

# Split SUSY at LHC and a 100 TeV collider

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With Hugues Beauchesne and  
Kevin Earl

1503.03099

GGI - 2015

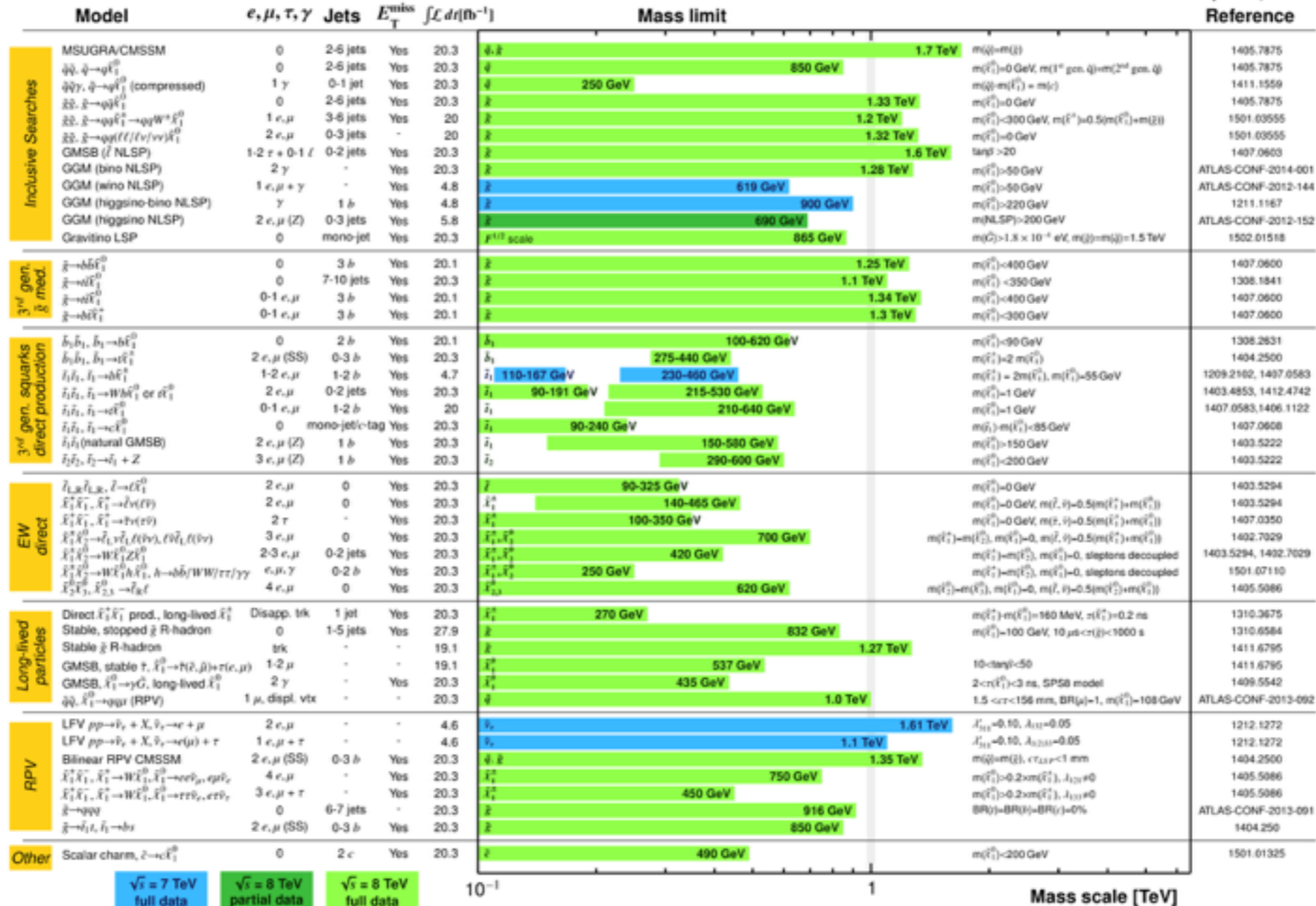
# Status of Supersymmetry

## ATLAS SUSY Searches\* - 95% CL Lower Limits

Status: Feb 2015

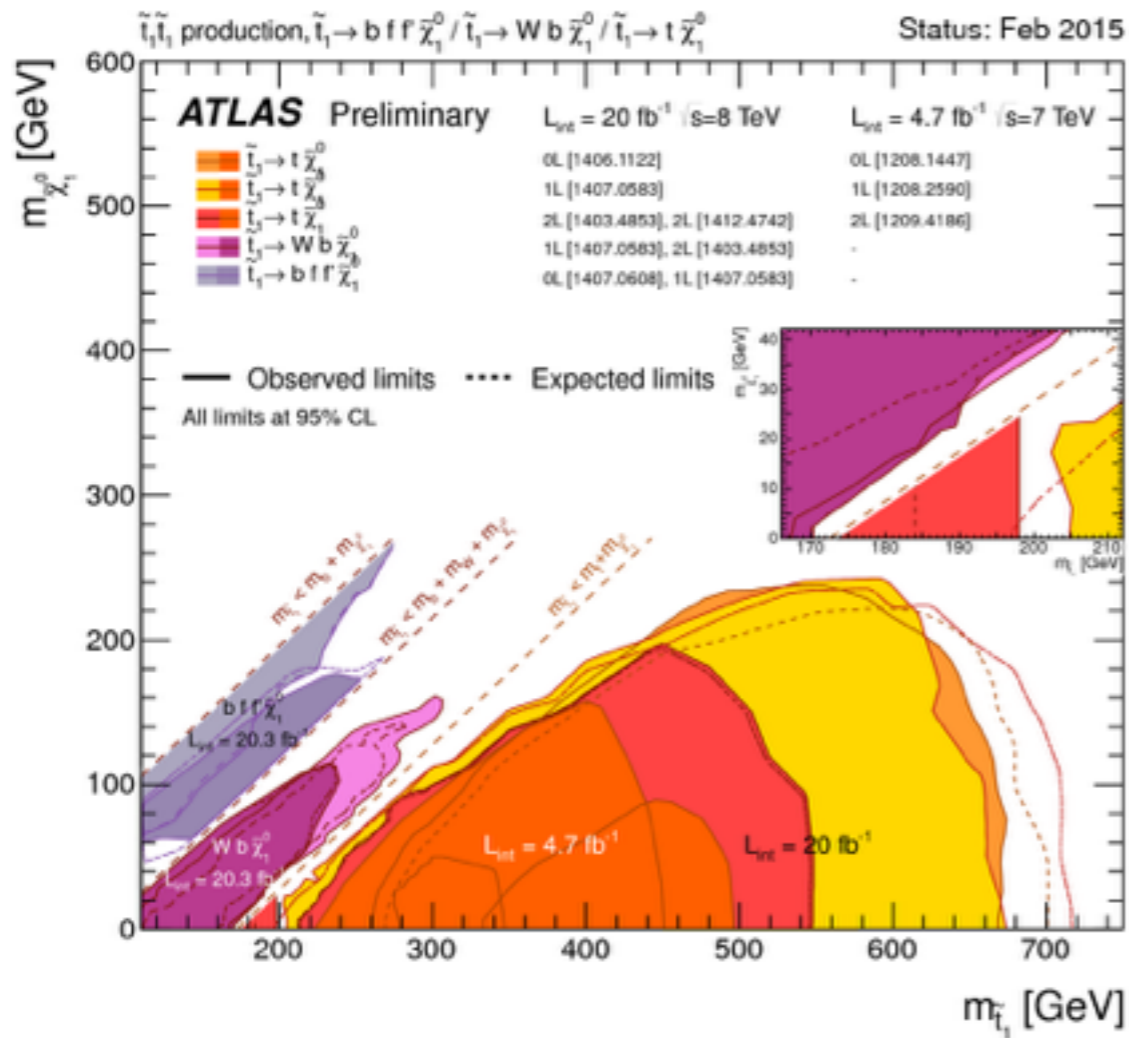
ATLAS Preliminary

$\sqrt{s} = 7, 8 \text{ TeV}$



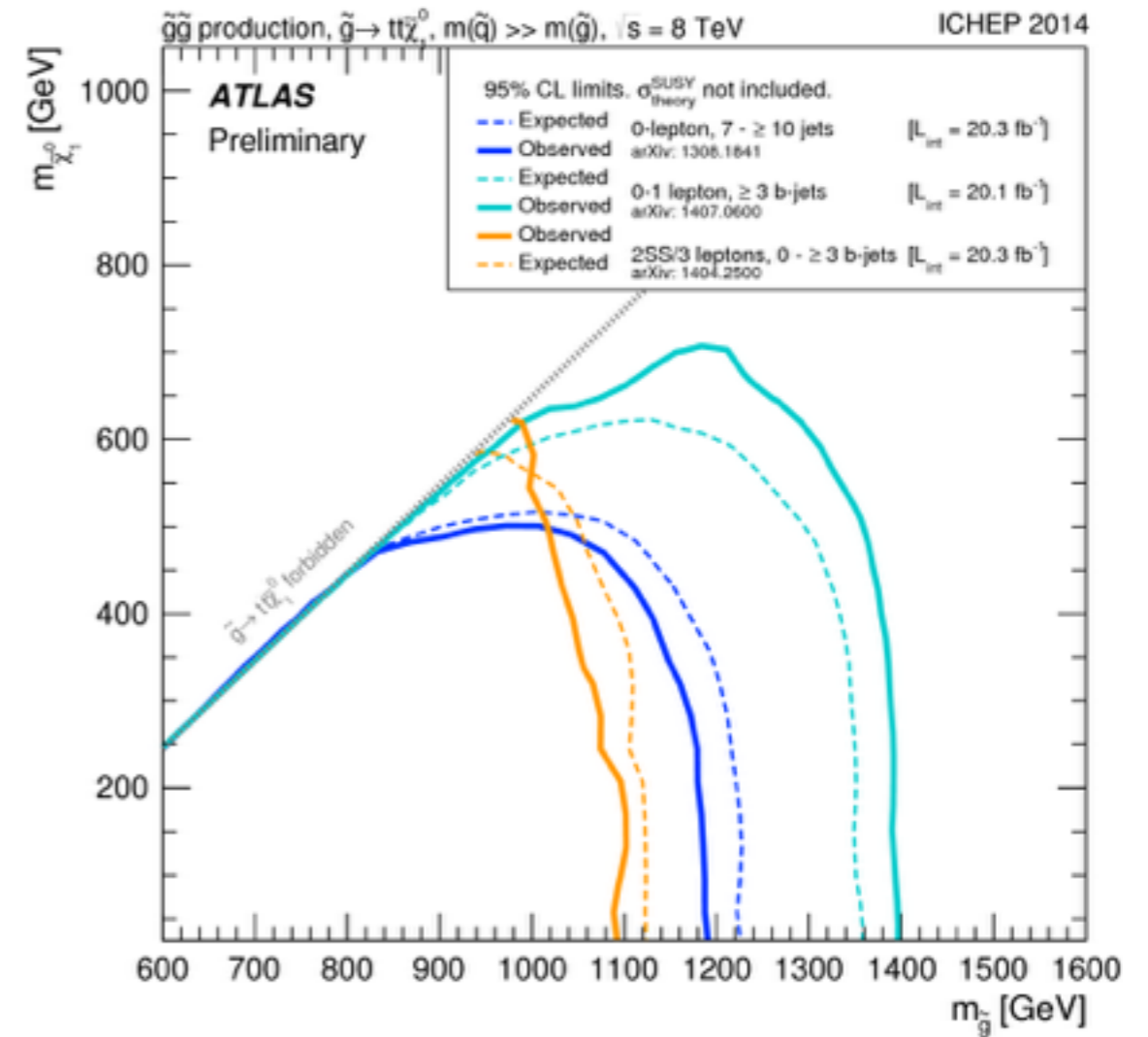
\*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus  $1\sigma$  theoretical signal cross section uncertainty.

# stop searches



$$m_{\tilde{t}} \gtrsim 700 \text{ GeV}$$

# gluino searches



$$m_{\tilde{g}} \gtrsim 1.4 \text{ TeV}$$

# What does it mean for naturalness?

‘natural SUSY’ (stop, gluino, higgsino)

Papucci, Ruderman, Weiler  
'11

stop

$$\delta m_h^2 = -\frac{3}{8\pi^2} y_t^2 \left( m_{Q_3}^2 + m_{u_3}^2 + |A_t|^2 \right) \log \left( \frac{\Lambda}{\text{TeV}} \right)$$

$$\sqrt{m_{\tilde{t}_1}^2 + m_{\tilde{t}_2}^2} \lesssim 600 \text{ GeV} \left( \frac{\Delta^{-1}}{20\%} \right)^{-1/2}$$

gluino

$$\delta m_h^2 = -\frac{2}{\pi^2} y_t^2 \left( \frac{\alpha_s}{\pi} \right) M_3^2 \log^2 \left( \frac{\Lambda}{\text{TeV}} \right)$$

$$M_3 \lesssim 900 \text{ GeV} \sin \beta \left( \frac{\Delta^{-1}}{20\%} \right)^{-1/2}$$

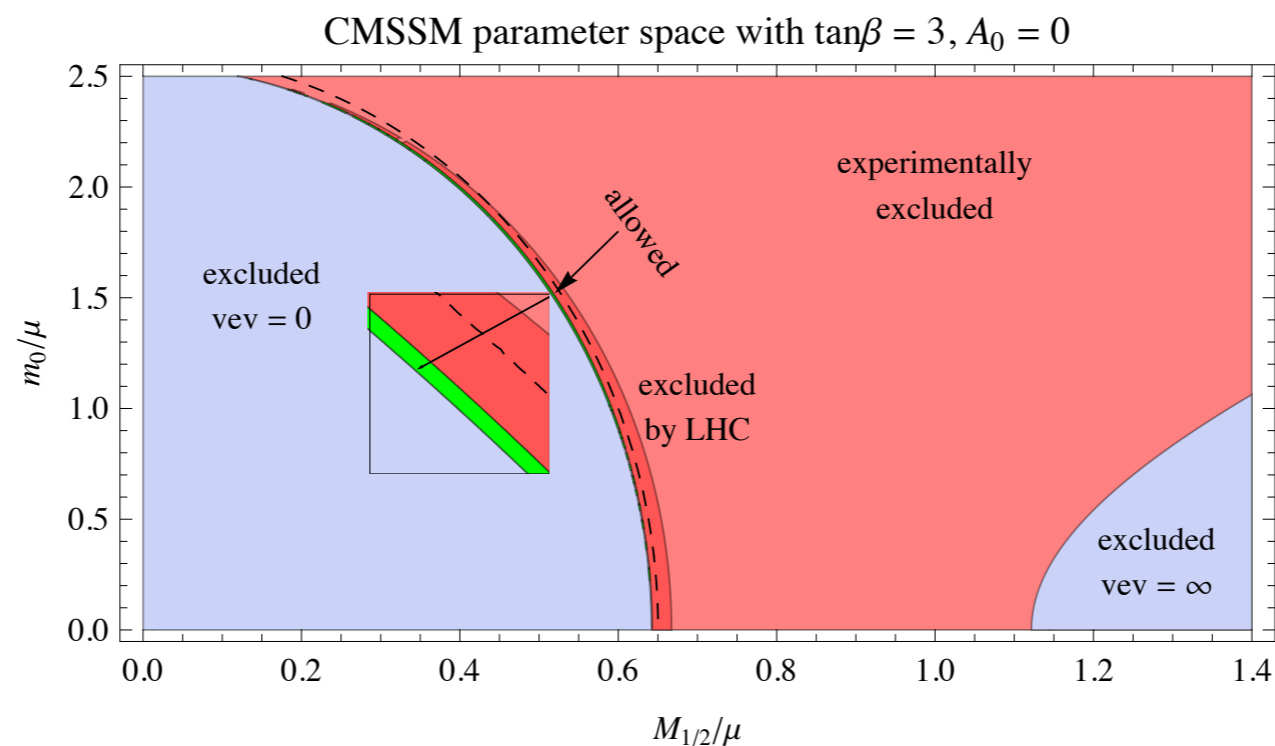
125 GeV Higgs : prefers heavy stops

$$\delta m_h^2 = \frac{3G_F}{\sqrt{2}\pi^2} m_t^4 \log \left( \frac{m_{\tilde{t}}^2}{m_t^2} \right)$$

~ 1-10 TeV stop depending on A-term

More complete models generally yield %-level fine-tuning

CMSSM



Strumia '11

NMSSM, split generation, low-scale mediation, Dirac gaugino,  
and R-parity breaking are **generically tuned**.

Arvanitaki, Baryakhtar,  
Huang, Tiburg, Villadoro '13

Better **model building** might save the day

**Scherk-Schwartz** SUSY breaking

Dimopoulos, March-Russell  
'14

LeComte, Martin '11

**Compressed** spectrum

Dimopoulos, March-Russell,  
Scoville '14

**Stealth** supersymmetry

Fan, Reece, Ruderman '11

**Twin** Higgs

Chacko, Goh, Harnik '05

Craig, Howe '13

# Dirac gauginos

In the MSSM gauginos are Majorana

$$M\lambda\lambda$$

$$F_X \theta^2 \leftarrow \int d^2\theta \textcircled{X} W_\alpha W^\alpha$$

Can be Dirac if new superfields are added

$$W_\alpha^1, W_\alpha^2, W_\alpha^3 \quad S, T, G$$

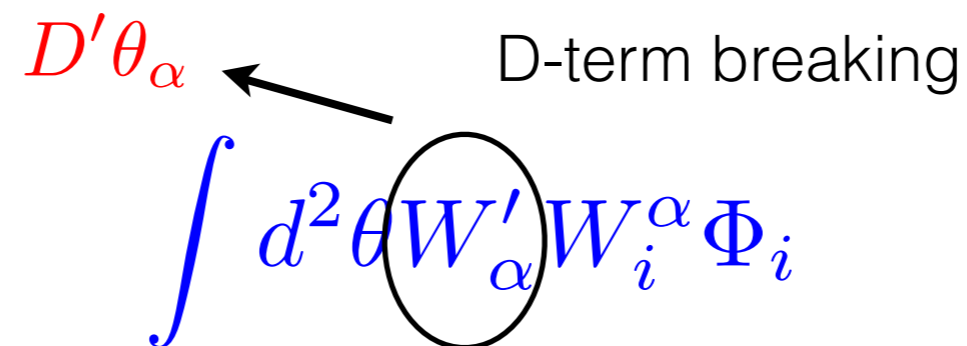
$$M_D \lambda \Psi$$

N=2 supersymmetry  
extra-dimension

# Supersoft SUSY breaking

Fox, Nelson, Weiner '02

$D'\theta_\alpha$  ← D-term breaking

$$\int d^2\theta W'_\alpha W_i^\alpha \Phi_i$$


Dirac gauginos do not feed into scalar masses through **renormalization**

$$m^2 = \frac{C_i(r) \alpha_i m_i^2}{\pi} \log \left( \frac{\delta^2}{m_i^2} \right)$$

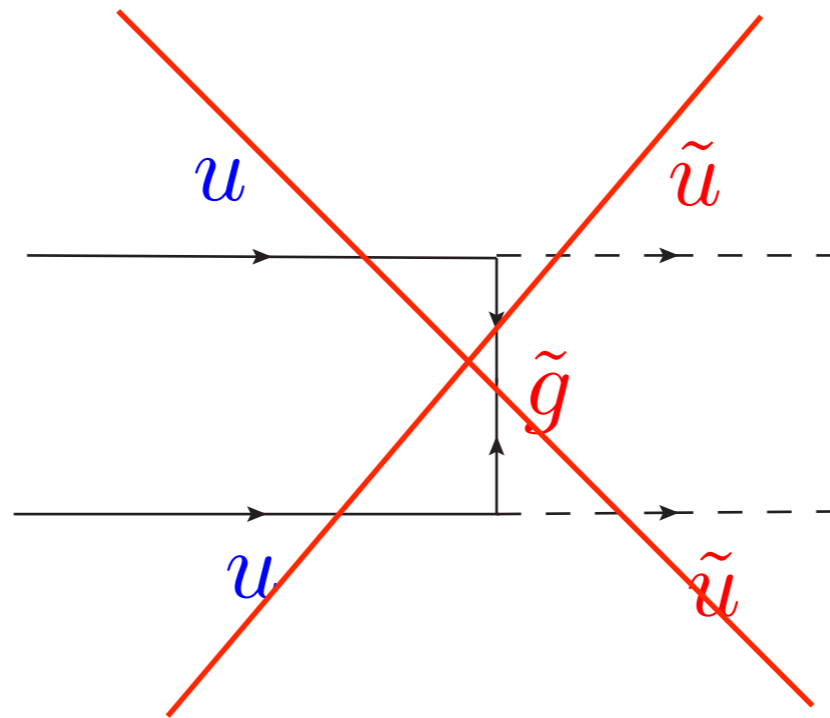


They can be naturally **heavier than scalars**

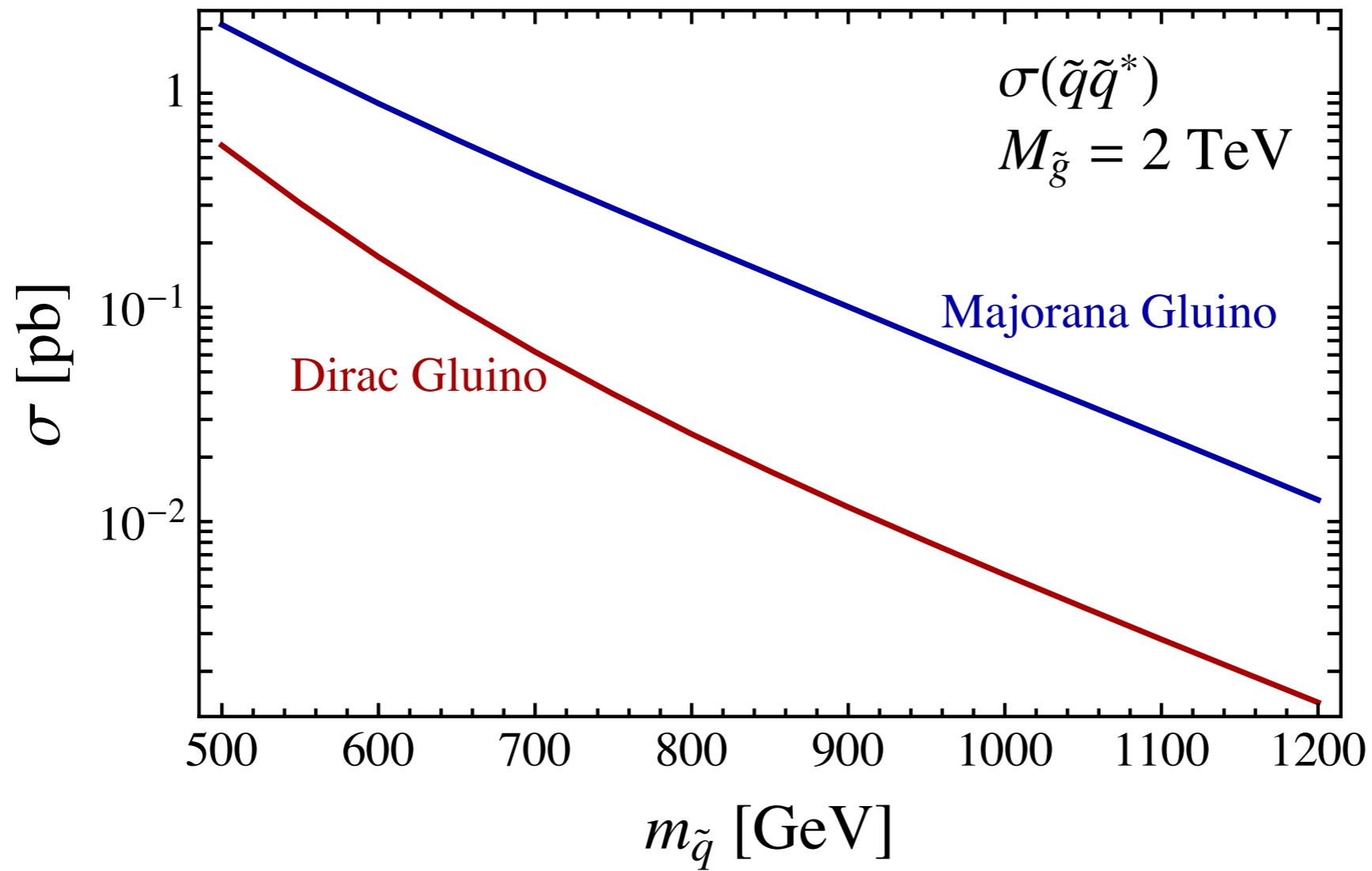
LHC will have a harder time seeing the gluino...

M. Heikinheimo, M. Kellerstein, V. Sanz '12  
Kribs, Martin '12

...and squarks



# Squark production



Frugieuele, T.G., Kumar, Ponton

# R-symmetry

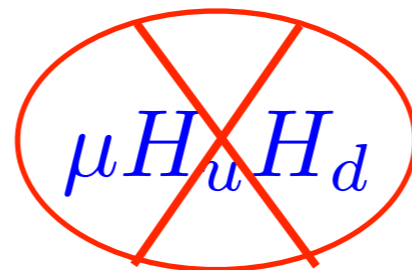
With Dirac gaugino: possible to impose an  
U(1) R-symmetry

$$M_D \lambda \Psi$$

Kribs, Poppitz, Weiner '02

- Bounds from FCNC are weaker: off diagonal  $m_{ij}$

$$R[Q, U^c, D^c, L, E^c] = 1 \quad R[H_u, H_d] = 0$$



~~$\mu H_u H_d$~~

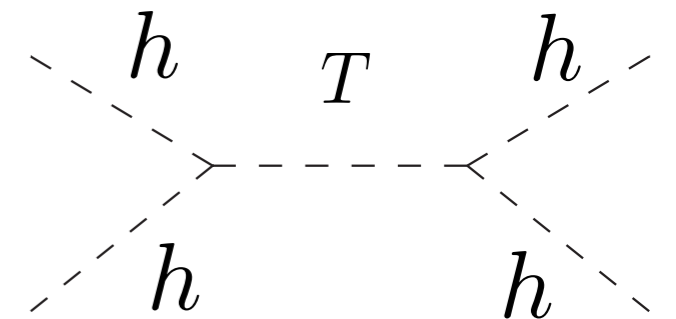
# Higgs mass

Tree-level:

Reduced quartic, usual of Dirac gauginos

$$\int d^2\theta W'_\alpha W_i^\alpha \Phi_i \quad \longrightarrow \quad D_2 = M_2 T^a + H_u^\dagger \sigma^a H_u + \dots$$

When the scalar  $T$  is integrated out:



$$\lambda \rightarrow 0$$

Higgs quartic

If the mass of  $T$  is set by  $M_2$  and  $\lambda_T = 0$

No help (at tree-level) from

$$\lambda_T H_u T(R_d) + \lambda_S H_u S(R_d)$$

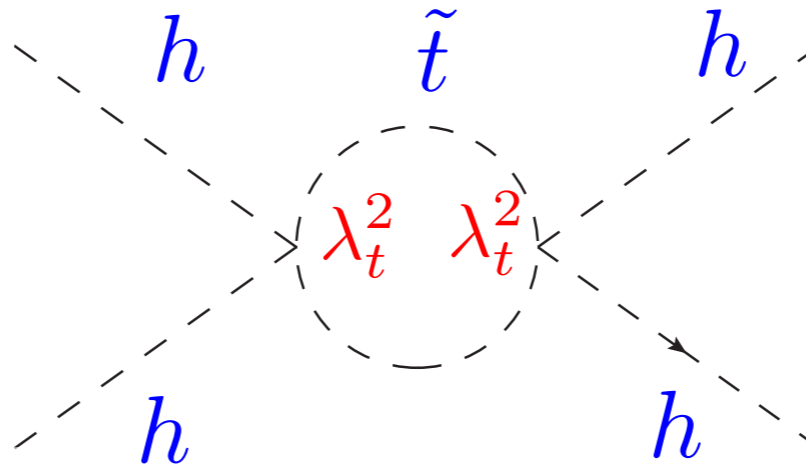
don't get a vev (In the limit of exact R-symmetry)

But do help in models *without an R-symmetry*

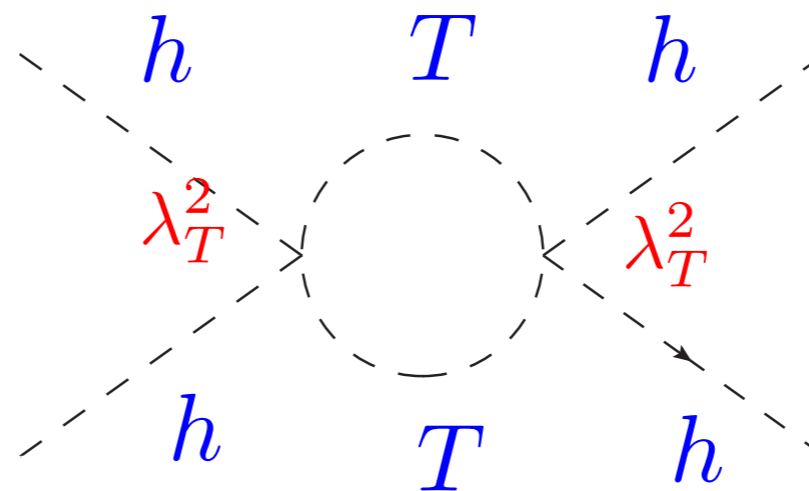
Benakli, Goodsell, Staub 1211.0552

# Loop-level

Usual stop correction (but A-terms are 0)



Similar loop from the triplet



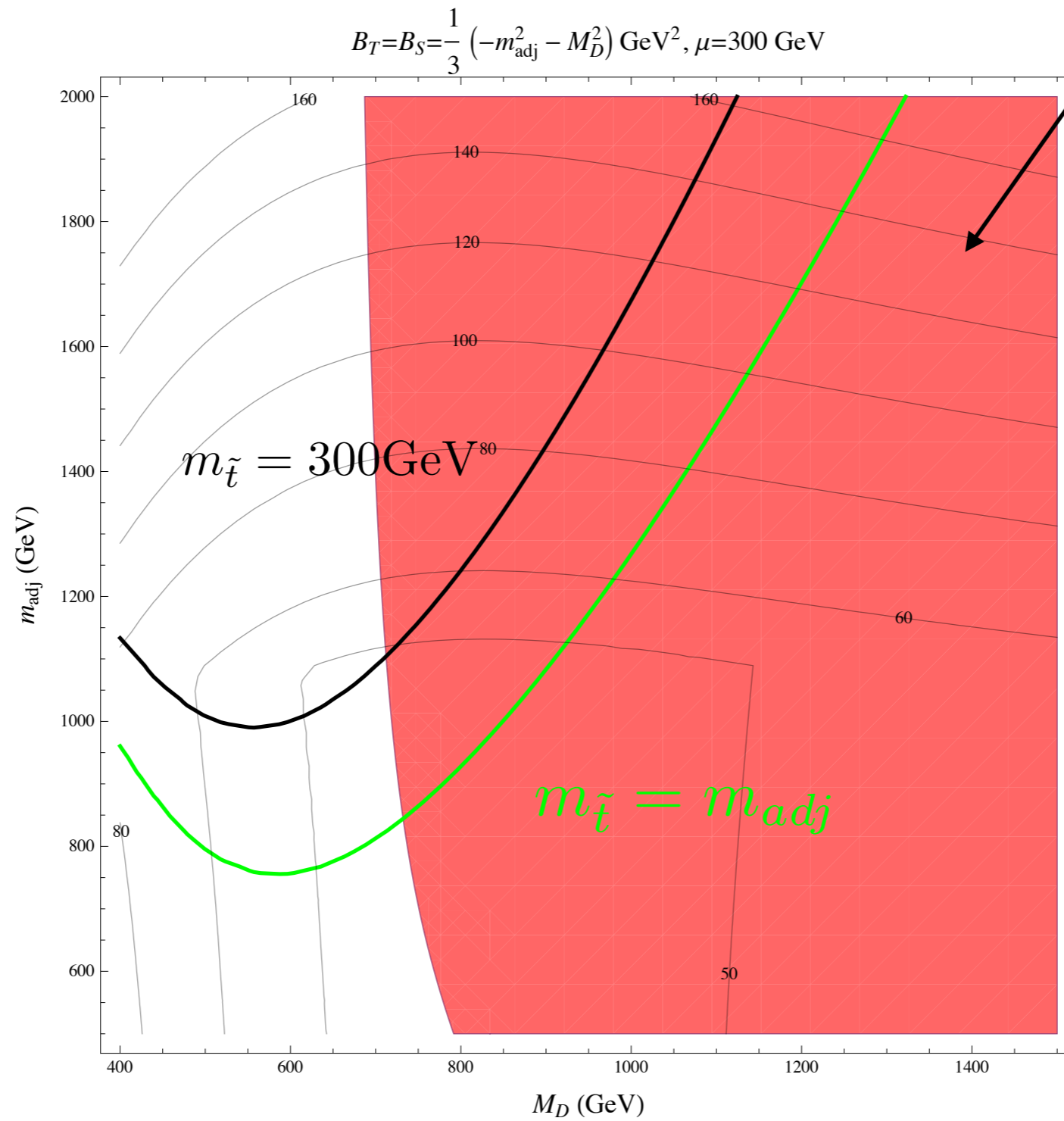
$$V_{\text{CW}} \sim \frac{1}{16\pi^2} \left( 5\lambda_T^4 \log \frac{m_T^2}{M_2^2} + 3\lambda_t^4 \log \frac{m_{\tilde{t}}^2}{m_t^2} \right)$$



Very sensitive to  $\lambda_T$

....but so are electroweak precision measurements

allowed by EWPT



Bertuzzo, Frugiuele, T.G., Ponton



# Split Supersymmetry

Naturalness **might not be a good guide**

SUSY might still be relevant

- Dark matter
- Gauge coupling unification
- 'UV' reasons

Gauginos and scalars might **not be at the same mass scale**

natural in for example anomaly  
mediation

## Prediction for the Higgs mass

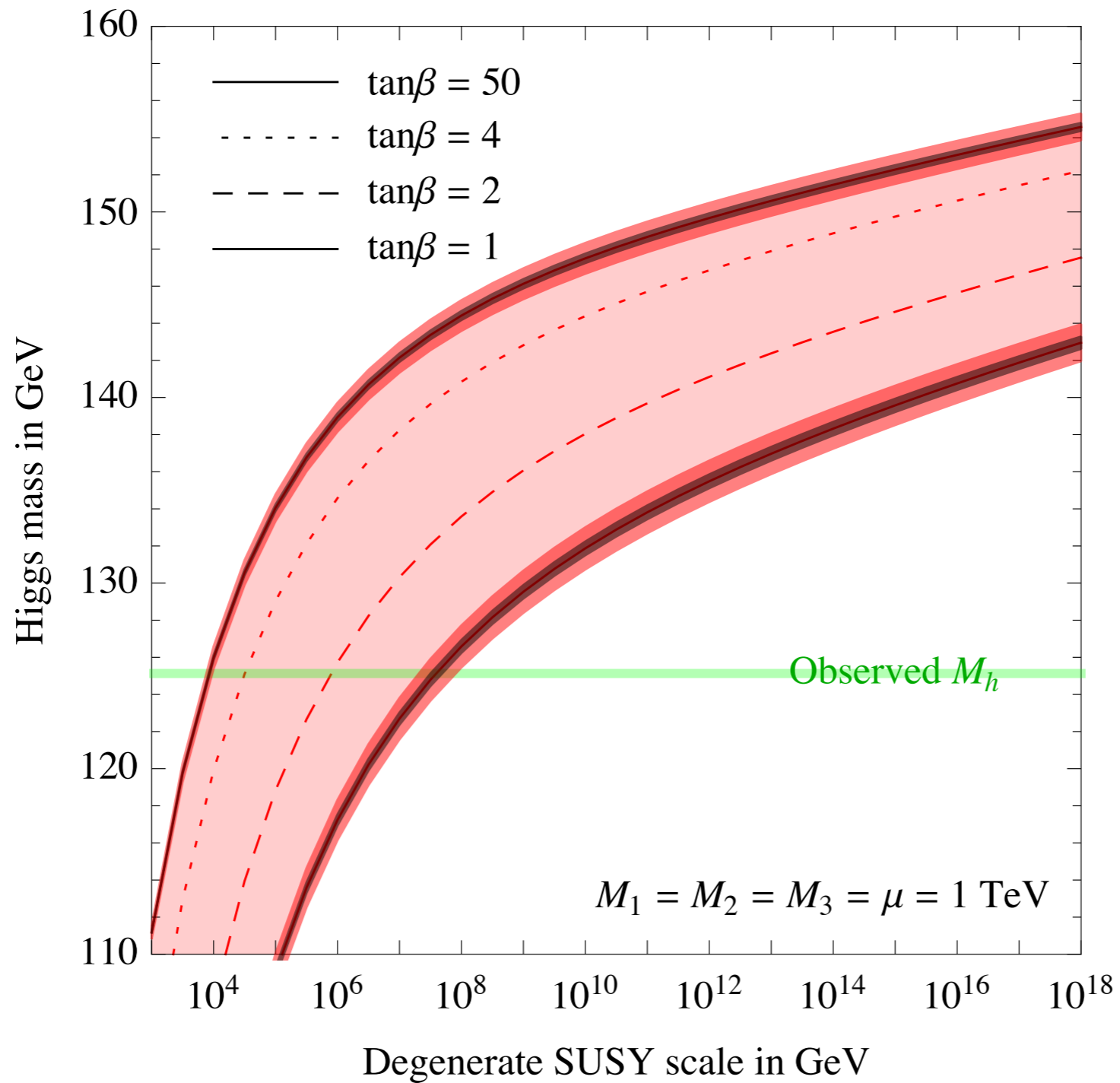
The Higgs quartic coupling is predicted at a high scale:

$$\lambda(m_{\text{scalar}}) = \frac{1}{4}(g^2 + g'^2) \cos^2 \beta \quad \text{tree-level}$$



Bagnaschi,  
Giudice, Slavich, Strumia

# Split-SUSY



Bagnaschi,  
Giudice, Slavich, Strumia  
'14

## Mini-split in **anomaly mediation**

scalar masses are generated by gravity mediation

$$\int d^4\theta \frac{X^\dagger X Q^\dagger Q}{M_{pl}^2} \quad m^2 = m_{3/2}^2$$

but the gaugino masses are generated by AMSB

$$M_i = \frac{b_i}{16\pi^2} g_i^2 m_{3/2}$$

~~$$\int d^2\theta \frac{X W_\alpha W^\alpha}{M_{pl}}$$~~

$\mu$ -term generated through Giudice-Masiero

conformal compensator

$$\int d^4\theta \phi^\dagger \phi H_u H_d$$

$$\mu \sim \frac{B_\mu}{\mu} \sim m_{3/2}$$

Heavy Higgsino

Similar spectrum could also arise in  
**gauge mediation**

Arvanitaki, Craig, Dimopoulos, Villadoro  
'12  
Buican, Meade, Seiberg, Shih  
'09

$$W = M_R (\phi_1 \bar{\phi}_1 + \phi_2 \bar{\phi}_2) + X \phi_1 \phi_2$$

$$M_i \sim \frac{g_i^2}{16\pi^2} \frac{M}{M_R} \frac{F^3}{M_R^3} = \frac{g_i^2}{16\pi^2} \Lambda$$

We take the scalars at  $\sim \Lambda$

Gaugino spectrum  
(deflected AMSB)

Higgsino  
threshold

$$M_{\tilde{B}} = M_1 \left[ 1 + \frac{C_\mu}{11} + \dots \right]$$

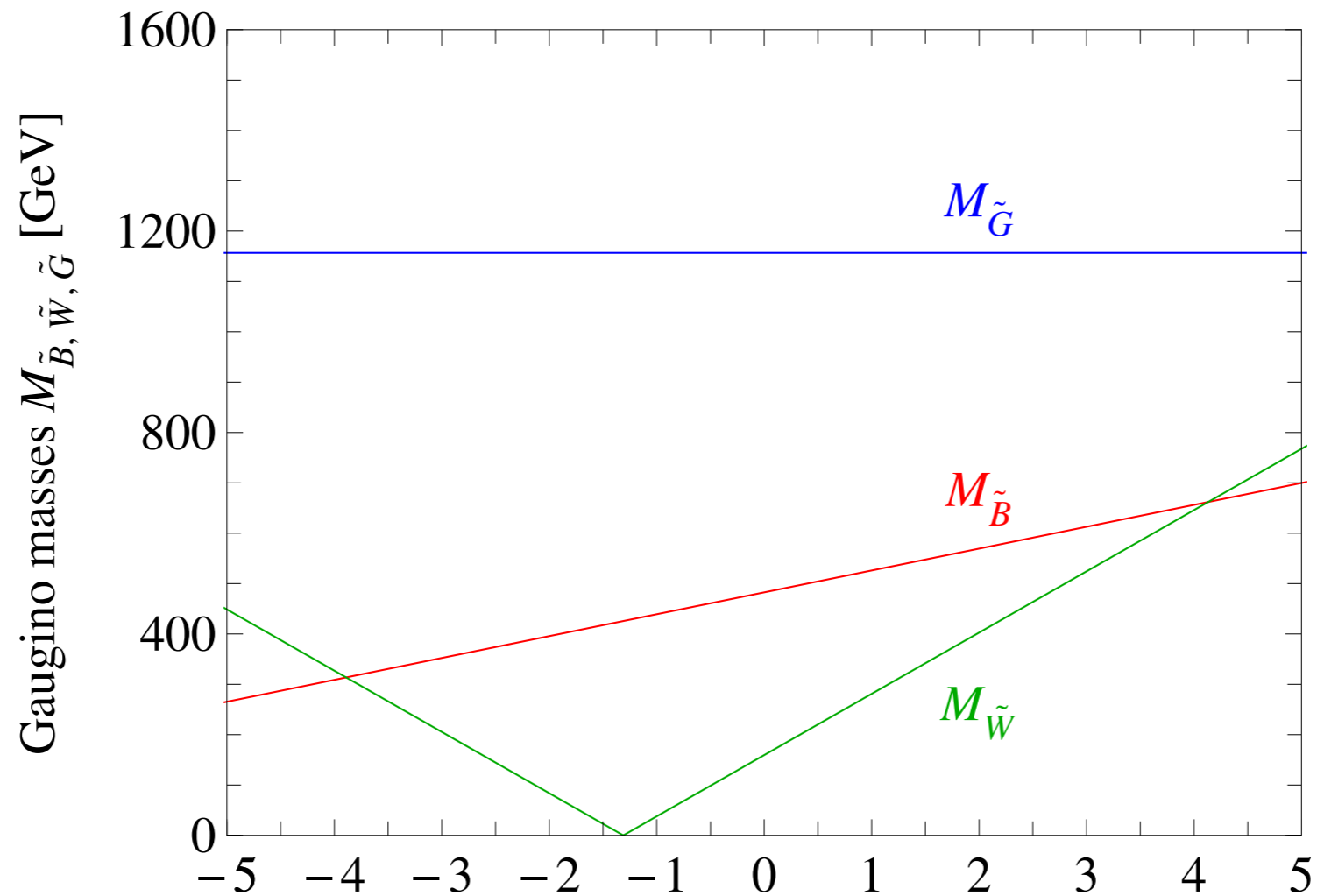
$$M_{\tilde{W}} = M_2 [1 + C_\mu + \dots]$$

$$M_{\tilde{g}} = M_3 [1 + \dots]$$

$$M_i = \frac{\beta_i}{g_i} m_{3/2} \quad C_\mu = \frac{\mu}{m_{3/2}} \frac{m_A^2 \sin^2 \beta}{m_A^2 - \mu^2} \ln \frac{m_A^2}{\mu^2}$$

# AMSB Gaugino spectrum

Beauchesne, Earl, T.G. '15



$c_\mu$

Parametrize deflection  
from AMSB spectrum



Similar expressions for  
gauge mediation

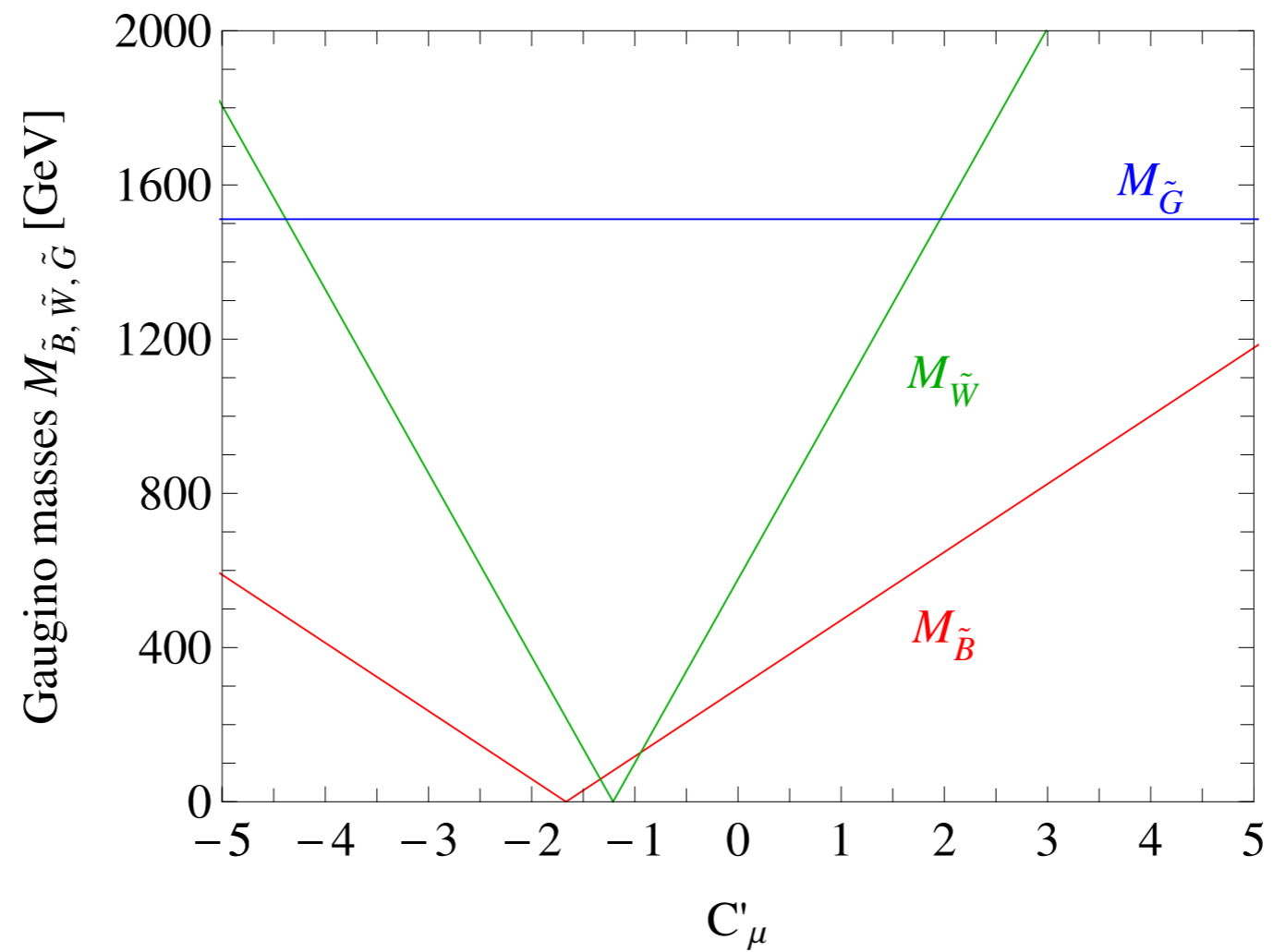
$$M_{\tilde{B}} = M_1 \left[ 1 + \frac{3C'_\mu}{5} + \dots \right]$$

$$M_{\tilde{W}} = M_2 \left[ 1 + C'_\mu + \dots \right]$$

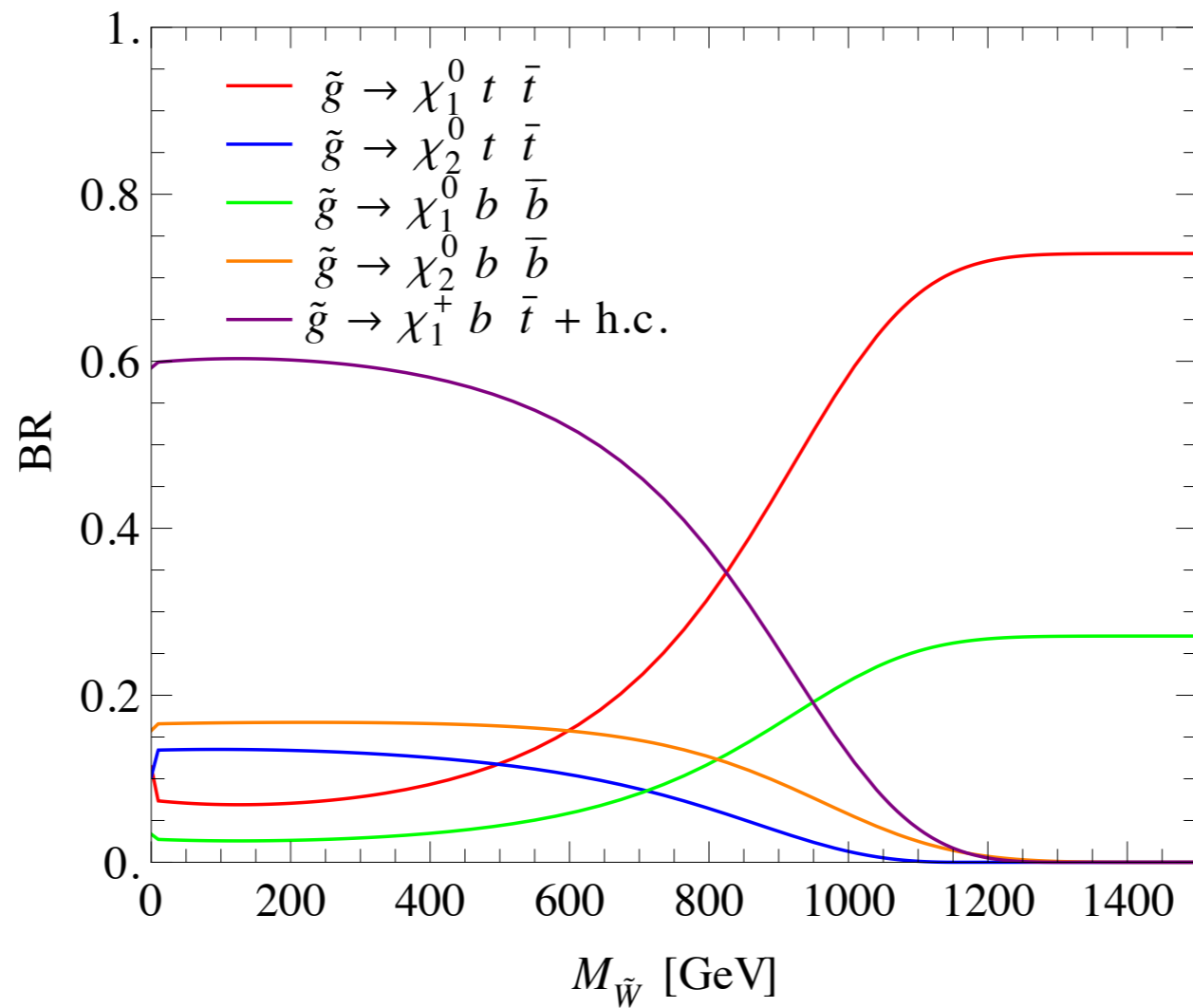
$$M_{\tilde{g}} = M_3 \left[ 1 + \dots \right]$$

$$M_i = \frac{g_i^2}{16\pi^2} \Lambda \quad C'_\mu = \frac{\mu}{\Lambda} \frac{m_A^2 \sin^2 \beta}{m_A^2 - \mu^2} \ln \frac{m_A^2}{\mu^2}$$

# Similar expressions for gauge mediation



Assume that gluino decay to 3rd generation quarks



$$m_{\tilde{B}} = 0 \text{ GeV}$$

$$m_{\tilde{G}} = 1.5 \text{ TeV}$$

## Electroweakino decays

$$\tilde{W}^0 \rightarrow \tilde{B}h$$

bino LSP

$$\tilde{W}^+ \rightarrow W^+ \tilde{B}$$

$$\tilde{B} \rightarrow \tilde{W}^0 h$$

wino LSP

$$\tilde{W}^+ \rightarrow W^0 + \text{soft}$$

## Parameter space

parameters of the model:

$$m_{3/2} \quad \mu \quad \tan \beta \quad m_{\text{scalar}}$$

choose:  $m_{\text{scalar}} \sim m_{3/2}$

set  $\tan \beta$  to reproduce the Higgs mass

results in term of

$$C_{\mu} \text{ and } m_{3/2}$$

## Recasting LHC bounds

Gaugino spectrum and branching ratios are obtained as a function of  $m_{3/2}$  and  $C_\mu$ .

we simulate the signal using MadGraph-Pythia-Delphes and **recast LHC searches**

ATLAS multi-leptons+b-jets

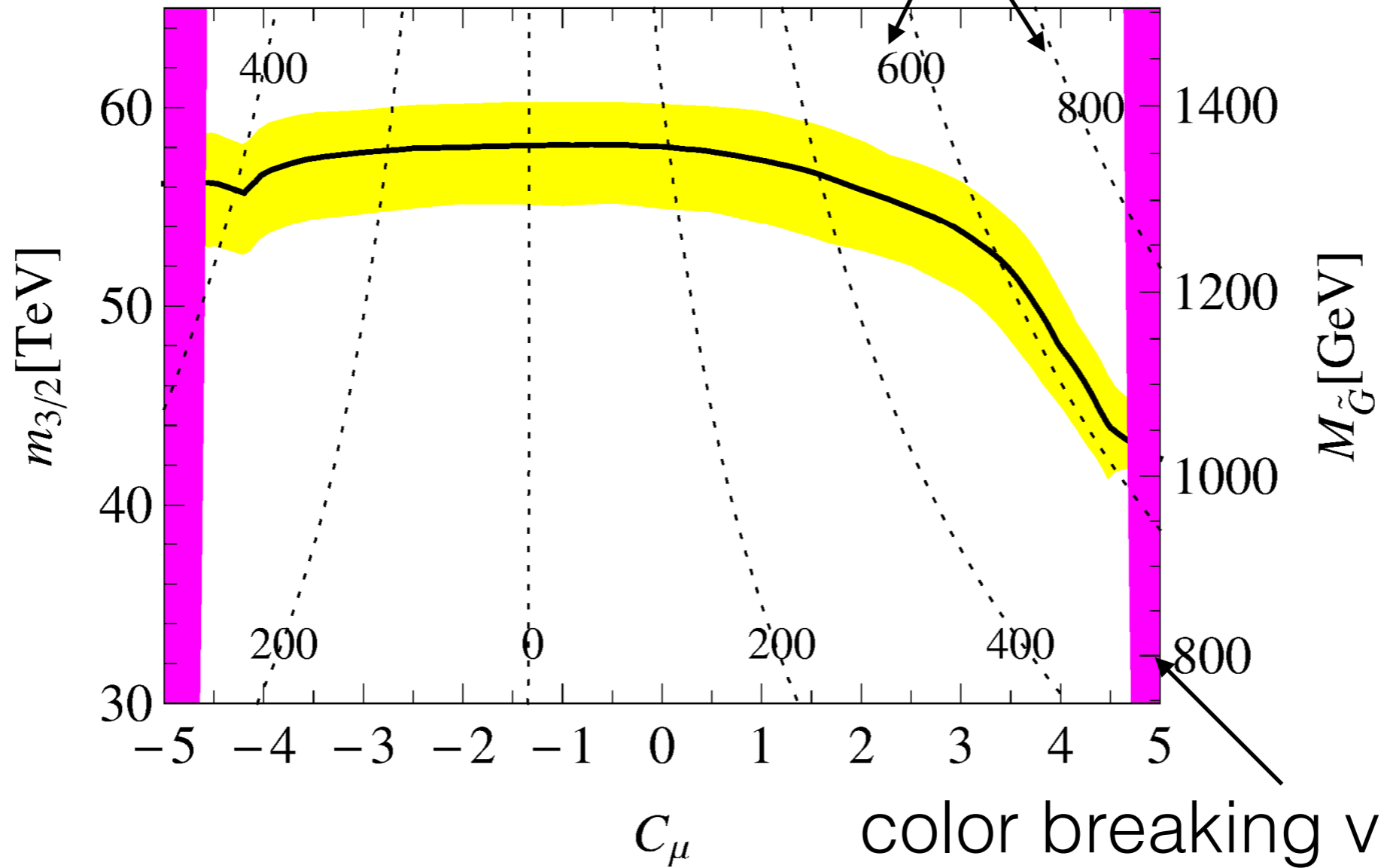
ATLAS 0-1 lepton+b-jets

CMS high jets multiplicity

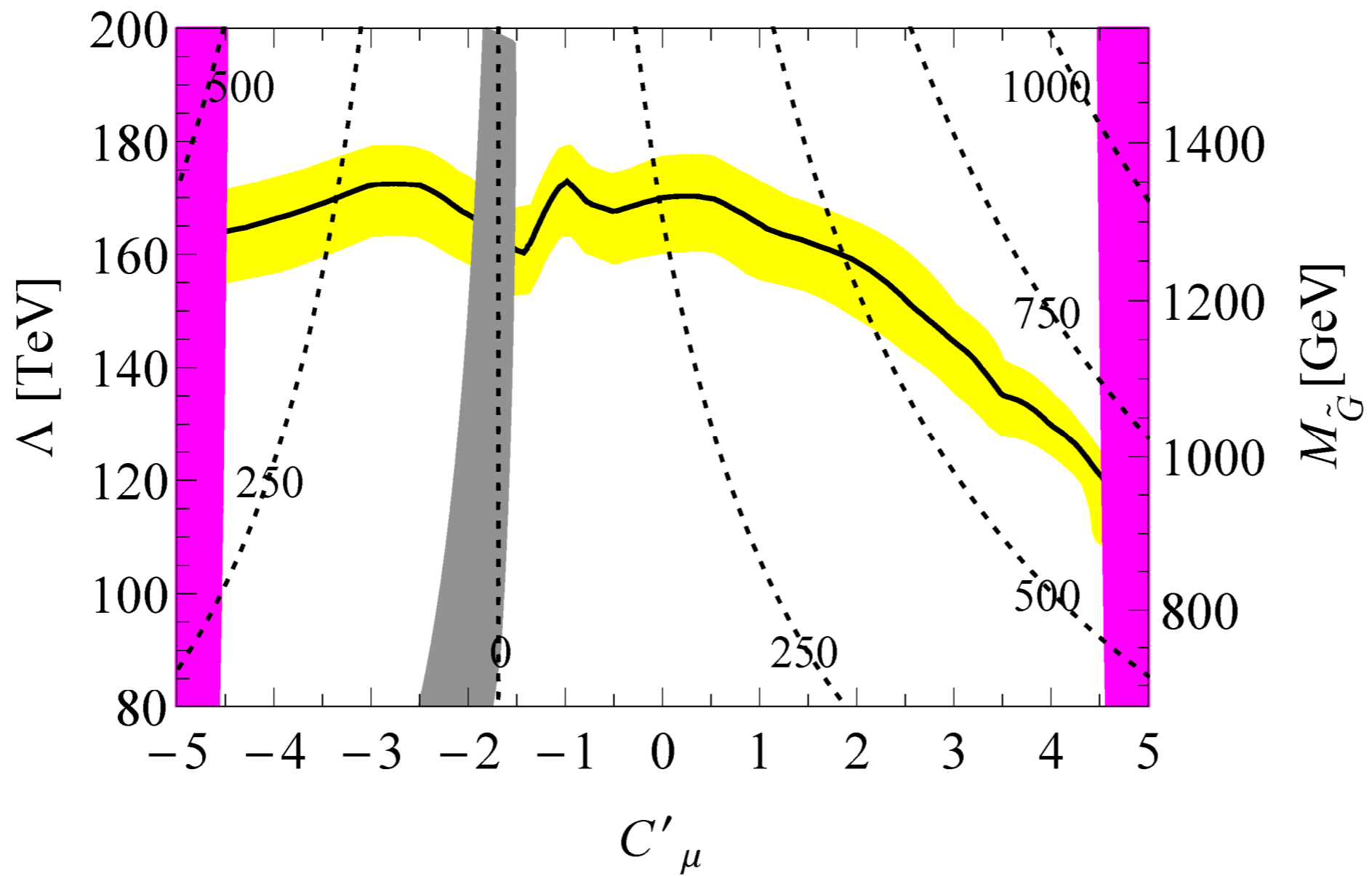
CMS 2 OS leptons+jets

AMSB

$M_{\tilde{W}}$  contours



# Gauge mediation





# LHC 14 prospects

looked at 2 sets of cuts

## same-sign dilepton

cohen et al. 1311.6480

- SSDL
- 2 b-jets or more
- 6 jets or more
- $H_T > 700\text{GeV}$
- $E_T^{\text{miss}} > 250\text{GeV}$

8 signal  
regions

## High missing energy

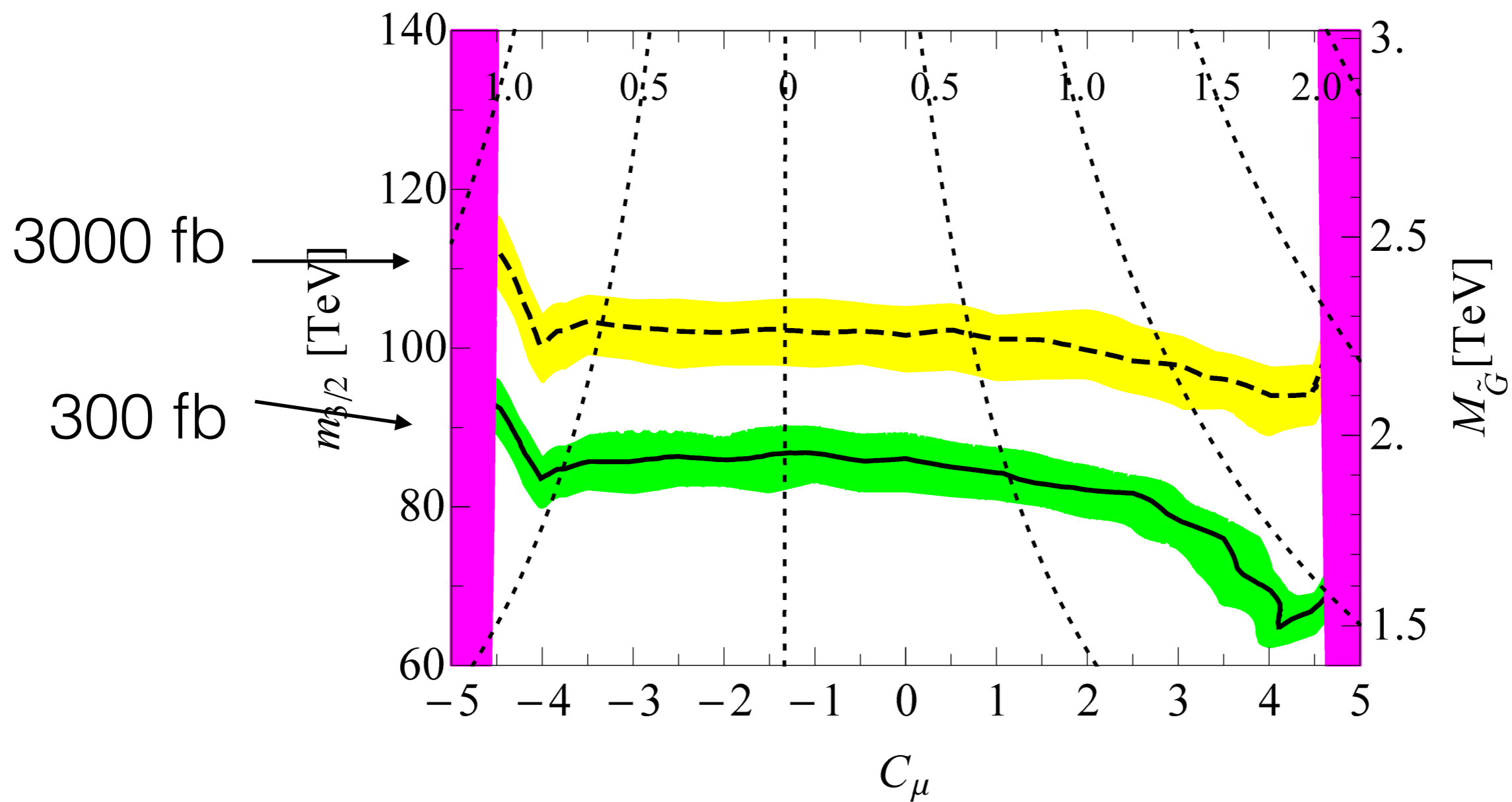
CMS-PAS-FTR-13-014

- 1 lepton
- 6 jets or more
- 1 b-jet
- $H_T > 500\text{GeV}$
- $E_T^{\text{miss}} + P_T^{\text{lep}} > 450\text{GeV}$

4 signal  
regions

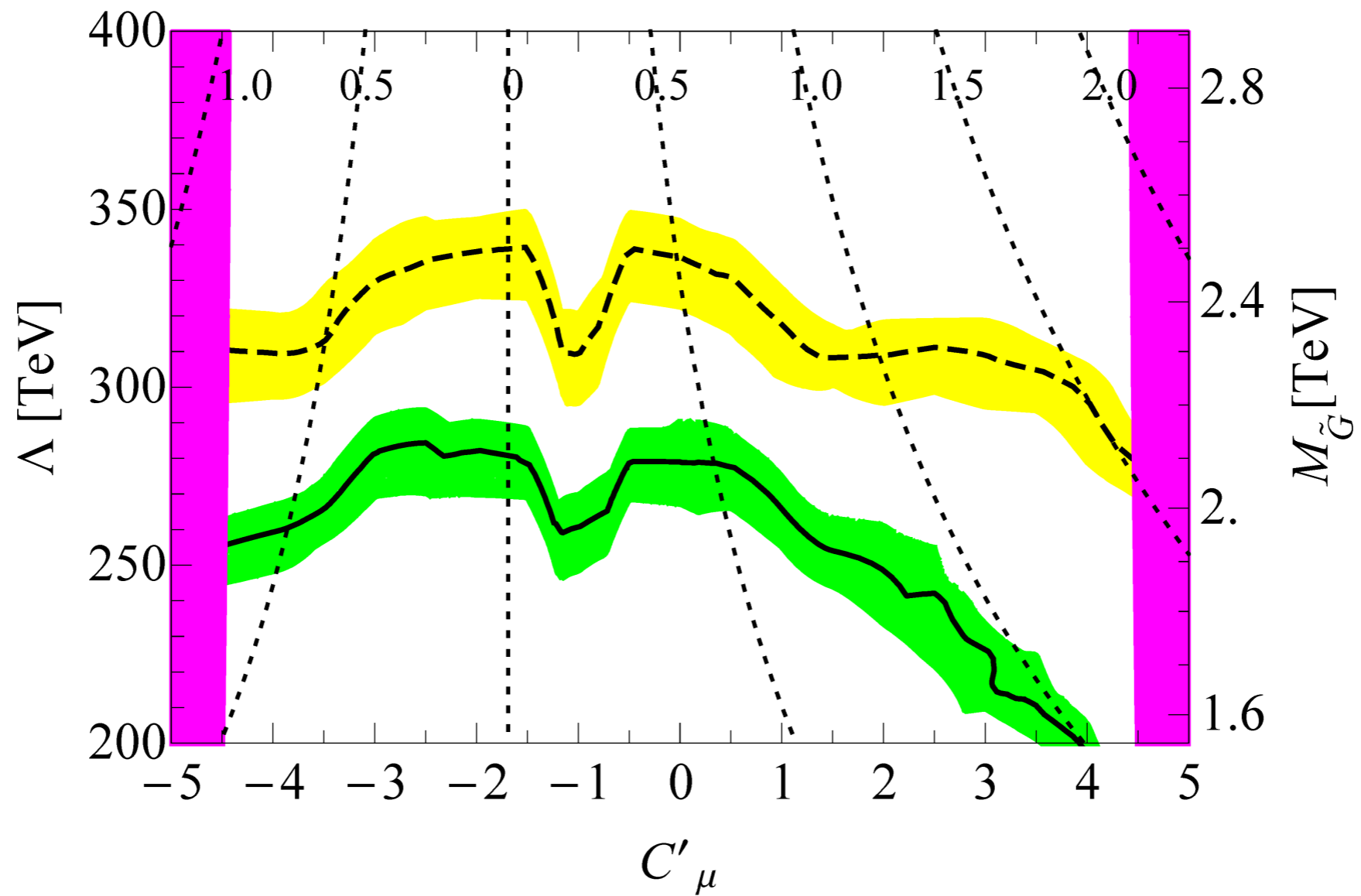
# AMSB (discovery) at LHC 14

$M_{\tilde{W}}$  contours



# GMSB (discovery) at LHC 14

$M_{\tilde{B}}$  contours



## Prospect for a 100 TeV collider

### same-sign dilepton

cohen et al. 1311.6480

- SSDL
- 3 b-jets or more
- 7 jets or more
- $H_T > 3000\text{GeV}$
- $E_T^{\text{miss}} > 800\text{GeV}$

8 signal  
regions

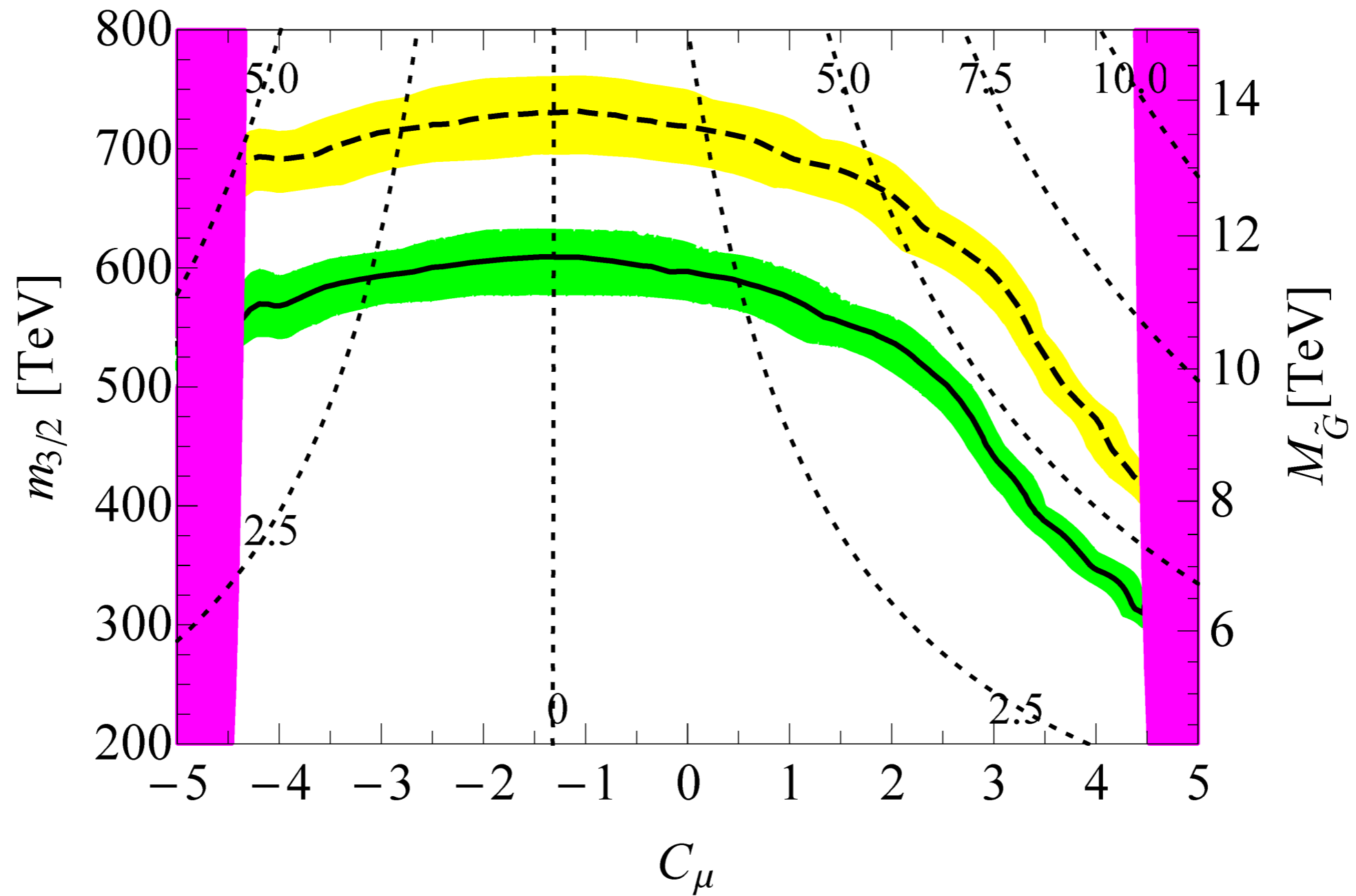
### High missing energy

adapted from Jung, Wells '13

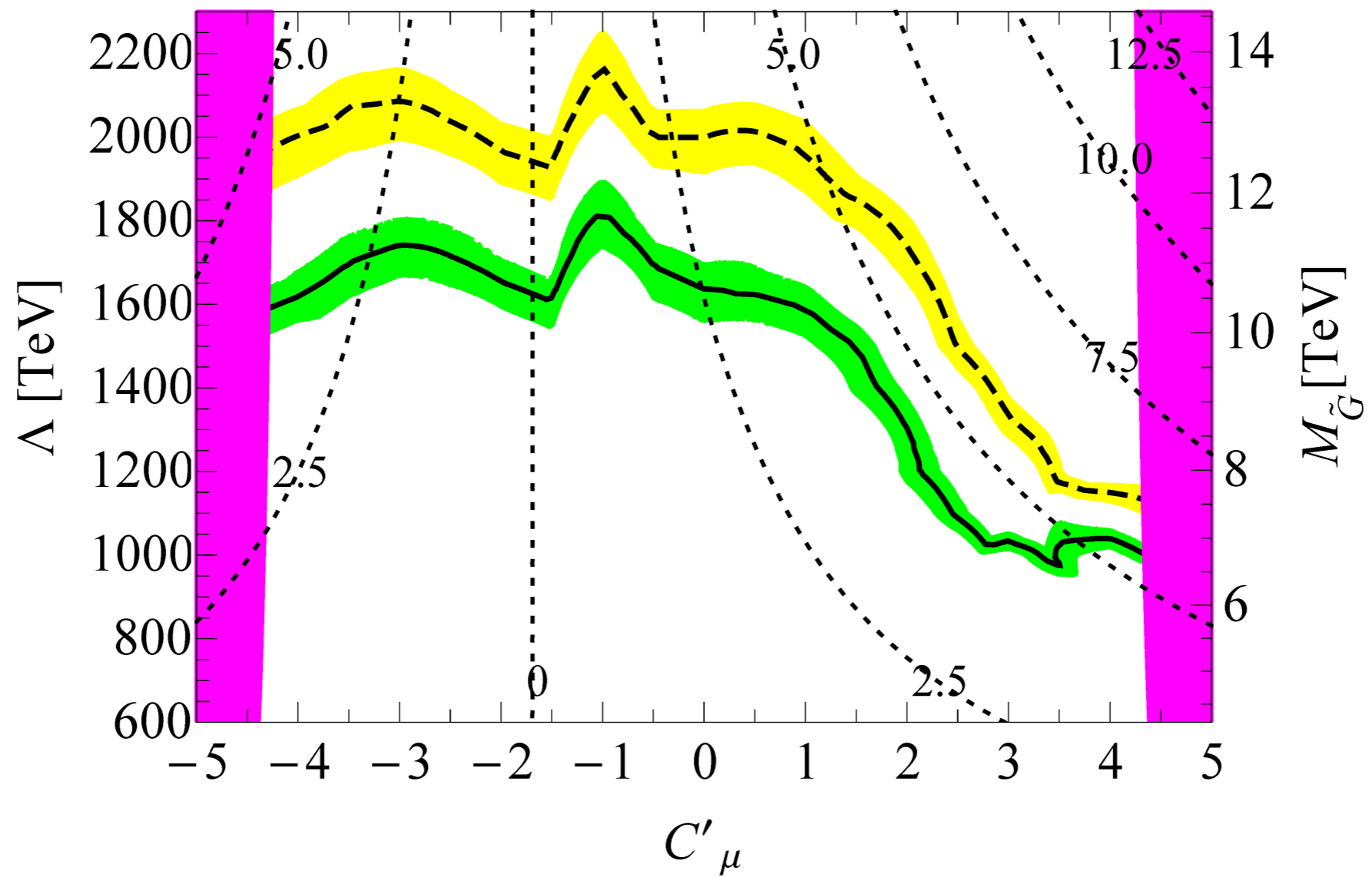
- 2 jets with  $p_T > 0.1M_{\text{eff}}$
- no lepton
- $E_T^{\text{miss}} > 0.2M_{\text{eff}}$
- $M_{\text{eff}} > 15\text{TeV}$
- 3 or more b-jest

5 signal  
regions

# AMSB at a 100 TeV collider ( $3ab^{-1}$ )



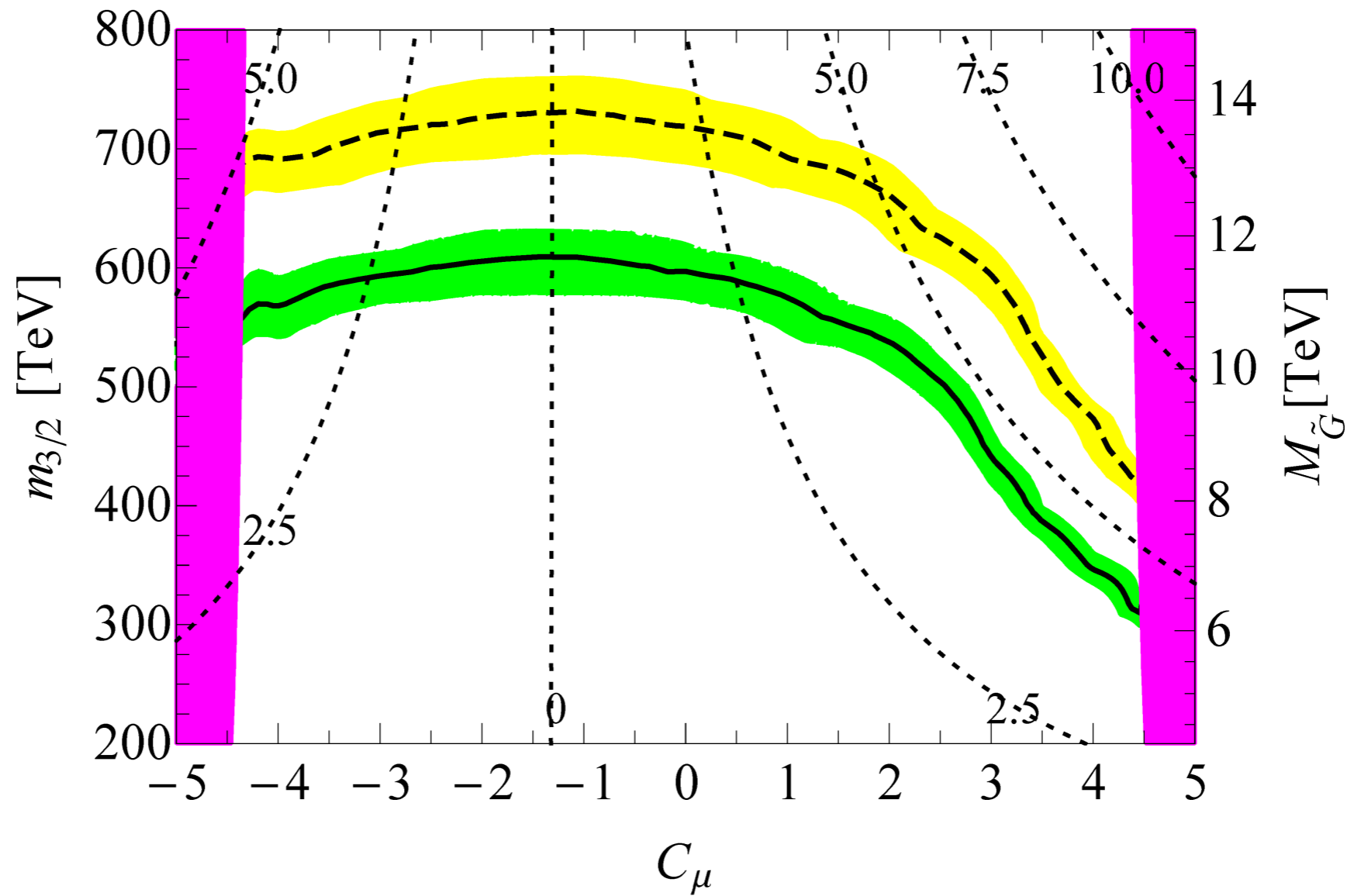
# GMSB at a 100 TeV collider



# Dark matter

If the LSP is a **wino**: need  $M_{\tilde{W}} \sim 2.7\text{TeV}$

100 TeV collider



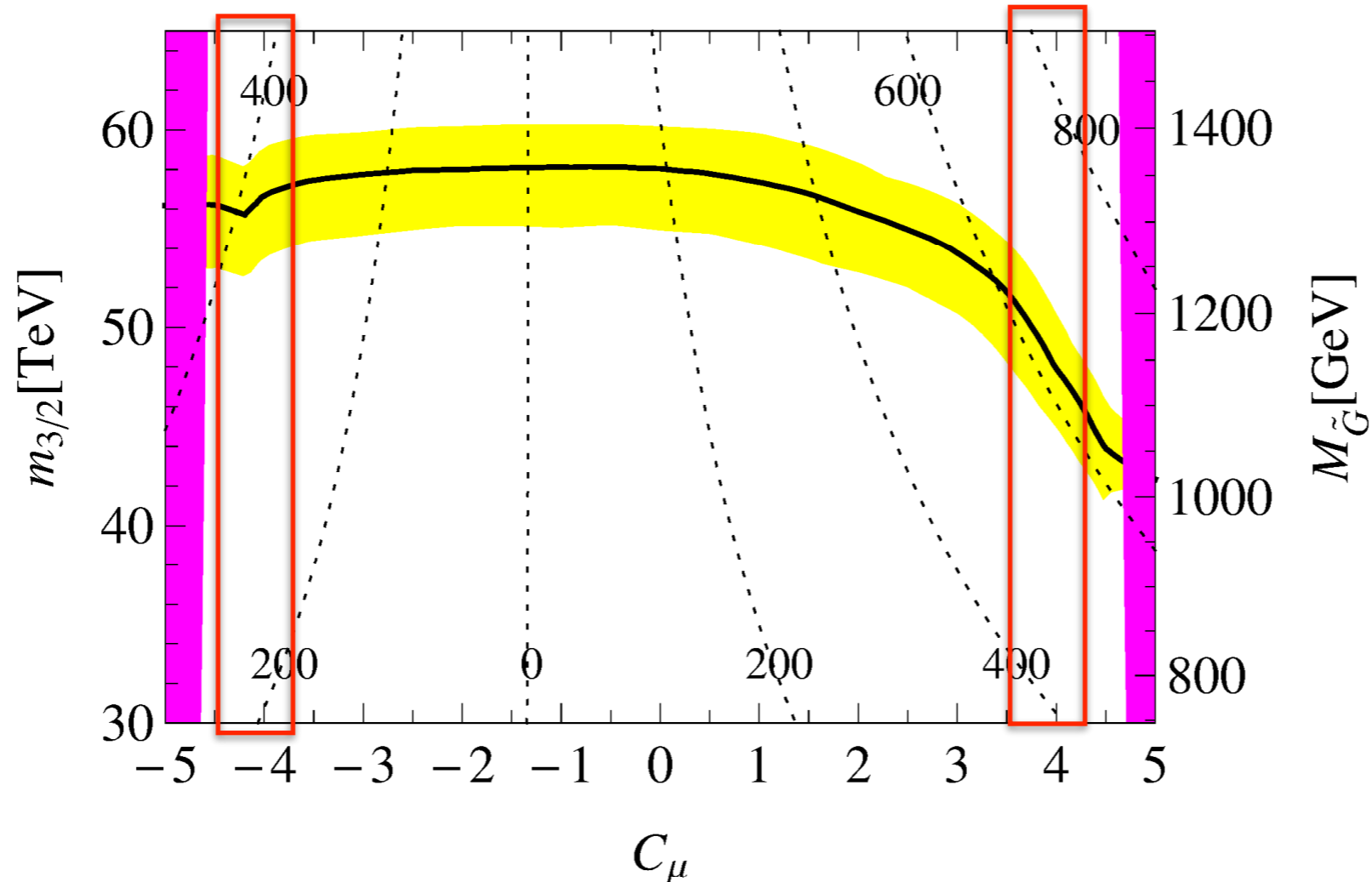
If the LSP is a **bino**: need dilution

If the LSP is a **bino-wino** mixture ( $|C_\mu| \sim 4$ ):

$M \sim$  several 100 GeV

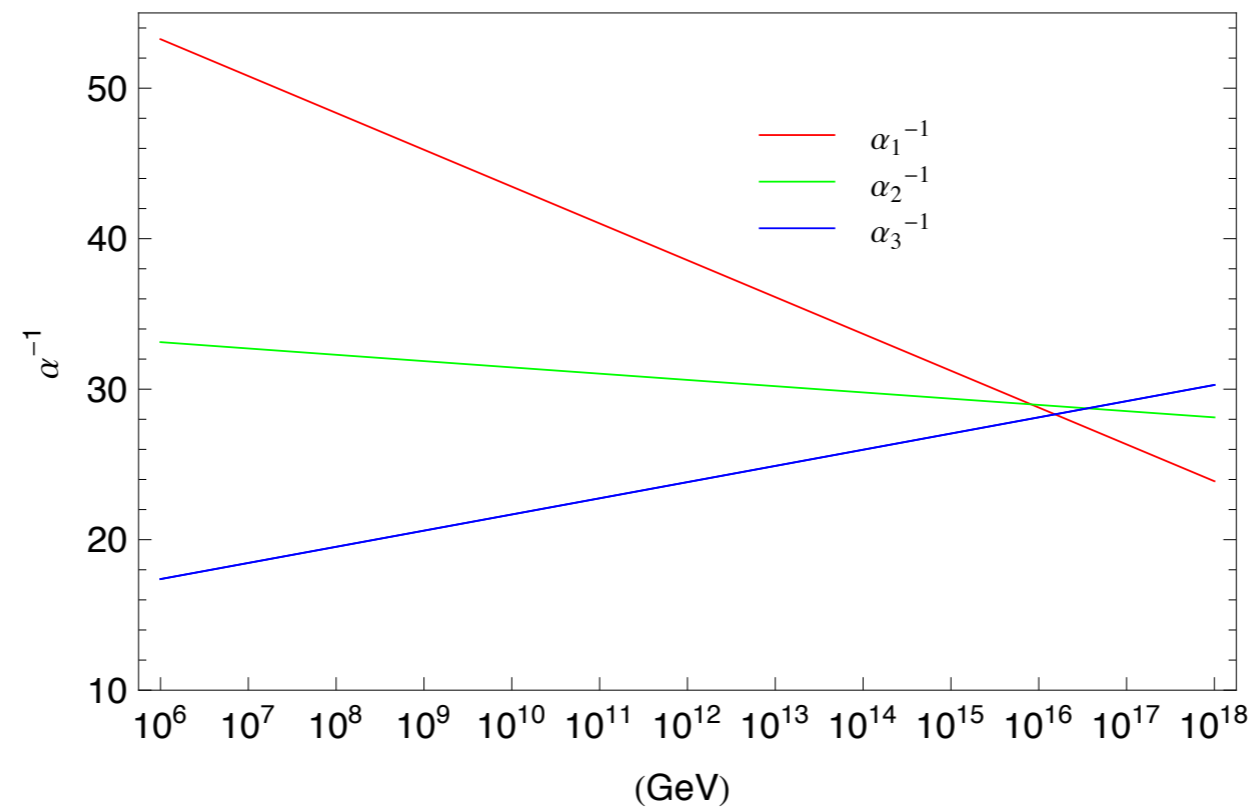
well-tempered neutralino

Arkani-  
Hamed, Delgado, Giudice





Gauge couplings unification  
modified by heavy Higgsino,  
but seems to work



Arkani-Hamed, Gupta,  
Kaplan, Weiner, Zowarski