

GGI, 28-30 September 2006

Advances in Precision Tests  
and  
Experimental Gravitation in Space

Alternative Theories of  
Gravity and Cosmology

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## *Outline*

- 1. Why alternatives?*
- 2. From quantum strings to classical "gravity"*
- 3. Light scalars?*
- 4. Large extra dimensions?*
- 5. Conclusions*

# Why alternatives?

Einstein's **General Relativity** (GR) is a very successful framework for describing gravity in most physical situations: our **Standard Model** of gravitational interactions, tested by now to  $O(10^{-3})$  accuracy

It seems to be applicable, equally well, to **isolated systems**, to **waves** in empty space, and to the **Universe** as a whole.

The universal attractive nature of gravity is responsible for the growth of small initial perturbations into the **large-scale structure** of our Universe.

There is no apparent reason for mistrusting GR in yet unexplored regimes...but

# Experimental puzzles abound

1. What is dark matter?
2. What is dark energy?
3. What's the origin of the initial inhomogeneities?
4. What's the origin of baryon asymmetry?
5. What's the origin of HECR and GRB?
6. ...

## Theoretical puzzles too!

Gravitational attraction is also responsible for gravitational collapse, formation of black holes, and of **singularities**.

**Theorems** by Hawking & Penrose imply that, under mild assumptions, smooth «initial conditions» lead, **inevitably**, to space-time **singularities**, e.g.

1. The singularity behind a **black-hole** horizon
2. The cosmological (**big bang**) singularity

Q<sub>1</sub>: What happens to singularities when we take quantum effects into account?

Answer not known: have you ever heard about **QGD**?

Q<sub>2</sub>: Can we reconcile **General Relativity & Quantum Mechanics**?

At present, the leading candidate for reconciling  
**GR** and **QM** is (Super)**String Theory**  
As such, it should provide answers to those hard  
questions

# From quantum strings to classical "gravity"

- Classical strings (e.g. cosmic strings) do **gravitate** but that's **not** what we are after. By contrast:
- Quantum (fundamental) strings **induce** gravity: how come?
- It's the consequence of some remarkable quantum miracles!

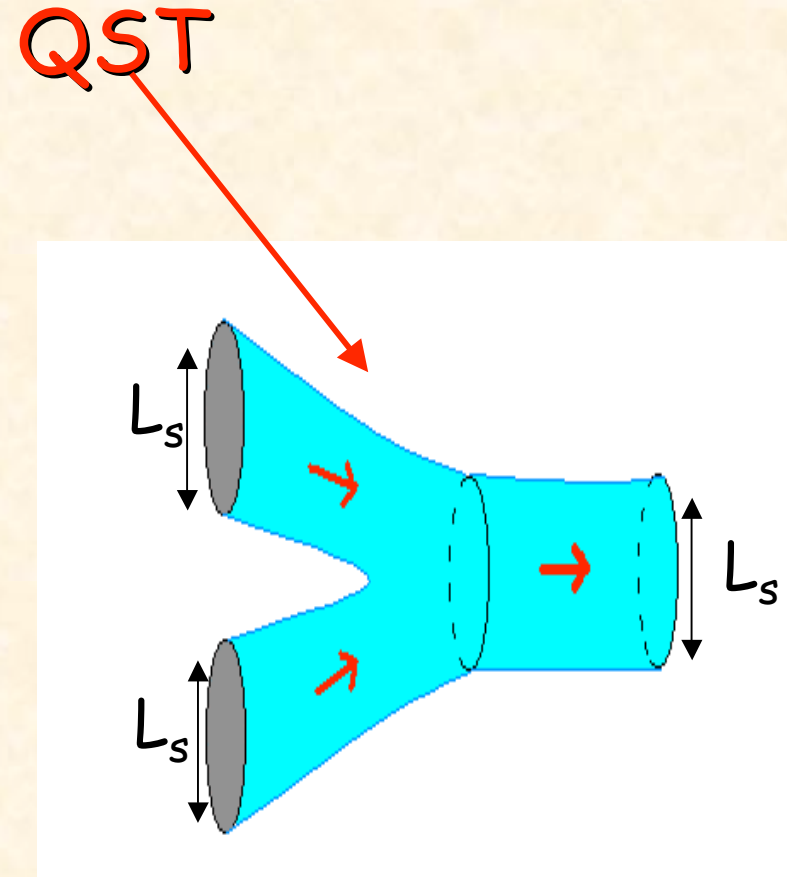
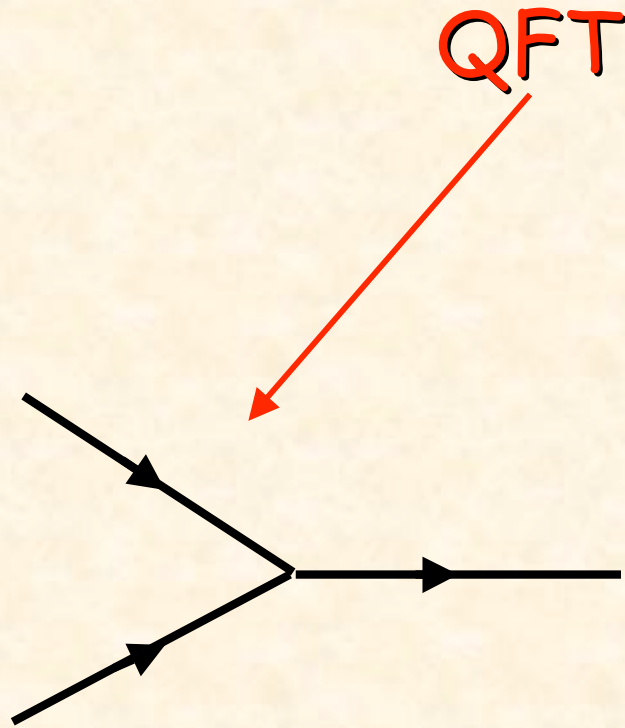
# Quantum miracles: I. Finite Size

- Classical relativistic strings with tension  $T$  may have any size  $L$  and any mass  $M \sim T L c^{-2}$ ;
- Quantum strings have a minimal (optimal) size  $L_s$  (Cf. Bohr radius), given by  $L_s^2 = hc/T$  \*).
- This length appears naturally in the (dimensionless, quantum) action of a string:

$$S_{class.} = T(\text{Area swept}) \Rightarrow \frac{1}{\hbar} S_{class.} = \frac{1}{L_s^2}(\text{Area swept})$$

\* ) Note analogy with  $L_p^2 = hG/c^3$  (if  $G \rightarrow c^4/T$ )





This finite string size,  $L_s$ , is responsible for the smearing of interactions over finite regions of spacetime and for the consequent disappearance of UV divergences

## Quantum miracles: II. Finite Spin

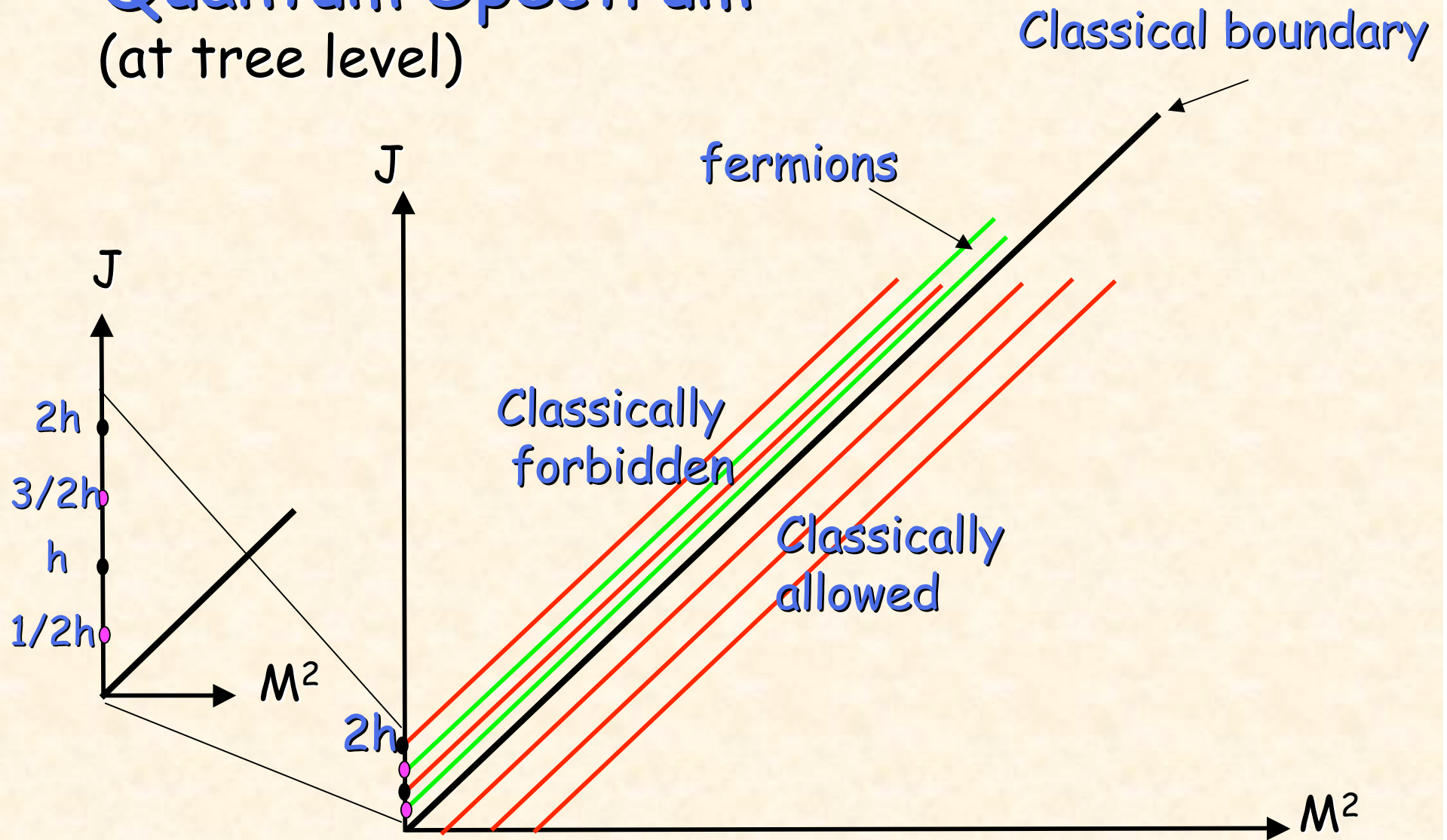
While classical string cannot have angular momentum without also having a finite size/mass, quantum strings may have up to **2 units of J without acquiring mass**:

$$\frac{M^2}{T} \geq J + \hbar \sum_1^{\infty} \frac{n}{2} = J - \alpha_0 \hbar$$

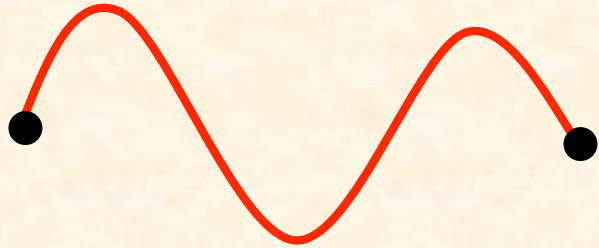
$$\alpha_0 = 0, 1/2, 1, 3/2, 2.$$

Cf. Casimir effect

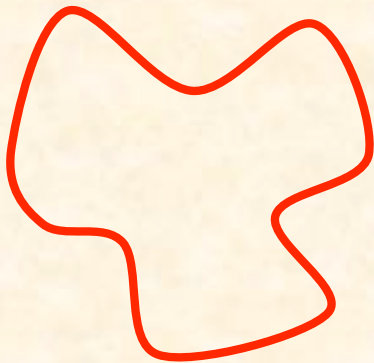
# Quantum Spectrum (at tree level)



In particular..



$\Rightarrow m=0, J = 1 \Rightarrow$  photon and other gauge bosons



$\Rightarrow m=0, J = 2 \Rightarrow$  graviton,

$\Rightarrow m=0, J = 0 \Rightarrow$  dilaton

Integer- $J$  massless states  $\Rightarrow$  **carriers of interactions;**  
1/2-integer- $J$  massless (light)states  $\Rightarrow$  **constituents of matter**

Combining both miracles provides

A **unified** and **finite** theory of elementary particles,  
and of their **gauge & gravitational** interactions

...an old challenge as we know from CSI...

Episode 101 - "Blink"

Detective Mac Taylor discovers the body of a missing woman in Brooklyn Heights. When he discovers a second victim on a garbage barge, his investigation leads him to a serial killer who "imprisons" his victims.

Grappling with memories of his own wife, Taylor follows the killer's trail to a live victim. A woman who cannot move, feel, or speak. A woman who **stretches Veneziano's theory of quantum physics to its outer limit**, challenging Taylor to prove that **"everything is connected."**

**But there are other quantum news..**

- Classical strings can move consistently in any ambient space-time; Quantum strings require **particular** background **space-times** in order to avoid lethal anomalies. E. g.: a Minkowskian space-time must have 1 time and 9 space dimensions. Six of them are presumably compact.
- **No free parameters**: replaced by **scalar fields** whose expectation values provide (dynamically?) the «Constants of Nature». For instance, the fine-structure constant  $\alpha$  and  $G_N T$  are fixed by the dilaton and by the various compactification radii.
- String theory goes one step further than GR by making **everything**, including microphysics, **soft** (T. Damour)

# Light scalars?

- Some  $J=0$  massless strings (at tree level) are there irrespectively of compactification: the **dilaton**  $\phi$  and its SUSY (pseudoscalar) partner, the (KR) **axion**  $\sigma$
- $\langle\phi\rangle$  controls the importance of loops (analogue of gauge coupling in QFT) but  $\phi$  itself is a bona-fide field associated with a spin 0 particle
- $\Rightarrow$  Gauge and gravitational couplings can be, in principle, functions of space and time.
- As such, the dilaton is responsible for an extra attractive force between two bodies A and B whose strength (in units of  $G_N$ ) is given by:

$$F_{A,B} \sim g_{\phi AA} g_{\phi BB} ; g_{\phi AA} = \left. \frac{\partial \log M_A}{\partial \phi} \right|_{\phi = \langle \phi \rangle}$$

This "5th force" violates the EP, universality of free fall

# The compactification moduli

- Sizes and shapes of the extra dimensions are controlled by a bunch of (pseudo)scalar fields called moduli, usually also massless at tree level (or even to all orders in PT)
- They **may** acquire a mass from higher order (or non-perturbative) effects (only "protected" by SUSY)
- If they end up being "heavy" they are not so interesting
- If they end up being very light (or massless) we may distinguish two cases:
  1. They have been light and coupled  $O(G_N)$  in the early universe, acquired a mass later => interesting for early cosmology, not for today's experiments
  2. They may be light and very weakly coupled ( $\ll G_N$ ) even today => interesting for dark energy, violations of Equivalence Principle, variations of "constants"



# 1. Light, gravitationally coupled scalar fields in EU

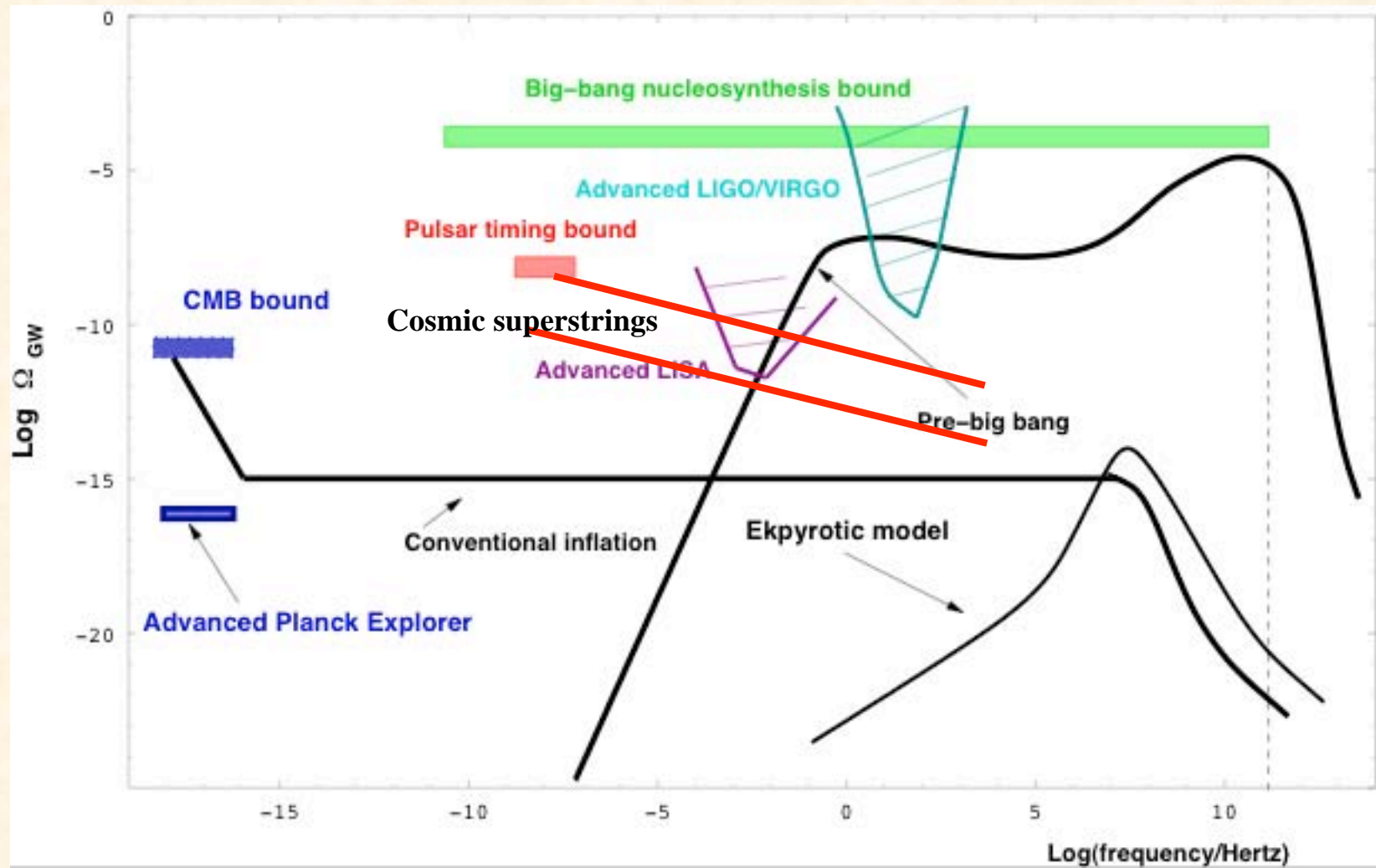
Early cosmology would **not** be described by **GR**, but by the appropriate (multidimensional) effective lagrangian of string theory, even in its classical regimes.

Basis of unconventional cosmologies such as the pre-big bang or expyrotic/cyclic scenarios

Typically, a classical (but not **GR**) pre-bang phase gets connected to a standard (**GR**) post-bang cosmology through a "**quantum bridge**", a high-curvature phase in which an effective field theory (let alone **GR**) description makes no sense

These cosmologies do not (seem to) give the right spectrum of adiabatic density perturbations that slow-roll inflation provides: blue, rather than nearly scale-invariant.

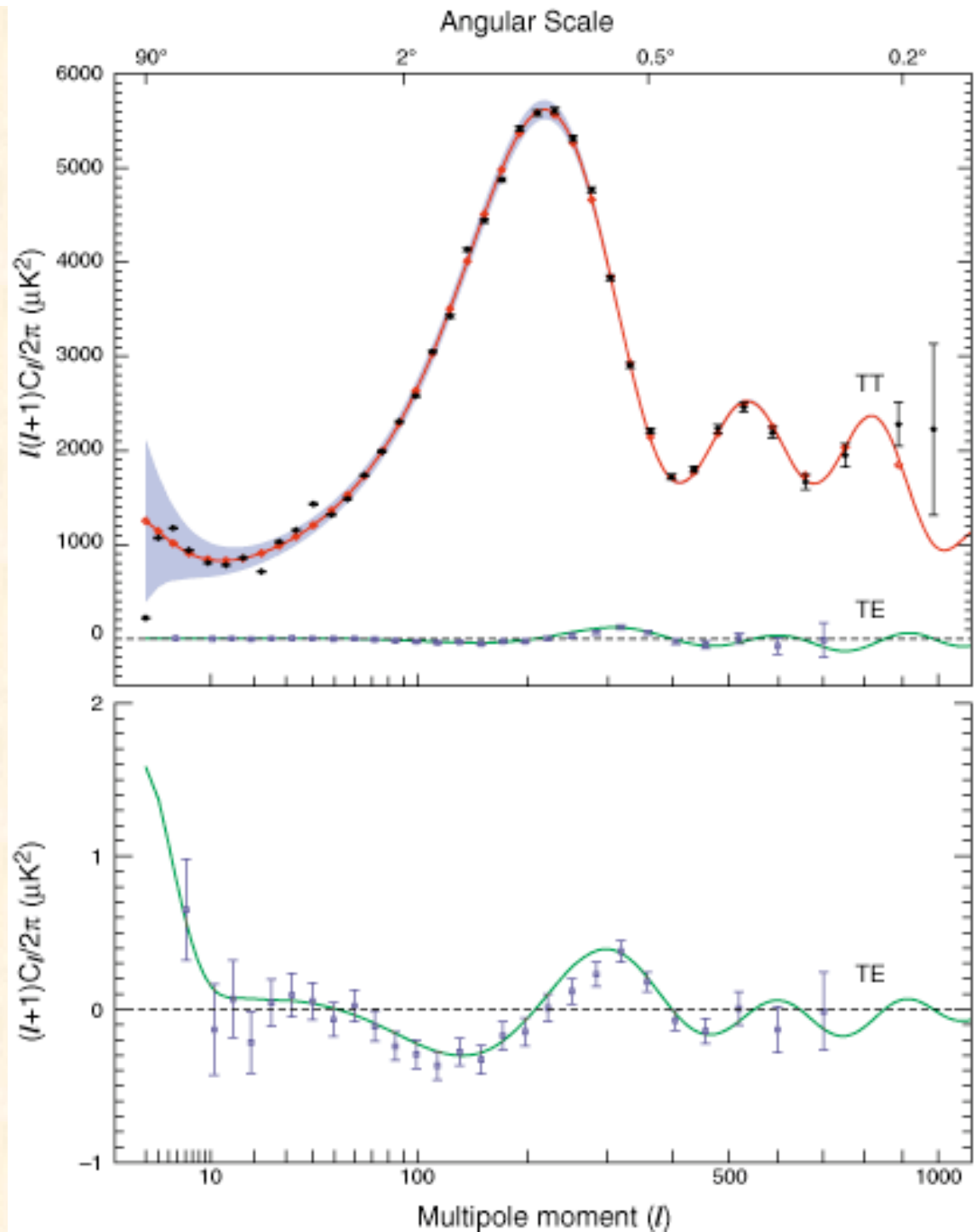
Example of tensor perturbations (**GW**)



# Can the axion save the day?

- The (KR) axion  $\sigma$  **can** have a scale-invariant spectrum. Its tilt,  $(n_\sigma - 1)$ , depends on evolution of the internal dimensions in pre-bang phase: SI spectrum corresponds to a "symmetric" evolution of **all nine** spatial-dimensions
- Unfortunately, axion perturbations do not talk, to first order, to metric perturbations (entropic, **isocurvature** fluctuations) => **bad** predictions for **acoustic peaks**
- The way to rescue these cosmologies is to have the axion play the role of the "**curvaton**" by first becoming a relevant fraction of the total energy density, and by then decaying (before Nucleosynthesis)
- Gives agreement with present CMB data and specific expectations on T-perturbations, non-gaussianity.

TT and TE  
(E-mode of  
polarization)  
correlations from  
WMAP  
B-mode needed to  
test different  
theories



2. Some scalar fields are light even today  
(and coupled  $\ll G_N$ )

Interesting for:

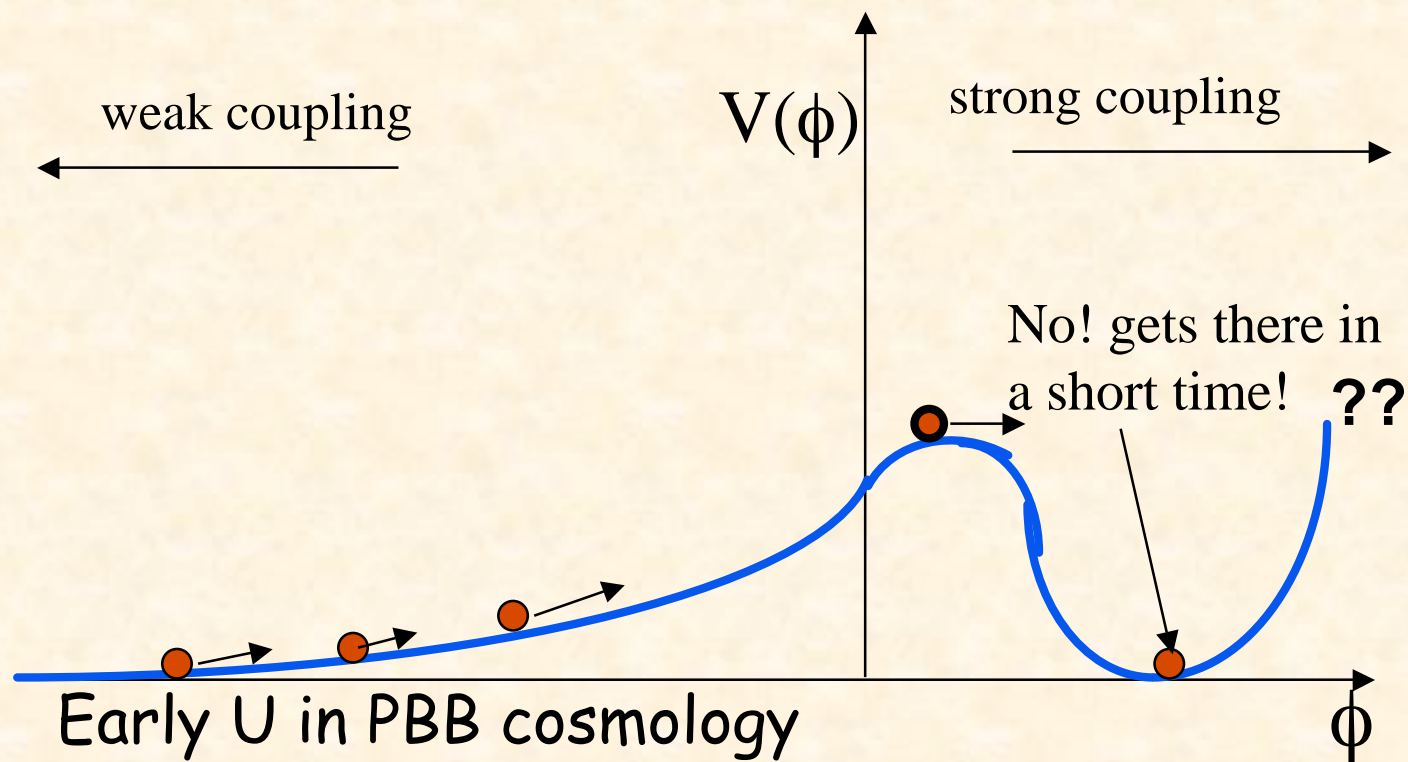
Dark energy,  
Violations of EP,  
Variations of "constants"

# Dilaton as dark energy?

- We have to settle first the question of its coupling to matter and of possible EP-violations
  - By supersymmetry,  $\phi$  is **massless to all orders** in perturbation theory
  - Its perturbative coupling to matter is **larger** than gravity and **non-universal**

This problem goes under the name of the  
“**Dilaton (moduli) Stabilization Problem**”  
in String Theory

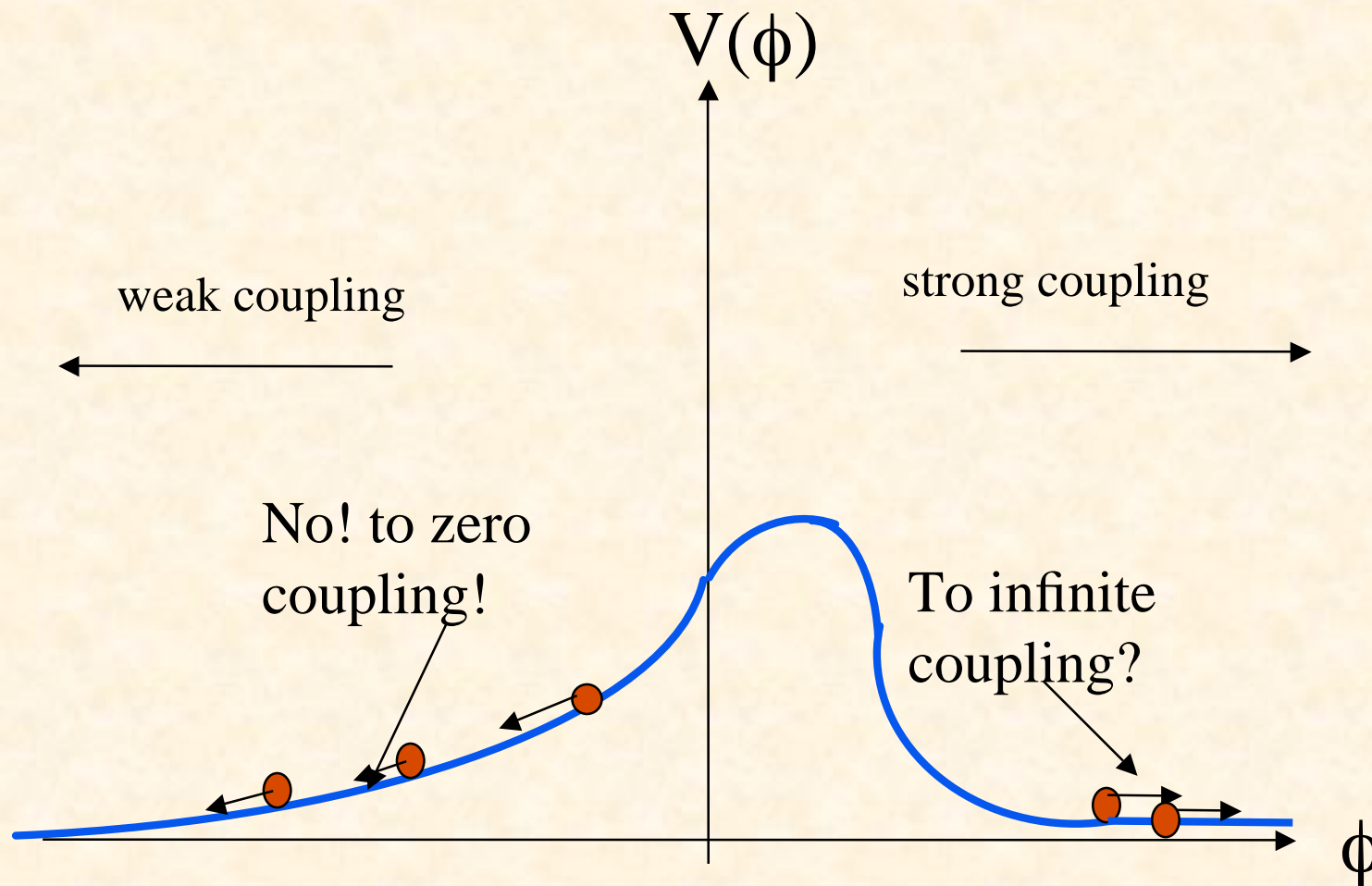
- Standard way out assumes that  $\phi$  acquires a NP mass and is **frozen today** at the bottom of its  $V$ .
- Solves EPV problem but cannot provide acceleration







**Another possibility: vacuum at infinity**  
Can provide dark energy but what about EPV problem?



# Induced gravity in String Theory?

- Is an **infinite bare** coupling ruled out in ST?
- So-called **compositeness** limit: kinetic terms of gauge and gravity fields generated by loops (Cf. **induced gravity** idea of Sakharov...)
- In some toy models one can argue that such **limit exists** and is even **interesting** if there are many matter fields ( $M_s/M_p$  lowered toward  $M_{GUT}$ ?)
- **Assuming** this to be the case, we can have **dilaton** induced acceleration & dilatonic couplings turning off as the dilaton slowly rolls to infinity

# Two cases

- ① If **infinite bare** coupling limit is same for ordinary matter and for non-baryonic dark matter, we fall into a rather conventional quintessence(Q)-model
- ② If the **infinite bare** coupling limit is not as smooth for non-baryonic dark matter then, at the cost of large EPV in the DM sector, we get a **non conventional** (so-called coupled-Q) **model**:  
Attractor towards  $\Omega_{DE} \sim \Omega_{DM}$  at **equality** followed by  
Acceleration with  $\Omega_{DE} / \Omega_{DM} \sim \text{constant}$

Solving coincidence problem?

Not quite: scale of  $V$  put in by hand!

## EP violations & varying $\alpha$

- Can we efficiently send the dilaton towards large values and get **small** enough -but perhaps **measurable**- EPV and/or **variations of  $\alpha$** ?
- This question was addressed and answered affirmatively a few years ago:  
(Damour, Piazza & GV, gr-qc/0204094,  
hep-th/0205111, Damour, gr-qc 0210059, 0306023)

- It is necessary to couple the dilaton to an **inflaton**  $\chi$  through a potential  $V(\chi, \phi)$  taken for simplicity of the form

$$V(\chi, \phi) = \lambda(\phi)\chi^n$$

Using standard chaotic-inflation results we can relate  $\phi$  at the end of inflation to the initial value of  $\chi$  and, eventually, to the amount  $\delta$  of primordial density fluctuations generated during inflation.

One gets:

$$e^{\phi_{end}} = O(10)(\delta)^{-\frac{4}{n+2}}$$

( $\delta \sim 5 \times 10^{-5}$  from CMB data)

At this point we can compute EPV,  $d\alpha / dt$

# EP violations

Instead of PT result,  $g_{\phi NN} / g_{gr NN} \sim 40$ , we get

$$\frac{g_{\phi NN}}{g_{gr NN}} \sim 40 e^{-\phi_{end}} \sim 4 (\delta)^{\frac{4}{n+2}}$$

for the **composition-independent** part of the new force, denoted by  $\alpha_{had}$ . As such it is quite safe

More interesting to look at the **composition-dependent** part, i.e. at violations of the universality of free fall (UFF).

- For two different bodies, A and B:

$$\left(\frac{\Delta a}{a}\right)_{AB} \equiv 2 \frac{a_A - a_B}{a_A + a_B} \sim \alpha_{had} (\alpha_A - \alpha_B)$$

- The small quantity  $(\alpha_A - \alpha_B)$  can be argued to be linear in baryon number, neutron excess and Coulomb energy. For pairs such as Be-Cu or Pt-Ti one finds

$$\frac{\Delta a}{a} \sim 5 \cdot 10^{-5} \alpha_{had}^2 \sim 5 \cdot 10^{-4} (\delta)^{\frac{8}{n+2}}$$

- For  $n=2$  (simplest chaotic inflation) this is compatible w/(but close to) present limits ( $10^{-12}$ ) while  $n=4$  looks already in trouble

**MICROSCOPE, STEP** aim at  $10^{-14}$  ,  $10^{-18}$  resp.

## A varying $\alpha$ ?

In general we expect a time-variation of  $\alpha$  given by

$$\frac{d \ln \alpha}{H dt} = e^{-\phi_{end}} \frac{d\phi}{d \ln a} \sim 2.5 \cdot 10^{-2} \alpha_{had} \frac{d\phi}{d \ln a}$$

The last factor is the main uncertainty. It depends on the coupling of the dilaton to DM and on its possible role as quintessence

If  $\phi$  has small coupling to DM, no role as Q, that factor kills any chance of appreciable variations of  $\alpha$ .



In scenarios of the Q type  $d\phi/d\ln a$  can easily be  $O(1)$  and we get a relation between  $d\ln\alpha/Hdt$  and UFFV

$$\frac{d \ln\alpha}{Hdt} \sim 3.5 \cdot 10^{-6} (1 + q_0 - 1.5\Omega_m)^{1/2} \left( \frac{10^{12} \Delta a}{a} \right)^{1/2}$$

i.e.  $d \ln\alpha/dt \sim 10^{-16} \text{ yr}^{-1}$  for  $\Delta a/a \sim 10^{-12}$

..below present sensitivity ( $10^{-14} \text{ yr}^{-1}$ ) but close to planned sensitivity of cold-atom clocks.

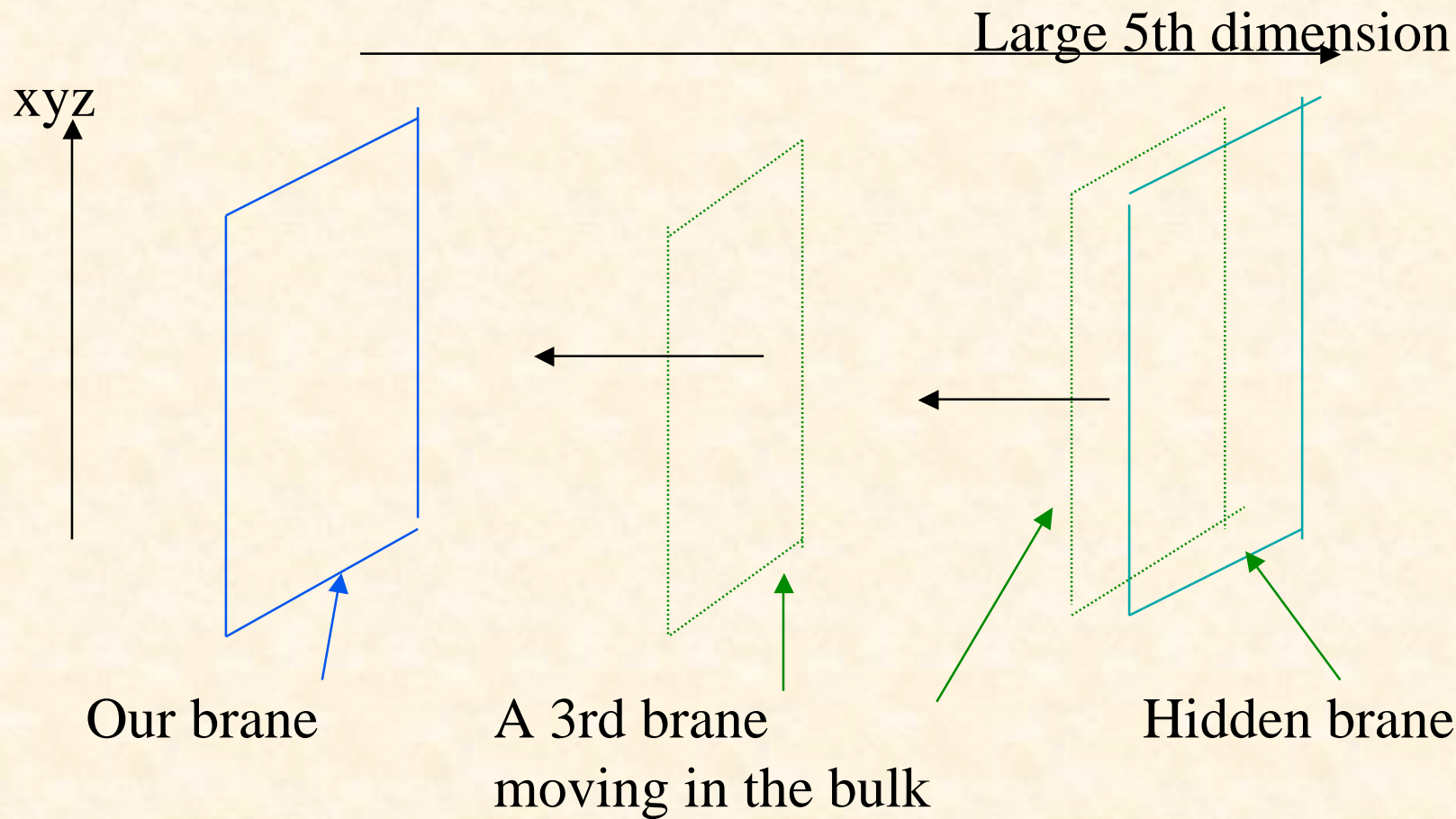
Upper limit from Oklo data ( $5 \times 10^{-17} \text{ yr}^{-1}$ ) would correspond to  $\Delta a/a \sim 10^{-13}$

Instead, difficult to get  $\Delta\alpha/\alpha \sim 10^{-6}$  at  $z \sim 0.5-3.5$  as claimed by Webb et al.

## Large extra dimensions?

- The VEVs of the moduli determine **sizes and shapes** of the internal dimensions
- If these are of string scale and frozen, we can only “see” them indirectly (e.g. gauge fields generated a la KK)
- If they are large and/or dynamical, we can distinguish again two cases:
  1. They were large/dynamical in the EU but not now
  2. They are large/dynamical even today

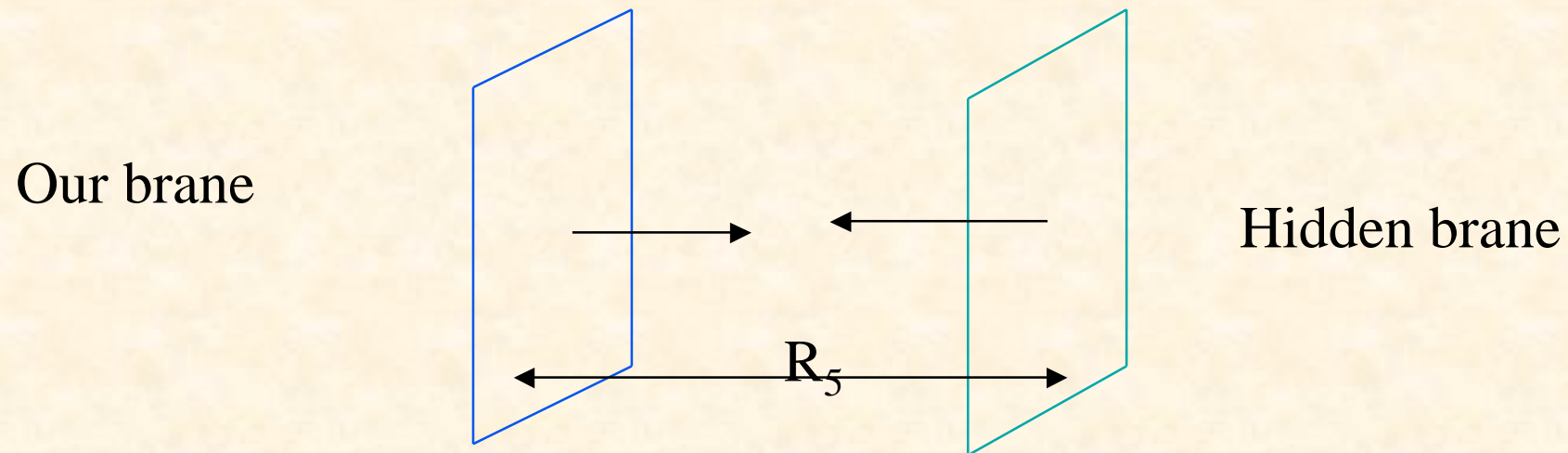
The first case is the one we have already discussed of some unconventional cosmologies of the pre-bang type. We may still distinguish the case of a multi-dimensional early universe (PBB) and that of an early brane-universe (EKP).



BB as result of impact of 3rd brane on ours

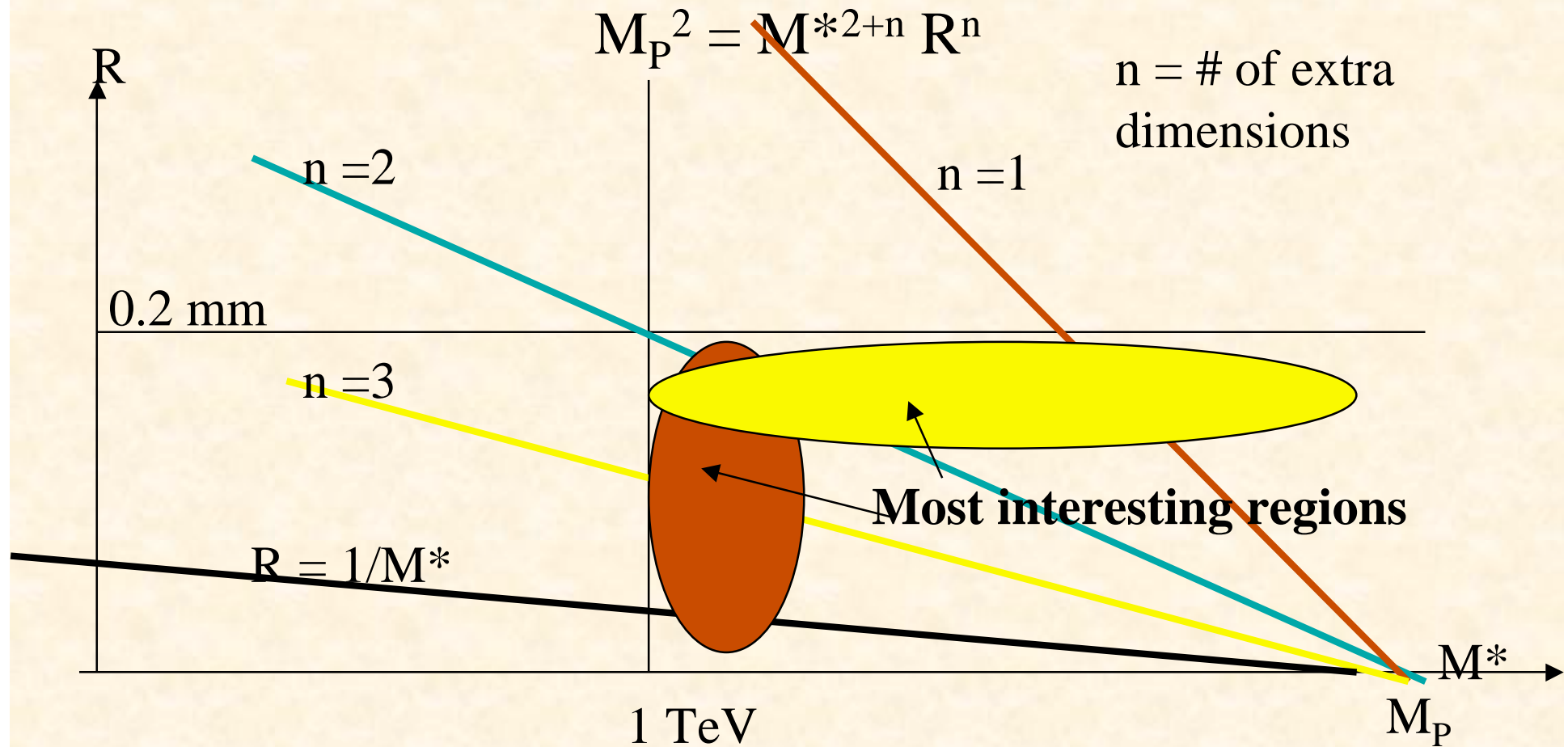
Alternatively:

Before the BB



BB is the collapse of the 5th dimension to zero size.

In the second case the extra dimensions can be small (but not infinitesimal) and may be seen through modifications of the  $1/r^2$  law at short distance, or via new strong-gravity phenomena that should occur at high energy accelerators (mini black holes etc.)



Or they are really macroscopic, even infinite **today**. In this case we have to accept that we (the SM) are (is) confined to a brane...

# Conclusions

Interesting **new phenomena** are usually expected in alternative theories of gravity and cosmology. Examples of BSGR physics:

1. Unconventional cosmic **perturbations** & **dark energy**
2. Violations of **EP**, UFF;
3. **Modifications** of Newton's law;
4. **Strong gravity** at accelerator energies;
5. Spacetime **variations** of "constants"

As with BSM physics, the obvious question making many theorists so nervous these days is:

Why have we not found **any evidence so far** for physics beyond the SM & GR?

Hopefully the answer will not be:  
because there isn't!