

SUSY INTEGRABLE SPIN CHAINS AND DISCRETE HIROTA DYNAMICS

A.Zabrodin (ITEP, Moscow)

GGI, Florence, 29 October 2008

based on

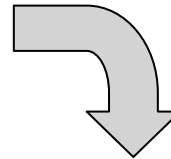
V.Kazakov, A.Sorin, A.Z. Nucl.Phys. B790 (2008) 345;

A.Z. Theor. Math.Phys. 155 (2008) 567

Motivation

- Classical versus quantum integrability:
Classical integrable equations take part in quantum problems even at $\hbar \neq 0$

Quantum spin chains



Discrete classical Hirota dynamics

(I.Krichever, O.Lipman, P.Wiegmann, A.Z. 1996)

- Clarifying SUSY Bethe ansatz,
explaining variety of Bethe equations,
“duality” transformations, etc.

Plan

1. Integrable SUSY GL(K|M)-invariant spin chains:
commuting transfer matrices
2. TT -relation (Hirota equation) for quantum transfer matrices
3. Boundary and analytic conditions
4. Solving the TT -relation by classical methods
 - Auxiliary linear problems
 - Backlund transformations
 - “Undressing” procedure: nested Bethe ansatz
as a chain of Backlund transformations
5. Generalized Baxter’s TQ -relations
6. QQ -relation (Hirota equation for Q ’s) and Bethe equations

$GL(K|M)$ spin chains

(P.Kulish 1985)

$GL(K|M)$ -invariant R -matrix acts in the tensor product $\pi_0 \otimes \pi_1$ of two irreps and

- obeys the graded Yang-Baxter equation

(P.Kulish, E.Sklyanin 1980)

- for any $g \in GL(K|M)$ obeys the condition

$$\pi_0(g) \otimes \pi_1(g) R_{01}(u) = R_{01}(u) \pi_0(g) \otimes \pi_1(g)$$

π_0 arbitrary, $\pi_1 = \pi_{\square}$

vector representation

in the space $\mathbb{C}^K \oplus \mathbb{C}^M$

spectral parameter

generators of $gl(K|M)$

parity

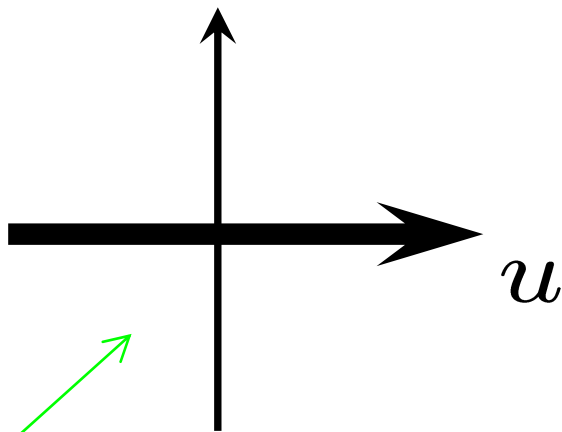
$$(E_{\alpha\beta})_{\alpha'\beta'} = \delta_{\alpha\alpha'}\delta_{\beta\beta'}$$

$$\alpha, \beta = 1, \dots, K+M$$

$$R_{01}(u) = u + 2 \sum_{\alpha\beta} (-)^{p(\beta)} \pi_0(E_{\alpha\beta}) \otimes E_{\beta\alpha}$$

=

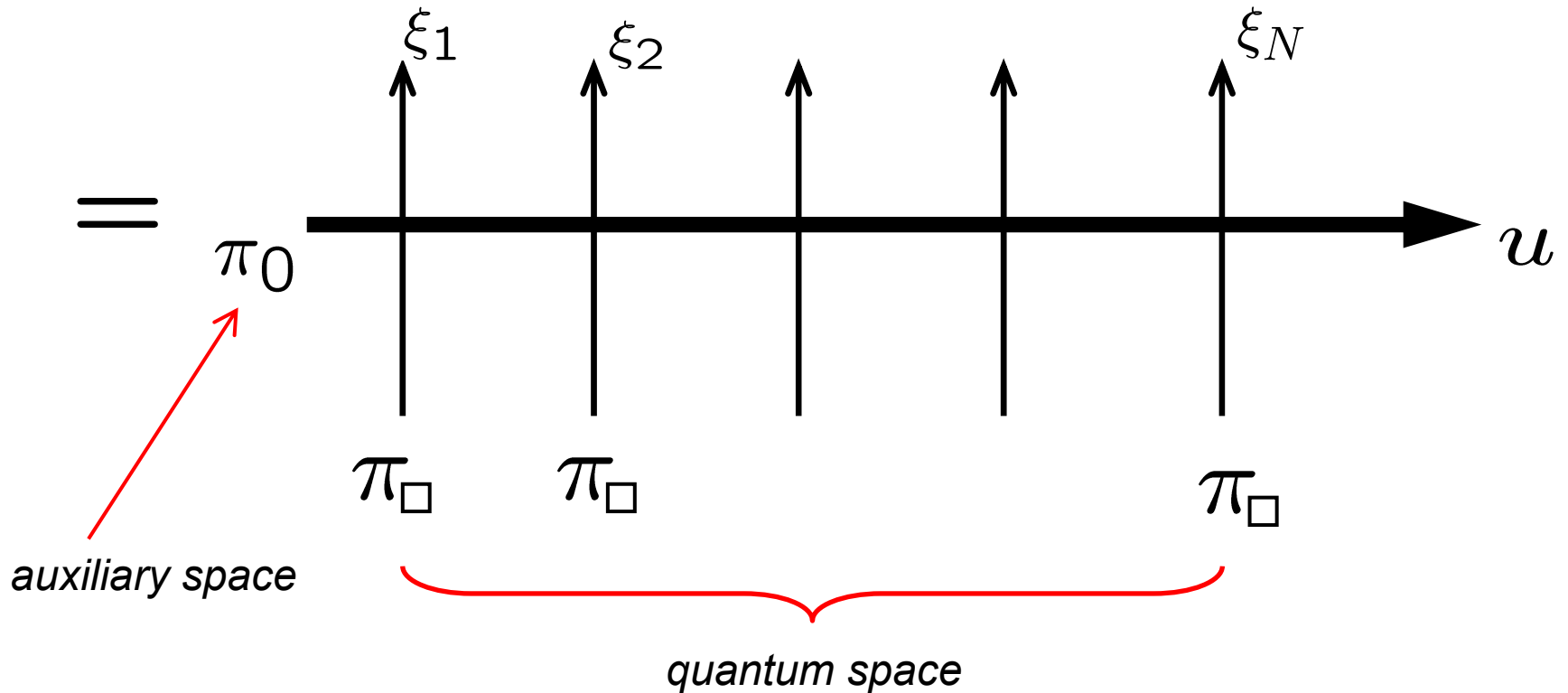
π_0



$\pi_1 = \pi_{\square}$

Quantum monodromy matrix

$$\mathcal{T}(u) = R_{01}(u - \xi_1)R_{02}(u - \xi_2) \dots R_{0N}(u - \xi_N)$$



$$\mathcal{T}(u) : \pi_0 \otimes (\pi_{\square})^{\otimes N} \rightarrow \pi_0 \otimes (\pi_{\square})^{\otimes N}$$

Quantum transfer matrix (periodic b.c.)

$$T^{(\pi_0)}(u) = \text{str}_{\pi_0} (\mathcal{T}(u))$$

Yang-Baxter equation  commutativity

$$[T^{(\pi_0)}(u), T^{(\pi'_0)}(u')] = 0$$

Generalization: twisted b.c.

$$g = \text{diag}(x_1, \dots, x_K, y_1, \dots, y_M) \in GL(K|M)$$

insert before taking **str**

$$T^{(\pi_0)}(u; g) = \text{str}_{\pi_0}(\pi_0(g) \mathcal{T}(u))$$

YB equation + $GL(K|M)$ -invariance



$$[T^{(\pi_0)}(u; g), T^{(\pi'_0)}(u'; g)] = 0$$

Rectangular representations

(We consider covariant representations only)

$$\pi_0 = \pi_s^a = \begin{array}{|c|c|c|c|} \hline & & & \\ \hline & & & \\ \hline & & & \\ \hline \end{array}$$

A 3x4 grid of squares. A green bracket above the grid spans the width of the four columns and is labeled 's'. A green bracket to the right of the grid spans the height of the three rows and is labeled 'a'.

$$\begin{aligned} T(a, s, u) &= T^{(\pi_s^a)}(u - s + a; g) \\ &= \text{str}_{\pi_s^a}(\pi_s^a(g) \mathcal{T}(u - s + a)) \end{aligned}$$

Functional relations for transfer matrices

*P.Kulish, N.Reshetikhin 1983;
V.Bazhanov, N.Reshetikhin 1990;
A.Klumper, P.Pearce 1992;
A.Kuniba, T.Nakanishi, J.Suzuki 1994;
Z.Tsuboi 1997 (for SUSY case)*

$T^{(\pi_0)}(u)$ for general irreps π_0 can be expressed

through $T(1, s, u)$ for the rows 

In particular,

$$T(a, s, u) \propto \det_{1 \leq i, j \leq a} T(1, s+i-j, u+a+1-i-j)$$

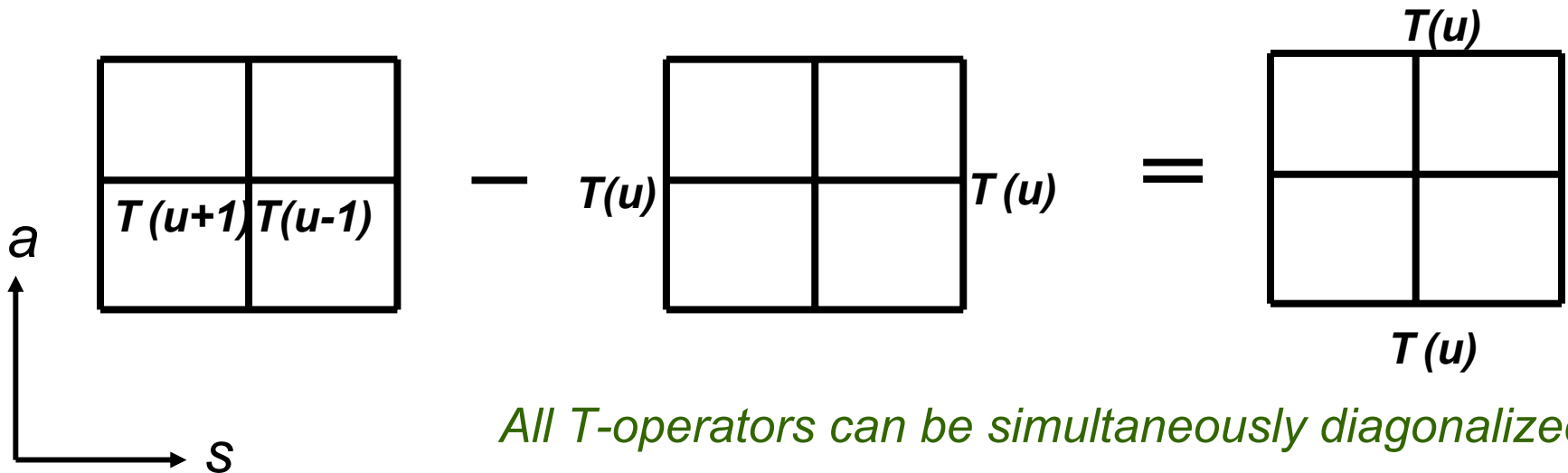
This is the Bazhanov-Reshetikhin det-formula for rectangular irreps.

- *Analog of Weyl formula for characters*
- *The same for SUSY and ordinary case*

TT-relation

- directly follows from the BR det-formula
- is identical to the **HIROTA EQUATION**

$$T(a, s, u+1)T(a, s, u-1) - T(a, s+1, u)T(a, s-1, u) = T(a+1, s, u)T(a-1, s, u)$$



All T-operators can be simultaneously diagonalized



The same relation holds for any of their eigenvalues

Hirota difference equation

(R.Hirota 1981)

- Integrable difference equation solvable by classical inverse scattering method
- A “master equation” of the soliton theory:
 - provides universal discretization of integrable PDE’s
 - generates infinite hierarchies of integrable PDE’s
- We use it to find all possible Baxter’s TQ-relations and nested Bethe ansatz equations for super spin chains

(similar method for ordinary spin chains:

I.Krichever, O.Lipan, P.Wiegmann, A.Zabrodin 1996)

Analytic and boundary conditions

$T(a, s, u)$ are polynomials in u of degree N for any a, s

length of spin chain

What is $T(0, s, u)$? And $T(a, 0, u)$?

Transfer matrix for the trivial irrep is proportional to the identity operator:

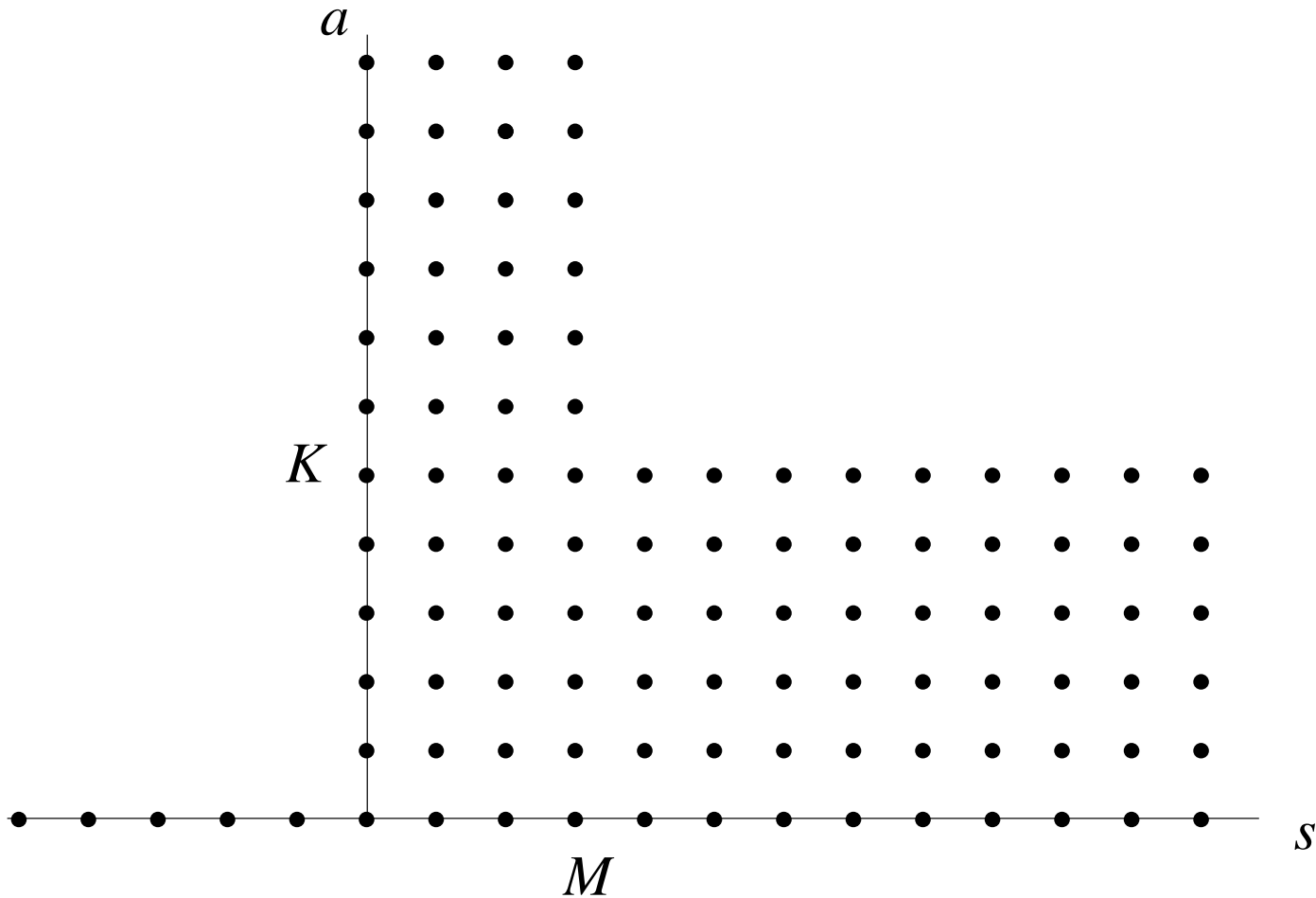
$$T(0, s, u) = \phi(u - s), \quad T(a, 0, u) = \phi(u + a)$$

$$\phi(u) = \prod_{i=1}^N (u - \xi_i)$$

*Here $a, s > 0$.
Extension to
negative values?*

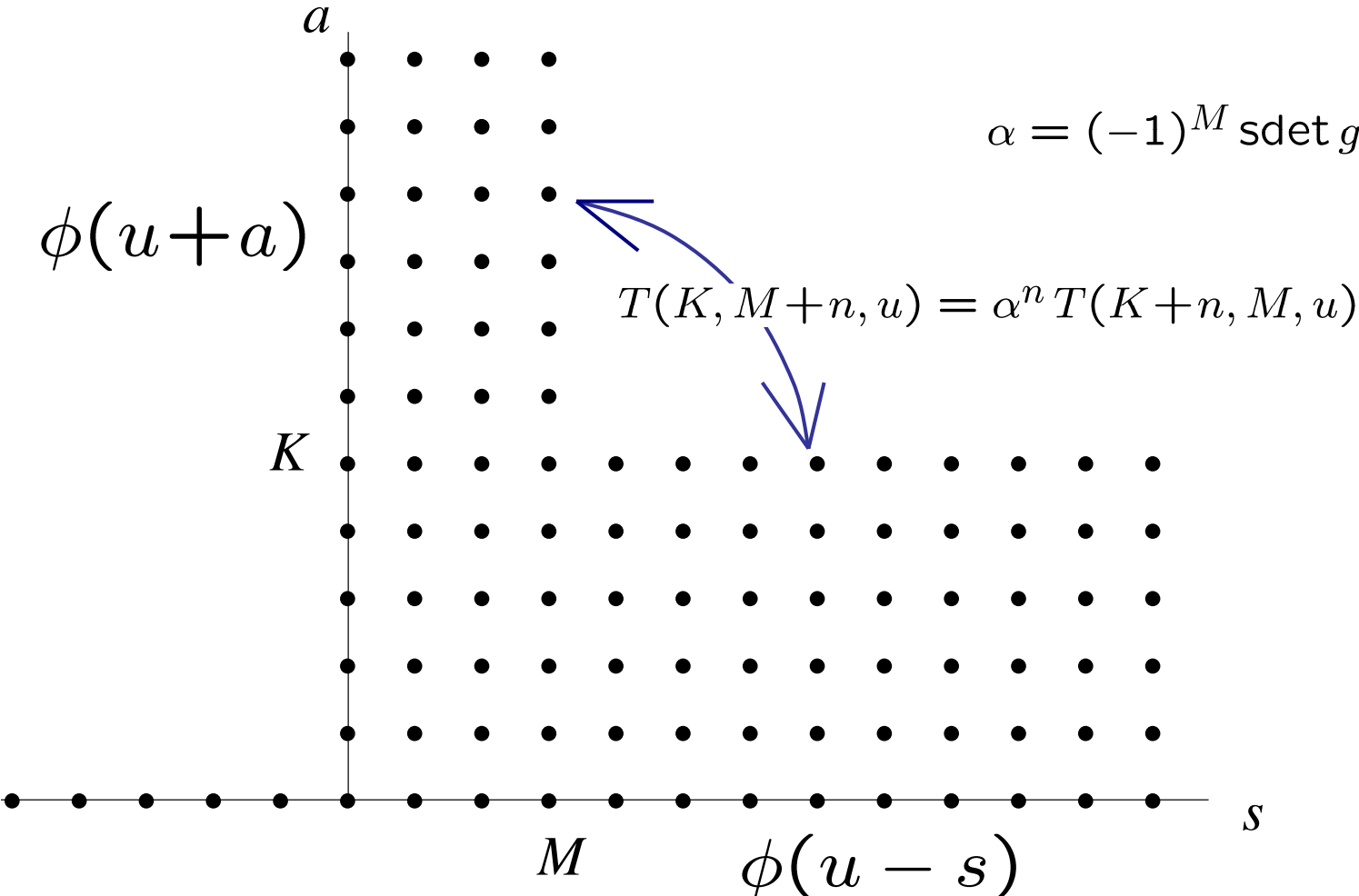
Fat hook

The domain where $T(a,s,u)$ for $GL(K|M)$ spin chains does not vanish identically

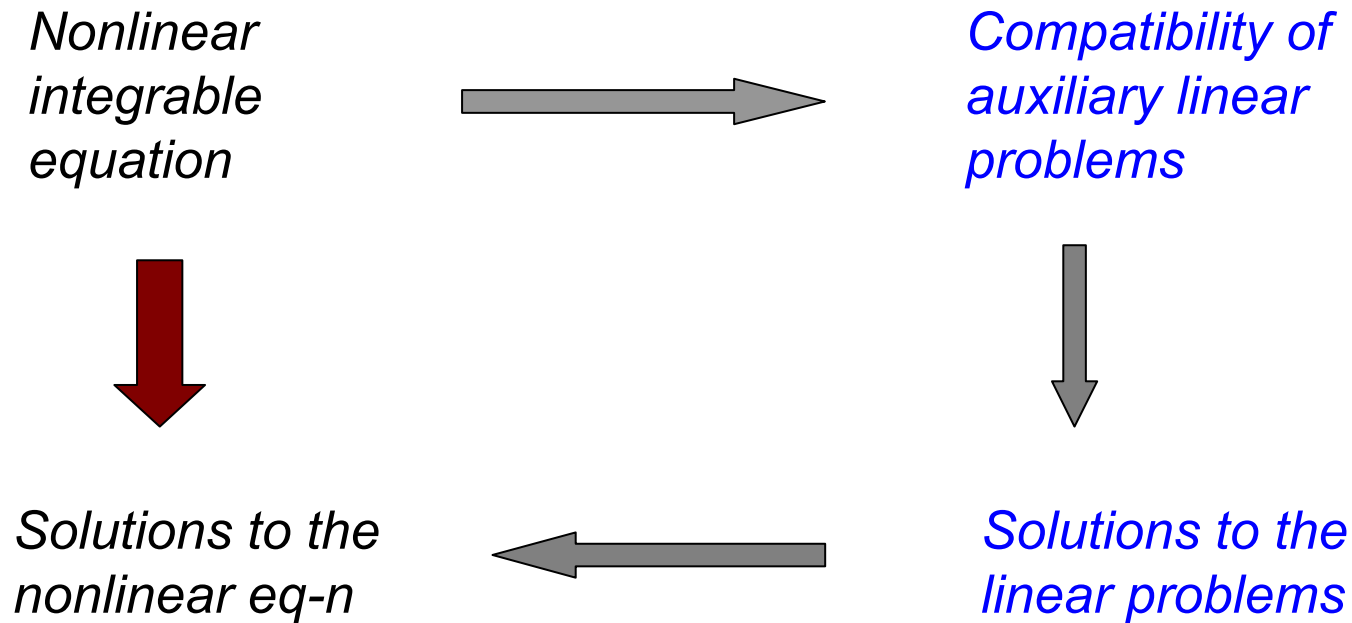


SUSY boundary conditions

Z.Tsuboi 1997
(at $g = id$)



The standard classical scheme

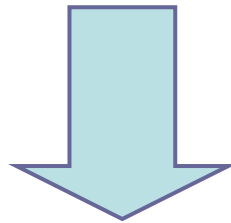


We are going to apply this scheme to the Hirota equation for transfer matrices

Auxiliary linear problems (Lax pair)

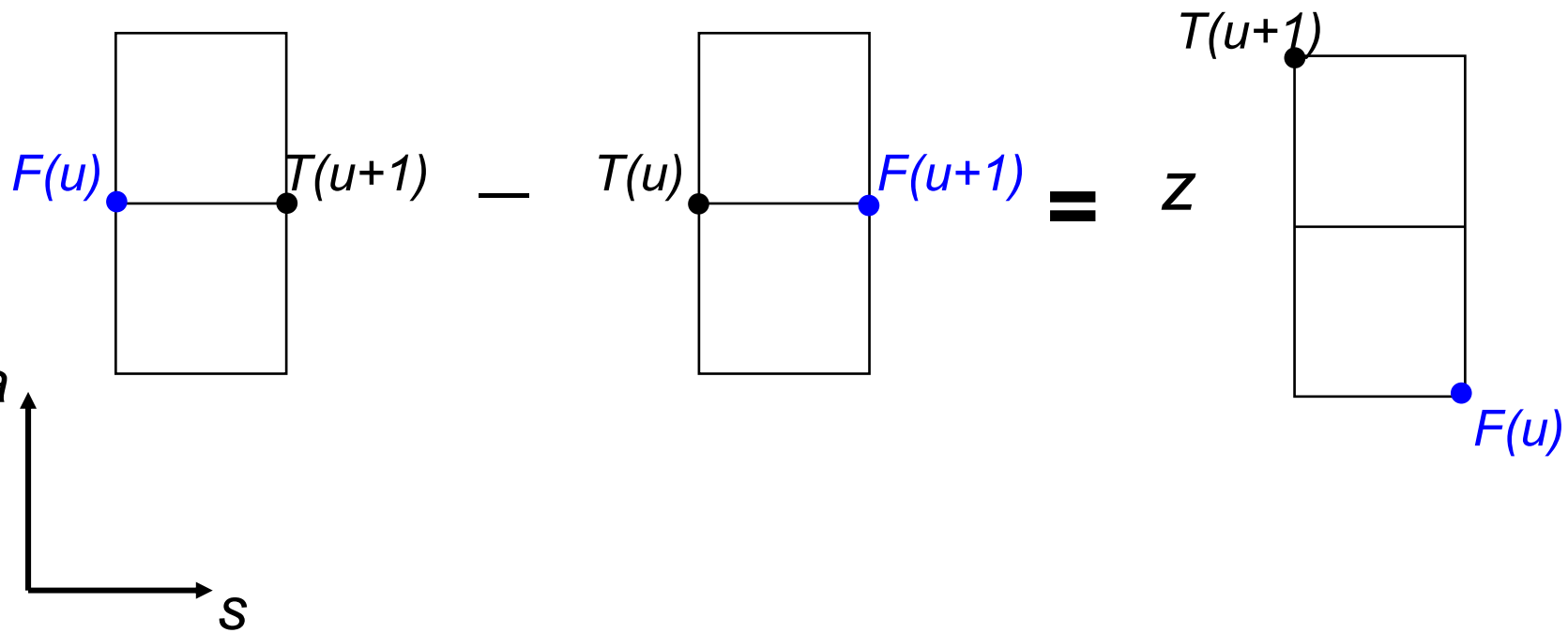
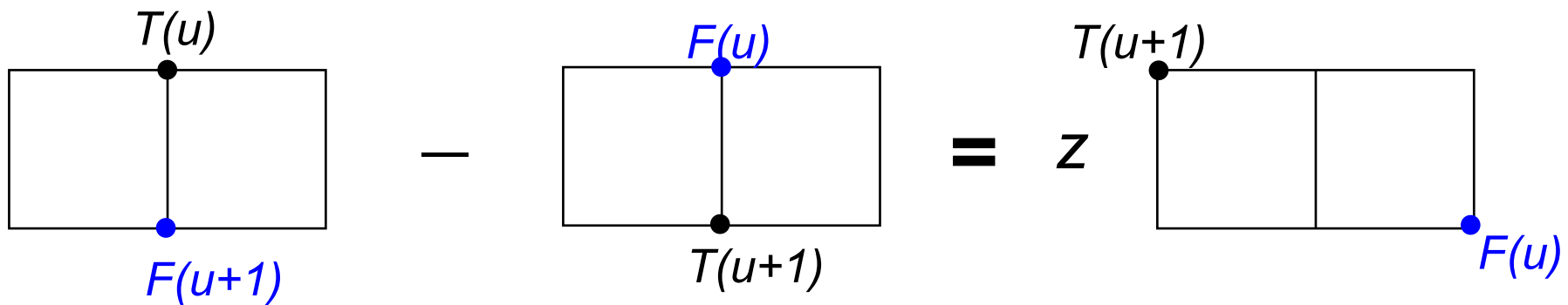
$$T(a+1, s, u)F(a, s, u+1) - T(a, s, u+1)F(a+1, s, u) = zT(a+1, s-1, u+1)F(a, s+1, u)$$

$$T(a, s+1, u+1)F(a, s, u) - T(a, s, u)F(a, s+1, u+1) = zT(a+1, s, u+1)F(a-1, s+1, u)$$



*for any z
(classical spectral parameter)*

The Hirota equation for $T(a, s, u)$



Backlund transformations

Simple but important fact:

F obeys the same Hirota equation

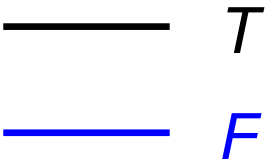
Therefore,

$$T \longrightarrow F$$

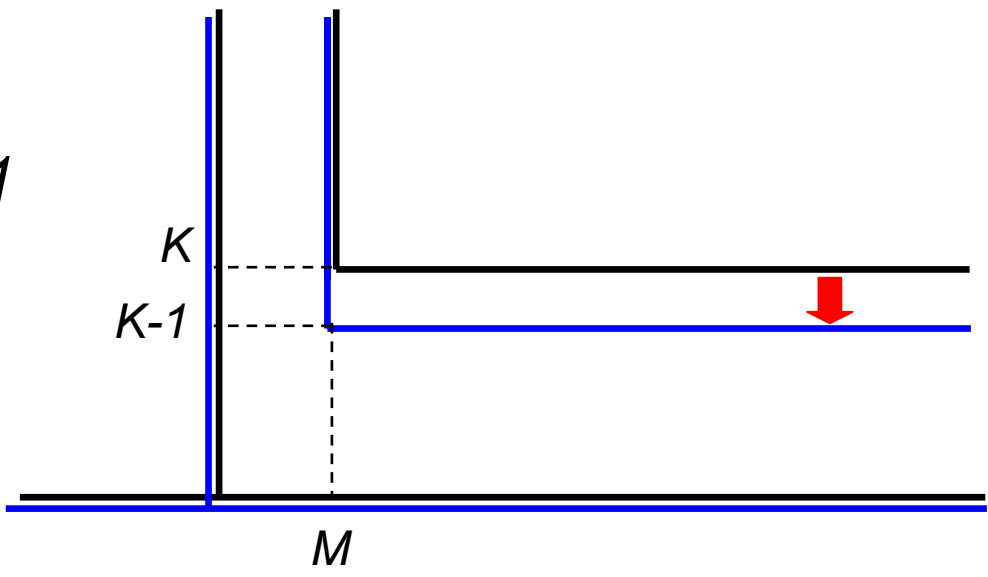
is a Backlund transformation

Boundary conditions for F - ?

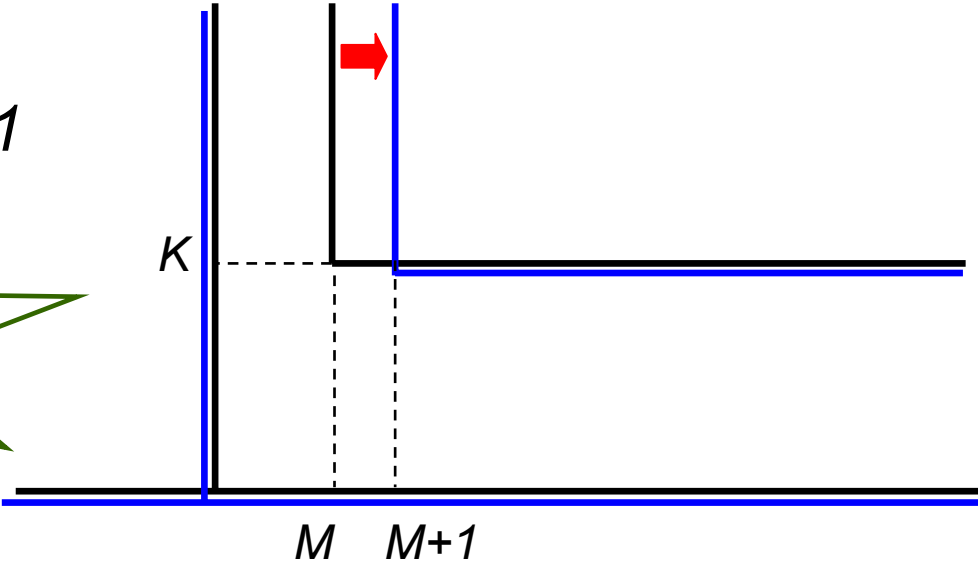
There are two possibilities:



1) $K \Rightarrow K-1$

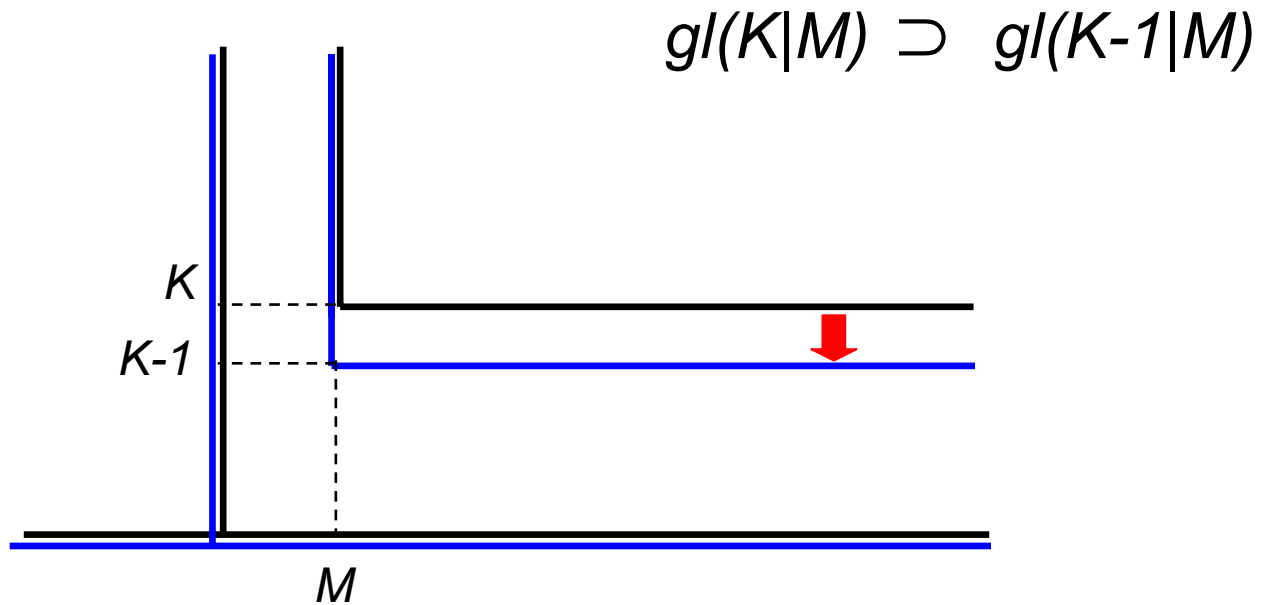


2) $M \Rightarrow M+1$



We need inverse transform!

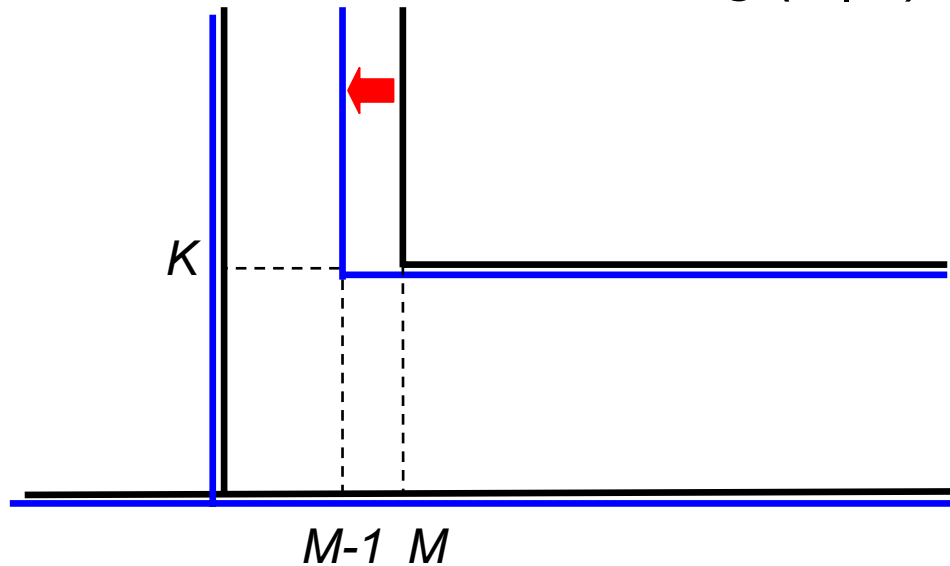
Undressing by BT-I: vertical move



Notation: $T(a,s,u) \equiv T_{K,M}(a,s,u) \rightarrow F(a,s,u) \equiv T_{K-1,M}(a,s,u)$

Undressing by BT-II: horizontal move

$$gl(K|M) \supset gl(K|M-1)$$



Notation: $T(a,s,u) \equiv T_{K,M}(a,s,u) \rightarrow F(a,s,u) \equiv T_{K,M-1}(a,s,u)$

Repeating BT-I and BT-II several times, we introduce the hierarchy of functions

$$T_{k,m}(a, s, u)$$

$$\begin{aligned} k &= 1, \dots, K; \\ m &= 1, \dots, M \end{aligned}$$

such that

- They obey the Hirota equation at any “level” k, m and are polynomials in u for any a, s, k, m

- They are connected by the Backlund transformations

$$\text{BT-I} : T_{k,m}(a, s, u) \rightarrow T_{k-1,m}(a, s, u)$$

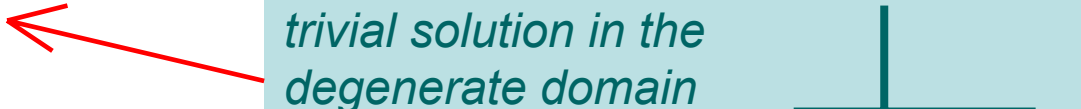
$$\text{BT-II} : T_{k,m}(a, s, u) \rightarrow T_{k,m-1}(a, s, u)$$

- At the highest level $T_{K,M}(a, s, u) = T(a, s, u)$

- At the lowest level $T_{0,0}(a, s, u) = 1$ at $a = 0$ or

at $s = 0, a \geq 0$, and 0 otherwise

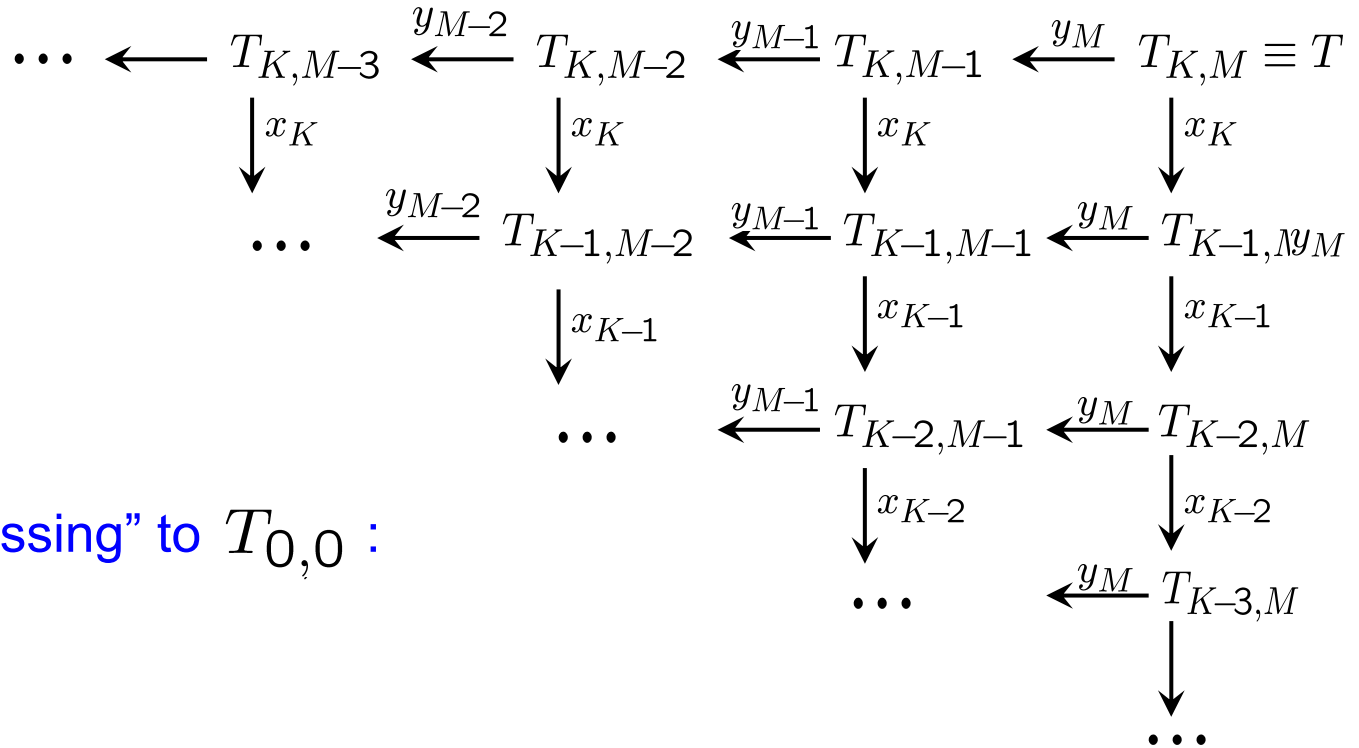
trivial solution in the degenerate domain



The idea is to “undress” the problem to the trivial one using a chain of Backlund transformations

At each step we have a continuous parameter z

Put $\begin{cases} z = x_k & \text{for BT-I } (k, m) \rightarrow (k - 1, m) \\ z = y_m & \text{for BT-II } (k, m) \rightarrow (k, m - 1) \end{cases}$



“Undressing” to $T_{0,0}$:

BT-I:

$$T_{k,m}(a+1, s, u)T_{k-1,m}(a, s, u+1) - T_{k,m}(a, s, u+1)T_{k-1,m}(a+1, s, u)$$

$$= x_k T_{k,m}(a+1, s-1, u+1)T_{k-1,m}(a, s+1, u),$$

$$T_{k,m}(a, s+1, u+1)T_{k-1,m}(a, s, u) - T_{k,m}(a, s, u)T_{k-1,m}(a, s+1, u+1)$$

$$= x_k T_{k,m}(a+1, s, u+1)T_{k-1,m}(a-1, s+1, u)$$

BT-II:

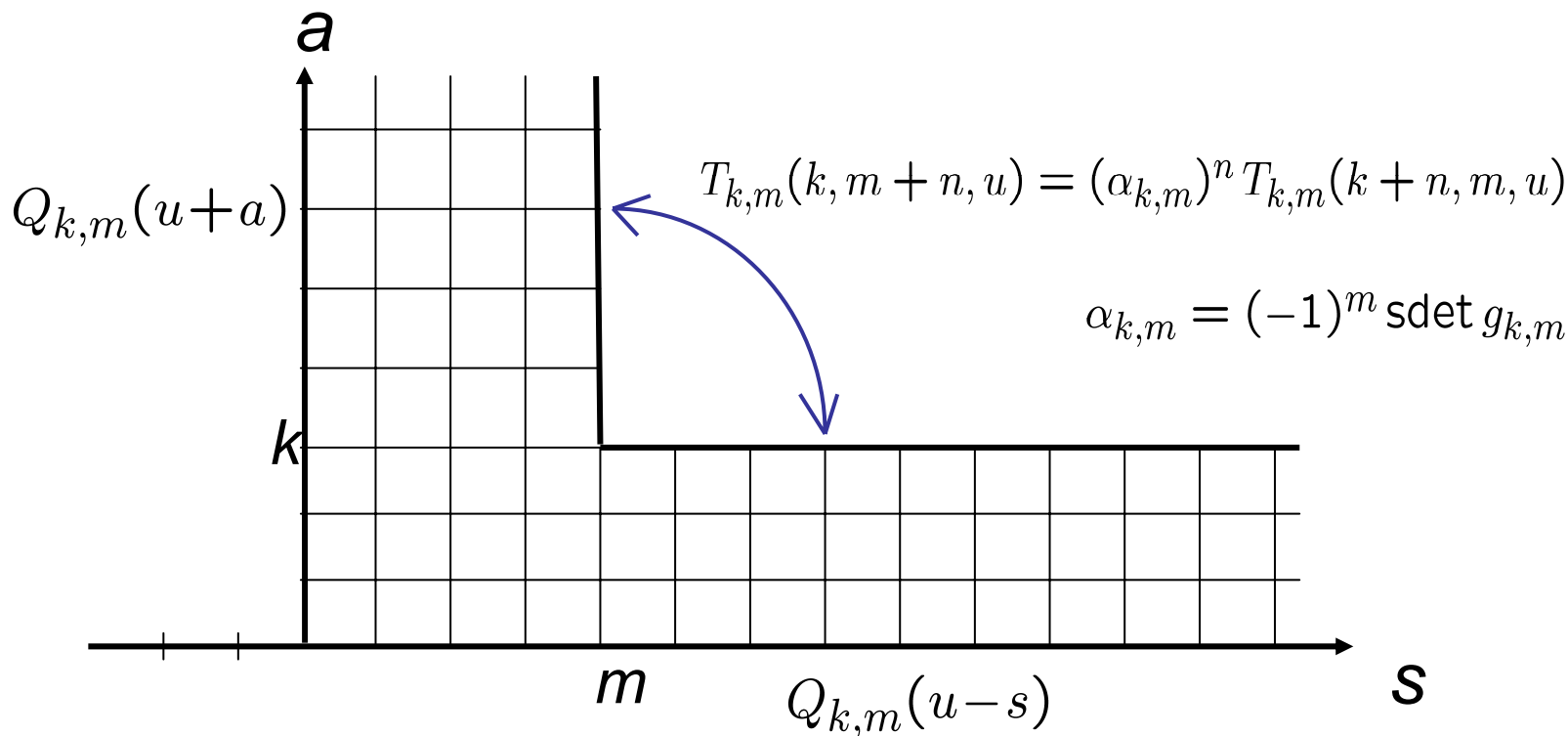
$$T_{k,m-1}(a+1, s, u)T_{k,m}(a, s, u+1) - T_{k,m-1}(a, s, u+1)T_{k,m}(a+1, s, u)$$

$$= y_m T_{k,m-1}(a+1, s-1, u+1)T_{k,m}(a, s+1, u),$$

$$T_{k,m-1}(a, s+1, u+1)T_{k,m}(a, s, u) - T_{k,m-1}(a, s, u)T_{k,m}(a, s+1, u+1)$$

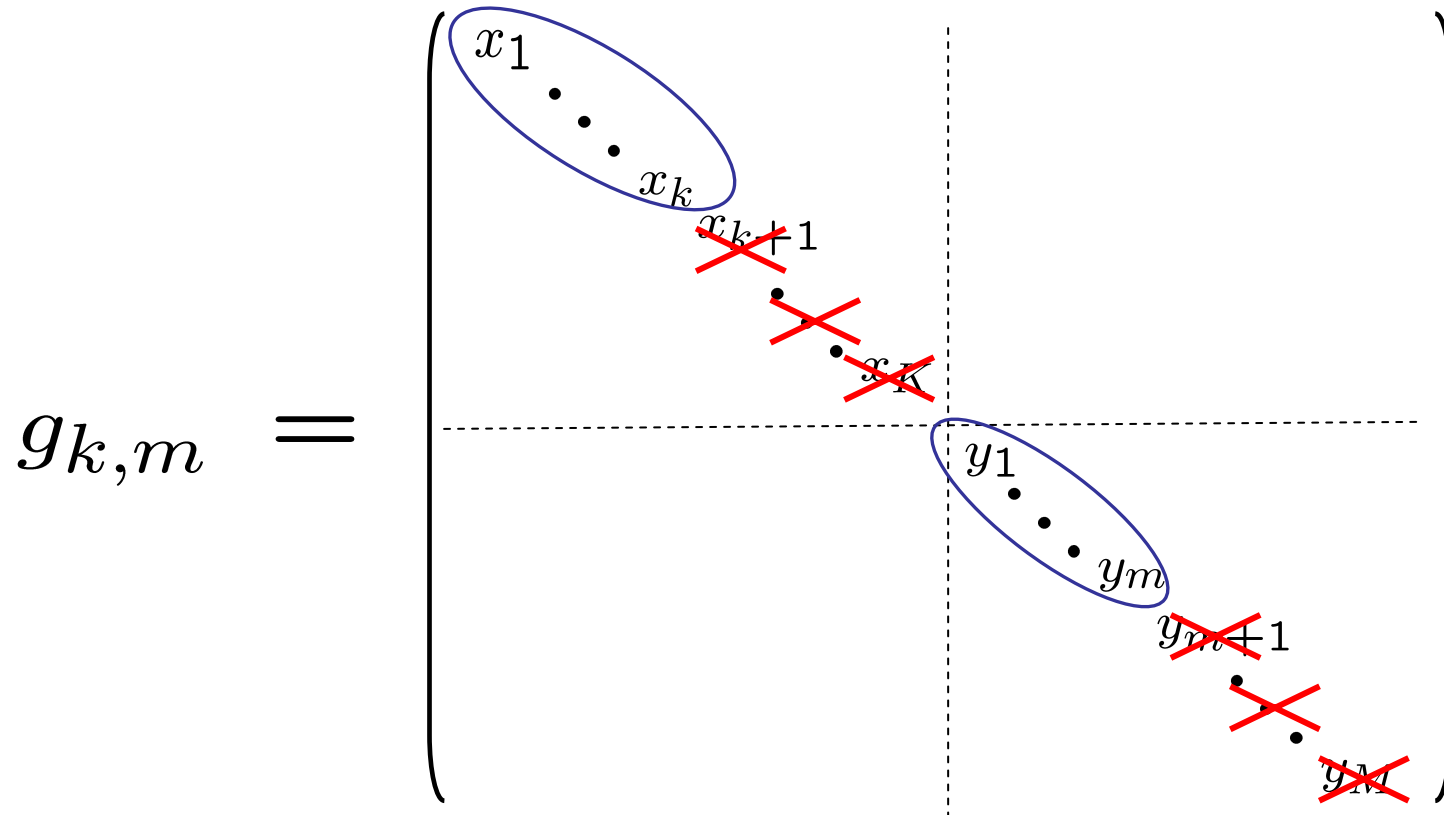
$$= y_m T_{k,m-1}(a+1, s, u+1)T_{k,m}(a-1, s+1, u)$$

Boundary conditions at intermediate levels $k=1, \dots, K, m=1, \dots, M$



$$Q_{K,M}(u) = \phi(u) = \prod_{j=1}^N (u - \xi_j) \quad \text{fixed polynomial}$$

$$Q_{k,m}(u) = \prod_{j=1}^{N_{k,m}} (u - u_j^{(k,m)}) \quad \text{eigenvalues of Baxter's Q-operators}$$




$$\text{sdet } g_{k,m} = \frac{x_1 \cdots x_k}{y_1 \cdots y_m}$$

Operator generating series

Pseudo-difference operator

$$\mathcal{W}(u) = \sum_{s \geq 0} \frac{T(1, s, u + s - 1)}{\phi(u)} e^{2s\partial_u}$$

 *for normalization*

Similar object at any level k, m :

$$\mathcal{W}_{k,m}(u) = \sum_{s \geq 0} \frac{T_{k,m}(1, s, u + s - 1)}{Q_{k,m}(u)} e^{2s\partial_u}$$

It is clear that

$$\mathcal{W}_{0,0}(u) = 1$$

Operator form of the Backlund transformations

Backlund transformations BT-I and BT-II can be represented as recurrence relations for the $\mathcal{W}_{k,m}(u)$

$$\mathcal{W}_{k-1,m}(u) = \left(1 - X_{k,m}(u)e^{2\partial u}\right) \mathcal{W}_{k,m}(u)$$

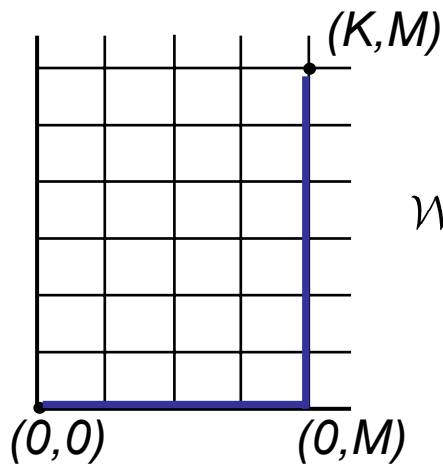
$$\mathcal{W}_{k,m+1}(u) = \left(1 - Y_{k,m+1}(u)e^{2\partial u}\right) \mathcal{W}_{k,m}(u)$$

where

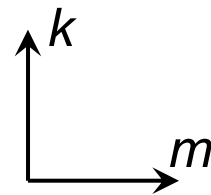
$$X_{k,m}(u) = x_k \frac{Q_{k,m}(u+2) Q_{k-1,m}(u-2)}{Q_{k,m}(u) Q_{k-1,m}(u)}$$

$$Y_{k,m}(u) = y_m \frac{Q_{k,m-1}(u+2) Q_{k,m}(u-2)}{Q_{k,m-1}(u) Q_{k,m}(u)}$$

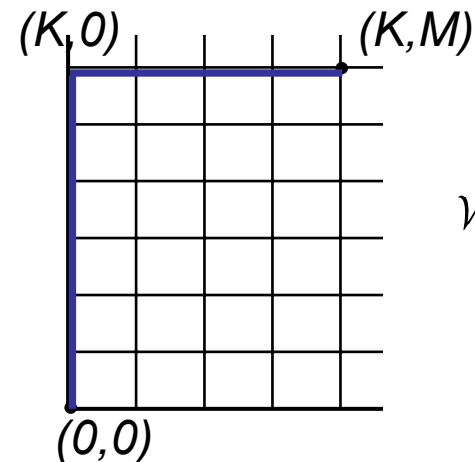
Factorization formulas



$$\mathcal{W}_{K,M}(u) = \overleftarrow{\prod}_{K \geq k \geq 1} (1 - X_{k,M}(u)e^{2\partial_u})^{-1} \cdot \overleftarrow{\prod}_{M \geq m \geq 1} (1 - Y_{0,m}(u)e^{2\partial_u})$$

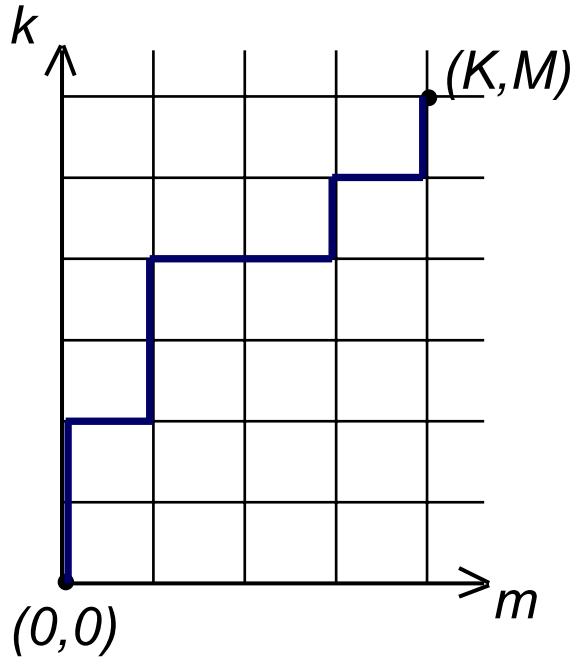


Ordered product: $\overleftarrow{\prod}_{J \geq i \geq I} A_i = A_J A_{J-1} \dots A_{I+1} A_I$



$$\mathcal{W}_{K,M}(u) = \overleftarrow{\prod}_{M \geq m \geq 1} (1 - Y_{K,m}(u)e^{2\partial_u}) \cdot \overleftarrow{\prod}_{K \geq k \geq 1} (1 - X_{k,0}(u)e^{2\partial_u})^{-1}$$

Arbitrary undressing path



Each step $(k, m) \rightarrow (k + 1, m)$ *brings*

$$\left(1 - X_{k+1,m}(u)e^{2\partial_u}\right)^{-1}$$

Each step $(k, m) \rightarrow (k + 1, m)$ *brings*

$$\left(1 - Y_{k,m+1}(u)e^{2\partial_u}\right)$$

$$\mathcal{W}_{K,M}(u) =$$

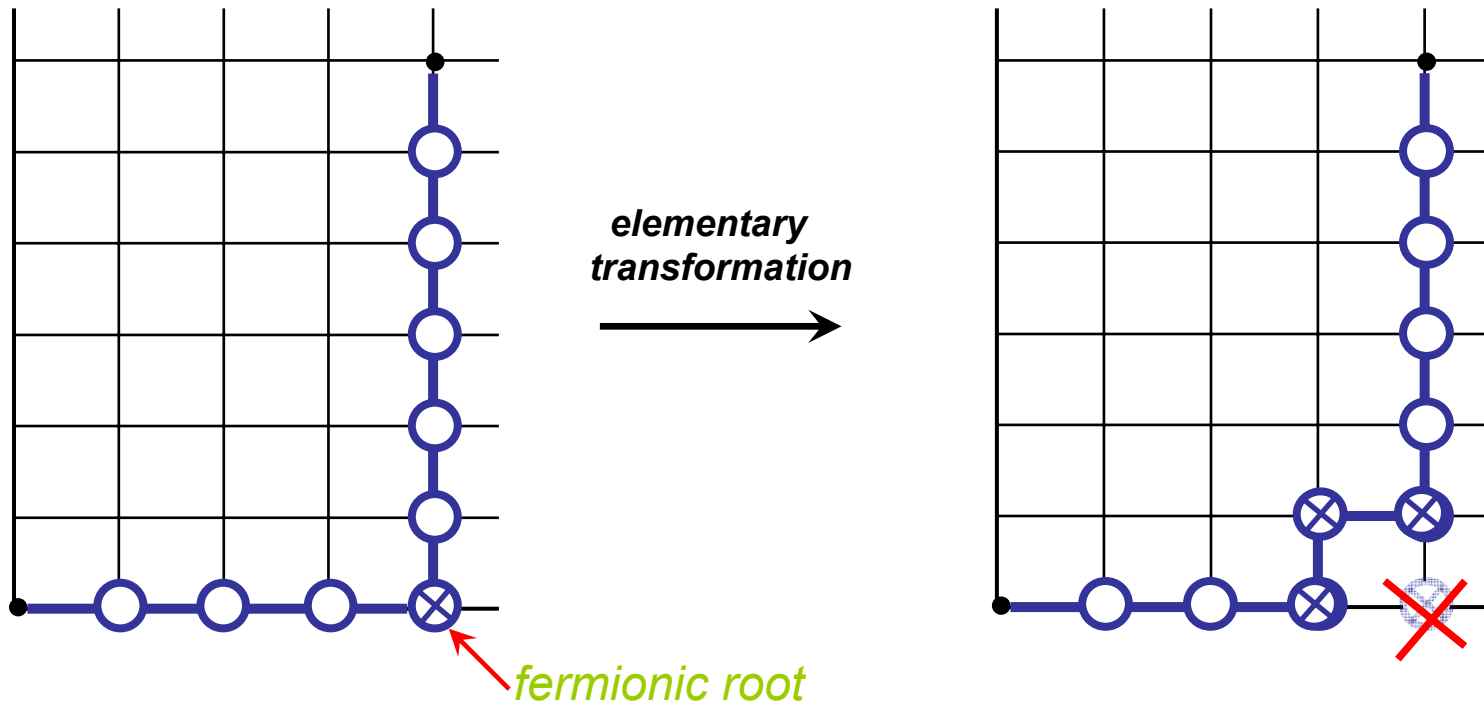
*product of these factors along the path
according to the order of the steps*

Equivalence of these representations follows from a discrete zero curvature condition

“Duality” transformations

Z.Tsuboi 1997;
N.Beisert, V.Kazakov,
K.Sakai, K.Zarembo 2005

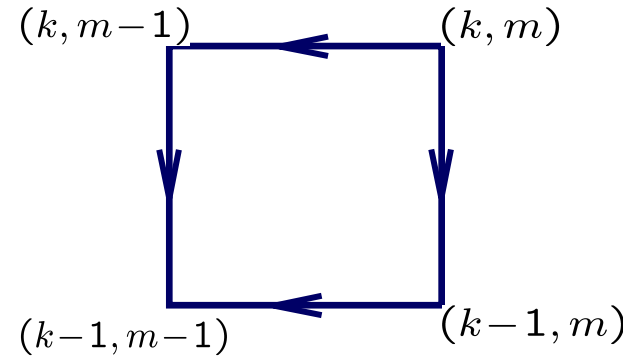
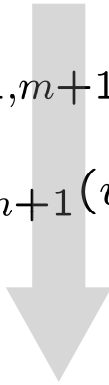
Undressing path \longleftrightarrow SUSY Dynkin diagram



The “duality” transformations can be most clearly understood in terms of **zero curvature condition** on the (k,m) lattice and the **QQ-relation** (V.Kazakov, A.Sorin, A.Zabrodin 2007)

“Zero curvature”

$$\begin{aligned} & (1 - Y_{k-1,m+1}(u)e^{2\partial u}) (1 - X_{k,m}(u)e^{2\partial u}) \\ = & (1 - X_{k,m+1}(u)e^{2\partial u}) (1 - Y_{k,m+1}(u)e^{2\partial u}) \end{aligned}$$



The QQ-relation

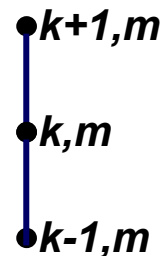
$$\begin{aligned} & x_k Q_{k-1,m-1}(u) Q_{k,m}(u+2) - y_m Q_{k,m}(u) Q_{k-1,m-1}(u+2) \\ = & (x_k - y_m) Q_{k-1,m}(u) Q_{k,m-1}(u+2) \end{aligned}$$

Again the Hirota equation!

We need polynomial solutions $Q_{k,m}(u) = \prod_{j=1}^{N_{k,m}} (u - u_j^{(k,m)})$
 with the “boundary conditions” $Q_{0,0}(u) = 1, Q_{K,M}(u) = \phi(u)$

Building blocks for Bethe equations

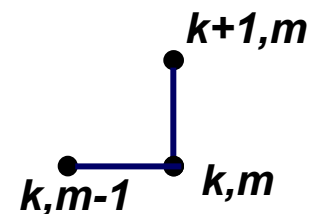
$$\frac{Q_{k-1,m} \left(u_j^{(k,m)} \right) Q_{k,m} \left(u_j^{(k,m)} - 2 \right) Q_{k+1,m} \left(u_j^{(k,m)} + 2 \right)}{Q_{k-1,m} \left(u_j^{(k,m)} - 2 \right) Q_{k,m} \left(u_j^{(k,m)} + 2 \right) Q_{k+1,m} \left(u_j^{(k,m)} \right)} = - \frac{x_k}{x_{k+1}}$$



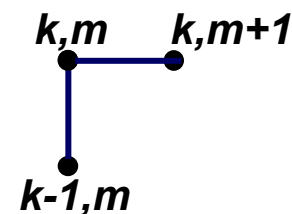
$$\frac{Q_{k,m+1} \left(u_j^{(k,m)} \right) Q_{k,m} \left(u_j^{(k,m)} - 2 \right) Q_{k,m-1} \left(u_j^{(k,m)} + 2 \right)}{Q_{k,m+1} \left(u_j^{(k,m)} - 2 \right) Q_{k,m} \left(u_j^{(k,m)} + 2 \right) Q_{k,m-1} \left(u_j^{(k,m)} \right)} = - \frac{y_{m+1}}{y_m}$$



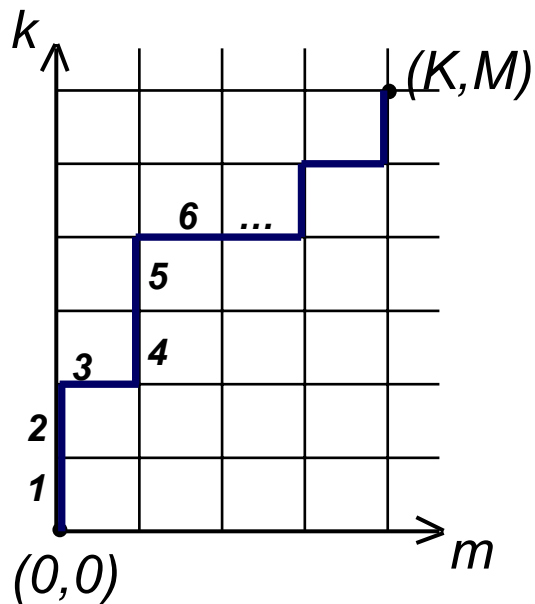
$$\frac{Q_{k+1,m} \left(u_j^{(k,m)} \right) Q_{k,m-1} \left(u_j^{(k,m)} + 2 \right)}{Q_{k+1,m} \left(u_j^{(k,m)} + 2 \right) Q_{k,m-1} \left(u_j^{(k,m)} \right)} = \frac{x_{k+1}}{y_m}$$



$$\frac{Q_{k,m+1} \left(u_j^{(k,m)} \right) Q_{k-1,m} \left(u_j^{(k,m)} - 2 \right)}{Q_{k,m+1} \left(u_j^{(k,m)} - 2 \right) Q_{k-1,m} \left(u_j^{(k,m)} \right)} = \frac{y_{m+1}}{x_k}$$



Bethe equations for arbitrary path (with $x_k = y_m = 1$ for simplicity)



$$p_n = \begin{cases} 1 & \text{for } \uparrow n \\ -1 & \text{for } \underline{n} \end{cases}$$

$$Q_{k,m}(u) = \check{Q}_{k+m}(u - k + m)$$

$$\check{u}_j^{(n)} = u_j^{(n)} - k + m$$

$$\prod_{b=1}^{K+M} \frac{\check{Q}_b(\check{u}_j^{(a)} - K_{ab})}{\check{Q}_b(\check{u}_j^{(a)} + K_{ab})} = (-1)^{\frac{1+p_a p_{a+1}}{2}} \quad a=1, \dots, K+M-1$$

$K_{ab} = (p_a + p_{a+1})\delta_{a,b} - p_{a+1}\delta_{a+1,b} - p_a\delta_{a,b+1}$ is the Cartan matrix

Conclusion

- A lesson: SUSY case looks more natural and transparent
- Our approach provides an alternative to algebraic Bethe ansatz
- Generalizations and problems:
 - mixed (covariant + contravariant) irreps, infinite dimensional (non-compact) irreps
 - models with other types of R -matrices including non-standard ones like Hubbard or $su(2|2)$ R -matrix in AdS/CFT