

INDUCED EWSB

GGI, Florence
September 21, 2015



Jamison Galloway

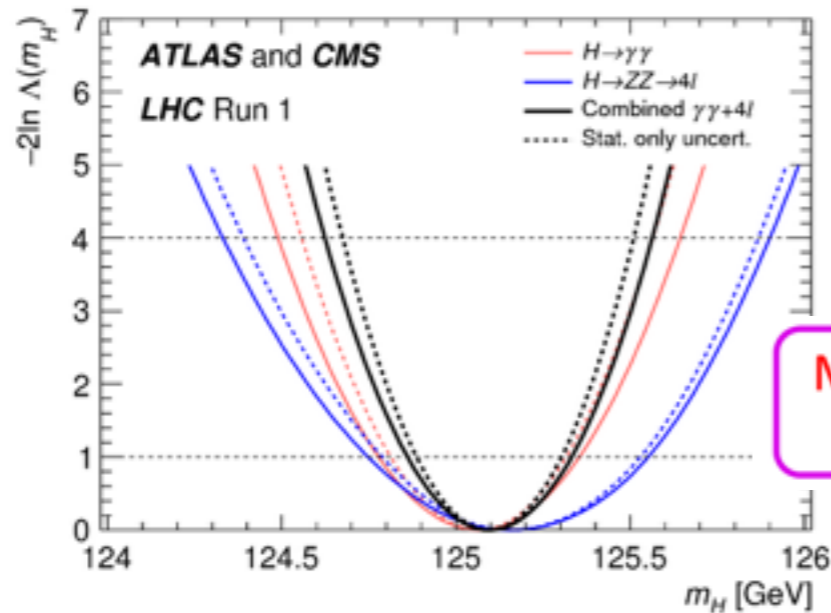
Based on work with:
A. Azatov, S. Chang, M. Luty, E. Salvioni, Y. Tsai, Y. Zhao

outline

- o Introduction/Motivation
- o Modeling, take one: realization with strong dynamics
- o Modeling, take two: realization with perturbative dynamics
- o Implications: phenomenology at the LHC
- o A conclusion or two

introduction

h mass:

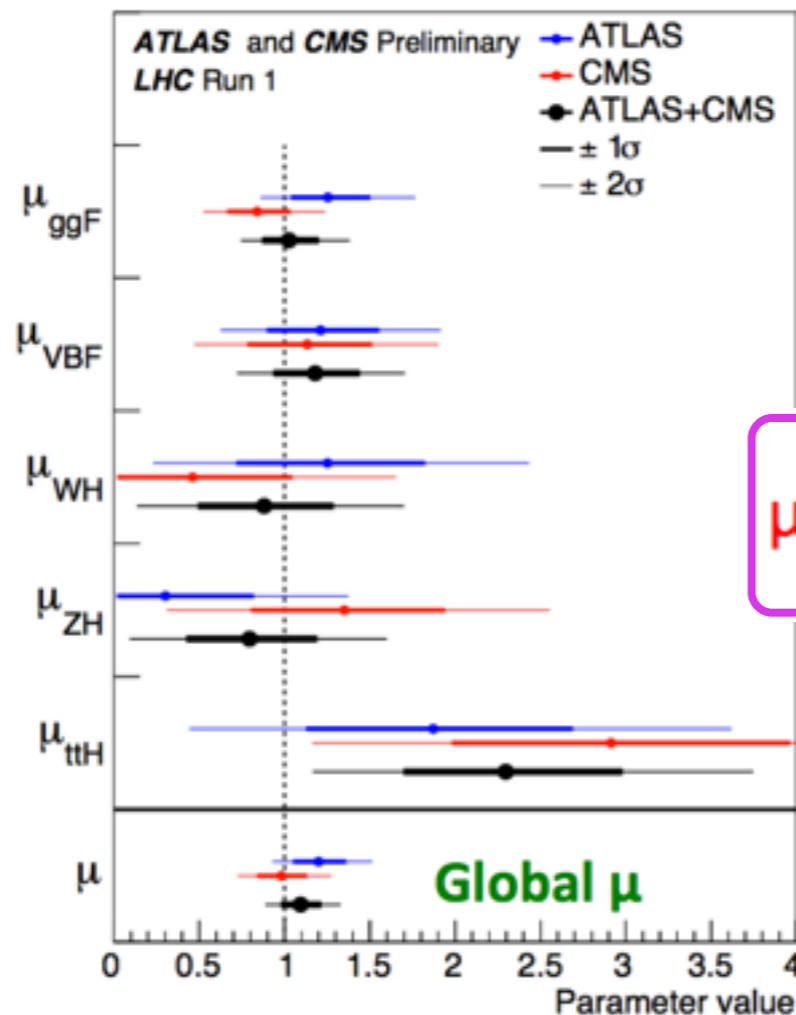


$$M_H = 125.09 \pm 0.24 \text{ GeV} \\ = \pm 0.21 \text{ (stat.)} \pm 0.11 \text{ (syst.) GeV}$$

Examples:

- v/f in comp. H (MCHM4)
- ‘non-alignment’ in type-1 2HDM at large $\tan \beta$
- Higgs mixing with singlet

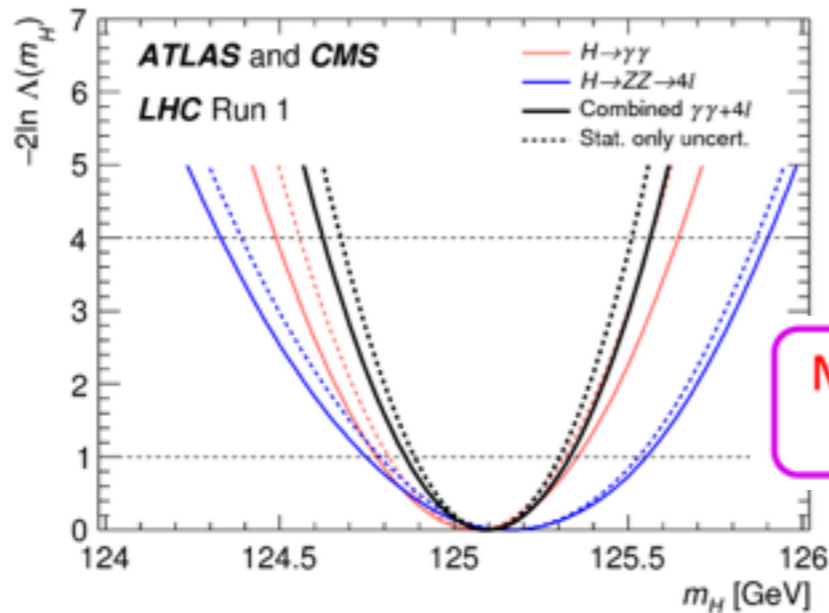
h couplings:
(tree-level)



$$\mu = 1.09^{+0.11}_{-0.10}$$

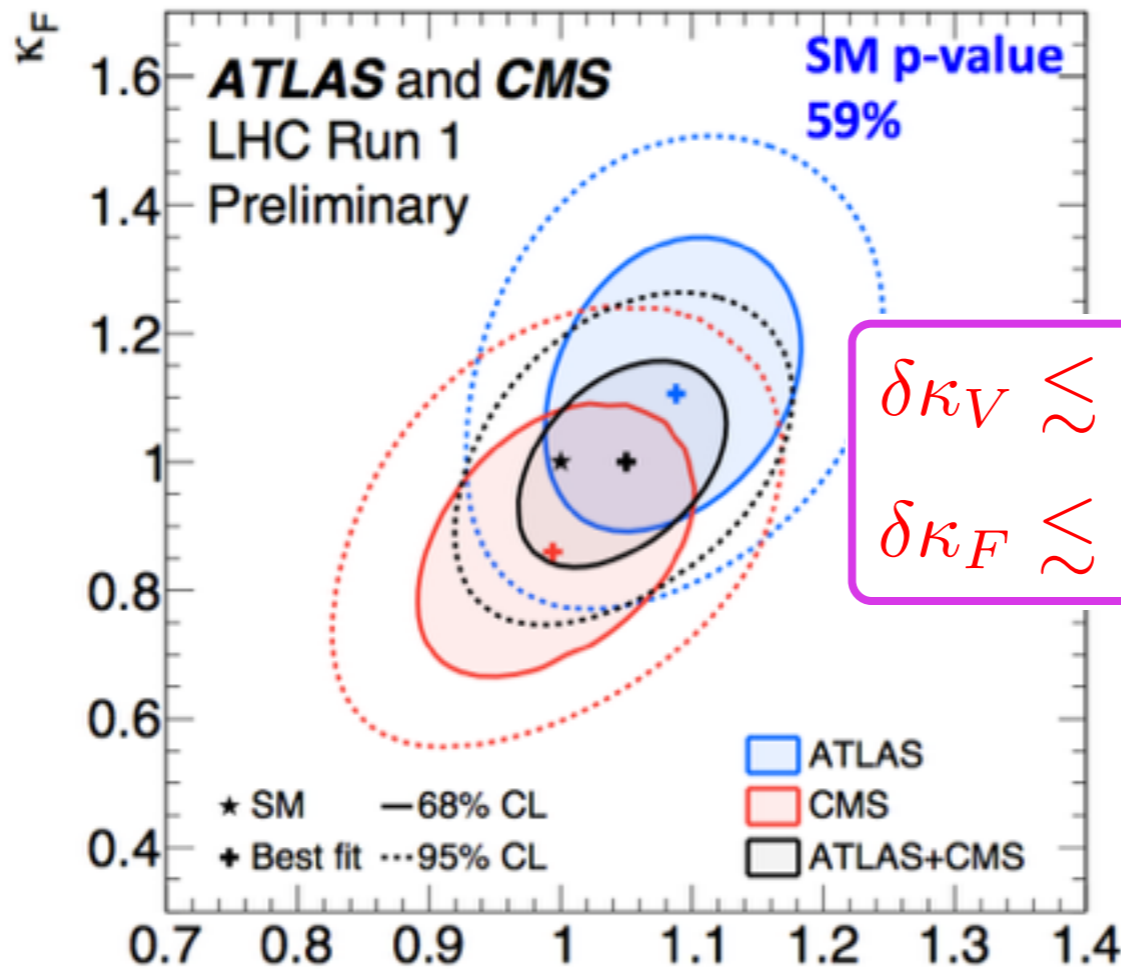
introduction

h mass:



$$M_H = 125.09 \pm 0.24 \text{ GeV} \\ = \pm 0.21 \text{ (stat.)} \pm 0.11 \text{ (syst.) GeV}$$

h couplings:
(tree-level)



$$\delta \kappa_V \lesssim 0.08 \\ \delta \kappa_F \lesssim 0.16$$

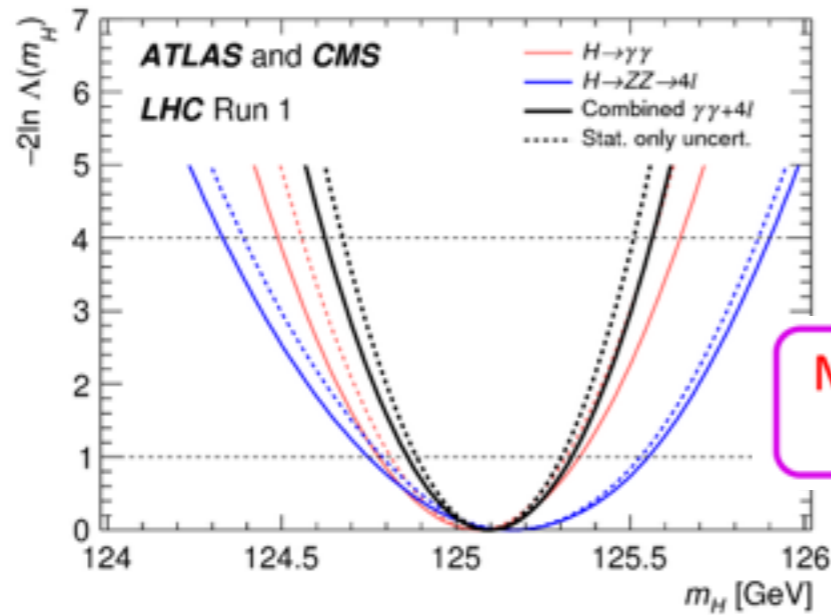
Examples:

- o type-1 2HDM at variable $\tan \beta$

*ref: M. Pieri @ LHCP 2015

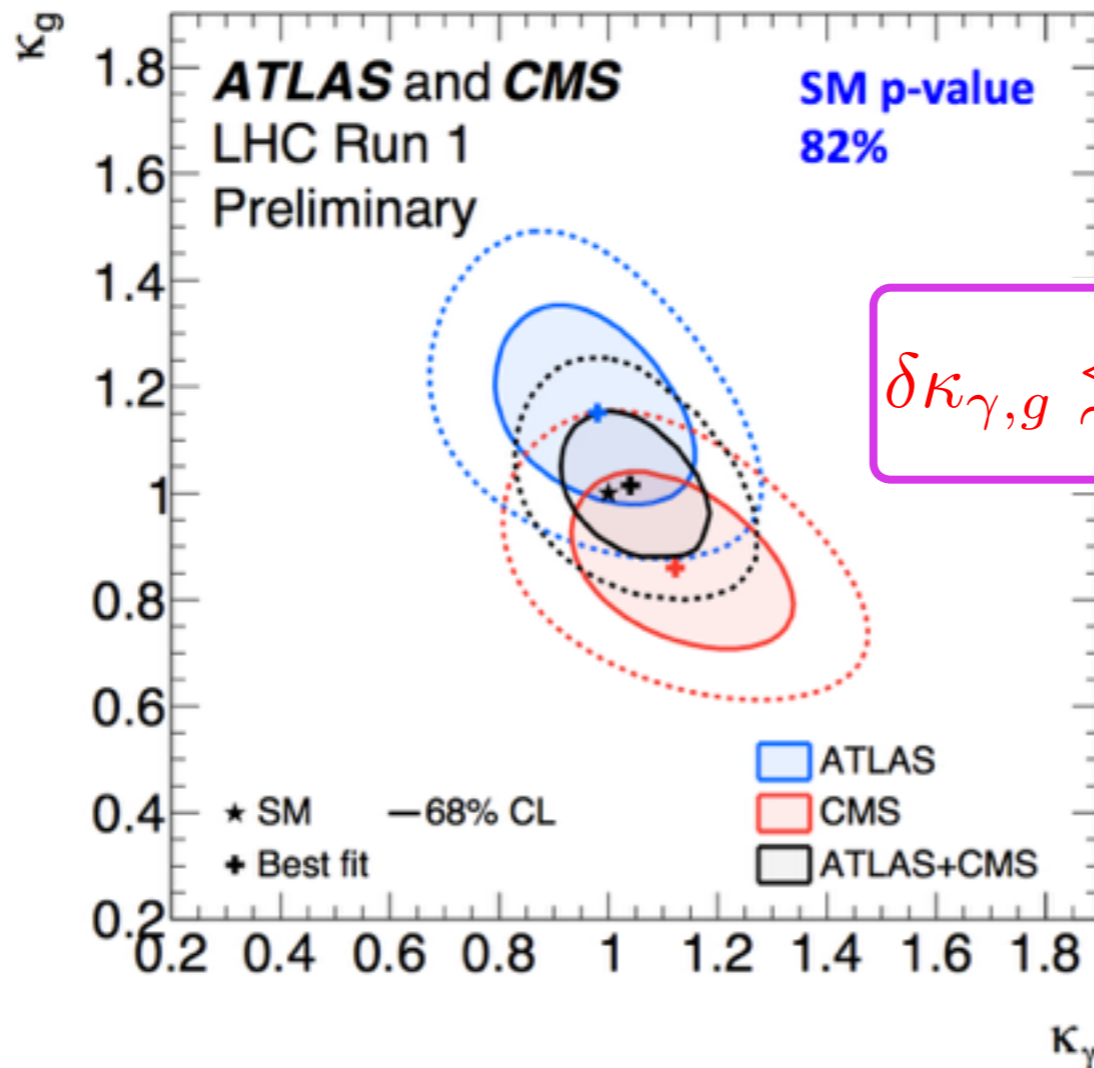
introduction

h mass:



$$M_H = 125.09 \pm 0.24 \text{ GeV} \\ = \pm 0.21 \text{ (stat.)} \pm 0.11 \text{ (syst.) GeV}$$

h couplings:
(loop-level)



$$\delta\kappa_{\gamma,g} \lesssim 0.14$$

Examples:

- light top partners
- charged Higgses

*ref: M. Pieri @ LHCP 2015

takeaway (and assumptions going forward)

- o b confirmed at 125 GeV
- o Tree-level couplings are already at $\sim\text{SM} \pm 10\%$
- o Loop-level at $\sim\text{SM} \pm 15\%$, consistent with...
- o ...null results from partner searches up to ~ 600 GeV

takeaway (and assumptions going forward)

- o b confirmed at 125 GeV

****~ SUSYish****

- o Tree-level couplings are already at $\sim\text{SM} \pm 10\%$

****composite H tuned at least at 10%****

- o Loop-level at $\sim\text{SM} \pm 15\%$, consistent with...

****~inconclusive****

- o ...null results from partner searches up to ~ 600 GeV

****still room for (somewhat) natural elementary H ****

I'll take this circumstantial evidence for an *elementary* Higgs seriously;

assume SUSY stabilization and focus on the question of **mass**

SUSY Higgs and its mass

$$m_h^2 = 2\lambda v^2$$

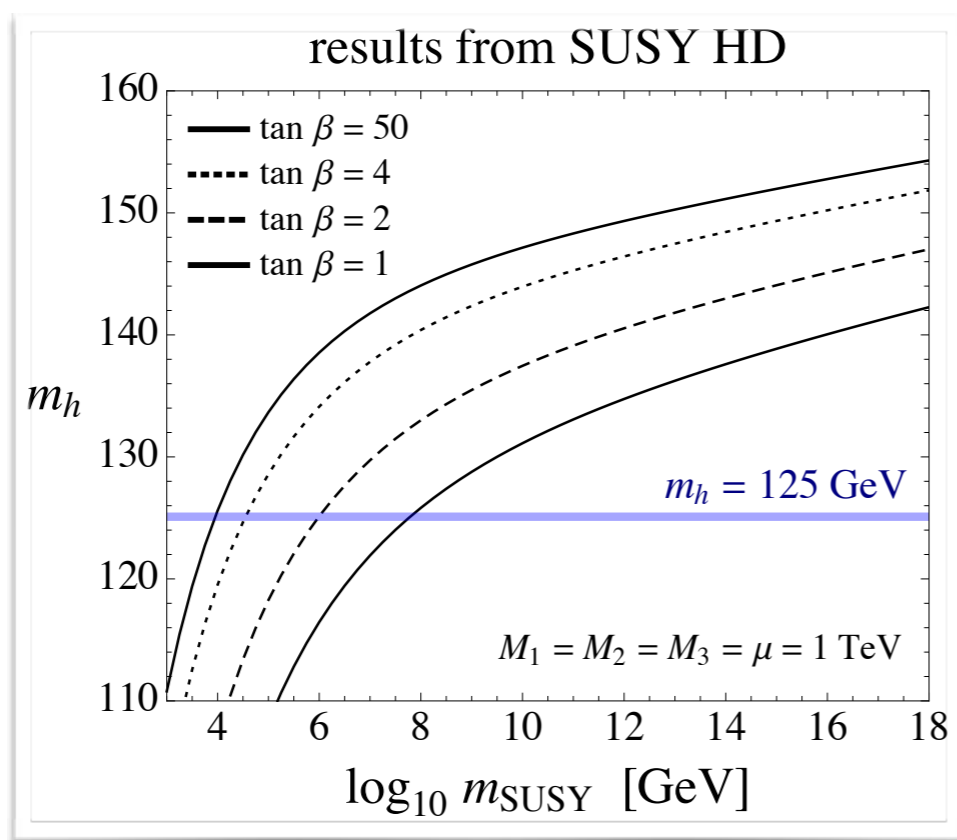
$$\lambda \leq \frac{1}{8}(g^2 + g'^2)$$

$$m_h \leq m_Z$$

SUS'ic relation for H quartic is too small by a factor of 2...

⇒ need an order one breaking!

Can be done with spectrum, but not very naturally



*ref: Vega, Villadoro (JHEP 2015)

| $\tan \beta$ | $m_{\text{SUSY}} \text{ (GeV)}$ |
|--------------|---------------------------------|
| 1 | 3×10^7 |
| 2 | 10^6 |
| 4 | 35×10^3 |
| 50 | 10^4 |

problem: $\delta m_H^2 \propto m_{\text{SUSY}}^2$

High SUSY scale (and thus pressure on naturalness) boils down to very special role of quartic (and very 'special' smallness of it)
 unavoidable consequence of $V(H)$ with negative quadratic

induced EWSB: a strong model

[turn that frown upside down]

what if... $V(H) \sim \oplus (125)^2 |H|^2 + \lambda |H|^4$

naively: $\left\{ \begin{array}{l} \circ \text{ EW intact } \text{👎} \\ \circ \text{ massless W, Z, fermions } \text{👎} \\ \circ \text{ physical mass approximately independent of quartic } \text{👍} \end{array} \right.$

less naively: $\left\{ \begin{array}{l} \circ \text{ EW broken by QCD } \text{👍} \\ \circ \text{ W, Z acquire mass } \sim gf_\pi/2 \sim 50 \text{ MeV } \text{👎} \\ \circ \text{ electron mass } m_e \sim y_e y_q \times 4\pi f_\pi^3 / m_h^2 \sim 10^{-5} \text{ eV } \text{👎} \end{array} \right.$

4 strong NO votes
~1.5 yes votes

(consensus may well be misguided;
cf. Trump leading GOP)

induced EWSB: a strong model

Higgs 'VEV' in previous example fixed by QCD...
...consider instead a TC-like sector

$$\left. \begin{array}{l} \text{"}v\text{"} \sim \frac{y_u \Lambda_{\text{QCD}}^3}{16\pi^2 m_h^2} \rightarrow \frac{\lambda \Lambda_{\text{TC}}^3}{16\pi^2 m_h^2} \end{array} \right\} \begin{array}{l} \Lambda_{\text{TC}} = \text{TeV} \\ \text{"}v\text{"} \rightarrow \lambda \times \text{TeV} \end{array}$$

$$\Delta V_{\text{UV}} = m_H^2 |H|^2 - (\lambda H \psi \psi' + \text{h.c.}); \quad \psi = (\square, 2)_0, \quad \psi' = (\bar{\square}, 1)_{-1/2}$$


$$\longrightarrow \Delta V(\mu < \text{TeV}) = m_H^2 |H|^2 - c_1 \left(\frac{\lambda \Lambda_{\text{TC}}^3}{16\pi^2} \textcircled{H} + \text{h.c.} \right) \quad \begin{array}{l} \text{*contrast} \\ \text{e.g. SILH} \end{array}$$

induced EWSB: a strong model

Higgs 'VEV' in previous example fixed by QCD...
...consider instead a TC-like sector

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 $\Delta V(\mu < \text{TeV}) = m_H^2 |H|^2 - c_1 \left(\frac{\lambda \Lambda_{\text{TC}}^3}{16\pi^2} H + \text{h.c.} \right)$

Upshot:

- o Confining dynamics **induces** $\langle H \rangle \neq 0$
- o Elementary Higgs VEV naturally right size.
- o Elementary Higgs mass is **independent of quartic.***
- o New isotriplet (minimally) of scalars exists below $\sim \text{TeV}$.

* Corrections from quartic $< 20\%$

induced EWSB: a strong model

EFT coupling H to TeV scale strong sector (w/ nonlinear sigma field)

*ref: Azatov, Galloway, Luty (PRL 2012)

kinetic:
$$\Delta\mathcal{L} = \frac{f_{\text{TC}}^2}{4} \text{tr} [(D_\mu \Sigma)^\dagger (D^\mu \Sigma)] + \frac{1}{2} \text{tr} [(D_\mu \mathcal{H})^\dagger (D^\mu \mathcal{H})]$$

$$\Rightarrow m_W^2 = \frac{g^2}{4} (f_{\text{TC}}^2 + v_h^2)$$

“bipartisan EWSB”

interaction:
$$\Delta\mathcal{L} = \sum_{n \geq 1} c_n \frac{\Lambda^{4-n}}{16\pi^2} \text{tr} (\mathcal{H}^\dagger \lambda \Sigma)^n$$

controlled expansion parameter $\epsilon \equiv \frac{\lambda v_h}{\Lambda} \ll 1$

$\epsilon \lesssim 0.1$ with $v_h = 230 \text{ GeV}$, $\lambda = 0.5$

constraint from Higgs @ LHC:

$$\frac{\delta g_{VVH}}{g_{VVH}^{(\text{SM})}} \lesssim 0.08 \quad \Rightarrow \quad f < v \times \sqrt{(2 - \delta)\delta} \approx 95 \text{ GeV}$$

induced EWSB: a strong model

Recap: an 'induced' VEV for the elementary field

$$v_h \sim \frac{\lambda}{4\pi} \frac{\Lambda^2}{m_h^2} \times f$$

[sensible for $\lambda = \mathcal{O}(1)$, $\Lambda \sim \text{TeV}$]

...(recklessly) reimagine as a linear sigma model

$$v_h \sim \frac{\lambda}{4\pi} \frac{m_{\sigma_{\text{TC}}}^2}{m_h^2} \times f_{\sigma_{\text{TC}}}$$

[i.e. treat $\Lambda \sim 4\pi f \rightarrow m_\sigma$]

$$\Rightarrow \epsilon \equiv \frac{\lambda v_h}{\Lambda} \rightarrow \boxed{\frac{v_h^2}{f^2} \frac{m_h^2}{m_\sigma^2}}$$

induced EWSB: a strong model

Recap:

Criteria for generalized induced EWSB

- o H in isolation does NOT break EW
[sensible for $\lambda = 0$, $\Lambda \sim 10^4$ GeV]
- o EW broken appreciably by heavy fields
 $(f \sim v_h)$ $(m_\sigma \gg m_h)$
- o i.e. EW nonlinearly realized at scales > 125
- o coupling H to heavy EWSB source induces $\langle H \rangle$
- o can be realized in 2HDM (with weaker couplings)...

$$\Rightarrow \epsilon \equiv \frac{\lambda v_h}{\Lambda} \rightarrow \frac{v_h^2}{f^2} \frac{m_h^2}{m_\sigma^2}$$

induced EWSB: a perturbative model

[considering a single H doublet]

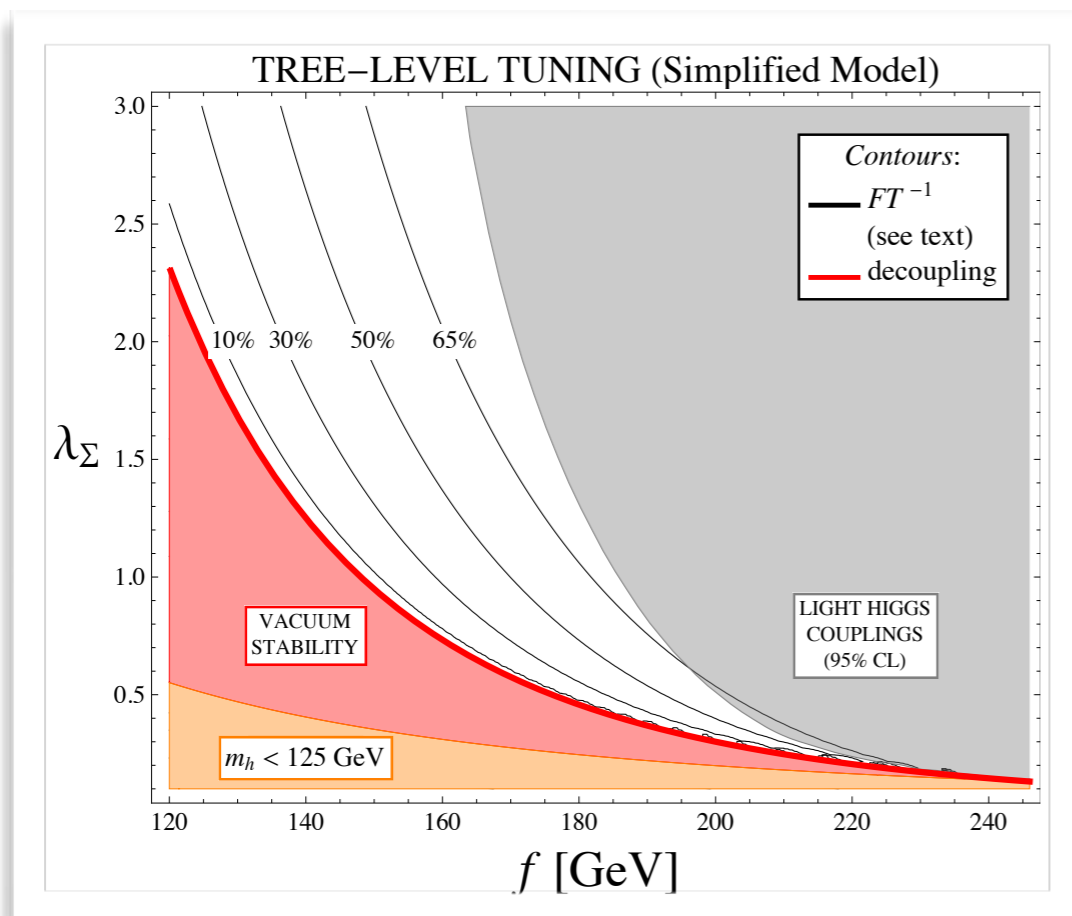
$$V = m_H^2 |H|^2 - m_\Sigma^2 |\Sigma|^2 - \kappa^2 (H^\dagger \Sigma + \text{h.c.}) + \lambda_\Sigma |\Sigma|^4$$

'auxiliary Higgs' $\lambda_\Sigma \gg \lambda_H$

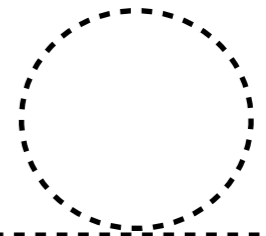
small mixing $\Rightarrow \langle \Sigma \rangle = f \propto \frac{|m_\Sigma|}{\sqrt{\lambda_\Sigma}}, m_\sigma^2 \propto \lambda_\Sigma f^2 \gg m_h^2$

$$\hookrightarrow V_{\text{eff}}(h) = \frac{1}{2} m_H^2 h^2 - \kappa^2 f h + \mathcal{O}(\kappa^4) \Rightarrow v_h \propto \frac{\kappa^2 f}{m_H^2}$$

*ref: Galloway, Luty, Tsai, Zhao (PRD 2014); Alves, Fox, Weiner (PRD 2015)



Higgs mass corrected
via coupling to Σ :



$$\Rightarrow \delta m_h^2 \propto \frac{\lambda}{16\pi^2} m_\sigma^2 \quad (\times \text{mixing angles})$$

(reminiscent of corrections from stops with important distinction that σ needn't be pushed to $\gg \text{TeV}$ scales)

phenomenology: TC-like model

$(H_u, H_d, \Sigma) \Rightarrow 8$ physical scalars:

$$\text{MSSM} \left\{ \begin{array}{l} H_2^\pm, A_2^0 \\ H_1^0, H_1^\pm, A_1^0 \\ h \\ G^\pm, G^0 \end{array} \right\} \pi_{\text{TC}}^{(1,2,3)}$$

$$\Delta\mathcal{L} \supset \lambda(v_h + H)\psi\psi' \Rightarrow m_\pi^2 \sim (\lambda_u v_u + \lambda_d v_d) \times \Lambda$$
$$\equiv (\epsilon_u + \epsilon_d) \times \Lambda^2 \approx (500 \text{ GeV})^2$$

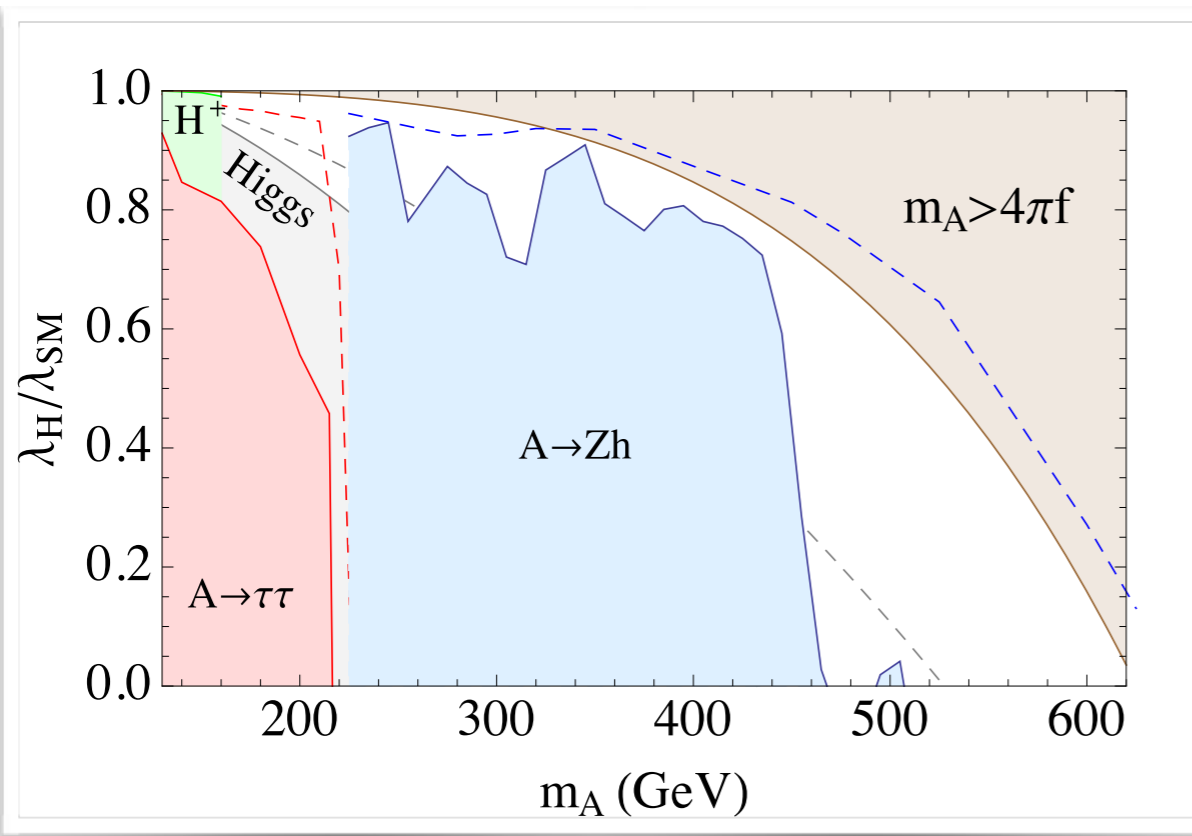
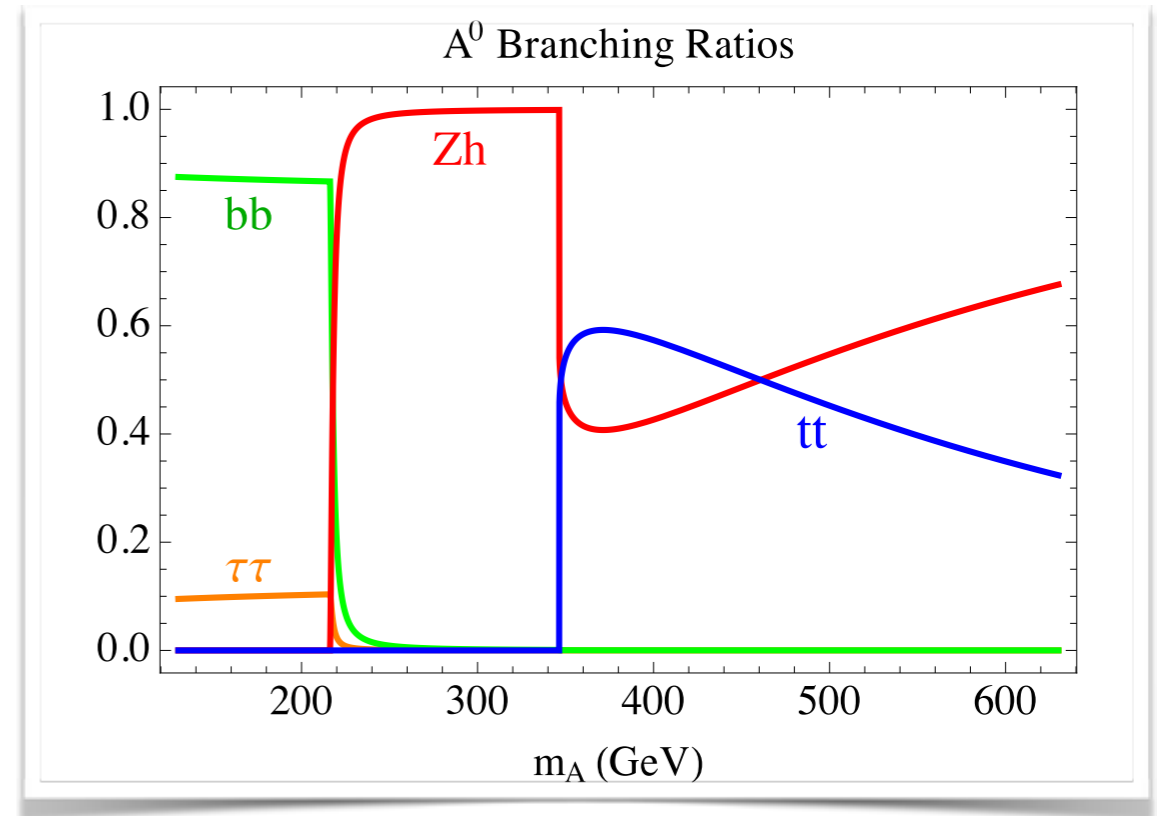
Heavy Higgses (pions) produced by, decay to, SM via mixing
or through auxiliary fields' SU(2) couplings

UNIQUE signals: compare with MSSM (H couples to f),
NMSSM (" S " inherits *all* quantum numbers from mixing),

phenomenology: TC-like model

[examples, exclusions]

- o Decouple second H (to simplify)
- o sub-TeV pseudoscalar remains from TC sector
- o couples to fermions only through mixing: Zh persists even at $m > 350$

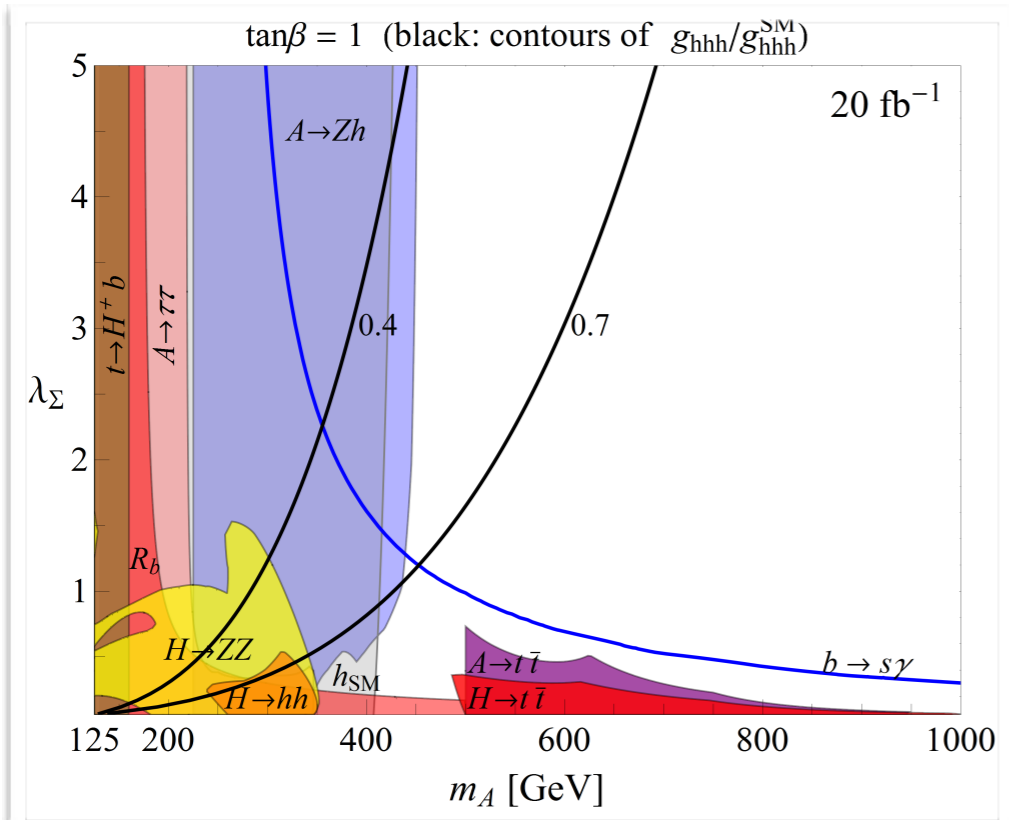


- o $A > Zh$ to cover most space @ LHC
- o powerful exclusion for strong model due to small ff couplings

*ref: Chang, Galloway, Luty, Salvioni, Tsai (JHEP 2015)

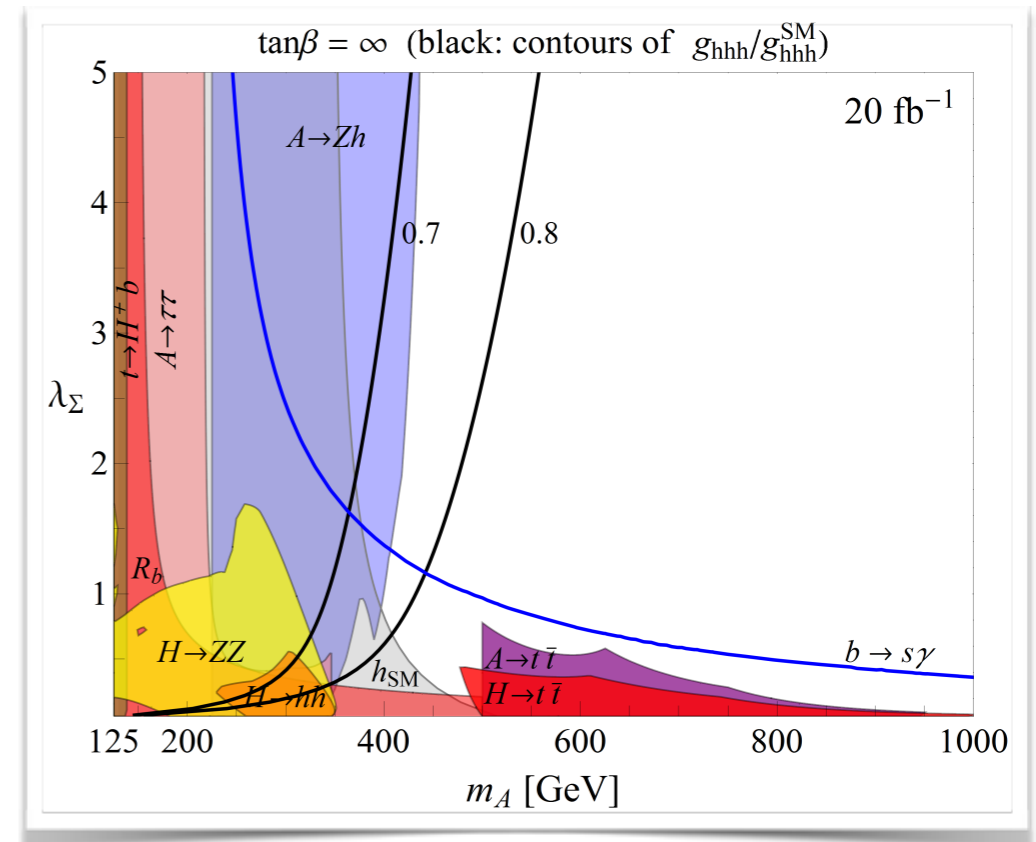
phenomenology: 2HDM-like model

[illustrating possibility of reduced trilinear]



- o direct searches exclude up to $m \sim 400$
- o 50% reduction in H trilinear remains possible (!)

- o even at large $\tan\beta$ significant reduction persists ($\sim 30\%$) for large 'auxiliary' self-interactions



conclusions

- o h @ LHC still allows $\sim 1/3$ of W mass to be generated elsewhere
- o If excitations of this other EWSB source are heavy and couple to H , the Higgs EFT contains a tadpole
- o non-zero Higgs VEV may not require negative quadratic; H may not break EW at all *in isolation*
- o Higgs quartic is consequently untethered from mass
 - > may provide breathing room in SUSY theories especially
 - > can generate large deviations in Higgs self-couplings
- o appearance of light stops will require explanation of Higgs mass; physical mass is essentially free parameter in induced EWSB
- o rich spectrum contains sub-TeV scalars with unique (i.e. non-MSSM) footprints
- o Nonstandard Higgs and add'l scalars still in play...
...any surprise is welcome; many nicely motivated, *and* still viable!

BACKUP

coincidence issue

The active participants:

| | $SU(2)_{TC}$ | $SU(2)_L$ | $SU(2)_R$ |
|---------------|--------------|-----------|-----------|
| ψ | □ | □ | 1 |
| ψ' | □ | 1 | □ |
| \mathcal{H} | 1 | □ | □ |

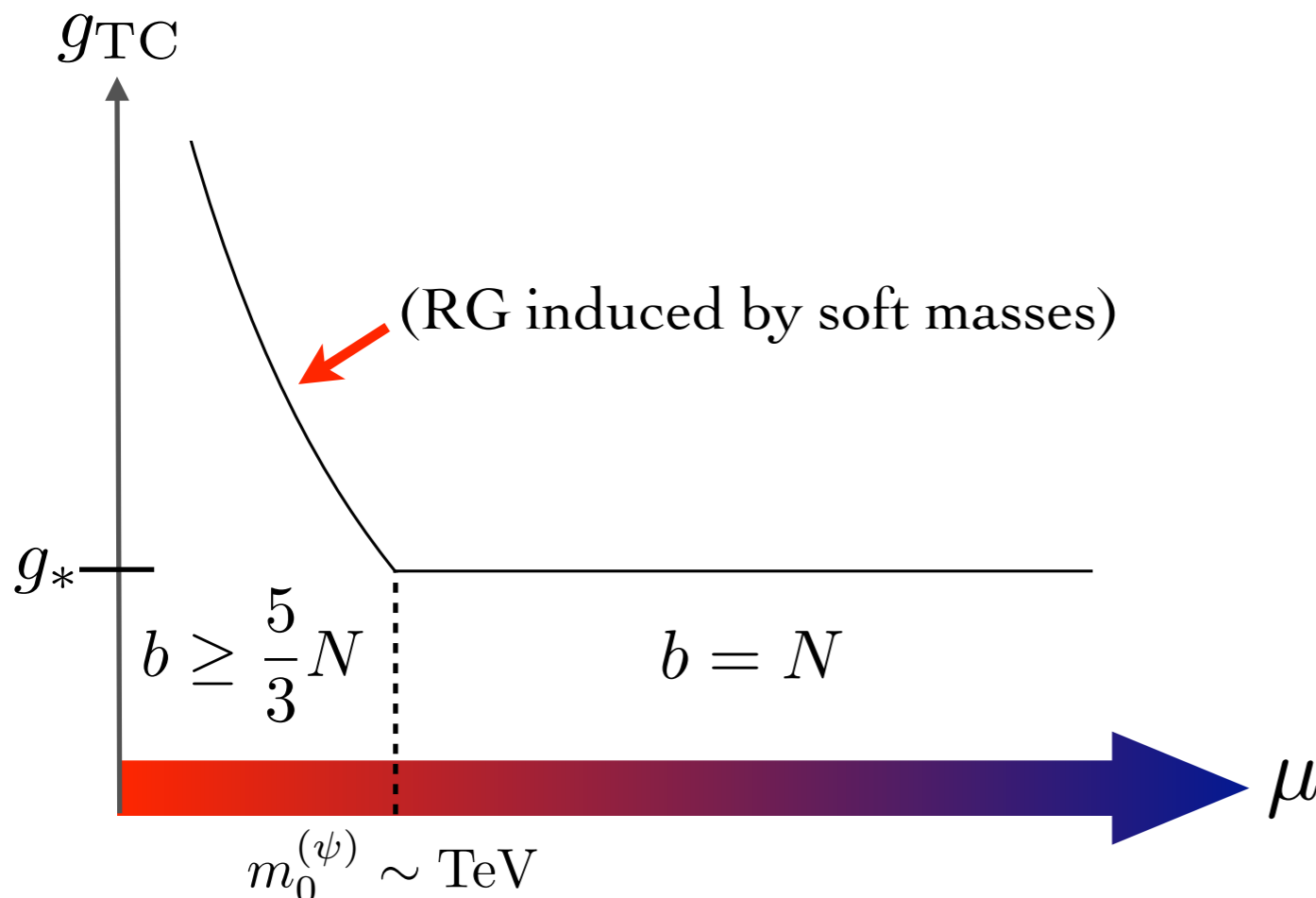
with potential $\Delta W = \lambda \mathcal{H} \psi \psi'$

plus two sterile flavors:

$$b = N$$

\implies IR fixed point

- o Strongly coupled
- o Self-dual



- o Phase transition induced by SUSY breaking
- o **STRONG** fixed point above m_0 sets $\Lambda_{TC} \sim m_h$ without conspiracy
- o As in the QCD toy case: m_h is *free*, **independent of quartic** at leading order

EWPT and strong model

S Parameter:

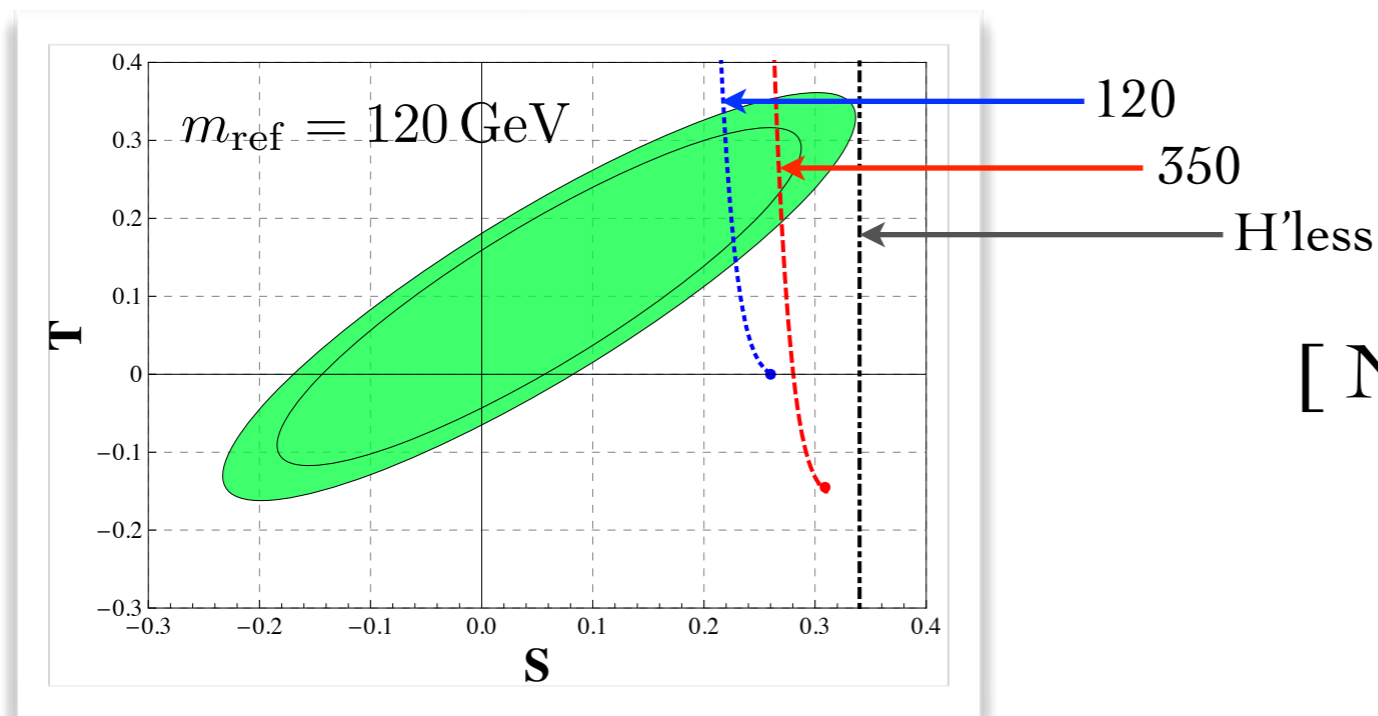
$$\Delta S_{\text{TC}}^{(\text{IR})} \simeq \frac{1}{12\pi} \log \left(\underbrace{\frac{\Lambda^2}{m_h^2}}_{\sim 16\pi^2} \right)$$

vs.

$$\Delta S_{\text{Ind.}}^{(\text{IR})} \simeq \frac{1}{12\pi} \log \left(\underbrace{\frac{\Lambda^2}{m_\pi^2}}_{\sim 4\pi} \right)$$

T Parameter: $\lambda_u v_u = \lambda_d v_d \Rightarrow$ custodial limit

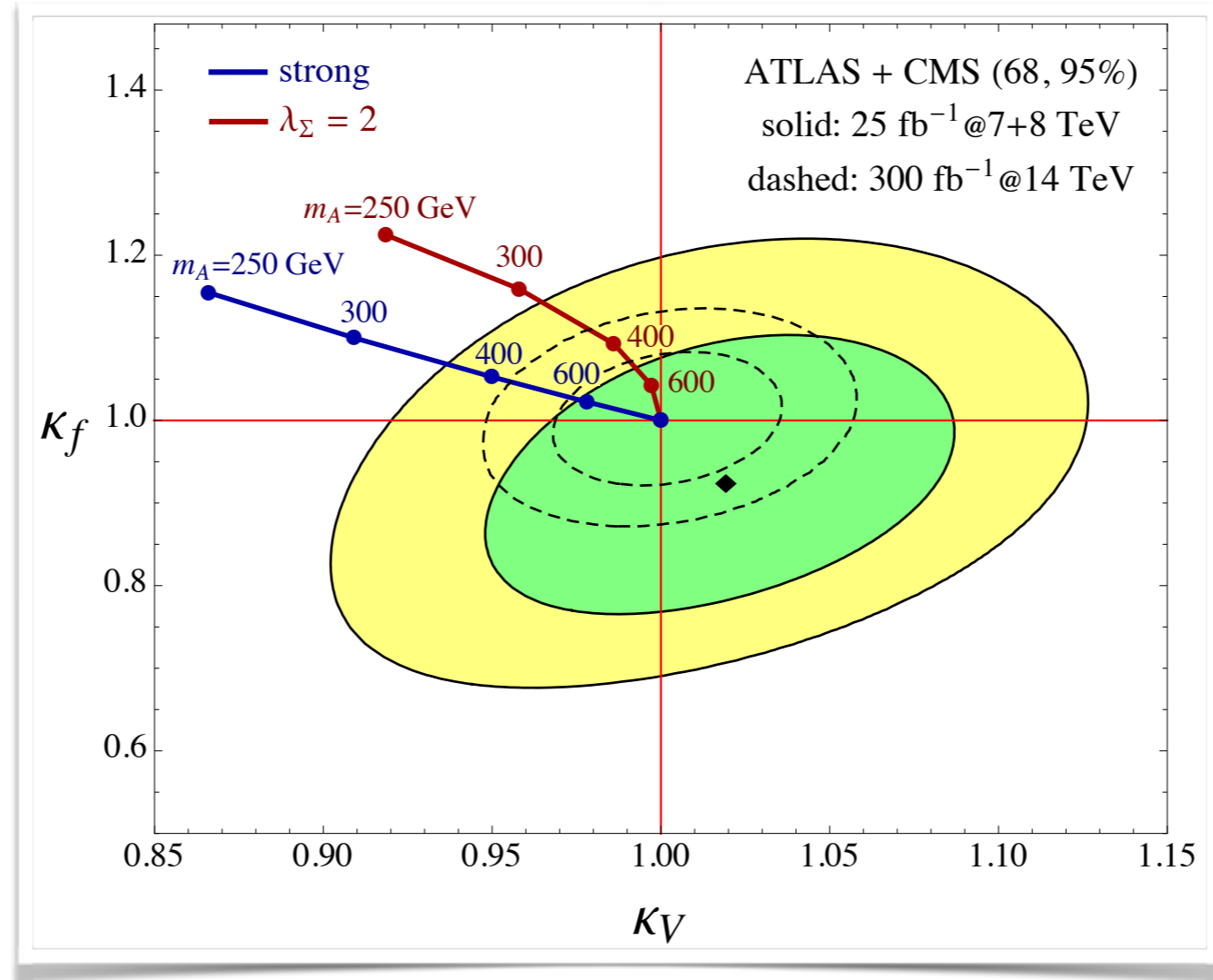
$\hookrightarrow T$ corresponds to a variable parameter of the theory



$$\alpha T = (\epsilon_u - \epsilon_d)^2$$

[Notice that increasing T tends to decrease S above]

'universal' phenomenology: H couplings



generation of auxiliary quartics

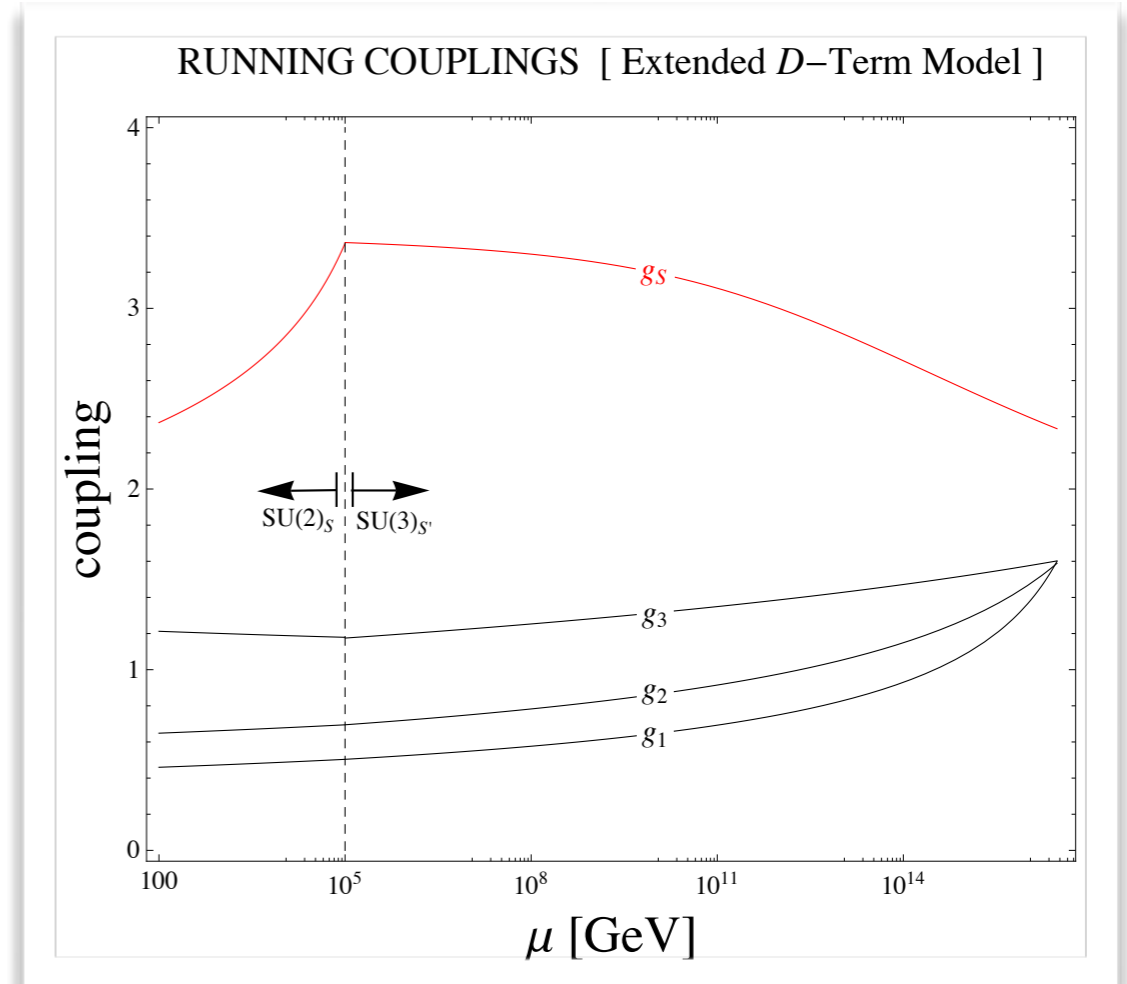
$$D\text{-Terms} \left\{ \begin{array}{l} \Delta K = \Sigma_{u,d}^\dagger \exp(g_S V^a T^a) \Sigma_{u,d}; \quad T^a \in SU(2)_S \\ \Sigma \in \Psi_5; \quad \Psi_5 = (T, \Sigma) \quad [SU(2)_S \text{ broken by } \langle \Phi \rangle] \end{array} \right.$$



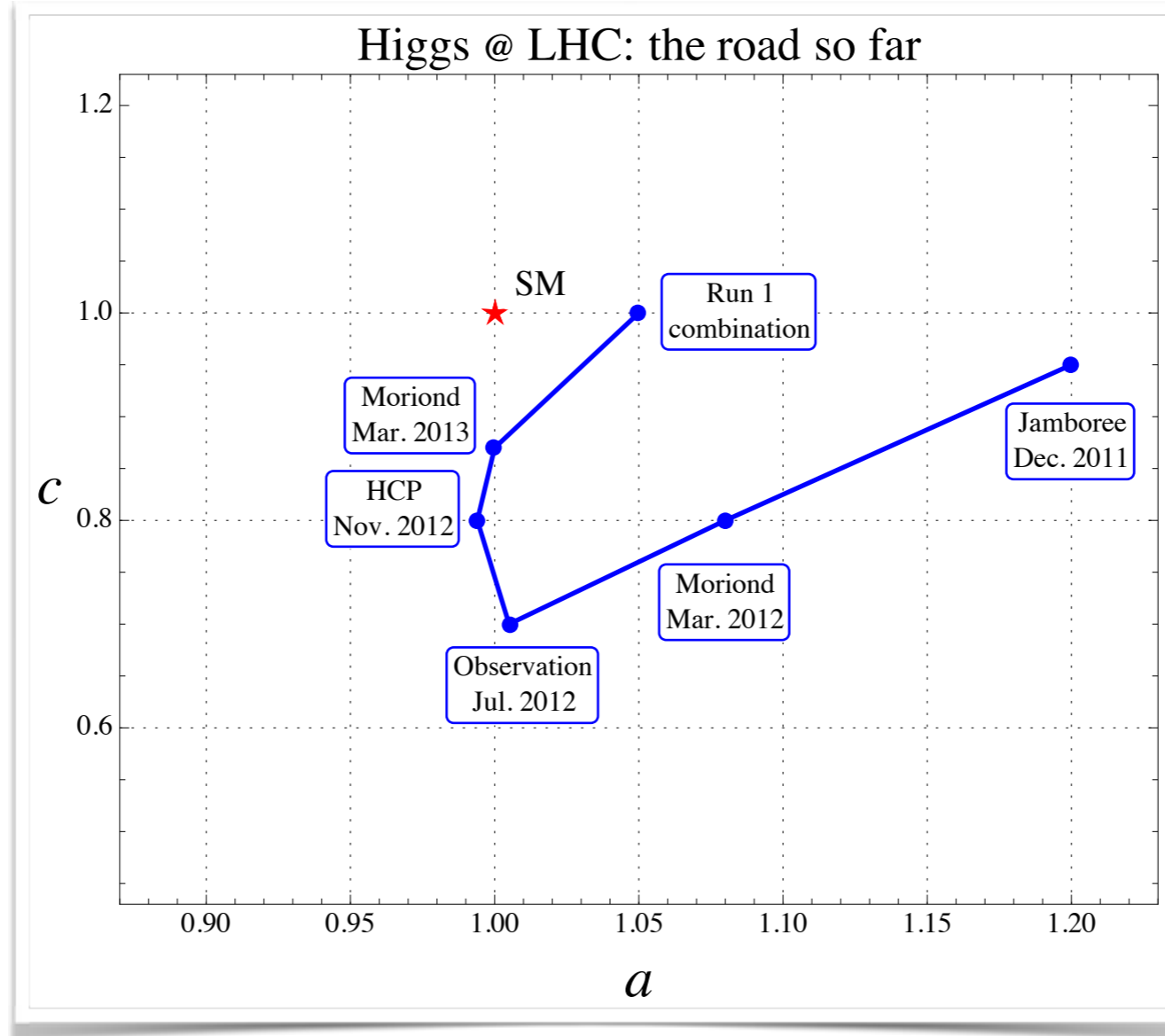
$(\Psi, \bar{\Psi}, \Phi, \bar{\Phi}) = 6$ flavors



running starts at two loops;
some completion still required



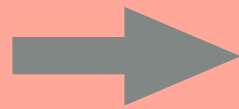
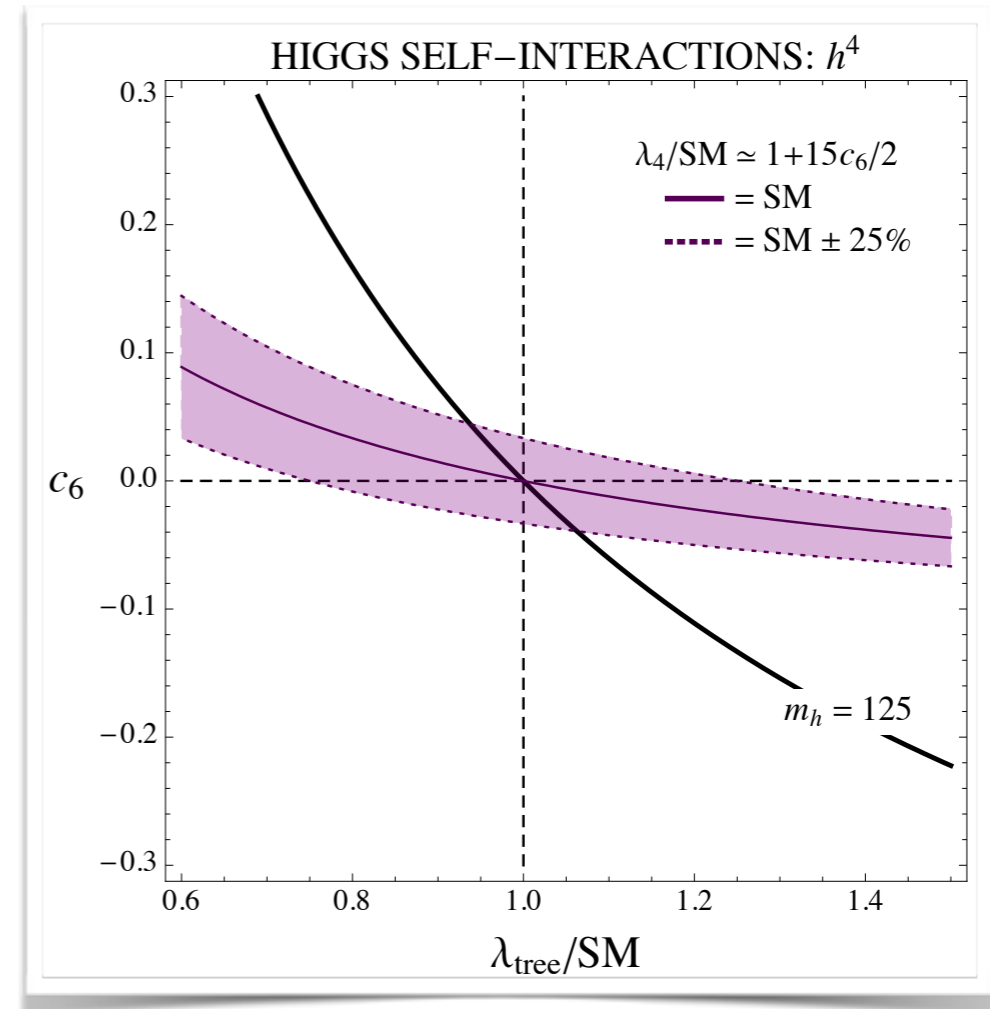
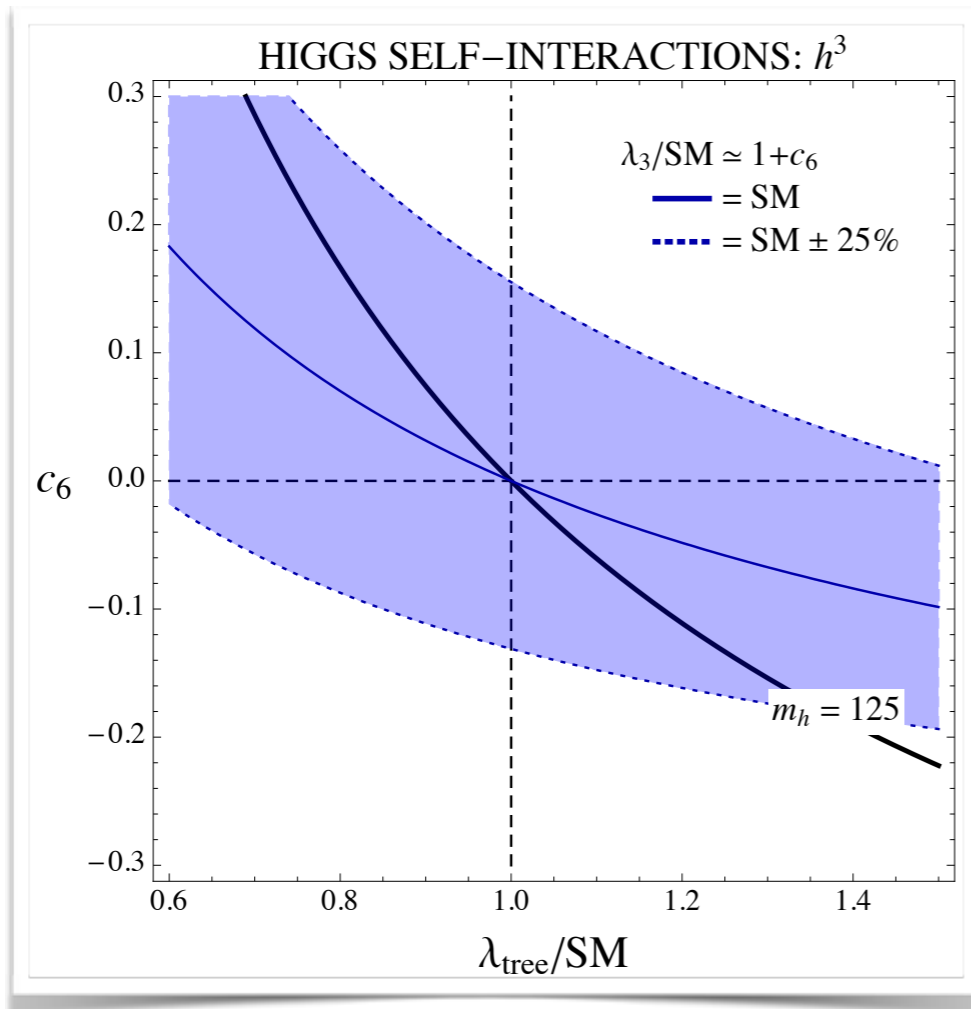
history of Higgs @ LHC



Higgs self-interaction: comparison with SILH

H carries all light scalars of the theory.

Self-interactions modified by $(H^\dagger H)^3$, $(\partial_\mu(H^\dagger H))^2$:



Small cubic implausible: $M < 500$ GeV for vectors!
 hVV and $hff \sim \text{SM}$ *does* paint us into a corner