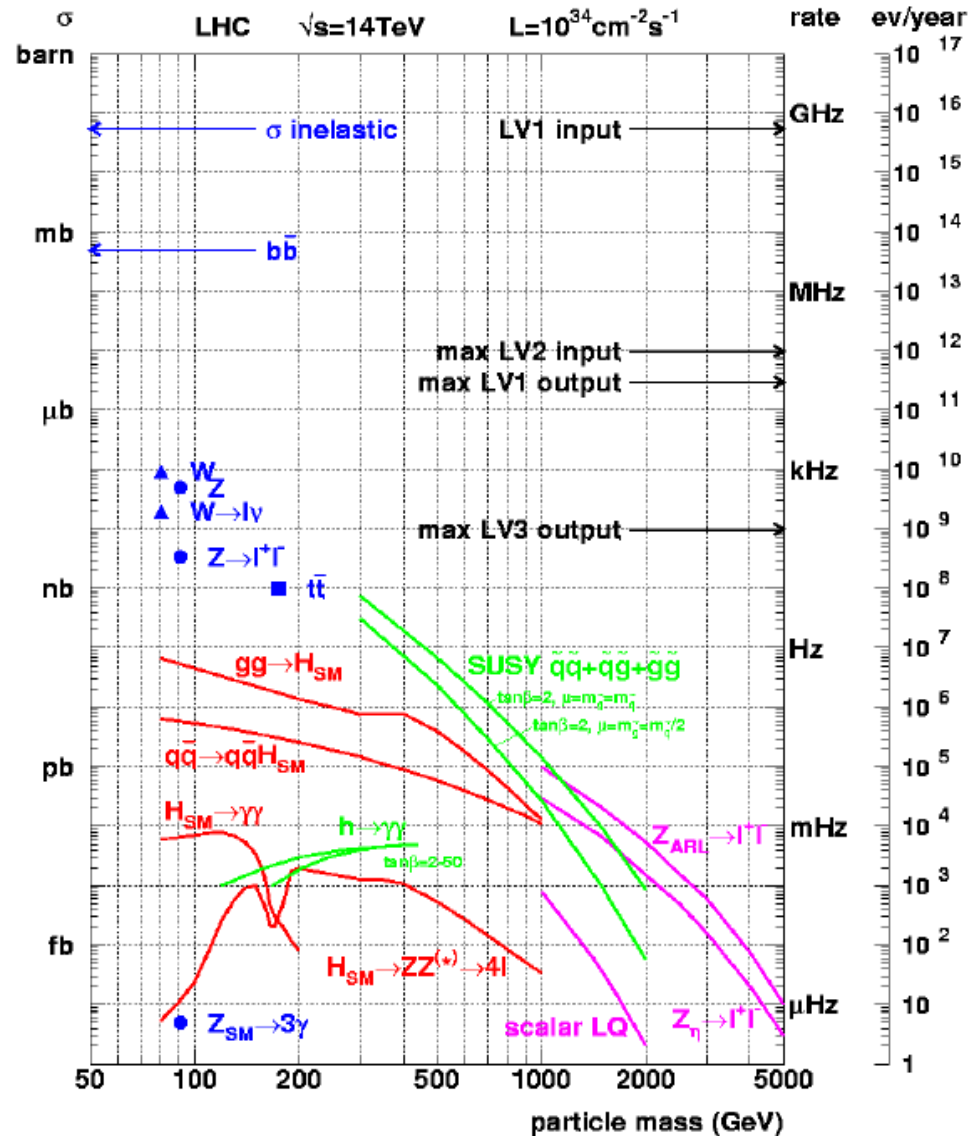
No exotic source of bosons/top in excess of 10-20% of SM
But this is only the start of the story

The problem: signal much smaller than bkg

For each signal need to devise selections reducing background by several orders of magnitude:

Need to predict SM in extreme corners of kinematic space

Necessary to complement MC with data-driven estimate

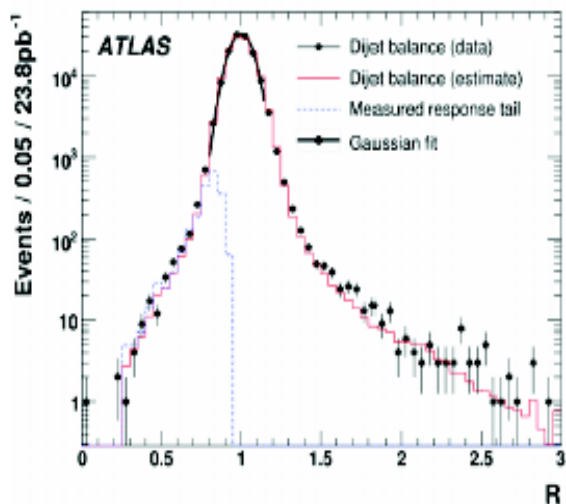


Selections and backgrounds

- QCD jet production overwhelming at LHC, need to add something else
- Signatures classified in terms of
 - non-QCD objects: leptons (e, μ), E_{miss} , τ -jets, b-jets
 - Number of QCD jets
- For each signature two types of backgrounds
 - Irreducible backgrounds: basic signature identical to signal
 - Reducible backgrounds: mimic signature because of detector effects – examples:
 - Fake E_{miss} in multijet events
 - Fake leptons
- For each type of background need to develop specific strategies

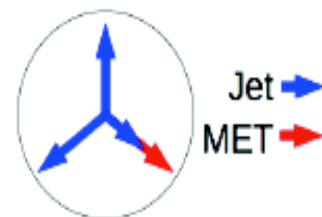
Fake E_T^{miss} estimate

- Large E_T^{miss} can be induced by a jet mis-measurement.
- Relevant for processes with high cross section and no “real” E_T^{miss} (multi-jet, Z→ll)



- Derive a “jet response function” from MC and adapt it to data:

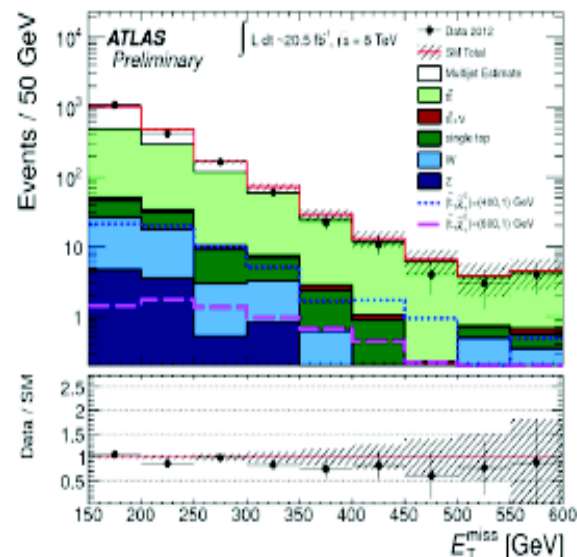
- **core:** p_T balance in di-jet events
- **tail:** three-jet (Mercedes) events



- Use response function to smear jets in real data events with low MET:

- Obtain events with large “fake” E_T^{miss}

- Validate the estimation in a dedicated control region



Fake lepton estimate

- General approach to **fake lepton background estimation** based on a **loose/tight matrix method**
- Example with 1 lepton (easily extendable to multi-lepton signatures):
- Strategy: **define a “loose”** (pre-selected) and a **“tight”** (signal) lepton selection.
- Then, solve the following system of equations

$$N^{loose} = N_{real}^{loose} + N_{fake}^{loose}$$

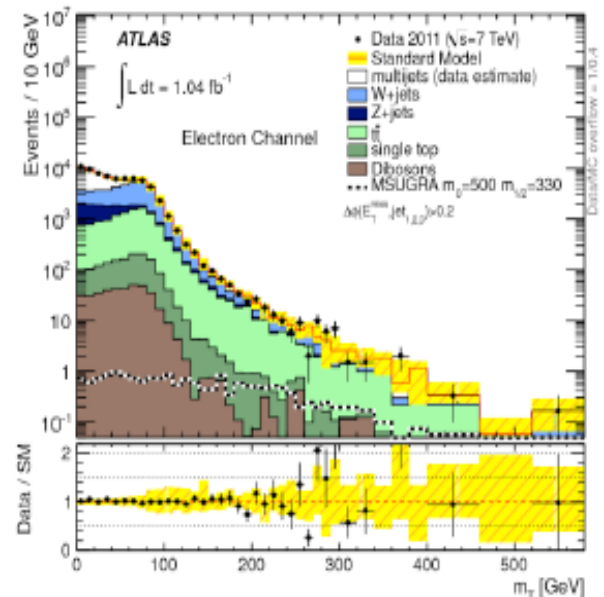
$$N^{tight} = \epsilon_{real} N_{real}^{loose} + \epsilon_{fake} N_{fake}^{loose}$$

- A fake lepton can arise from:
 - Off-axis HF semileptonic decays
 - Photon conversion

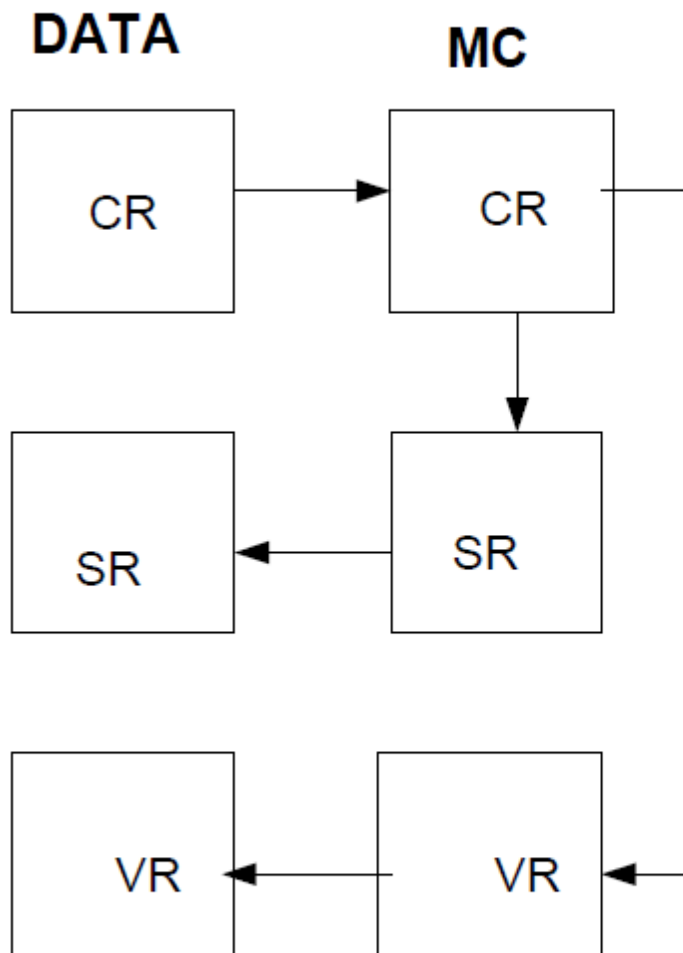
Need to be measured independently from data

Simply count how many of them

$$N_{fake}^{tight} = \frac{\epsilon_{fake}}{\epsilon_{real} - \epsilon_{fake}} (N_{real}^{loose} \epsilon_{real} - N^{tight})$$



Background predictions: CR method



Control Region: CR
No signal expected

Signal Region: SR
Where signal is expected

Validation Region (VR)
No signal expected

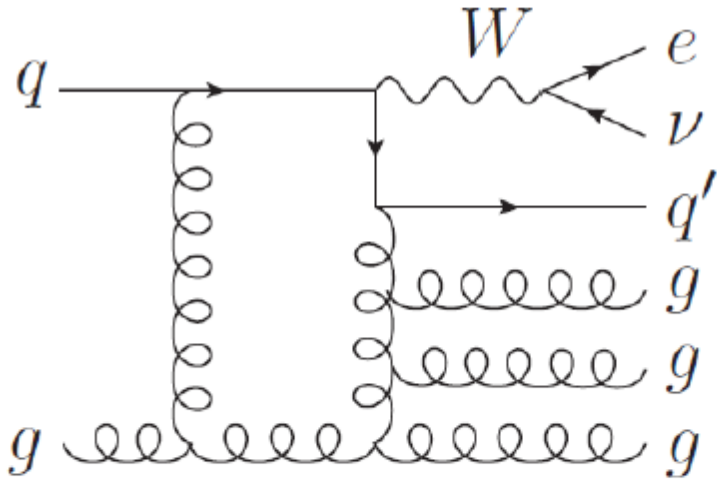
Transfer factor (TF)=
 $N(\text{MC}, \text{SR}) / N(\text{MC}, \text{CR})$

$N(\text{DATA}, \text{SR}) = N(\text{DATA}, \text{CR}) * \text{TF}$

Estimate relies on the MC to predict TF correctly
Use VR to give you confidence on prediction

Key is clever
Choice of CR

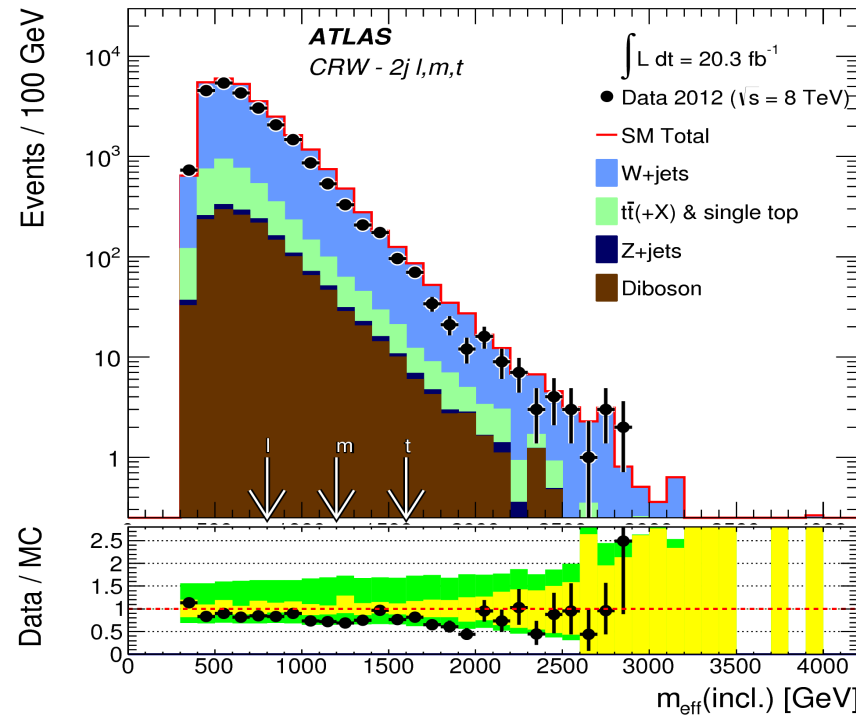
Example



- Multijet+ Emiss:
Background from W+jets
Where lepton is not detected
- Too low pt
 - Outside acceptance
 - Id inefficiency

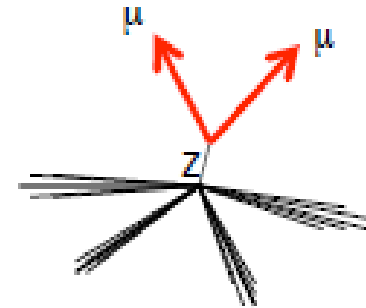
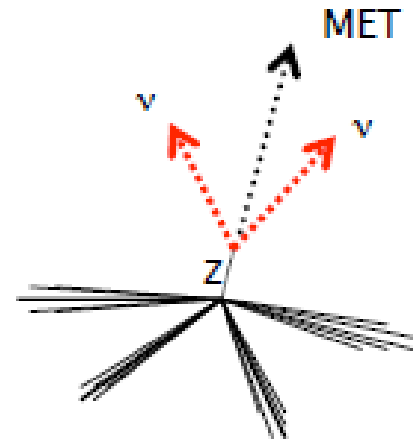
Discriminant variable:
Meff: sum of pt of jets and Emiss

Leptonic decay of W well
predicted by MC
Measure Meff in a sample with
One reconstructed lepton
Predict with MC when leptons
get lost

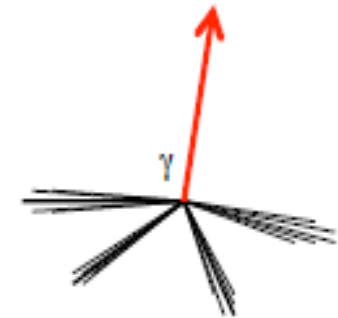


Example 2

- Replacement Method: $Z \rightarrow \nu\nu + \text{jets}$
- Main irreducible background to multijets+ E_{Tmiss}
- Apply the analysis cuts except E_{Tmiss} to a replacement process
 - Take $Z \rightarrow \mu\mu$ and replace leptons with E_{Tmiss}
 - Take prompt photon events and replace photon with E_{Tmiss}
- Transfer the measured E_{Tmiss} spectrum in replacement process to the original process via MC



$Z \rightarrow \mu\mu + \text{jets}$

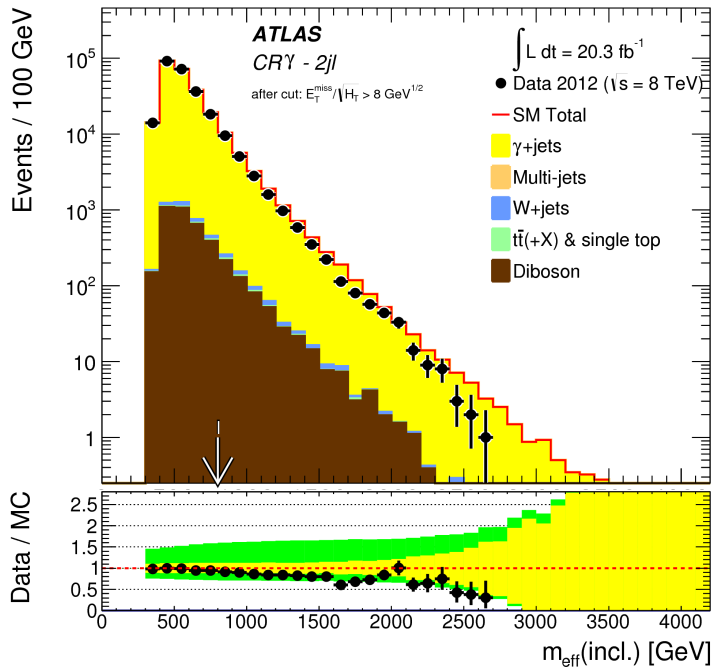


$\gamma + \text{jets}$

MC still has a key role in transferring the result from the Replacement process to the original one

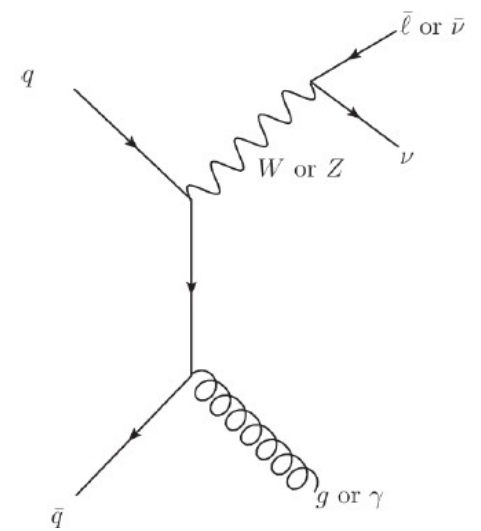
Transfer is 'easy' for $Z \rightarrow \mu\mu$,
And more complex for prompt Photon \rightarrow Larger systematics

Statistical error much bigger
For $Z \rightarrow \mu\mu$



ABCD Method

In a search for mono-photon+E_{miss}, background from W/Z+jets where the jet is identified as a photon
 Use CR with one or two lepton+E_{miss} recoiling against a jet + estimate transfer factor from jet to fake photon



Photons separated from jets with two criteria:

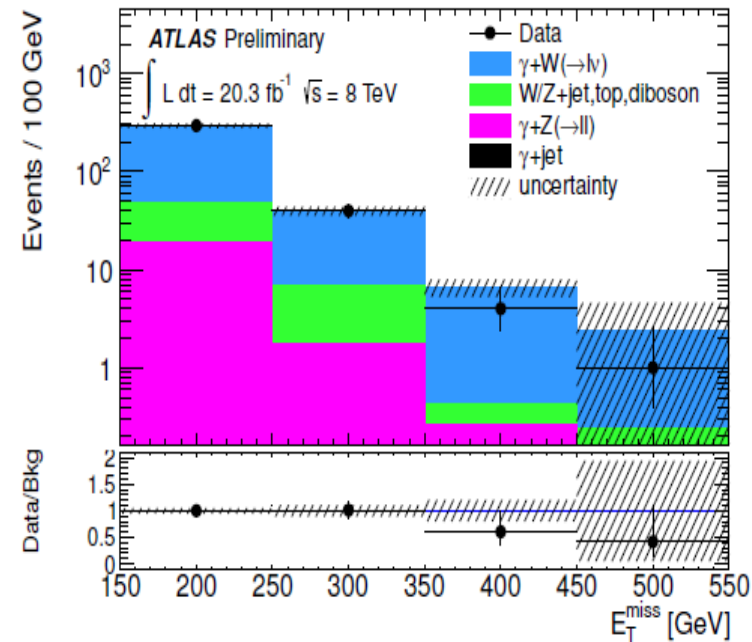
- Shower shape and track veto
- Isolation: no activity in cone around photon

By releasing one or both of these criteria

Create 3 control regions

If the two criteria are independent:

| | | | |
|----------------|---------|--------------|----------|
| Identification | nominal | CR A | SR |
| | loose | CR B | CR C |
| | | non-isolated | isolated |

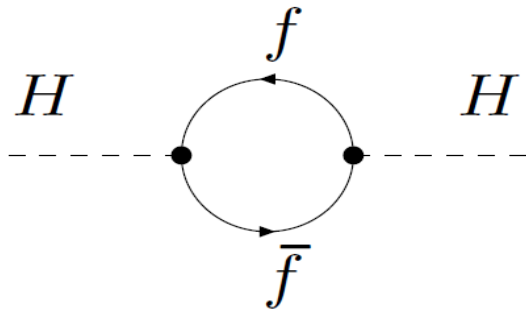


$$N_{SR} = N_{CR C} \times \frac{N_{CR A}}{N_{CR B}}$$

SUSY

SUSY solution to naturalness problem

Correction to higgs mass from fermion loop:



$$\Delta m_H^2 \sim \frac{\lambda_f^2}{4\pi^2} (\Lambda^2 + m_f^2) +$$

Where Λ high energy cutoff

For $\Lambda \sim M_{\text{Planck}} \sim 10^{18}$ GeV corrections explode

Correction from scalar \tilde{f}

$$\Delta m_H^2 \sim -\frac{\lambda_{\tilde{f}}^2}{4\pi^2} (\Lambda^2 + m_{\tilde{f}}^2) + \dots$$

Corrections have opposite sign. Cancellations if for each fermion degree of freedom one has scalars such that:

$$\lambda_{\tilde{f}}^2 = \lambda_f^2 \quad m_{\tilde{f}} = m_f$$

Achieved in theory invariant under transformation Q:

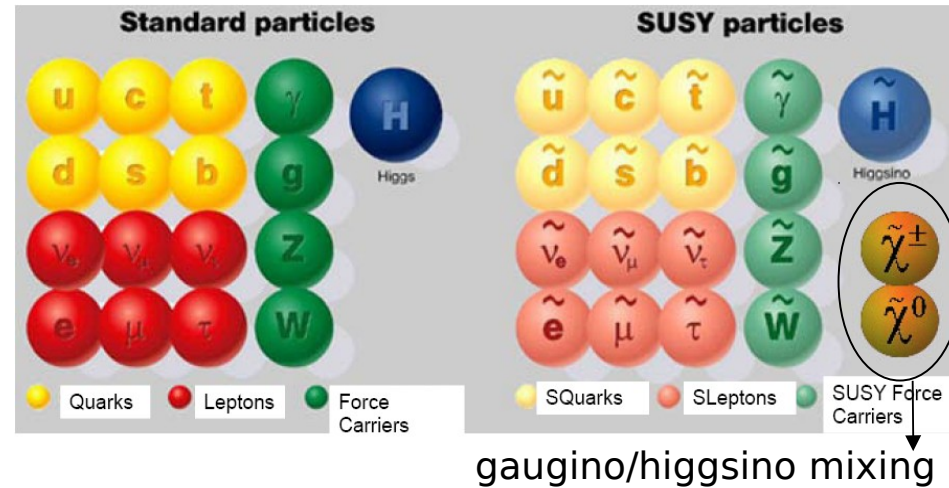
$$Q|\text{boson}\rangle = |\text{fermion}\rangle \quad Q|\text{fermion}\rangle = |\text{boson}\rangle \quad \text{Supersymmetry}$$

Very general class of theories, specialize to minimal model: **MSSM**

Minimal Supersymmetric Standard Model (MSSM)

Minimal particle content:

- A superpartner for each SM particle
- Two Higgs doublets and spartners:
5 Higgs bosons: h, H, A, H^+, H^-



- Insert in Lagrangian all soft breaking terms: **105 parameters**.
- If we assume that flavour matrices are aligned with SM ones (minimal flavour violation): **19 parameters**

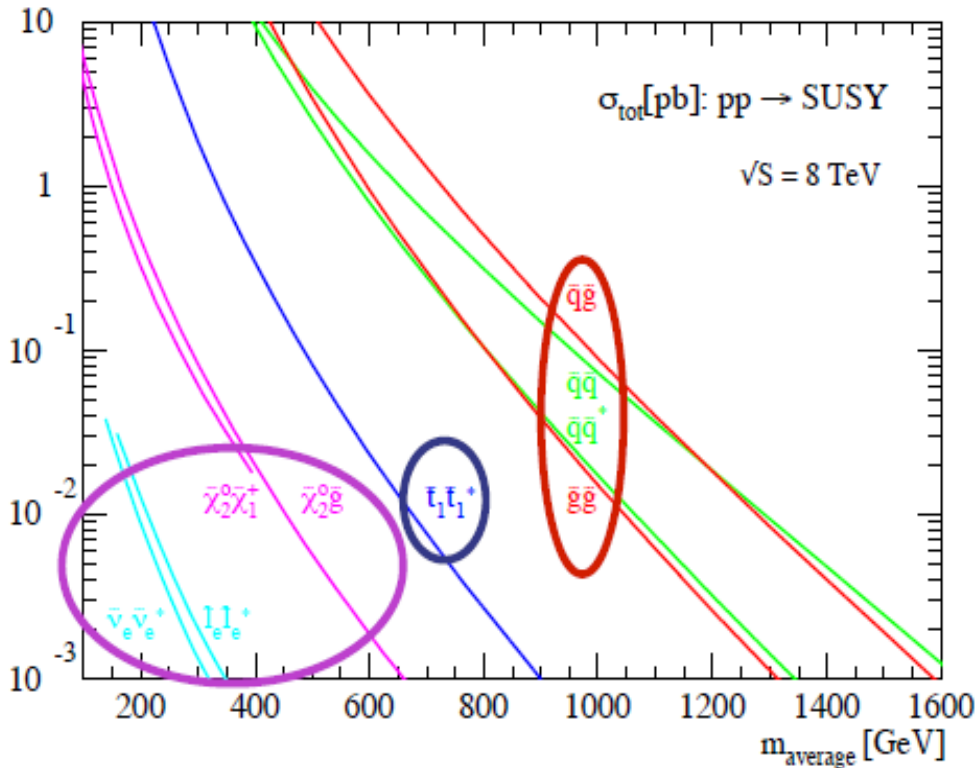
Additional ingredient: R-parity conservation: $R = (-1)^{3(B-L)+2S}$

- Sparticles are produced in pairs
- The Lightest SUSY particle (LSP) is stable, neutral weakly interacting
 - Excellent dark matter candidate
 - It will escape collider detectors providing E_{T}^{miss} signature

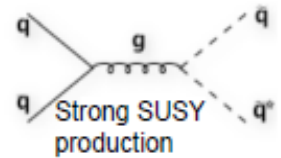
Models with R-parity violating terms are also studied: no E_{T}^{miss} signature, but often 'easier' kinematic signatures

SUSY search strategy

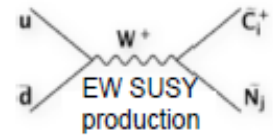
- SUSY pair productions at the LHC:



- **gluinos and 1st and 2nd gen. squarks** with high cross-section, reachable up to $> 1 \text{ TeV}$ mass
- **3rd gen. squarks with moderate cross-section**, up to $\sim 0.5 \text{ TeV}$

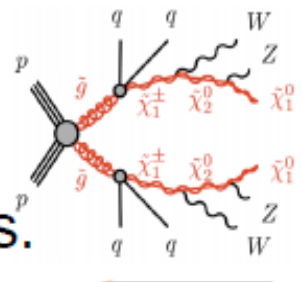


- **charginos, neutralinos and sleptons** with small cross-section, becoming feasible with the current dataset.



- Sparticles decay into characteristic signatures:
 e.g. E_T^{miss} , (b/c-) jets, leptons, and photons.

→ Designed various analyses to cover many signatures.



All hadronic signature optimisation

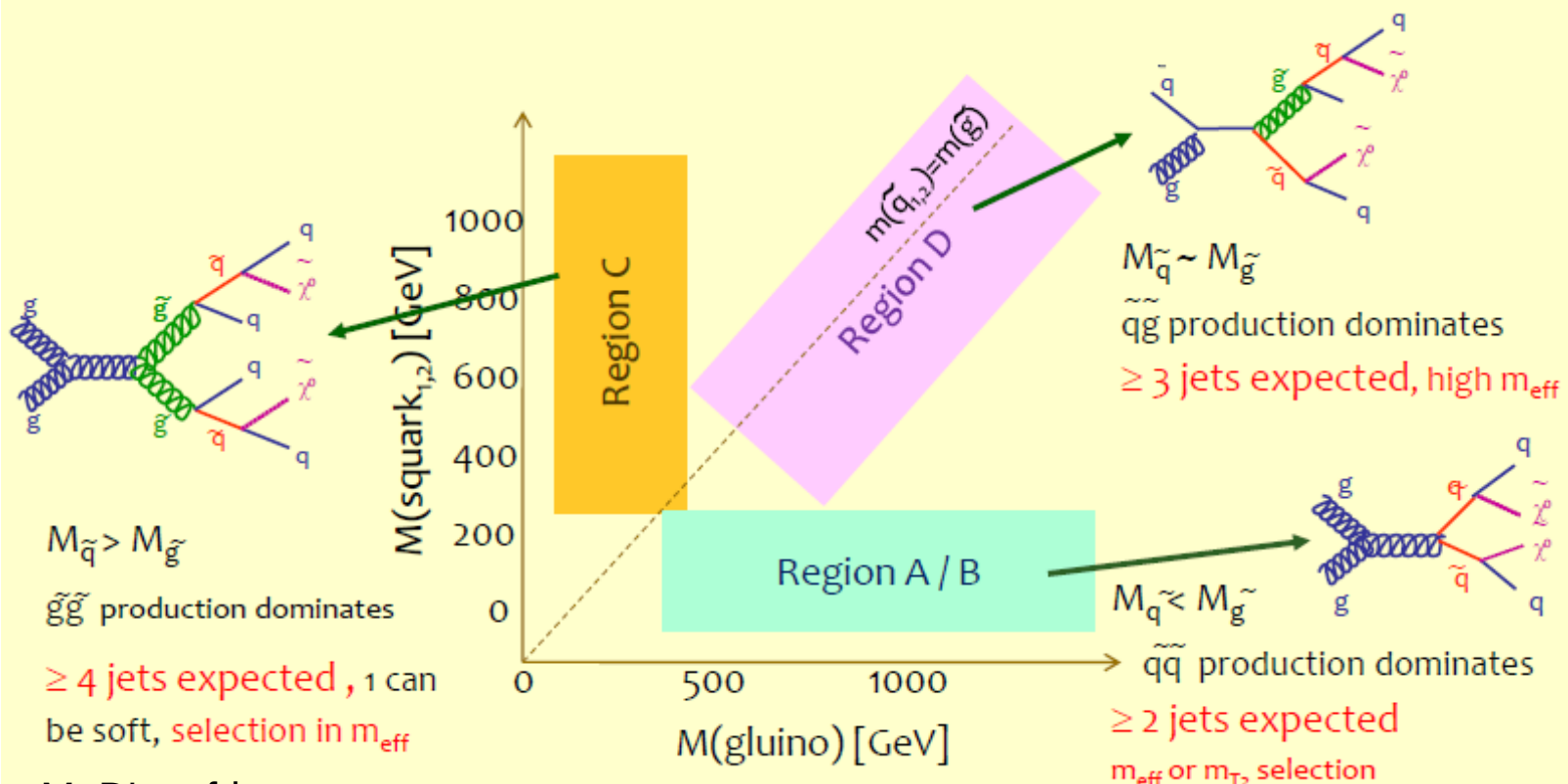


Figure by M. D'onofrio

Require 2 to ≥ 6 (8) Jets and E_{tmiss} . Signal regions classified according to:

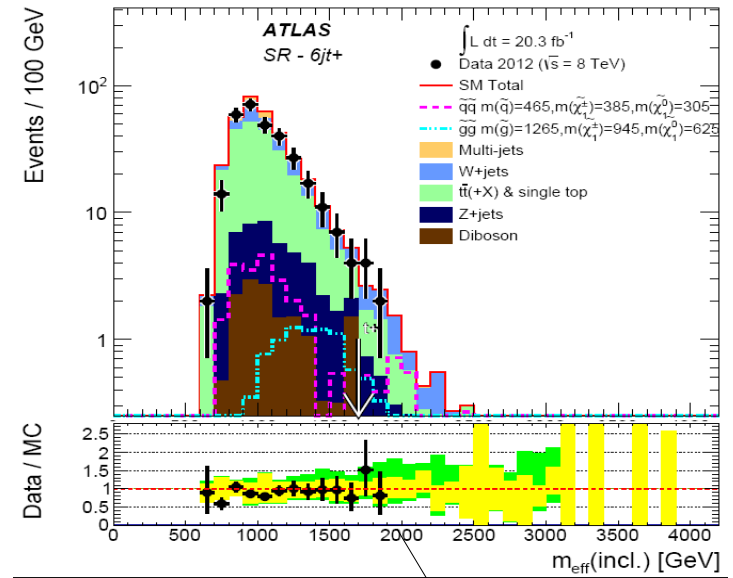
- Number of jets (ATLAS and CMS)
- E_{Tmiss} (ATLAS) HT_{miss} (-vector sum of jet pT) (CMS)
- $M_{\text{eff}} = E_{\text{Tmiss}} + \text{scalar sum of jet pT}$ (ATLAS)
- $HT = \text{scalar sum of jet pT}$ (CMS)

Results

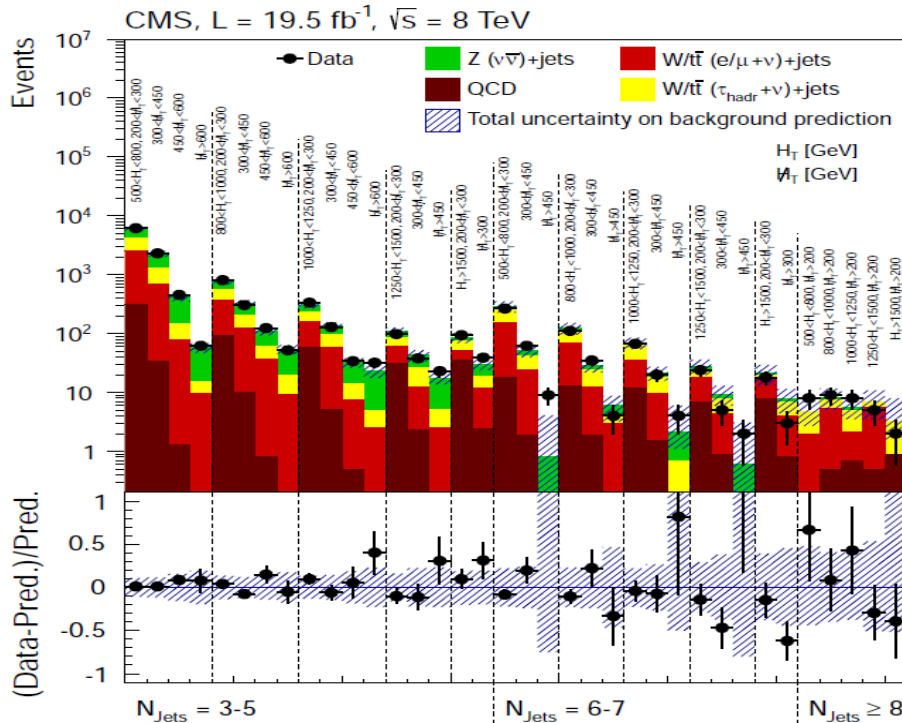
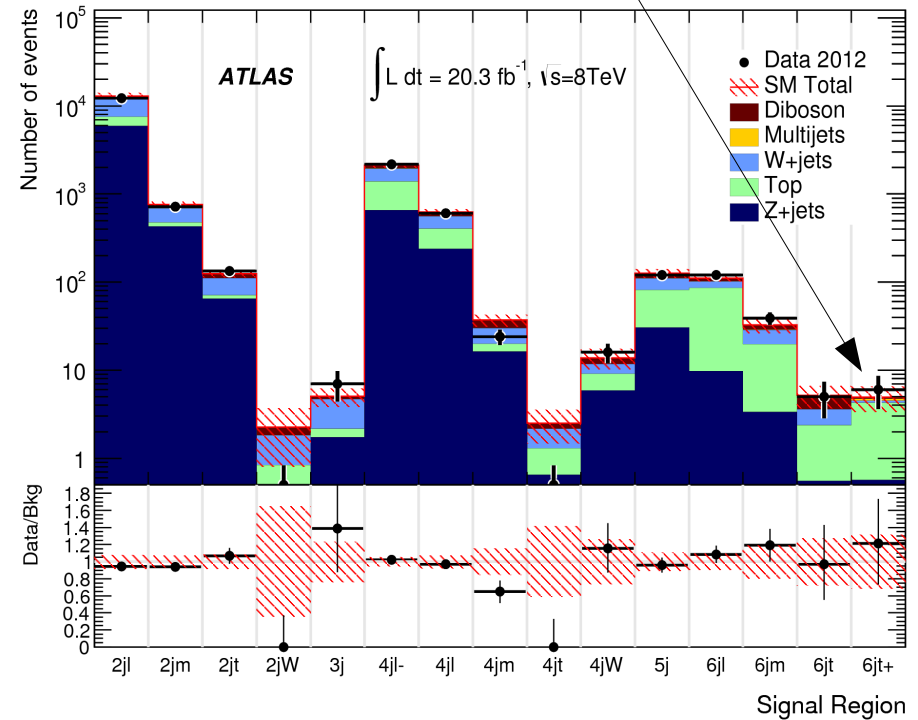
Good agreement between data and prediction in all signal regions

→ Interpret in term of coverage of SUSY space

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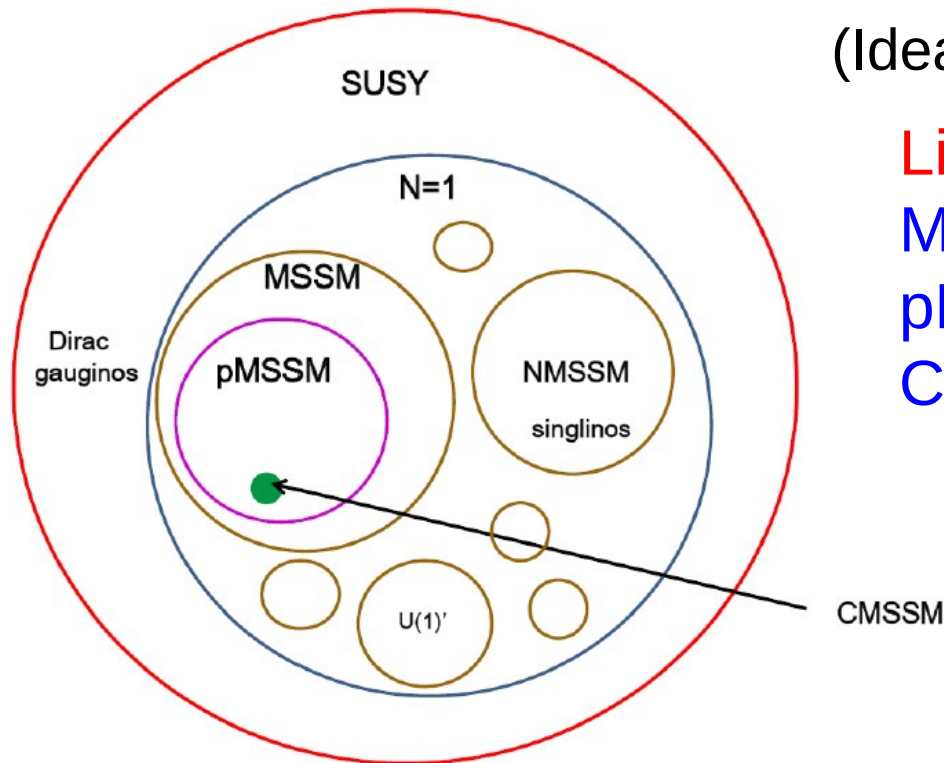


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Interpretation

SUSY theory space



For interpretations need to reduce
To small parameter dimensionality
(Ideally 2)

Limiting to MSSM:

MSSM: ~109 parameters

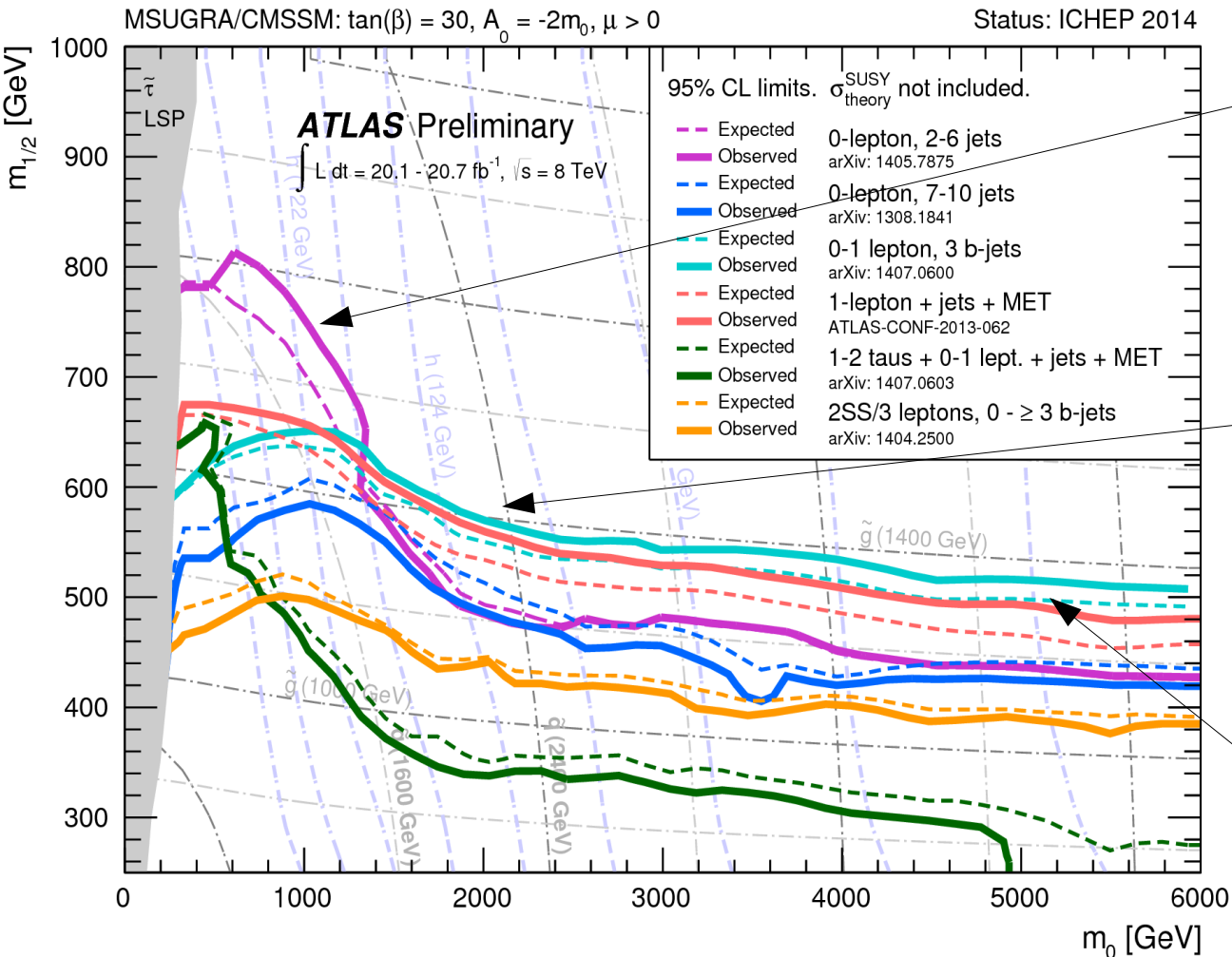
pMSSM: 19 parameters

CMSSM: 4 parameters

The smaller the number
Of parameters, the smaller
The fraction of SUSY space
explored

CMSSM interpretation

CMSSM has 4 parameters. For fixed $\tan\beta$ phenomenology essentially only dependent on the mass of the scalars (M_0) and of the fermions ($M_{1/2}$) at SUSY breaking scale. Useful benchmark of different topologies



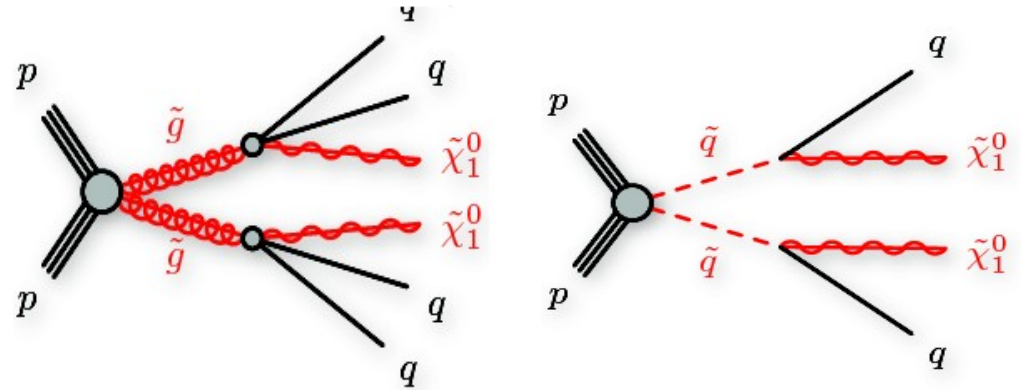
Low jet multiplicity
 0 lepton analysis:
 Excellent coverage
 Where squark
 Production dominant

Intermediate m_0 :
 1l+jets gives large
 contribution

High m_0 : only gluino
 production, decay mainly
 into 3rd generation:
 0l + 3b best analysis

pMSSM interpretation

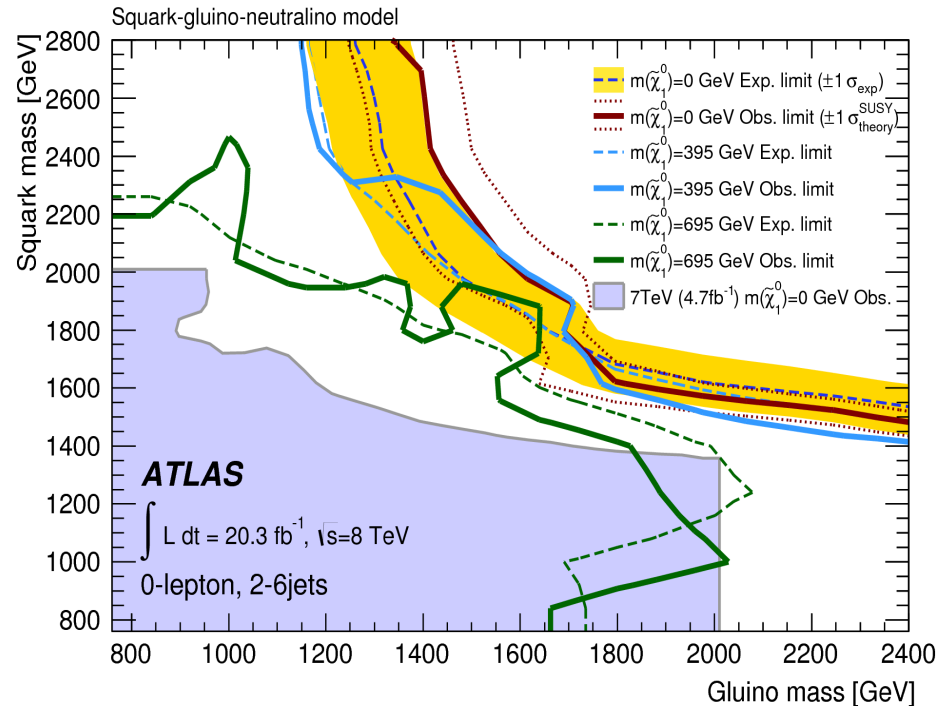
pMSSM: slice: fix all but two parameters, and choose Signature where reach mostly determined by free parameters
 Example: 1-step decays of squark and gluinos: 0 lepton signature
 All other sparticles decoupled
 Except LSP: only two decays allowed



Squark-gluino excluded up to ~ 1.5 TeV

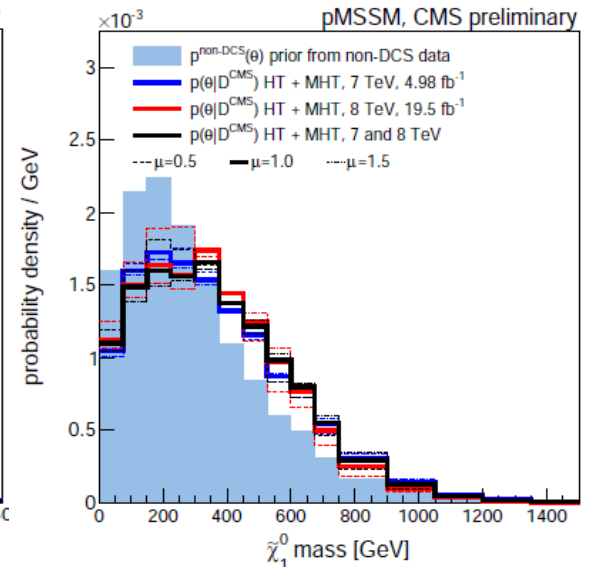
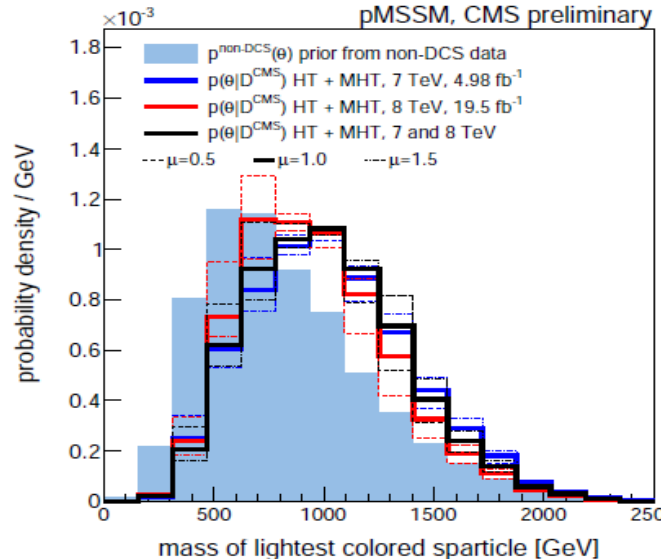
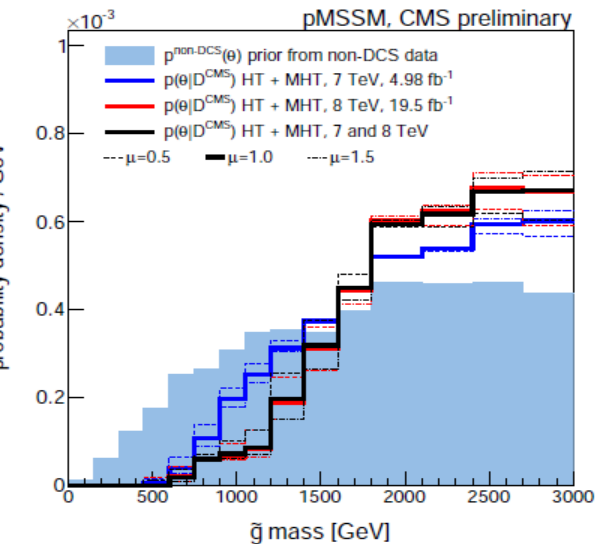
BUT

Dependence on neutralino mass



pMSSM interpretation (CMS)

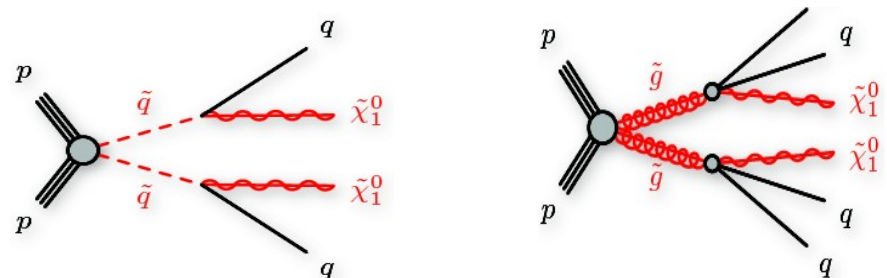
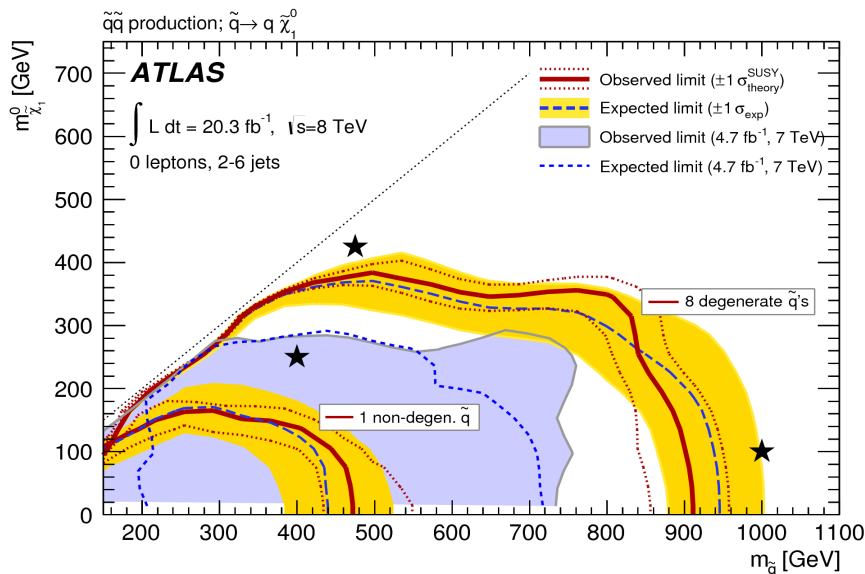
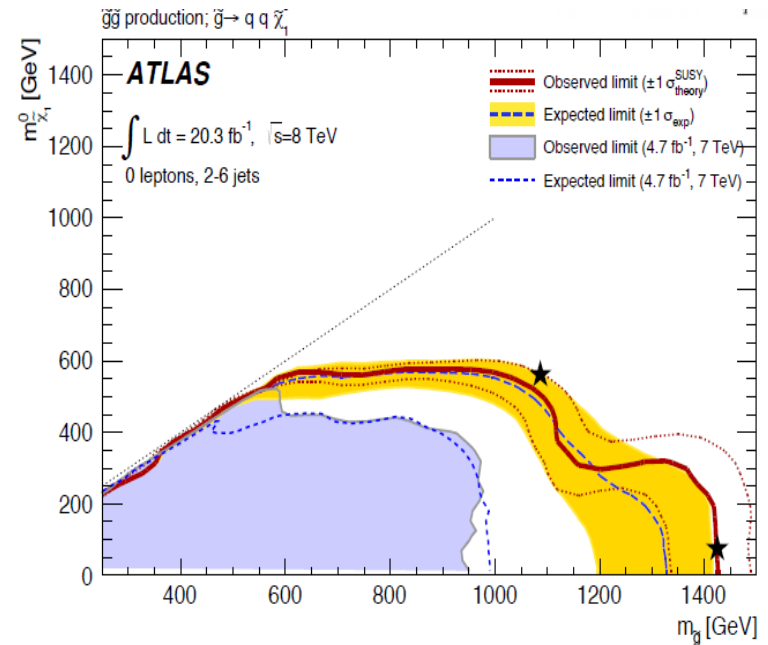
- Select large grid of points in 19-parameters space compatible with LEP and flavour constraints, neutralino LSP and sparticles lighter than 3 TeV
- Build likelihood with results of CMS EW and inclusive Ht + Emiss (+b-jets) searches
- Show marginalized distributions for sparticle masses
 - Blue are prior distributions
 - Lines are posteriors from CMS searches



“simplified model” interpretation

Simplified models as a tool for analysis optimisation and display:

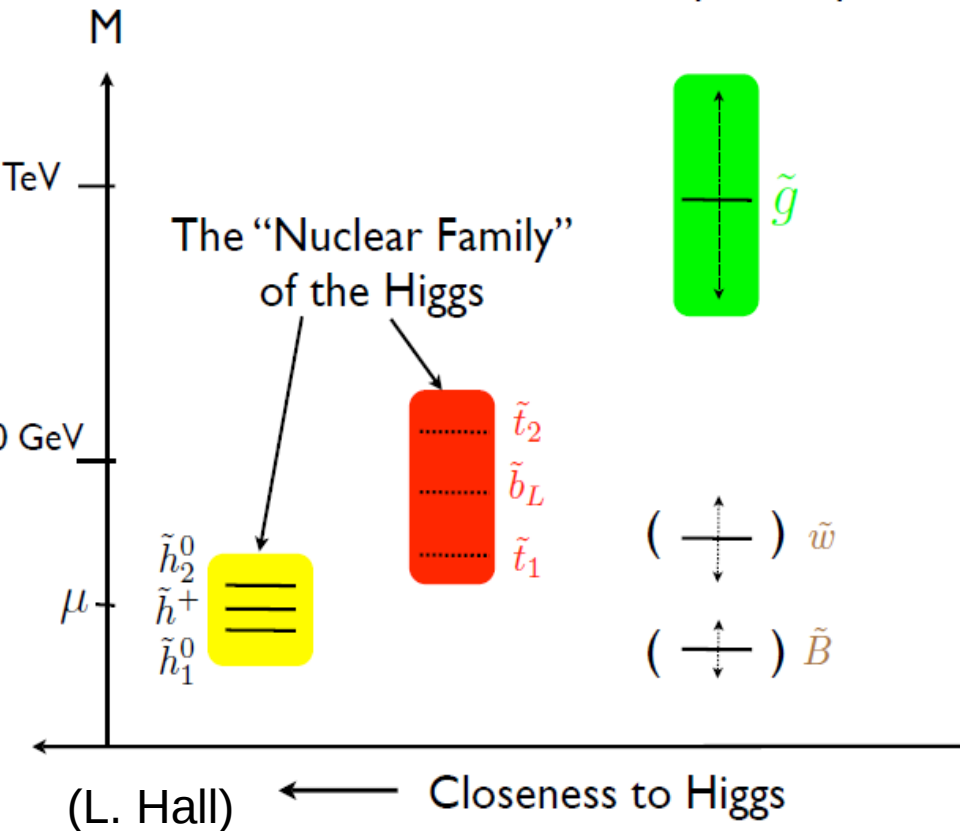
- Generate events with given decay chain on both legs
- Assume 100% BR in both legs and the SUSY production cross-section
- Express reach in 2d mass plane
- No statement on theory but very clear Representation of our potential for a specific kinematics



For low LSP mass, exclude gluinos with mass below ~ 1.4 TeV
 And squarks with mass below ~ 900 GeV

'Natural' SUSY

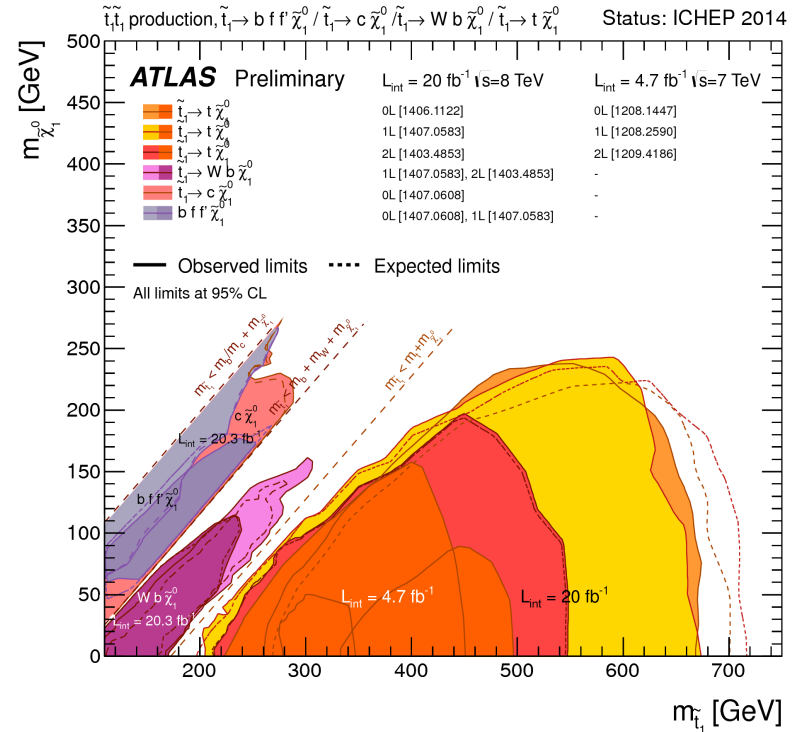
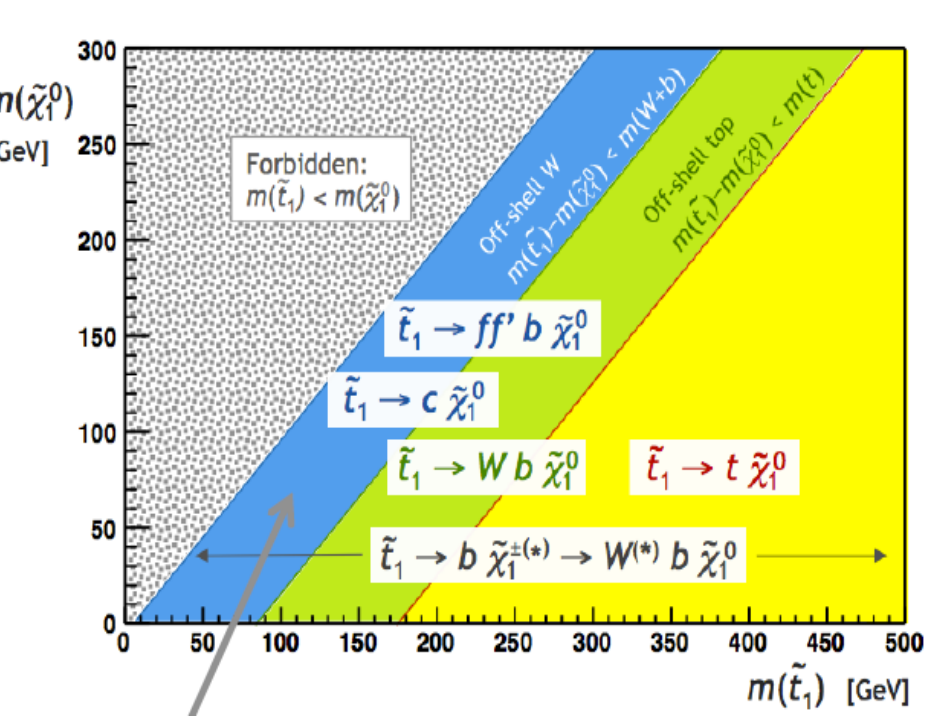
Inclusive searches with multijet+ E_{miss} + (0-2) leptons push masses of squarks of first two generations and gluinos uncomfortably high
 → dedicated searches for part of SUSY spectrum most relevant to naturalness



Assume other squarks too heavy
 Three steps:

- Search for gluino decay through real/virtual 3rd generation quarks
 - b-jets in decay
 - high multiplicity
- Search for direct production of stop/sbottom
- Try to cover all possible phenomenology in terms of decay patterns
- Search for direct production of Ewkinos (4 parameters + slepton sector)

Search for direct stop pair production



Extensive search in all possible decay channels:

2-body stop \rightarrow top LSP, stop \rightarrow chargino b, stop \rightarrow charm LSP

3-body stop \rightarrow W b LSP **4-body:** stop \rightarrow ffbar b LSP

Up to ~ 700 GeV stop mass in configurations with large visible energy

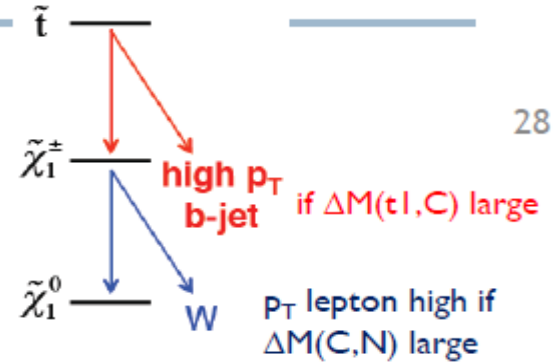
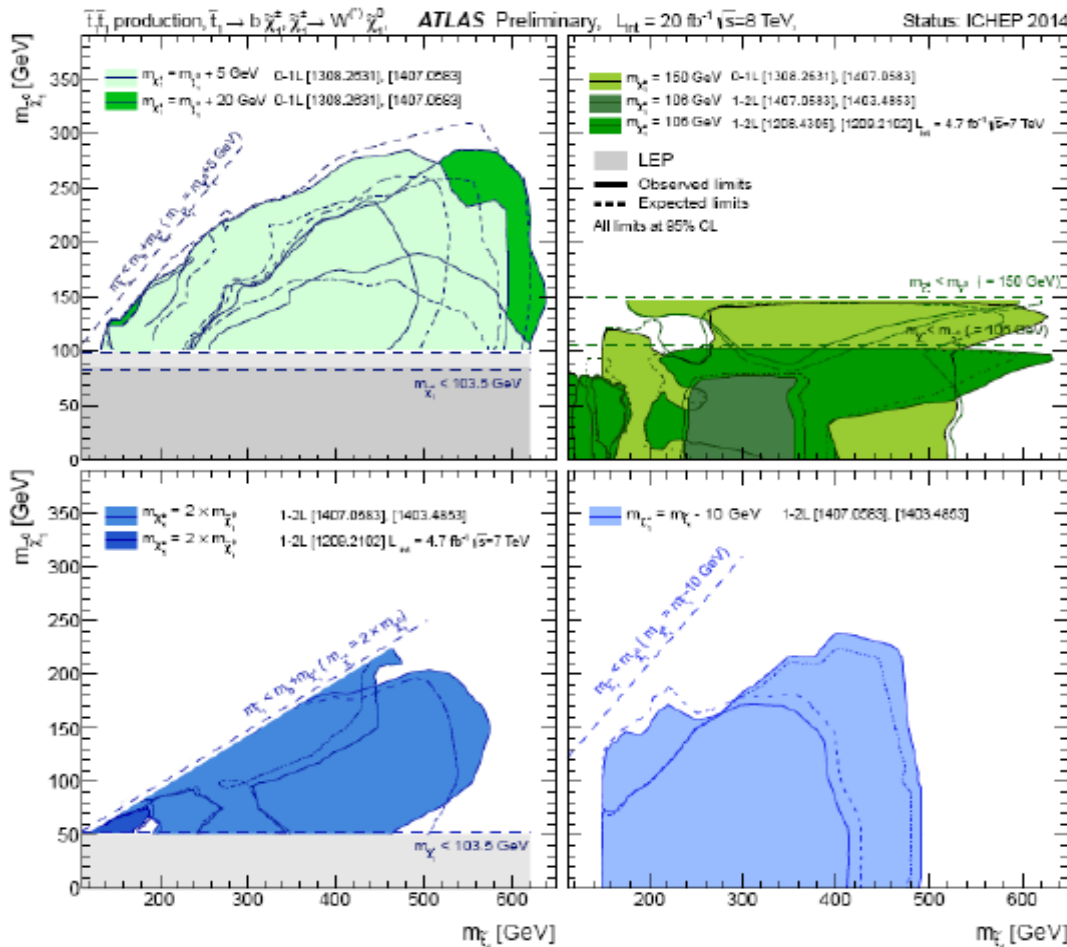
Difficult region for $m(\text{stop})=m(\text{top})+m(\text{LSP})$

For compressed topologies reach up to ~ 250 GeV with some remaining holes

Direct stop to chargino

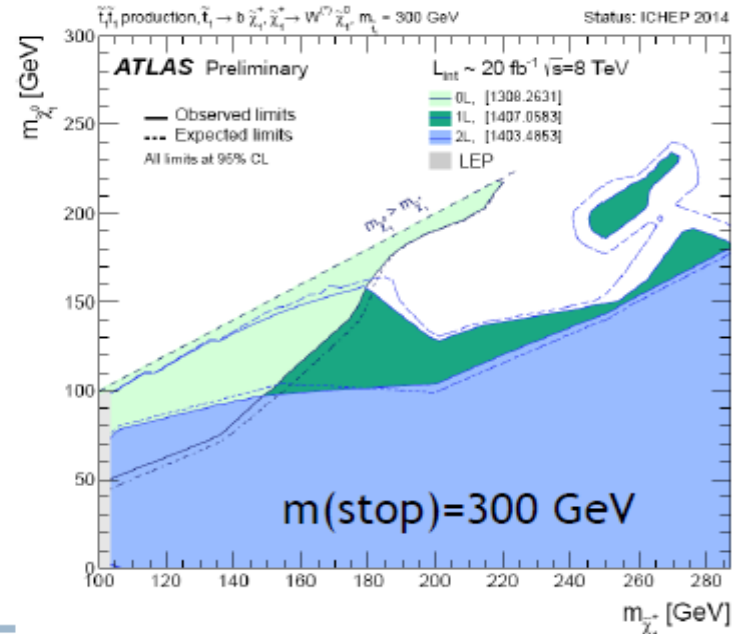
3 parameters: $m(\text{stop})$, $m(\text{chargino})$, $m(\text{LSP})$, show 2-d slices

- Various assumptions of $\Delta M(\text{stop-chargino})$ and $\Delta M(\text{chargino-neutralino})$



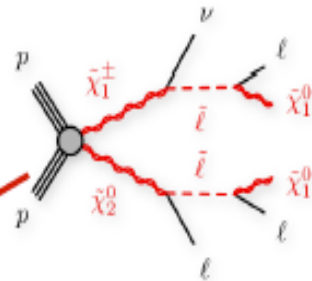
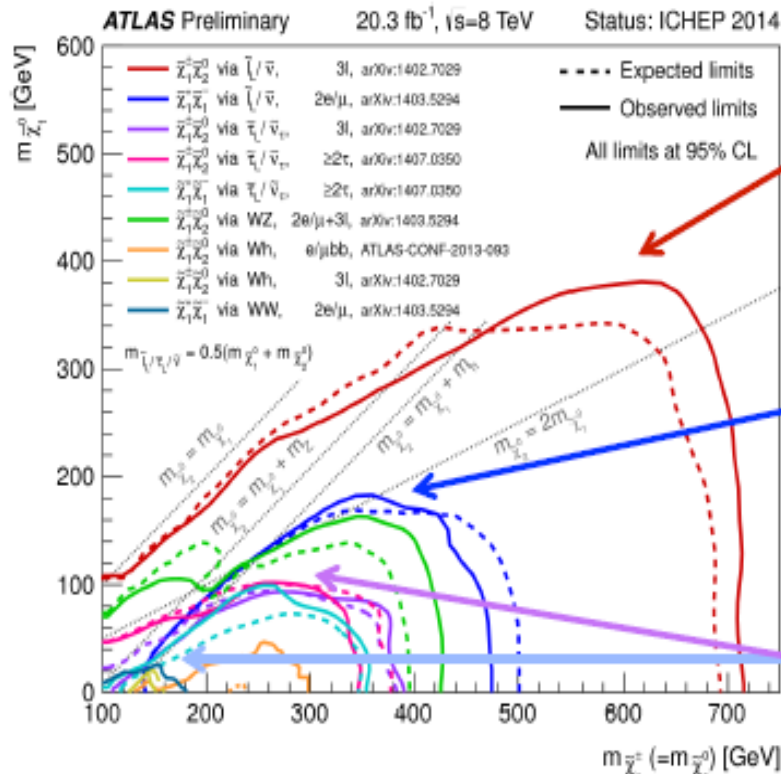
28

Fixed mass stop, function of chargino-neutralino mass

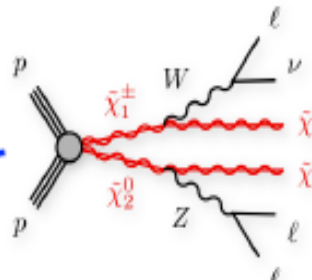


Electroweak SUSY production

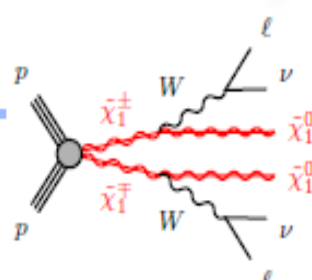
- Clean signature: multi leptons, depending on slepton masses and gauge mixture.
- Many possible models are covered by several comprehensive analyses.



Large cross-section.
 3 lepton final states if light sleptons:
 exclusion up to $m_{LSP} < \sim 380$ GeV.



If heavy slepton, smaller cross-section to lepton WZ/h with 3/2 leptons:
 up to $m_{LSP} < \sim 180$ GeV



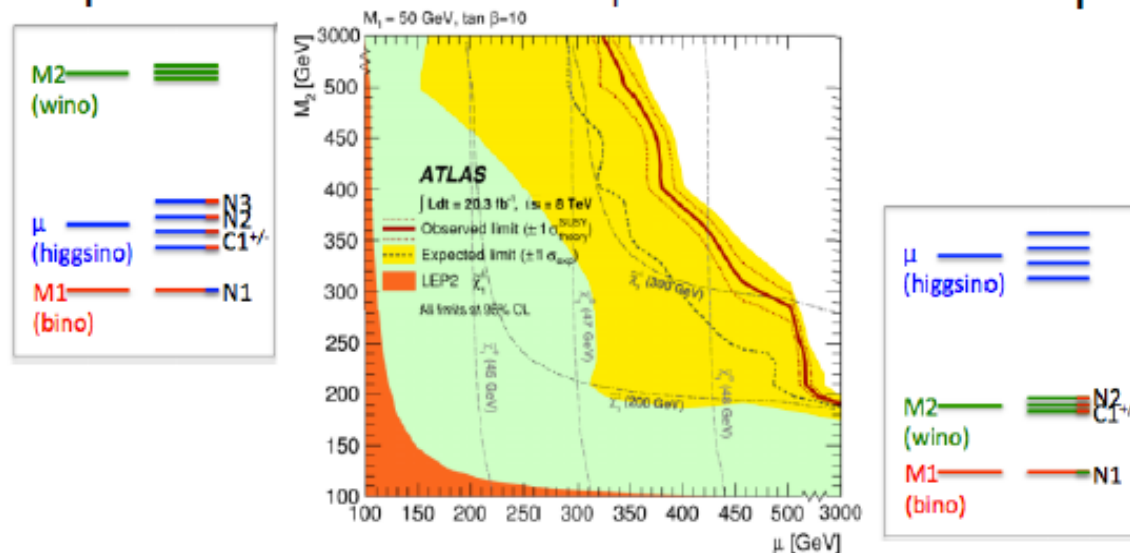
Tiny cross-section of chargino pair production via WW with 2 leptons.
 Sensitivity increased if light stau.

Interpretation of EW production in pMSSM



JHEP 05 (2014) 071

- Phenomenological MSSM (pMSSM) can put generic constraints on most of the phenomenological features of the RPC MSSM.
- Interpretation of 2-lepton+3-lepton analyses in pMSSM:
 - on higgsino μ -wino M_2 mass-plane also with very large slepton masses.
 - Assumption is bino mass $M_1 = 50\text{ GeV}$ and $\tan\beta = 10$.

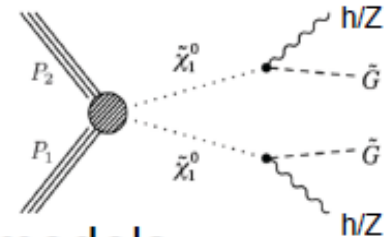




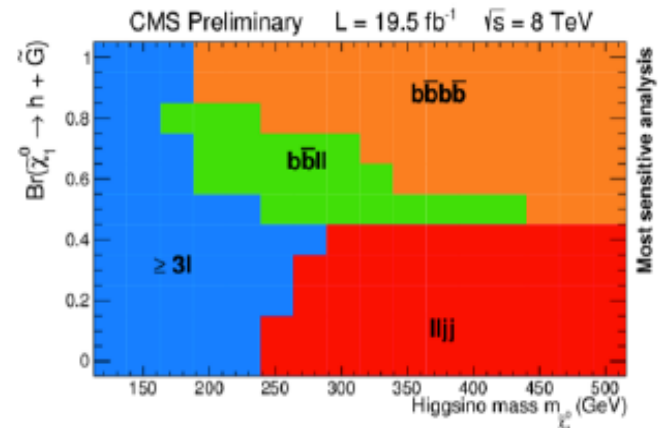
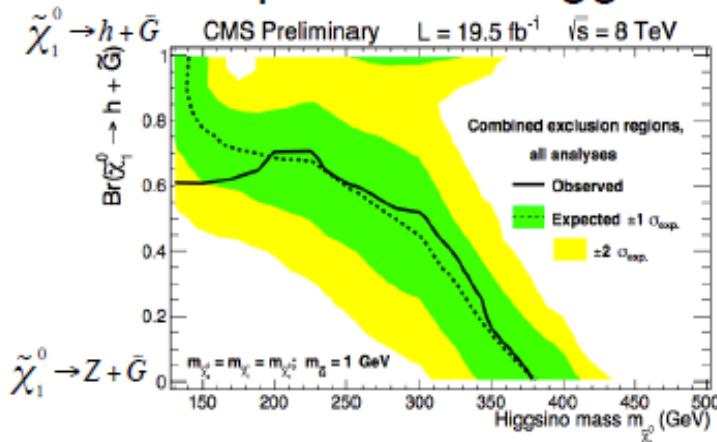
Search for EW production with Higgs

CMS-SUS-14-002, arXiv:1409.3168

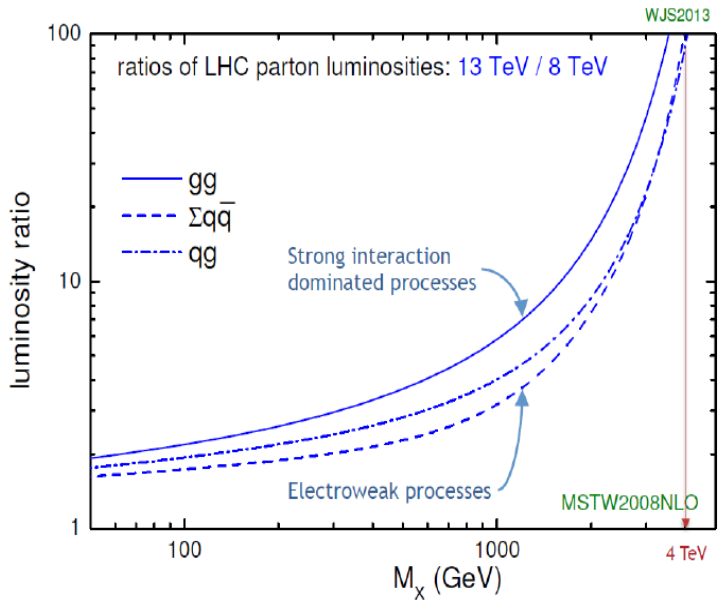
- Higgs discovery opens up new SUSY searches:
 - Lightest neutral CP -even Higgs (h) expected to be SM-like, if others heavy.
 - Neutralino can decay to h/Z +LSP.
- CMS performed comprehensive search program with di-boson + E_T^{miss} including hh, Zh, Wh .
 - $h \rightarrow ZZ, WW, \gamma\gamma, bb$



- Interpretation: Higgsino NLSP in GMSB models



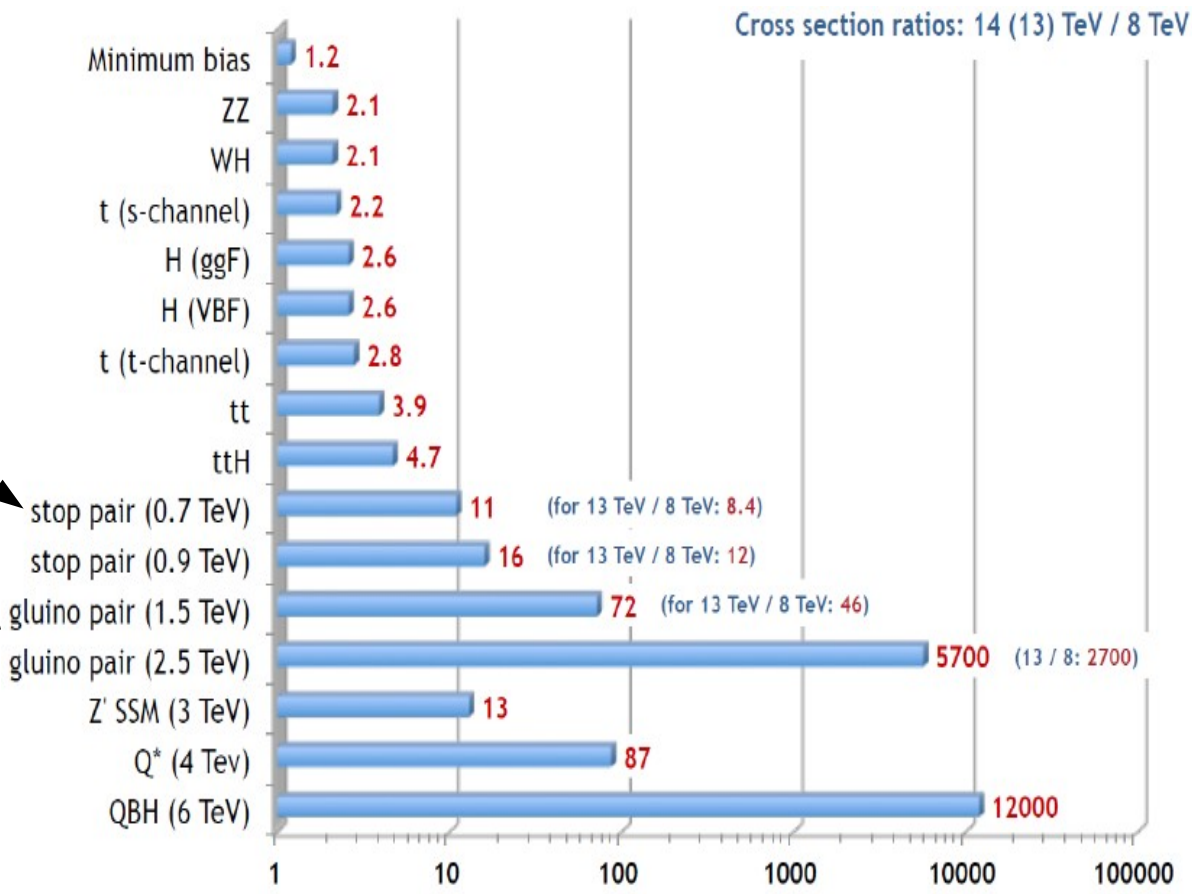
Prospects for SUSY in Run2



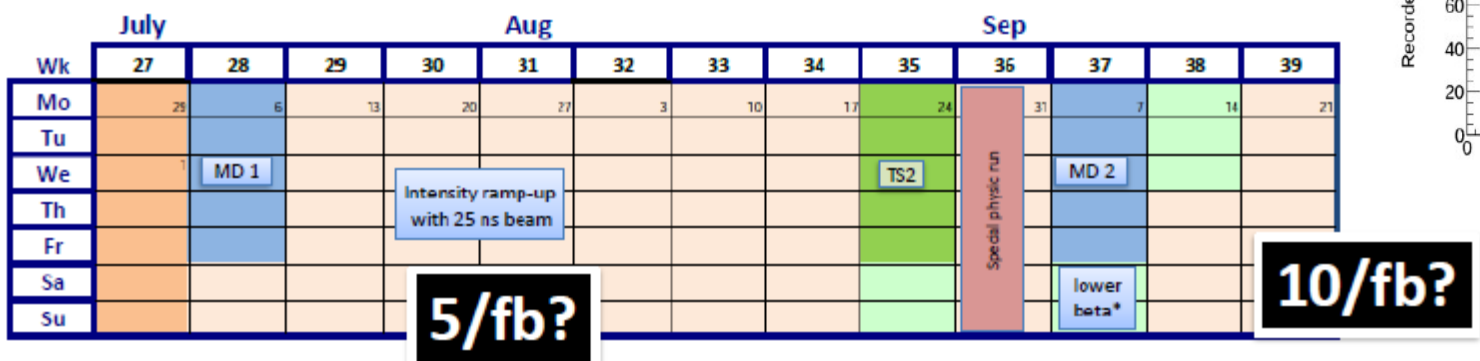
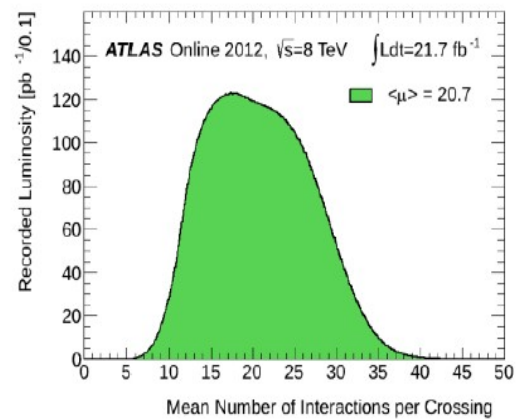
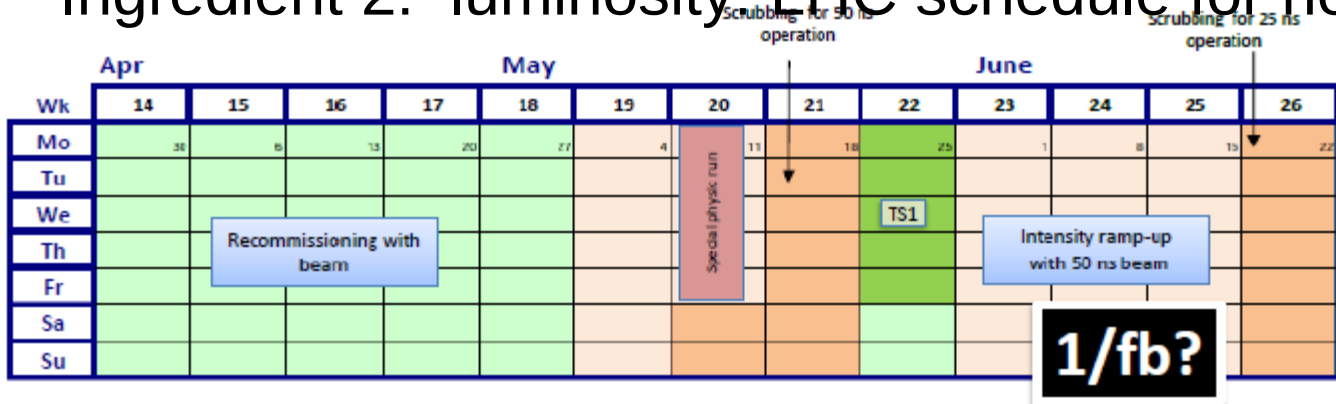
Ingredient number 1: CMS Energy

~Reach we had with 8 TeV
20 fb⁻¹

With 1 fb⁻¹ we will produce
~twice as many gluino pairs
at 1.5 TeV as in full Run 1
With 5 fb⁻¹ we will produce
~twice as many stop pairs at
0.7 TeV as in full Run 1



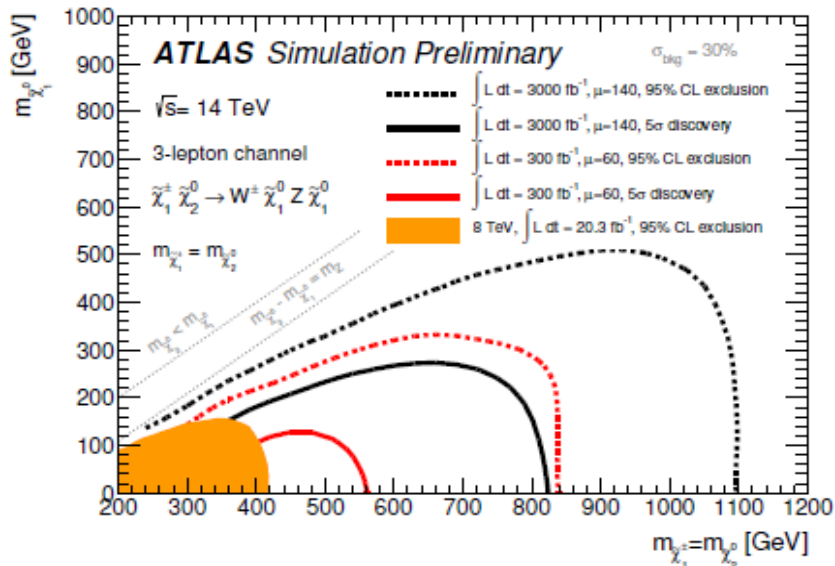
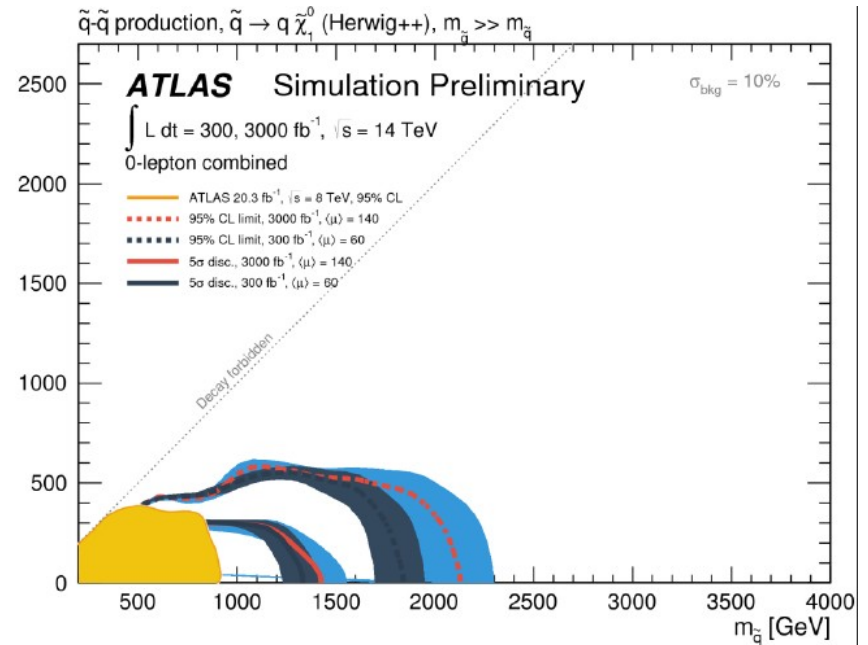
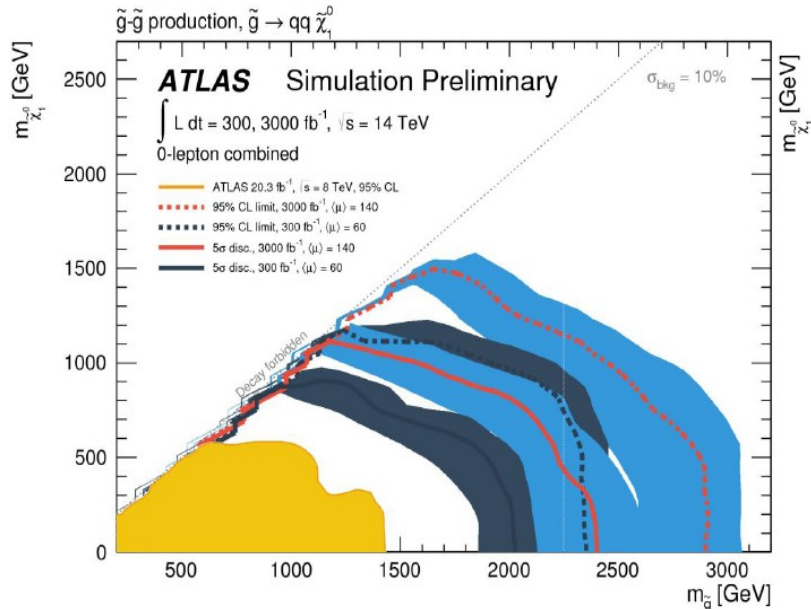
Ingredient 2: luminosity: LHC schedule for next year



https://espace.cern.ch/pe-aep/bc/departmentalDocuments/BE/LHC_Schedule_2015.pdf

| | <u>Nc</u> | Beta * | ppb | <u>EmitN</u> | <u>Lumi</u> | Days (<u>approx</u>) | <u>Int lumi</u> | Pileup |
|--------|-----------|-----------|--------|--------------|-------------|---------------------------|----------------------|--------|
| 50 ns | 1300 | 80 | 1.2e11 | 2.5 | 4.6e33 | 21 | ~1 fb ⁻¹ | 27 |
| 2015.1 | 2496 | 80 | 1.1e11 | 2.5 | 7.4e33 | 44 | 5.1 fb ⁻¹ | 22 |
| 2015.2 | 2496 | 40 | 1.1e11 | 2.5 | 1.3e34 | 46 | 9.2 fb ⁻¹ | 39 |

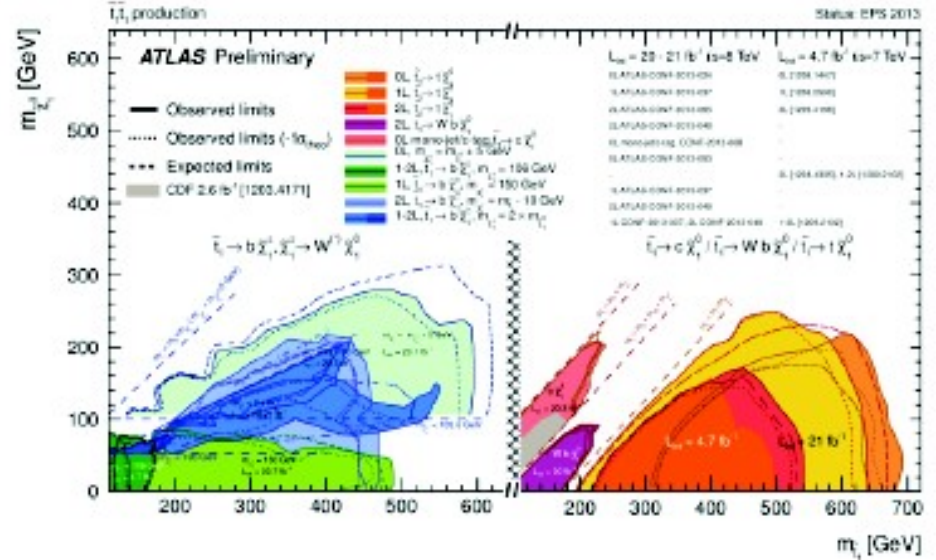
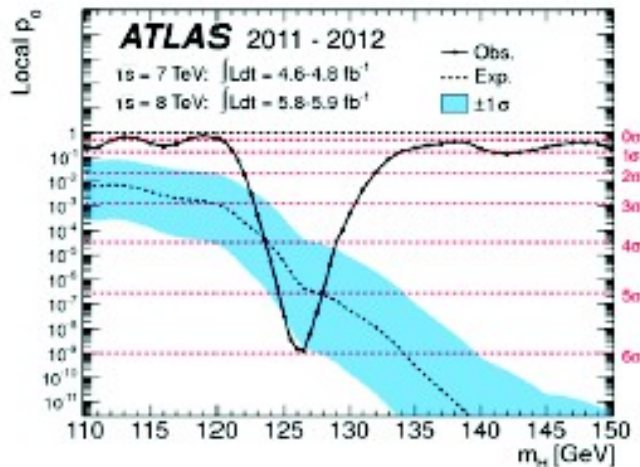
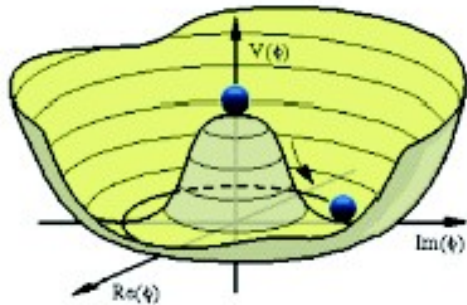
Longer term perspective



For high luminosity running need
 To take into account large pileup
 Which will smear Emiss.
 Simulation done in two scenarios:
 $\langle \mu \rangle = 60$ for 300 fb^{-1}
 $\langle \mu \rangle = 140$ for 3000 fb^{-1}

Exotic searches

- No precise model to guide us.
- No unified parameter phase space to map results

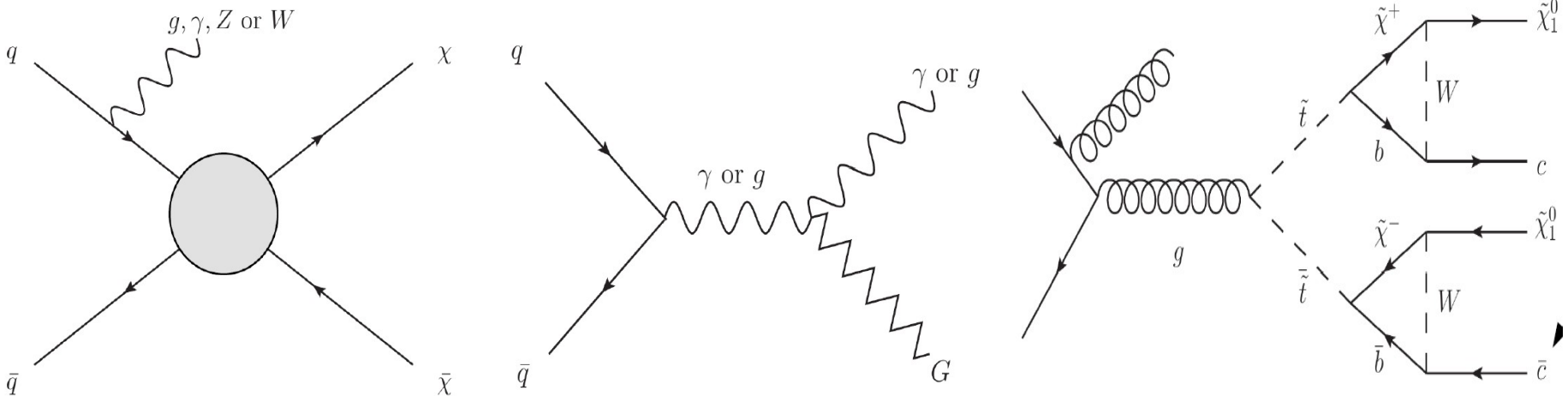


Strategy

- Address wider range of final state topologies
- Concentrate on topologies:
 - Giving easily identifiable signature
 - Largely model independent or predicted by several classes of models. Examples
 - Mono-object+Emiss
 - Resonances
 - High multiplicity final states
 - Predicted by well motivated theoretical speculation. Examples from naturalness:
 - Top partner
 - Contact interactions
- Concentrate in the following on Mono-X, most recent and hottest topic

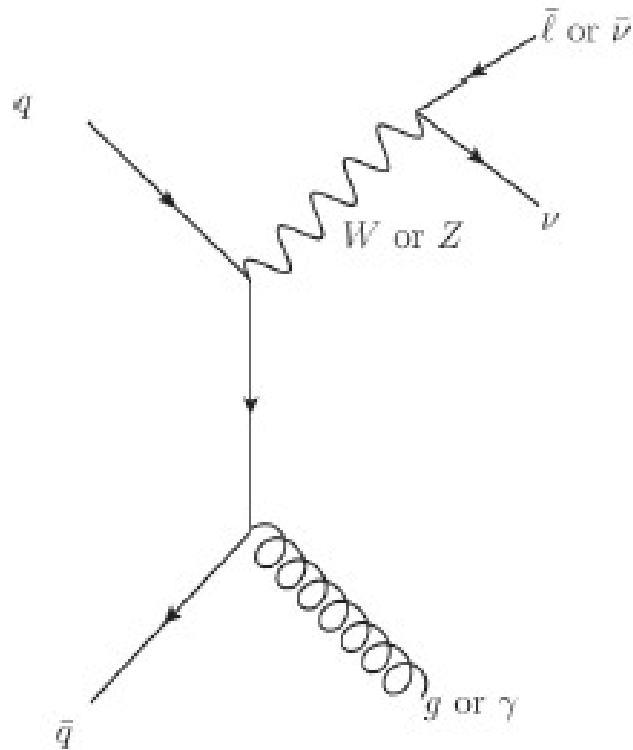
The mono-X signature

- A single high pt object (jet, photon, W, Z) associated with large E_{miss} can be produced by several different BSM processes such as:
 - Invisible particles produced in association with QCD or EWK initial state radiation (ISR). Example: Dark Matter
 - Two-body production of gravitino/stop recoiling against photon/gluon
 - Production of particles decaying into an almost degenerate invisible particle: need to rely on ISR to extract visible signal.



The mono-X signature

- Simple final state with well-known backgrounds from electroweak processes



Use same estimation techniques as described for multijet+MET SUSY searches,

Main differences:

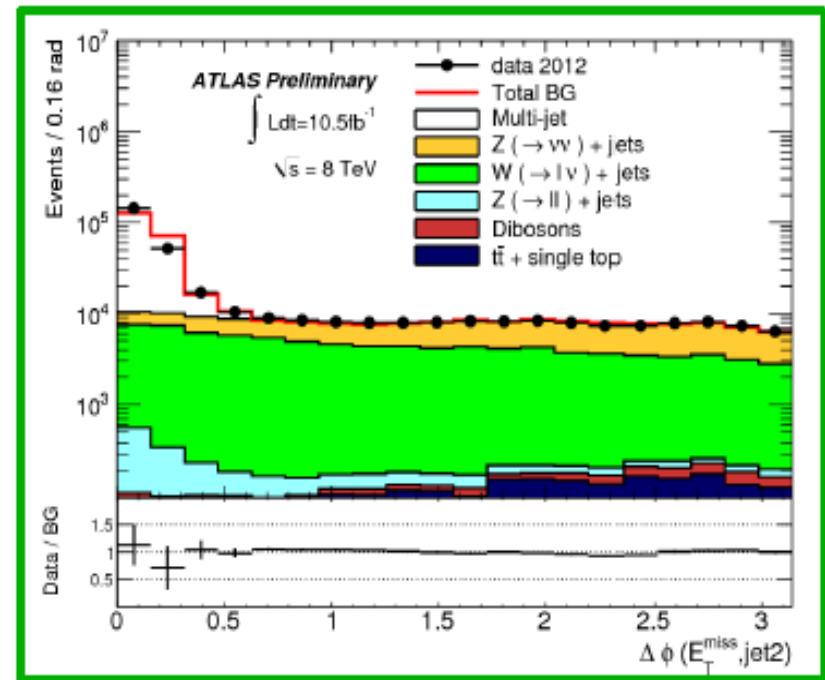
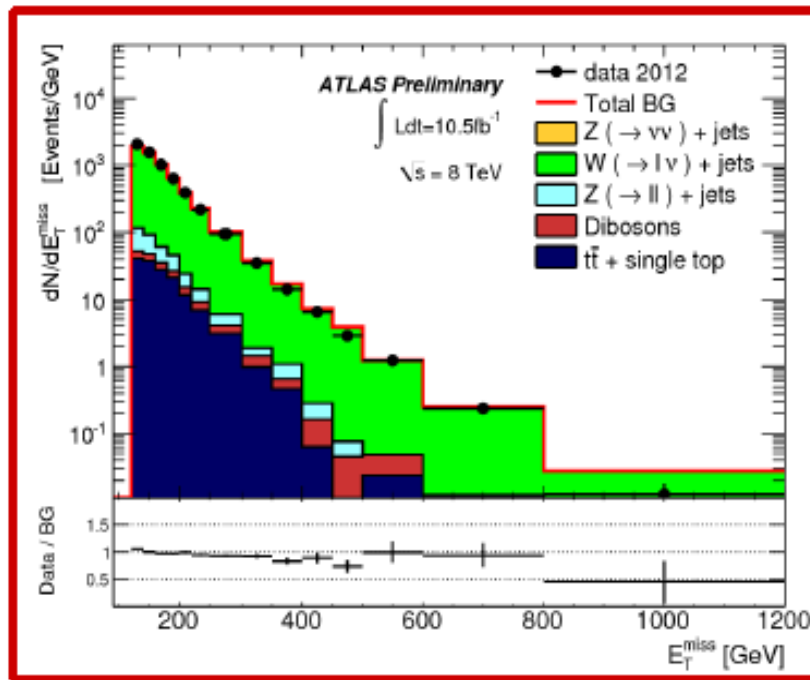
- Low jet multiplicity
- Hard kinematics

The mono-x analyses

- Select events with a high pt object (jet, photon, lepton hadronically decaying W/Z) and large MET
- Veto events in which:
 - A lepton is identified: *remove electroweak background*
 - There are more than 2 jets: *remove top or multijets*
 - MET is pointing along an jet: *remove fake MET from mismeasured jets*
- Estimate from data main backgrounds:
 - $(Z \rightarrow \nu\nu) + X$ (irreducible)
 - $(W \rightarrow l\nu) + X$, $(Z \rightarrow ll) + X$, with lost lepton
 - Multi-jets, γ +jets with fake MET
 - Non-collision events
- Estimate from MC smaller backgrounds: top, diboson

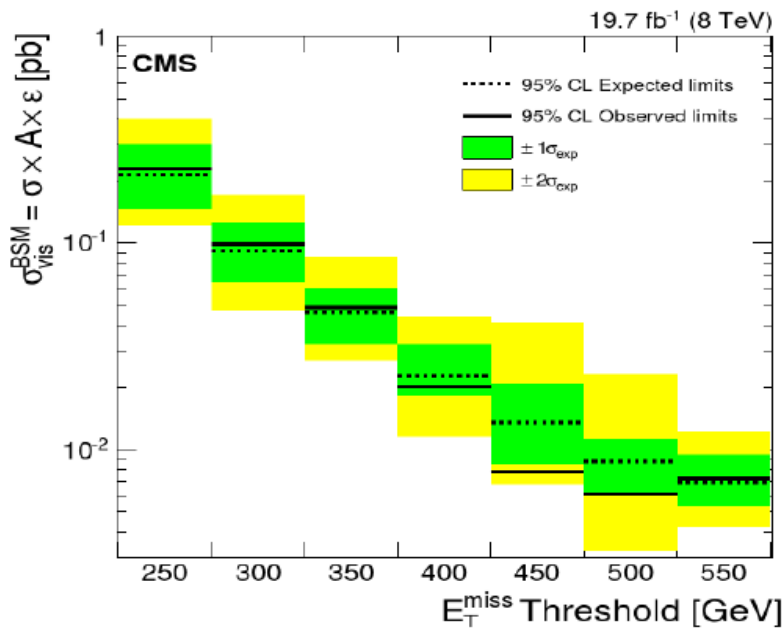
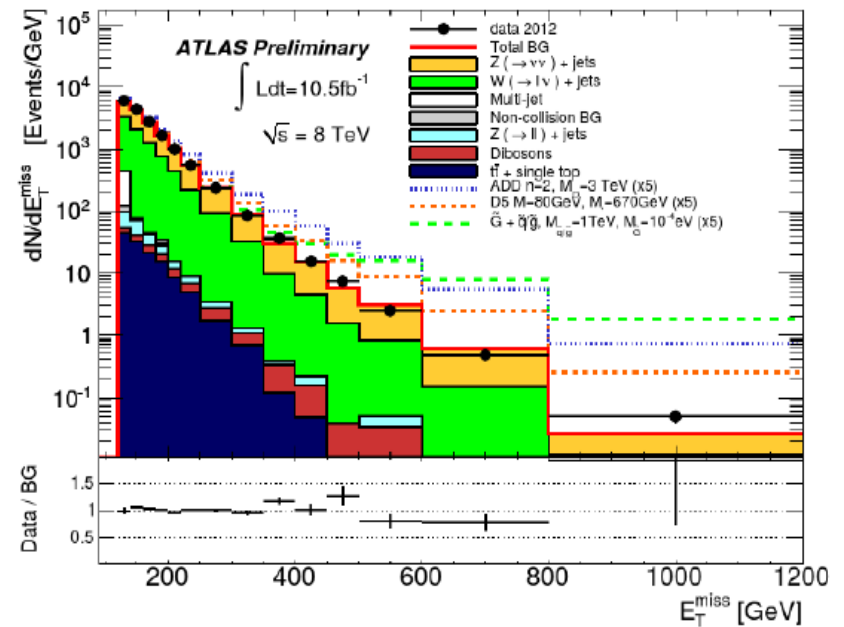
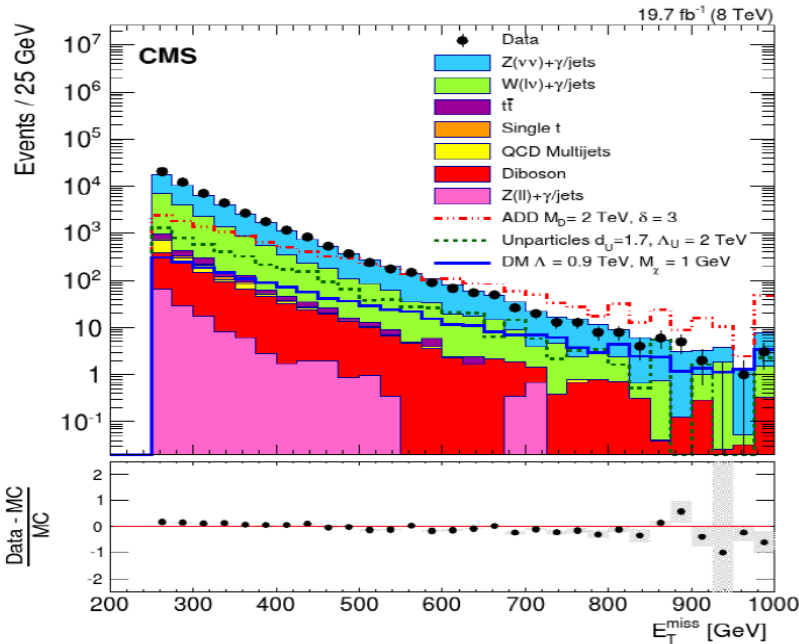
Monojet Analysis: Backgrounds

$Z \rightarrow \nu\bar{\nu} (\mu\bar{\mu}, e\bar{e}, \tau\bar{\tau}) + \text{jets}$ and $W \rightarrow e\nu (\mu\nu, \tau\nu) + \text{jets}$ are estimated using CR with leptons enriched in $W \rightarrow \mu\nu + \text{jets}$ and $W \rightarrow e\nu + \text{jets}$ (also cross checked with a $Z \rightarrow \mu\bar{\mu} + \text{jets}$ CR)



Multijet background estimated using a MET+ 2(3) jets CR where $\Delta\phi(\text{MET}, \text{jet}2(3)) < 0.5$

Monojet Analysis: Results



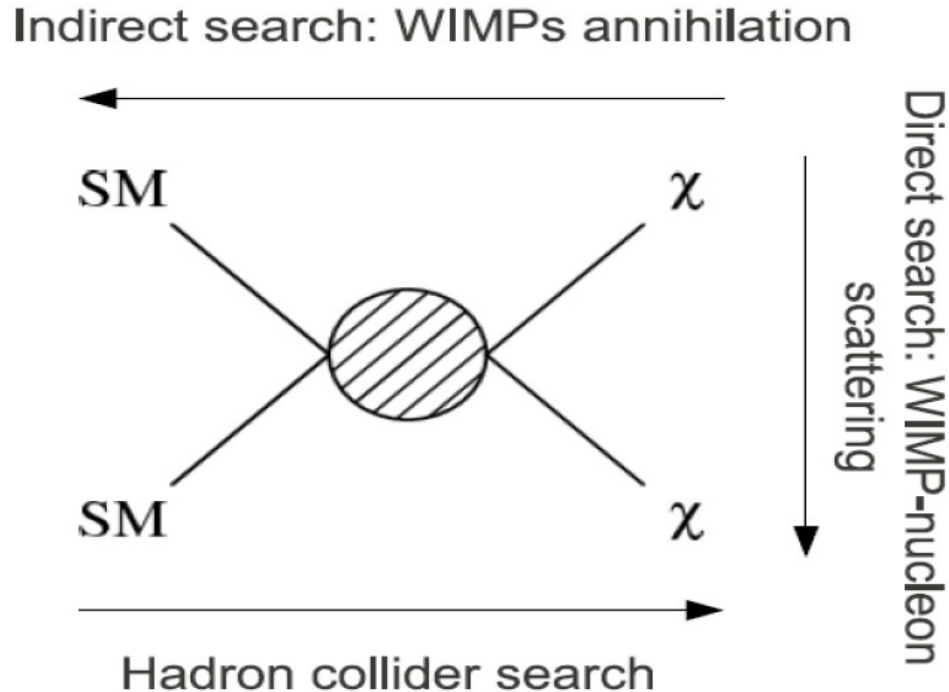
ATLAS results for 10 fb⁻¹ at 8 TeV
CMS results for 19.5 fb⁻¹

Good agreement of data with SM
Expectation used to set a
Model-independent limit on
Cross-section for new physics

Monojet/monophoton analysis: interpretations

- Dark Matter production
- Graviton production in Extra Dimensions
- Gravitino production in GMSB models
- Degenerate SUSY models:
 - Light stop
 - Higgsinos

Dark matter interpretation

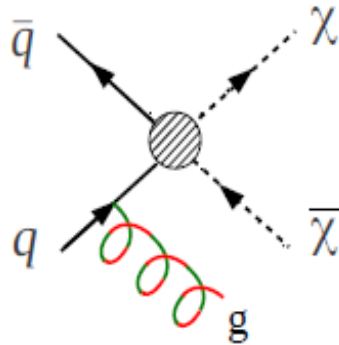


DM production at Colliders test same process as direct and indirect searches.

Need to put some theory in the blob to allow comparison

Two main model approaches

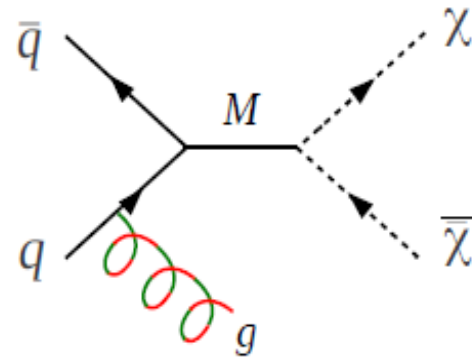
Can't resolve mediator \rightarrow use Effective Field Theory of contact interaction at scale Λ :



$$O_V \sim \frac{(\bar{\chi} \gamma_\mu \chi)(\bar{q} \gamma^\mu q)}{\Lambda^2}$$

$$O_A \sim \frac{(\bar{\chi} \gamma_\mu \gamma^5 \chi)(\bar{q} \gamma^\mu \gamma^5 q)}{\Lambda^2}$$

Light mediator of mass $M \rightarrow$ use Simplified Theory:



$$\mathcal{L}_V \sim \frac{i g_\chi g_q}{q^2 - M^2} (\bar{\chi} \gamma_\mu \chi)(\bar{q} \gamma^\mu q)$$

$$\mathcal{L}_A \sim \frac{i g_\chi g_q}{q^2 - M^2} (\bar{\chi} \gamma_\mu \gamma^5 \chi)(\bar{q} \gamma^\mu \gamma^5 q)$$

Nomenclature for the interactions:

"V" \rightarrow vector; "A" \rightarrow axial-vector; "S" \rightarrow scalar (describes gluon fusion with $O_S \sim \frac{\bar{\chi} \chi}{4\Lambda^3} \alpha_s (G_{\mu\nu}^a)^2$)

EFT and ST are equivalent for $q \ll M \rightarrow \Lambda = \frac{M}{\sqrt{g_\chi g_q}}$

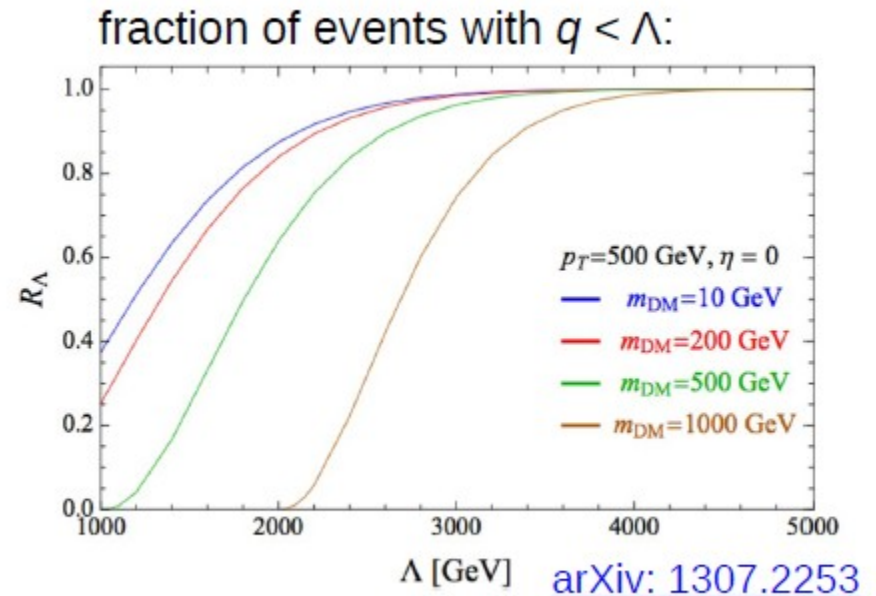
EFT vs simplified model

EFT

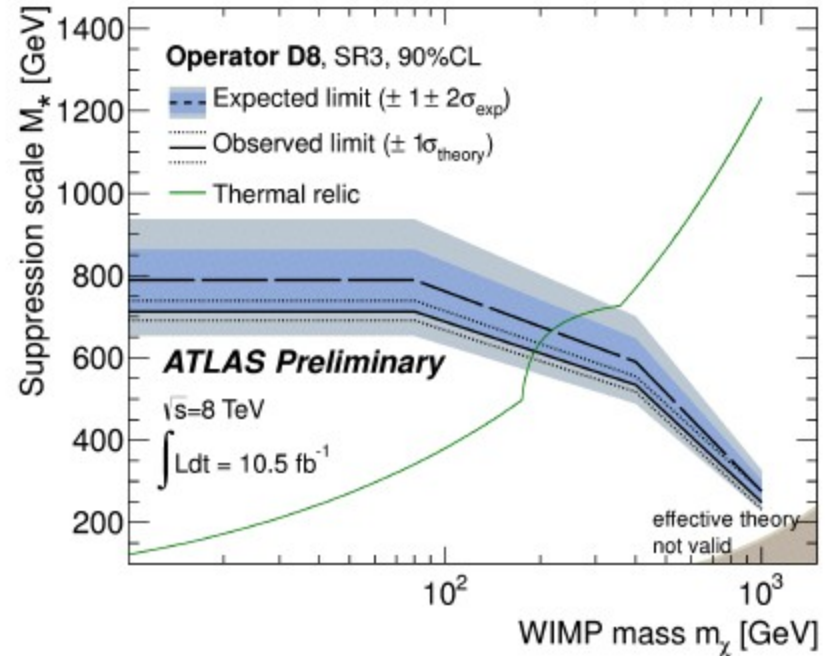
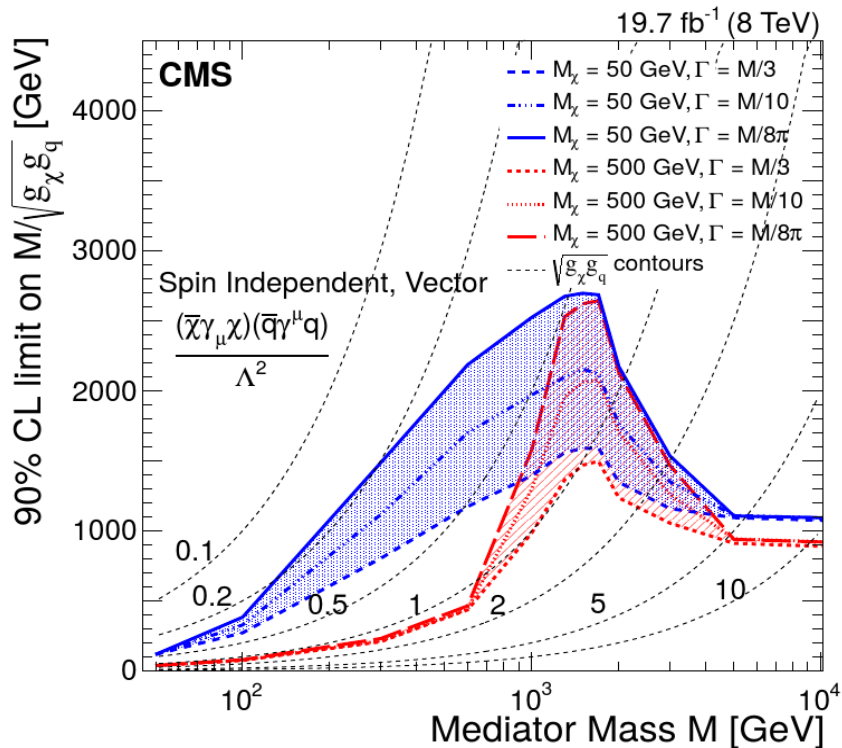
- Simple parameter space
 Λ and m_χ
- Breaks down when $q > \Lambda$

Simplified model

- UV complete
- Larger parameter space:
- M, m_χ, g_q, g_χ



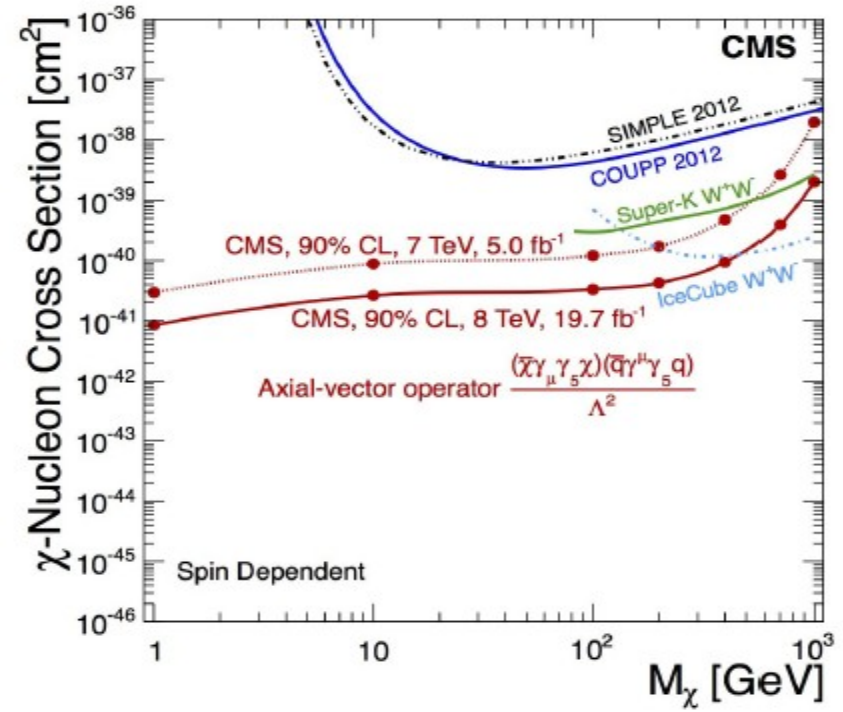
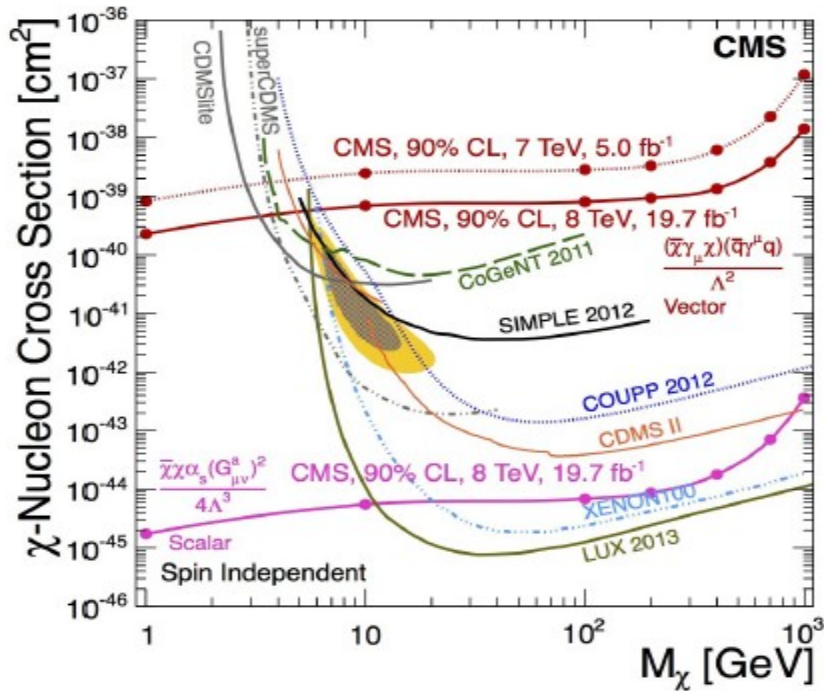
Example of limits in two approaches



- Light mediator or large couplings are ruled out
- Resonant structure
- Reduced to EFT for high M

D8 == Axial Vector
 For a wide range of χ masses
 The limit is order 0.8 TeV

EFT WIMP interpretation



Direct detection experiments use the same EFT

Limits can be translated on limits on χ -nucleon cross-section

EFT always valid for direct detection (low q).

For colliders, would need to integrate out high q events, depending on assumed mediator mass.

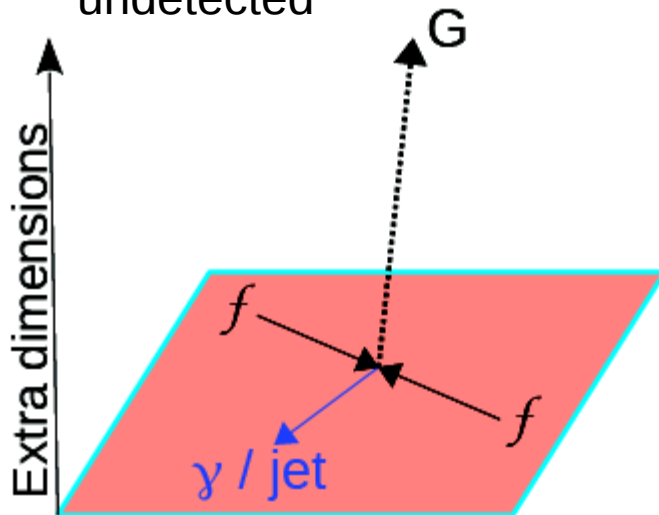
Interpretation: graviton in extra-dimensions

ADD model: gravity propagates in n Extra Dimension compactified on a radius R .

Characteristic scale of gravity is M_D given by

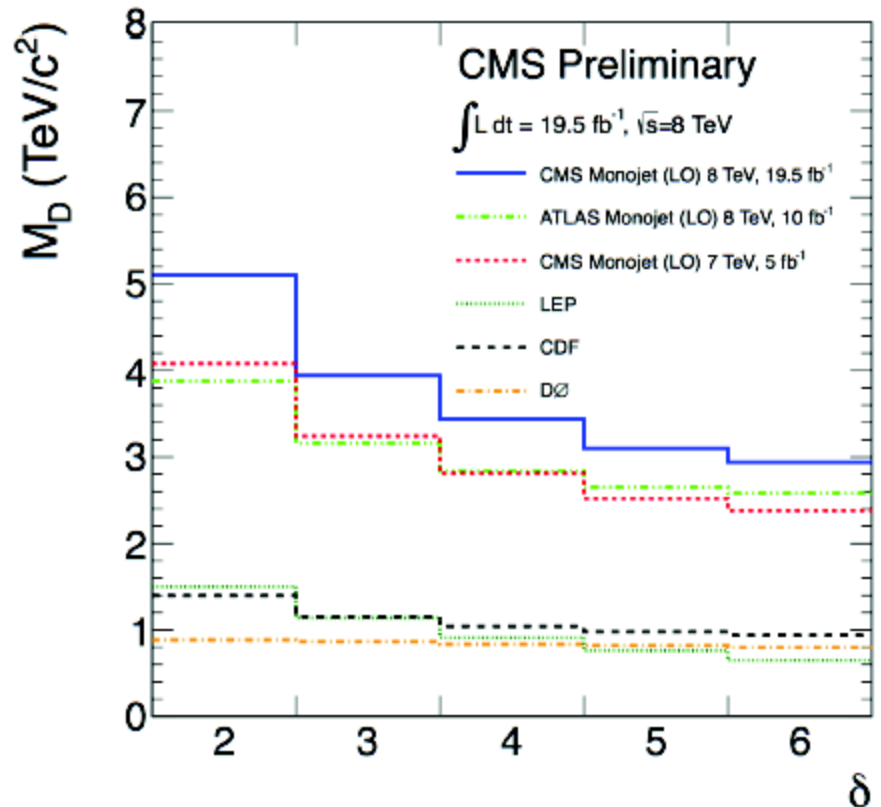
$$M_{Pl}^2 = M_D^{2+n} R^n$$

Graviton escapes in ED and goes undetected



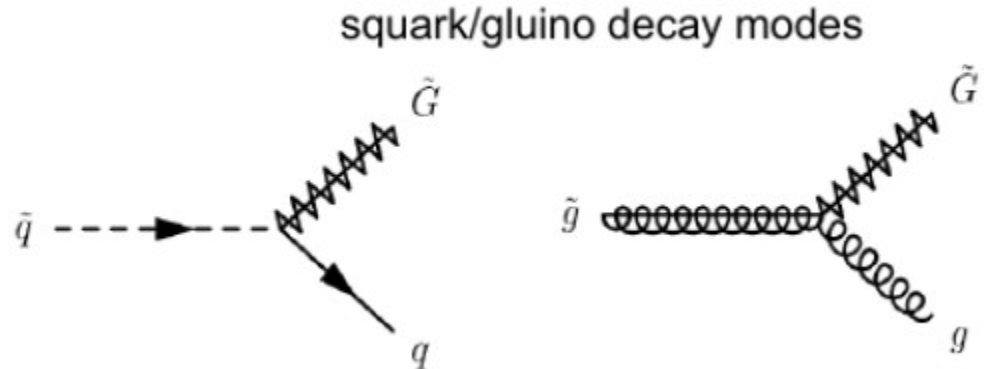
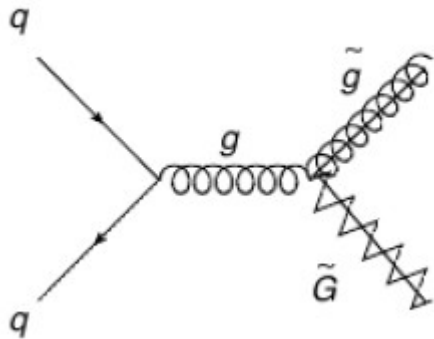
Can produce a KK tower of graviton states
Recoiling against a jet or a photon

$$\sigma_{KK} \sim \frac{1}{M_D^2} \left(\frac{\sqrt{s}}{M_D} \right)^n$$



Limit on M_D between 3 and 5 TeV depending on n

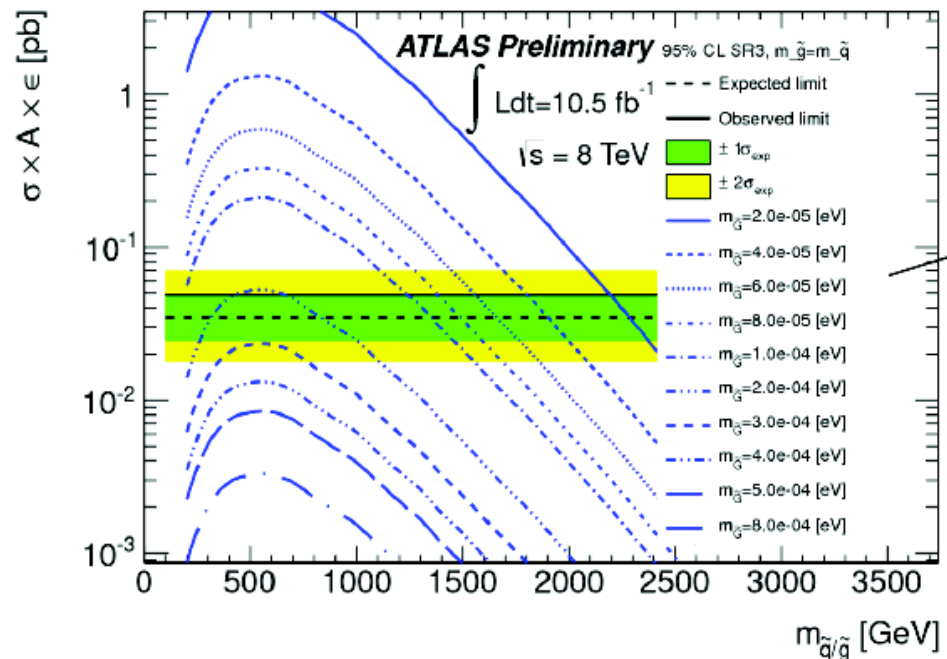
Interpretation: gravitinos in GMSB



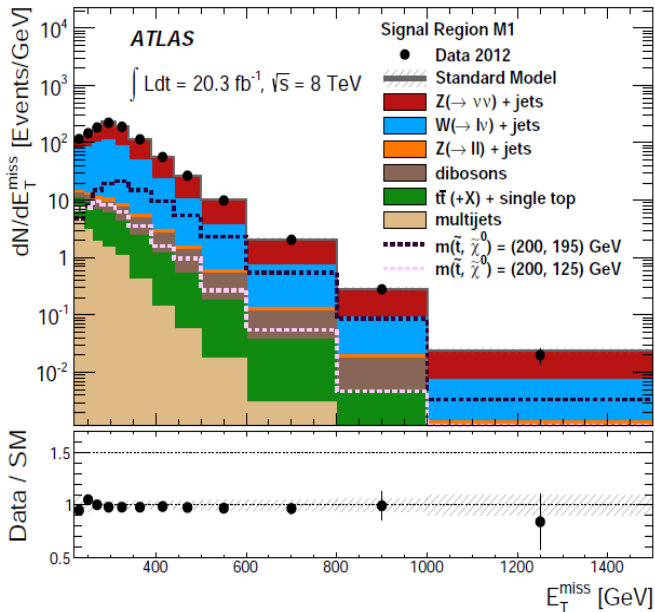
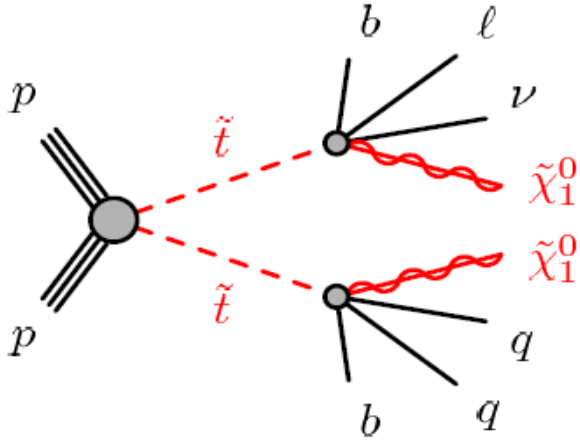
In GMSB model light gravitino often LSP

Study associated production of Gravitino with squark/gluino
Squark/gluino in turn decay into jet+Gravitino: monojet signature

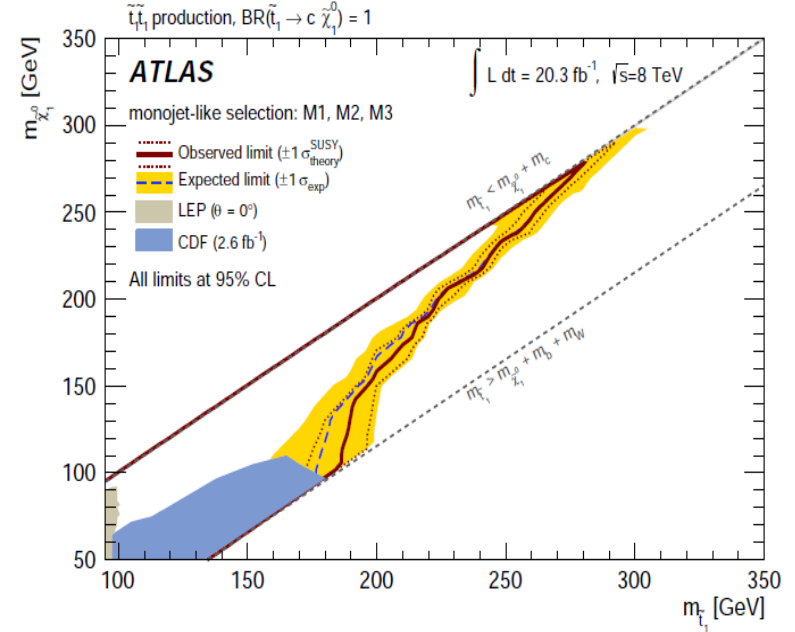
For a 1 TeV squark/gluino exclude A gravitino with mass above 1e-4eV



Interpretation: stop



M1: $P_{Tj} > 280$ GeV, $E_{Tmiss} > 220$ GeV



Search for 4-body decay of stop

Require:

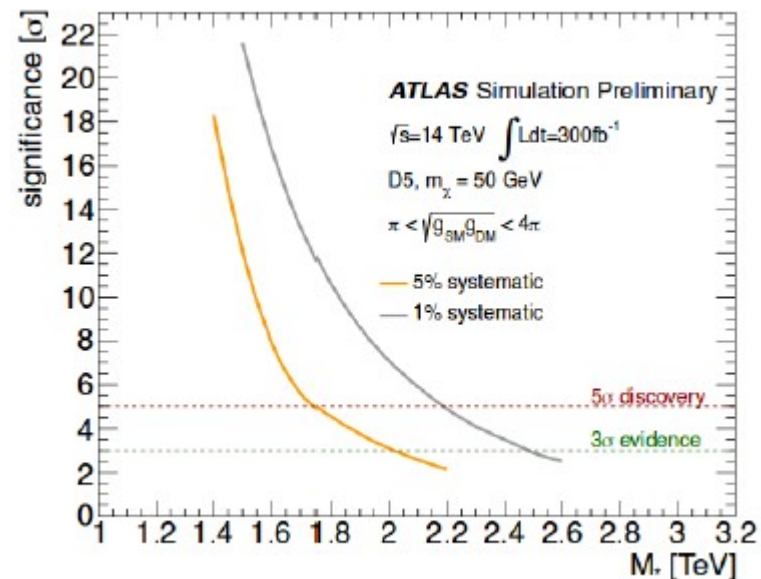
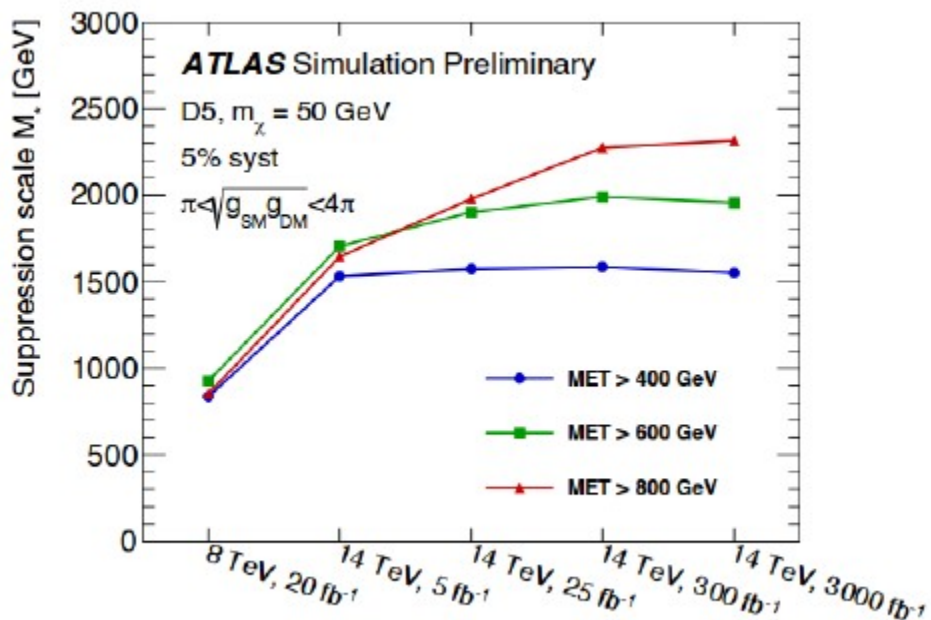
One high p_T jet and MET,

No more than 3 jets with $p_T > 30$ GeV

Lepton Veto

$\Delta\phi(\text{jets, MET}) > 0.4$

Outlook on monojet searches



Significant improvement in sensitivity expected with early run 2 data:
 Exclusion limit on mediator mass improved by a factor 2 with first few fb⁻¹
 5 σ discovery potential for $M^* \sim 1.7$ TeV with 300 fb⁻¹

Conclusions

- Searches for new physics performed on very broad range of signatures, addressing many BSM models on Run 1 LHC data
- Null results strongly constrain BSM model space
- Squarks of first two generations and gluinos heavy $> \sim \text{TeV}$
- Good Run 1 coverage also for production of stop and EWKinos
- Through mono-X analysis constraints on production of Dark Matter
- Run 2 will open a further kinematic region, experiments are ready to take advantage of the opportunity

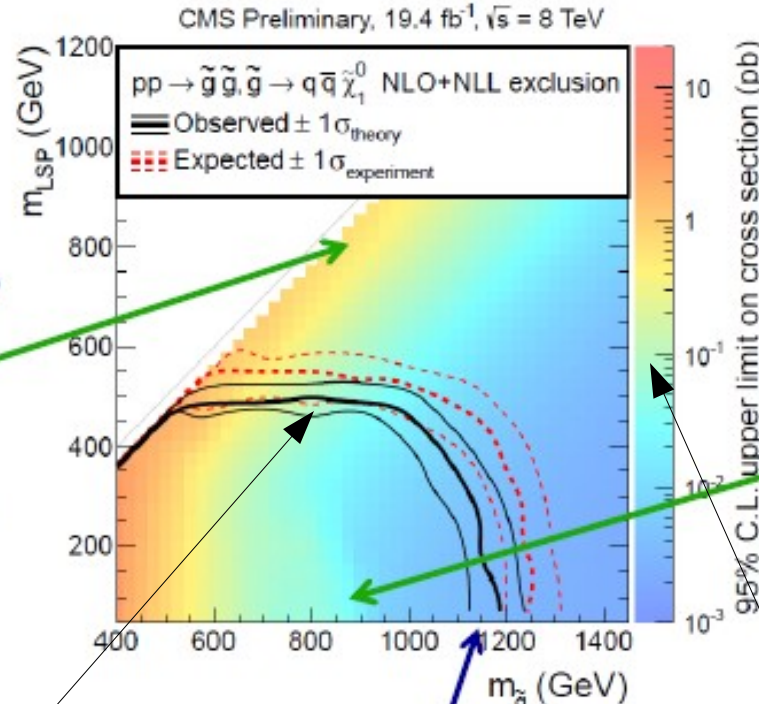
How to read a simplified model plot

“Compressed case”
Worse cross-section limit

Small Δm :
Softer jets,
less MET

— Gluino mass
— $\tilde{\chi}_1^0$ mass

mass ↑



For light LSP, gluinos excluded to >1100 GeV

“Large Δm case”
Better cross-section limit

Large Δm :
Harder jets,
more MET

— Gluino mass
— $\tilde{\chi}_1^0$ mass

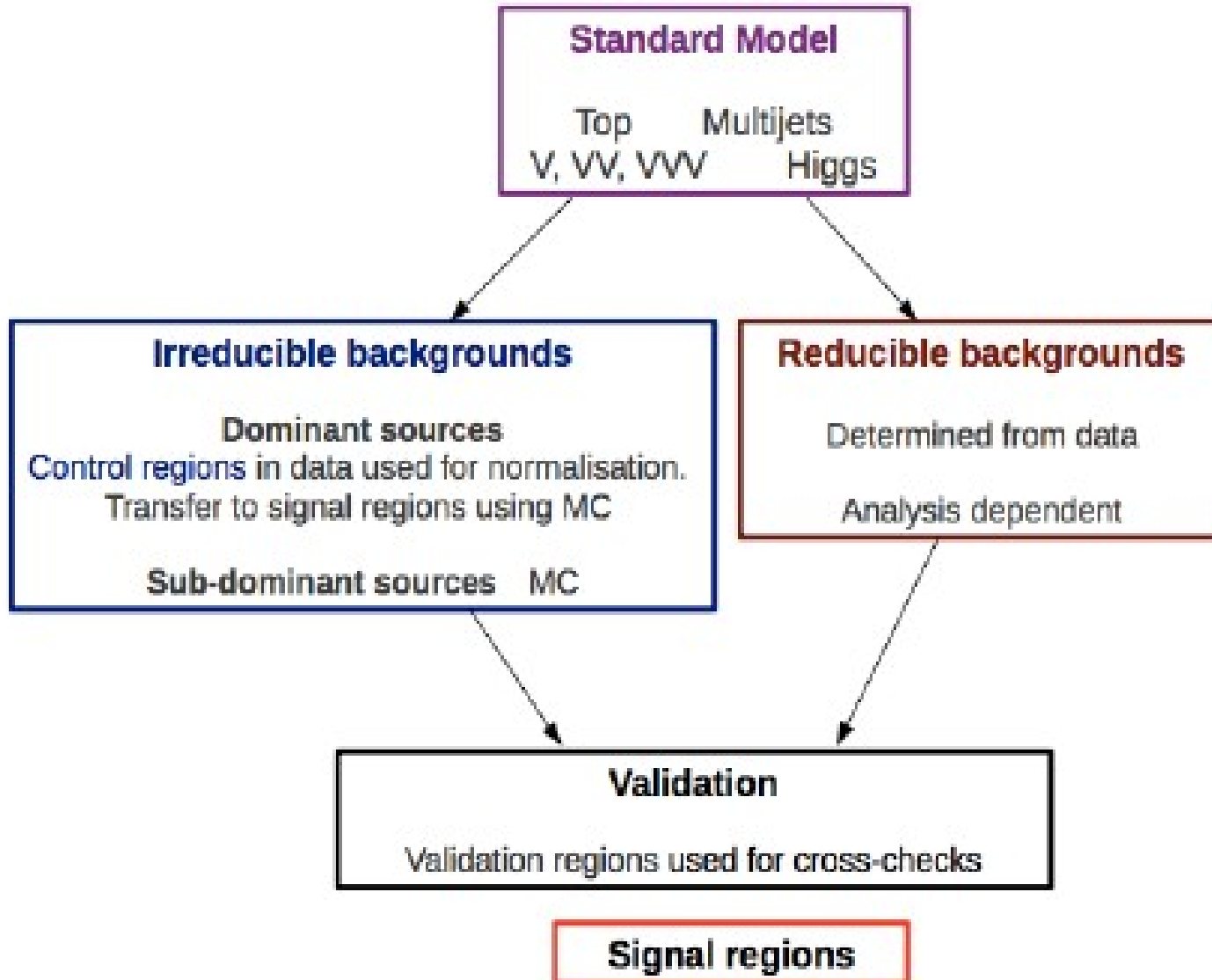
mass ↑

Lines are model dependent, assume

- Production cross-section for initial state
- Branching fraction for decay

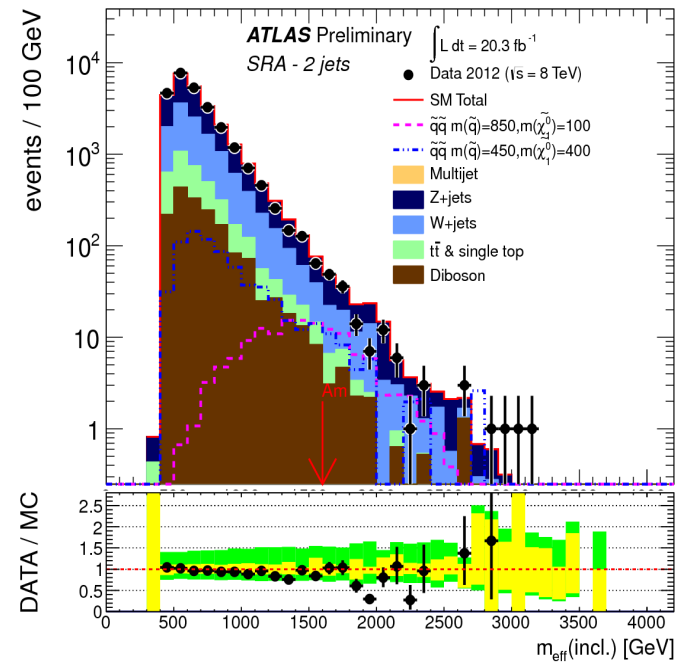
Color of plot is important:
It gives excluded cross section
In model-independent way,
Can be used to exclude different
Model with same topology

Flow of background evaluation

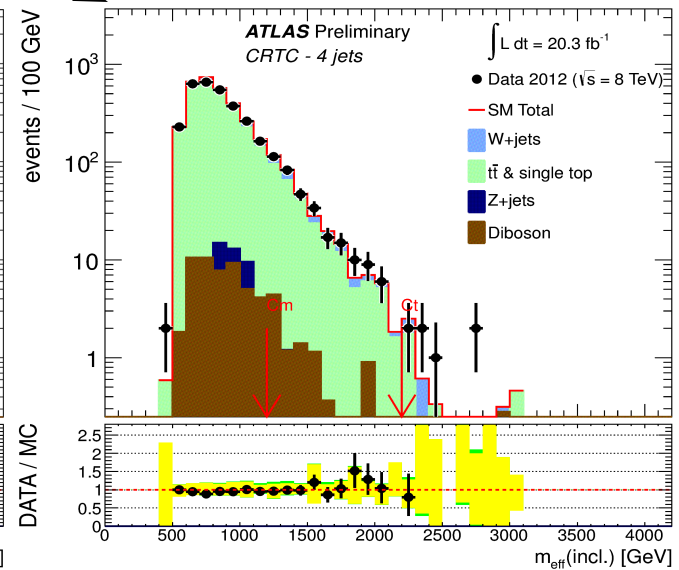
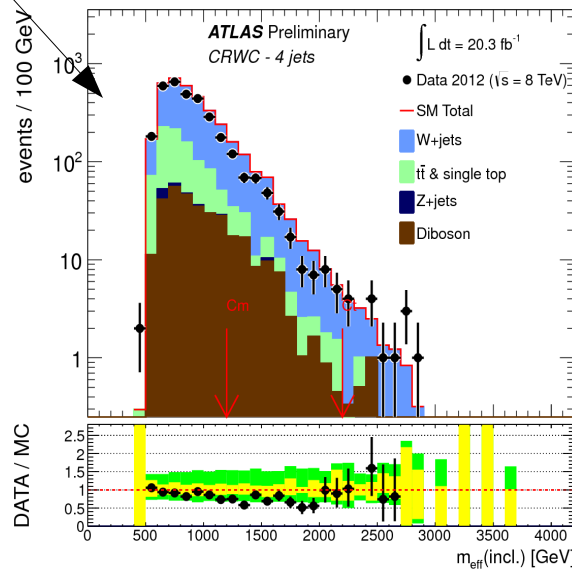
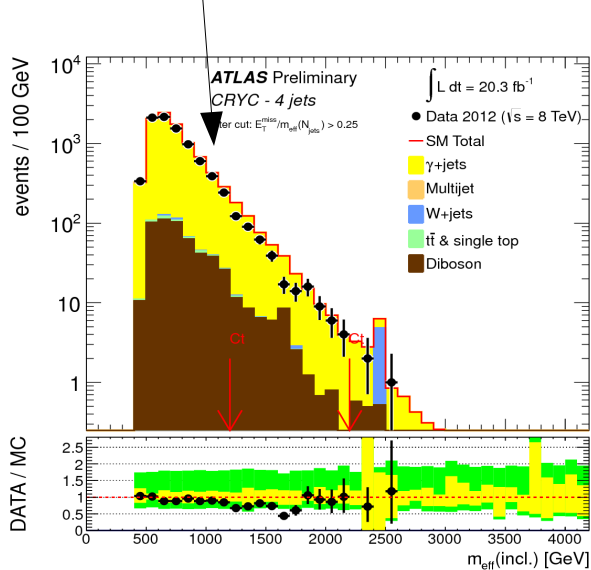


Background evaluation

| CR | SM Process | CR Process | CR Selection |
|-----|--------------------------------|--------------------------------|-----------------------|
| CRY | $Z \rightarrow \nu\nu + jets$ | $\gamma + jets$ | Isolated photon |
| CRW | $W \rightarrow \ell\nu + jets$ | $W \rightarrow \ell\nu + jets$ | trans. mass, b-veto |
| CRT | $t\bar{t}$ and single top | $t\bar{t}$ and single top | trans. mass, b-tag |
| CRQ | multijets | multijets | Reverse anti QCD cuts |



4 main backgrounds. For each signal region 4 control regions to constrain backgrounds



Dark Matter interpretation

- Need to assume model for DM interaction for connecting Collider data to DM experiments
- Use Effective Field (EFT) theory with contact interaction
- Ignore the nature of the mediator, write interaction as set of generic operators

Valid if the scale of interaction is less than the mediator mass M

| Name | Initial state | Type | Operator |
|------|---------------|--------------|---|
| D1 | qq | scalar | $\frac{m_q}{M_*^3} \bar{\chi} \chi \bar{q} q$ |
| D5 | qq | vector | $\frac{1}{M_*^2} \bar{\chi} \gamma^\mu \chi \bar{q} \gamma_\mu q$ |
| D8 | qq | axial-vector | $\frac{1}{M_*^2} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q} \gamma_\mu \gamma^5 q$ |
| D9 | qq | tensor | $\frac{1}{M_*^2} \bar{\chi} \sigma^{\mu\nu} \chi \bar{q} \sigma_{\mu\nu} q$ |
| D11 | gg | scalar | $\frac{1}{4M_*^3} \bar{\chi} \chi \alpha_s (G_{\mu\nu}^a)^2$ |